

# Fundamentals of Atomic Force Microscopy

## Part 2: Dynamic AFM Methods

Week 2, Lecture 4

Connections between amplitude reduction, phase,  
frequency shift, and excitation

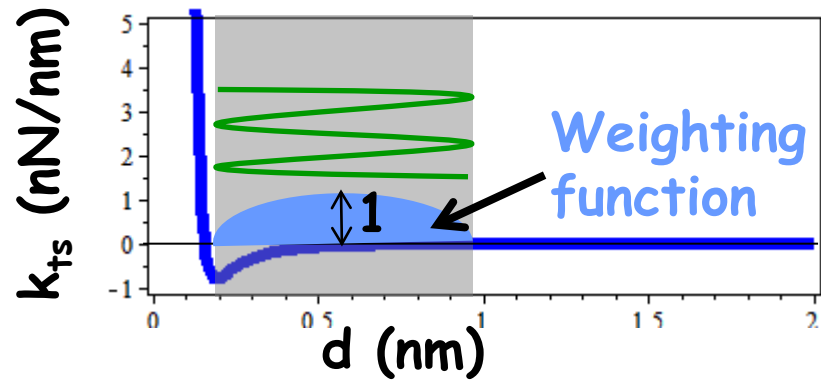
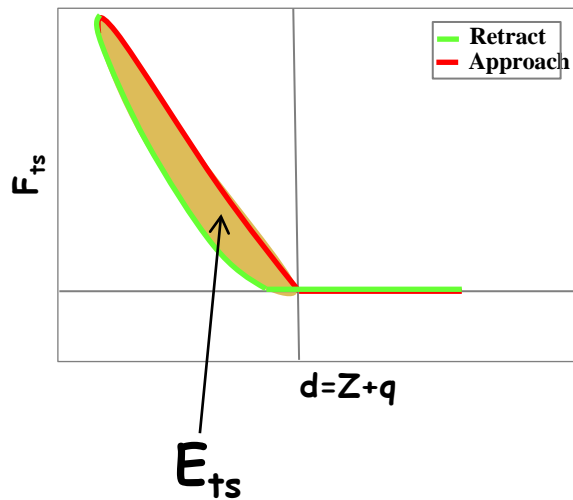
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# From the last lecture

$$V_{ts}(Z, A) = \langle F_{ts}^{CONS} \cdot q \rangle = \frac{kA^2}{2} \left( 1 - \frac{\omega^2}{\omega_0^2} \right) - \frac{AF_0}{2} \cos \phi \quad (\text{"Virial"}) \quad (1)$$

$$E_{ts}(Z, A) = - \int_0^{2\pi/\omega_0} F_{ts}^{DISS} \cdot \dot{q} dt = \pi AF_0 \sin \phi - \pi kA^2 \frac{\omega}{Q\omega_0} \quad (\text{"Dissipation"}) \quad (2)$$



# Connections between amplitude reduction, frequency shift and dissipation

Eliminating  $\phi$  from (1) and (2) (see Appendix) we arrive at

$$\frac{kA}{F_0} = \frac{1}{\sqrt{\left(\frac{\omega_0'^2 - \omega^2}{\omega_0^2}\right)^2 + \left(\frac{\omega}{\omega_0 Q'}\right)^2}} \quad (3)$$

$$\tan(\phi) = \frac{\left(\frac{\omega}{\omega_0 Q'}\right)}{\left(\frac{\omega_0'^2 - \omega^2}{\omega_0^2}\right)} \quad (4)$$

where

$$\omega_0' = \omega_0 - \omega_0 \frac{V_{ts}(Z, A)}{kA^2} \quad (5)$$

$$\frac{1}{Q'} = \frac{1}{Q} + \frac{\omega_0 E_{ts}}{\omega \pi k A^2} \quad (6)$$

In the limit of infinitesimal amplitude, recall  $V_{ts}(Z, A) \sim \frac{1}{2} k_{ts}(Z) A^2$ ,

leading to a classical result of the linearized analysis:

$$\omega_0' = \omega_0 \left(1 - \frac{k_{ts}}{2k}\right)$$

# AM-AFM or Tapping mode

Assume  $\omega = \omega_0$ , and  $F_0$  is constant

Recall,  $F_0 = kA^{\text{far}} / Q$  when  $\omega = \omega_0$

where  $A^{\text{far}}$ ,  $\phi^{\text{far}} = \pi / 2$  are the amplitude and phase far from the surface  
the equations reduce to

