

# From Atoms to Materials: Predictive Theory and Simulations

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Week 4: Connecting Atomic Processes to the Macroscopic World  
Lecture 4.3: Statistical Mechanics of the Harmonic Solid

**Ale Strachan**

[strachan@purdue.edu](mailto:strachan@purdue.edu)

School of Materials Engineering &  
Birck Nanotechnology Center  
Purdue University  
West Lafayette, Indiana USA

# Review of important ensembles

Microcanonical (NVE)	Canonical (NVT)	Isobaric/isothermal (NPT)
Probability distributions		
$P(\{P_i\}\{R_i\}) = \frac{1}{\Omega(E, V, N)}$	$P(\{R_i\}, \{P_i\}) = \frac{e^{-\frac{H(\{R_i\}, \{P_i\})}{kT}}}{Z(T, V, N)}$	$P(\{R_i\}, \{P_i\}, V) = \frac{e^{-\frac{1}{kT}[H(\{R_i\}, \{P_i\}) - PV]}}{Z_p(T, P, N)}$
$\Omega(E, V, N) = \sum_{micro} \delta(E - H(\{R_i\}\{P_i\}))$	$Z(T, V, N) = \sum_{micro} e^{-\frac{E}{kT}}$	$Z_p(T, P, N) = \sum_V \sum_{micro} e^{-\frac{E - PV}{kT}}$
Free energies (atomistic ↔ macroscopic thermodynamics)		
$S = k \log \Omega(E, V, N)$	$F(T, V, N) = -kT \log Z$	$G(T, P, N) = -kT \log Z_p$

# Some useful relationships + grand canonical

From the definition of canonical partition function:

$$Z(T, V, N) = \sum_{\text{microstates-}i} e^{-\beta E_i}$$

Average energy can be obtained from the partition function:

$$\langle E \rangle = \frac{1}{Z} \sum_{\text{microstates-}i} E_i e^{-\beta E_i} = - \frac{\partial \log(Z(T, V, N))}{\partial \beta}$$

Grand canonical ensemble (energy and number of particles fluctuate):

$$\Xi(\mu, V, T) = \sum_{N=0}^{\infty} \sum_{\substack{\text{microstates-}i \\ \text{with } N \text{ atoms}}} e^{-\beta(E_i - \mu N)} = \sum_{N=0}^{\infty} e^{\beta \mu N} Z(N, V, T)$$

→ chemical pot.

$$\langle N(\mu, V, T) \rangle = \frac{1}{\Xi} \sum_{N=0}^{\infty} N e^{\beta \mu N} Z(N, V, T) = kT \frac{\partial \log(\Xi)}{\partial \mu}$$

# Fluctuations

Fluctuations from equilibrium are also related to materials properties

$$\langle \delta A^2 \rangle = \frac{1}{\tau} \int_0^\tau [\underline{A(t) - \langle A \rangle}]^2 dt = \langle A^2 \rangle - \langle A \rangle^2$$

Specific heat:

$$C_V = \left. \frac{\partial E}{\partial T} \right|_{NVT}$$

$$\underline{\langle \delta H^2 \rangle}_{NVT} = kT^2 C_V$$

Compressibility:

$$\beta_V = \left. \frac{1}{V} \frac{\partial V}{\partial P} \right|_{NVT}$$

$$\langle \delta V^2 \rangle_{NPT} = k \langle V \rangle T \underline{\beta_T}$$

# Classical mechanics of the harmonic solid

Second order expansion of Hamiltonian around equilibrium  
(week 3 lecture 5)

$$H = \sum_{i=1}^{3N} \frac{P_i^2}{2M_i} + V(\{R_i^0\}) + \frac{1}{2} \sum_{i,j} \frac{\partial^2 V}{\partial R_i \partial R_j} \bigg|_{\{R_i^0\}} \Delta R_i \Delta R_j$$

After diagonalizing the Hessian matrix  $H$  can be written as a sum of  $3N$  de-coupled harmonic oscillators

$$H = V(\{R_i^0\}) + \sum_{\alpha=1}^{3N} \left( \frac{\pi_{\alpha}^2}{2} + \frac{1}{2} \omega_{\alpha}^2 \rho_{\alpha}^2 \right) + \sum_{i=1}^{3N} \frac{P_i^2}{2M_i}$$

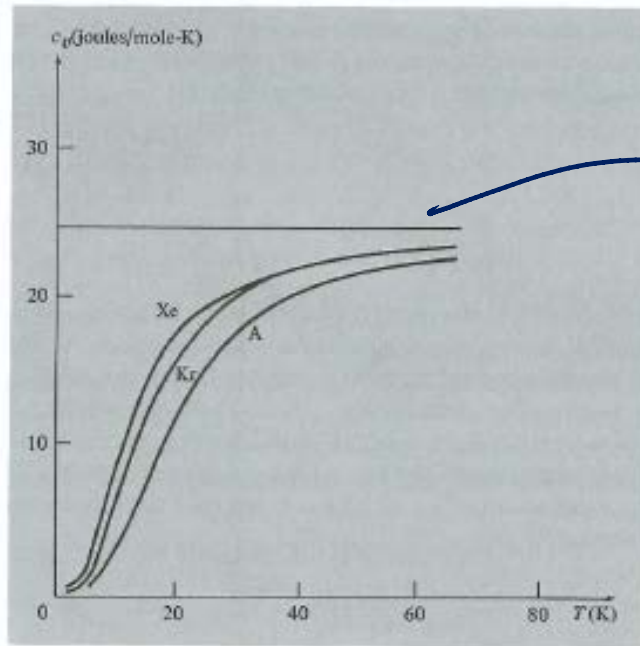
Equipartition: each degree of freedom contributes  $1/2kT$  of energy:

$$E(N, V, T) = V(\{R_0\}) + \frac{2 \times 3N}{2} kT = V(\{R_0\}) + 3NkT$$

Specific heat :

$$C_V(N, V, T) = \frac{\partial E}{\partial T} = 3Nk$$

# Comparing with experiments



**Figure 22.3**  
Measured specific heats of argon, xenon, and krypton. The horizontal line is the classical Dulong and Petit value. (Quoted in M. L. Klein, G. K. Horton, and J. L. Feldman, *Phys. Rev.* 184, 68 (1969).)

Classical harmonic result:

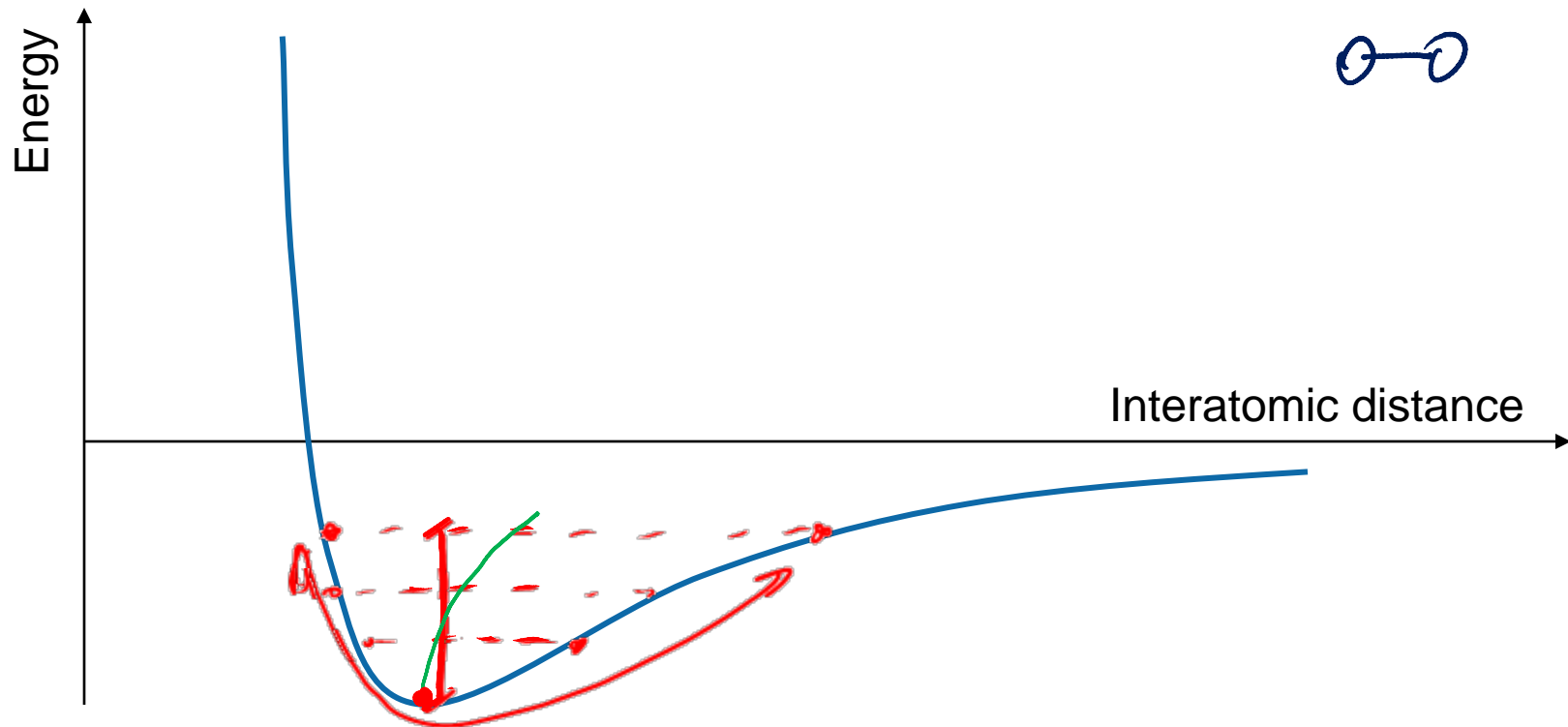
$$C_V(N, V, T) = \frac{\partial E}{\partial T} = \boxed{3Nk}$$

Solid-state Physics  
Ashcroft-Mermin

What error did we make?

# Quantum effects

When does classical mechanics for atoms stop working?



Temperature at which quantum effects kick in depends on frequency

$kT \gg \hbar\omega$       Classical regime

$kT \sim < \hbar\omega$       Quantum regime

# Quantum harmonic oscillator

How do we solve a harmonic oscillator using QM?

Write the Hamiltonian:  $H = \frac{P^2}{2M} + \frac{1}{2}KX^2 = -\frac{\hbar^2}{2M} \frac{\partial^2}{\partial X^2} + \frac{1}{2}KX^2$

Solve the Schrödinger Equation:

$$H\psi = E\psi$$

Family of eigenvalues:  $E_n = \hbar\omega \left( \frac{1}{2} + n \right)$

$n = 0, 1, 2, 3, \dots, \infty$

Compute the partition function:

$$Z = \sum_{\text{microstates}} e^{-\beta E_n} = \sum_{\text{microstates}} e^{-\beta \hbar \omega \left( \frac{1}{2} + n \right)}$$