$$V_{\text{ts}}(Z, A) = -\frac{kAA^{\text{far}}}{2Q} \cos \phi \quad (7)$$

$$E_{\text{ts}}(Z, A) = \frac{\pi kAA^{\text{far}}}{Q} \sin \phi - \frac{\pi kA^2}{Q} \quad (8)$$

$$\frac{A}{A^{\text{far}}} = \frac{1/Q}{\sqrt{\left(-2\frac{V_{\text{ts}}(Z, A)}{kA^2}\right)^2 + \left(\frac{1}{Q} + \frac{E_{\text{ts}}(Z, A)}{\pi kA^2}\right)^2}} \quad (9)$$

$$\tan(\phi) = \frac{\left(\frac{1}{Q} + \frac{E_{\text{ts}}(Z, A)}{\pi kA^2}\right)}{\left(-2\frac{V_{\text{ts}}(Z, A)}{kA^2}\right)} \quad (10)$$

$\frac{A}{A^{\text{far}}}$  called amplitude setpoint ratio

## AM-AFM or tapping mode

- While scanning  $Z$  adjusts so that  $E_{ts}$  and  $V_{ts}$  reach values where the desired amplitude reduction is achieved
- Topography is the set of those  $Z$  values which achieve a constant amplitude by adjusting  $E_{ts}$  and  $V_{ts}$  pointwise on the sample
- In tapping mode (AM-AFM)  $E_{ts}$  and  $V_{ts}$  are not independent, rather they are related through Eqs. (9, 10)
- Phase lag depends both on conservative and dissipative interactions

# FM-AFM

$\omega$  changes to keep phase fixed

From (4)

$$\omega = \omega'_0 = \omega_0 - \omega_0 \frac{V_{ts}(Z, A)}{kA^2} !!! \quad (11)$$

and  $F_0$  changes to keep  $A$  fixed

From (3)

$$F_0 = kA \frac{\omega}{\omega_0 Q'} = kA \frac{\omega}{\omega_0} \left( \frac{1}{Q} + \frac{\omega_0 E_{ts}}{\omega \pi k A^2} \right) \quad (12)$$

And  $Z$  changes to keep a constant frequency shift

$$\Delta\omega_0 = \omega'_0 - \omega_0 = -\omega_0 \frac{V_{ts}(Z, A)}{kA^2} \quad (13)$$

# FM-AFM

- Topography is set of those values of  $Z$  that keep  $V_{ts}/A^2$  constant
- Drive frequency  $\omega$  tracks effective natural frequency  $\omega'_0$
- $F_0$  picks up the dissipative interactions

# Appendix

Starting with

$$V_{ts}(Z, A) = \frac{kA^2}{2} \left( 1 - \frac{\omega^2}{\omega_0^2} \right) - \frac{AF_0}{2} \cos \phi \quad (\text{"Virial"}) \quad (14)$$

$$E_{ts}(Z, A) = \pi AF_0 \sin \phi - \pi kA^2 \frac{\omega}{Q\omega_0} \quad (\text{"Dissipation"}) \quad (15)$$

we can rewrite (14) and (15) as

$$\sin \phi = \frac{kA}{F_0} \frac{\omega}{\omega_0} \left( \frac{\omega_0}{\omega} \frac{E_{ts}(Z, A)}{\pi kA^2} + \frac{1}{Q} \right) \quad (16)$$

$$\cos \phi = \frac{kA}{F_0} \left( \frac{\omega_0^2 \left( 1 - \frac{2}{kA^2} V_{ts}(Z, A) \right) - \omega^2}{\omega_0^2} \right) \quad (17)$$

Define

$$\frac{1}{Q'} = \frac{1}{Q} + \frac{\omega_0}{\omega} \frac{E_{ts}(Z, A)}{\pi kA^2} \quad (18) \quad \text{and} \quad (\omega'_0)^2 = \omega_0^2 \left( 1 - \frac{2}{kA^2} V_{ts}(Z, A) \right) \quad (19)$$

If  $\frac{V_{ts}(Z, A)}{kA^2} \ll 1$  then

$$\omega'_0 = \omega_0 \left( 1 - \frac{V_{ts}(Z, A)}{kA^2} \right)$$

Eliminating  $j$  from (16) and (17) and using (18) and (19) we arrive at

$$\frac{kA}{F_0} = \frac{1}{\sqrt{\left( \frac{\omega'^2_0 - \omega^2}{\omega_0^2} \right)^2 + \left( \frac{\omega}{\omega_0 Q'} \right)^2}} \quad (20)$$

$$\tan(\phi) = \frac{\left( \frac{\omega}{\omega_0 Q'} \right)}{\left( \frac{\omega'^2_0 - \omega^2}{\omega_0^2} \right)} \quad (21)$$