

Principles of Electronic Nanobiosensors

Unit 3: Sensitivity

Lecture 3.2: Potentiometric Sensors:

Charge Screening for a Planar Sensor

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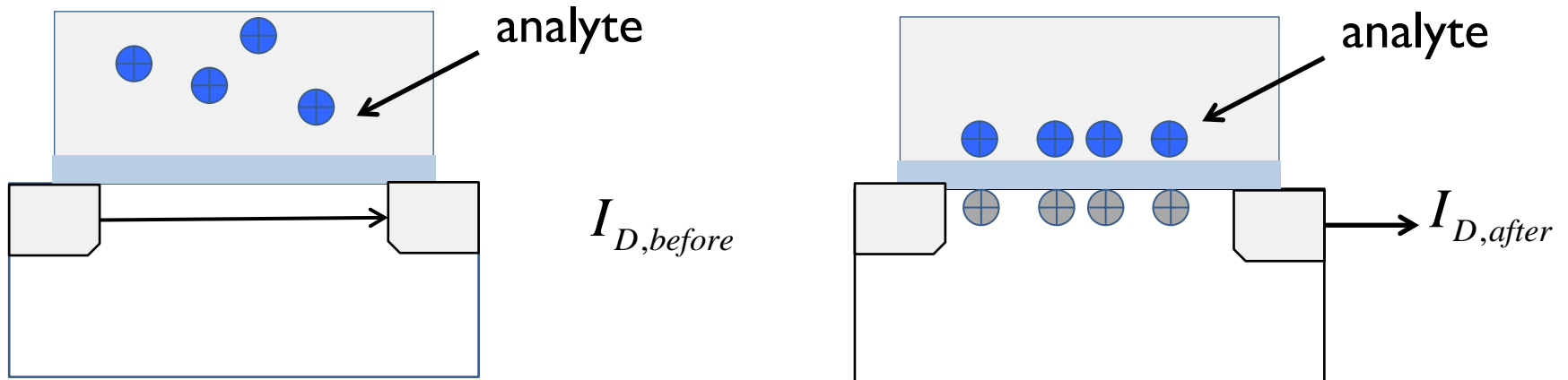
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Outline

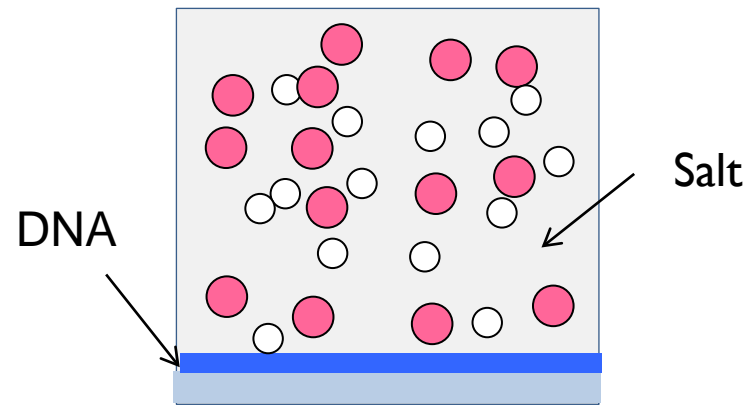
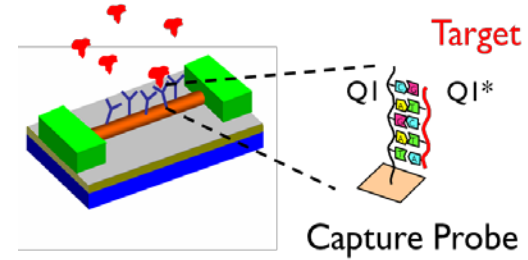
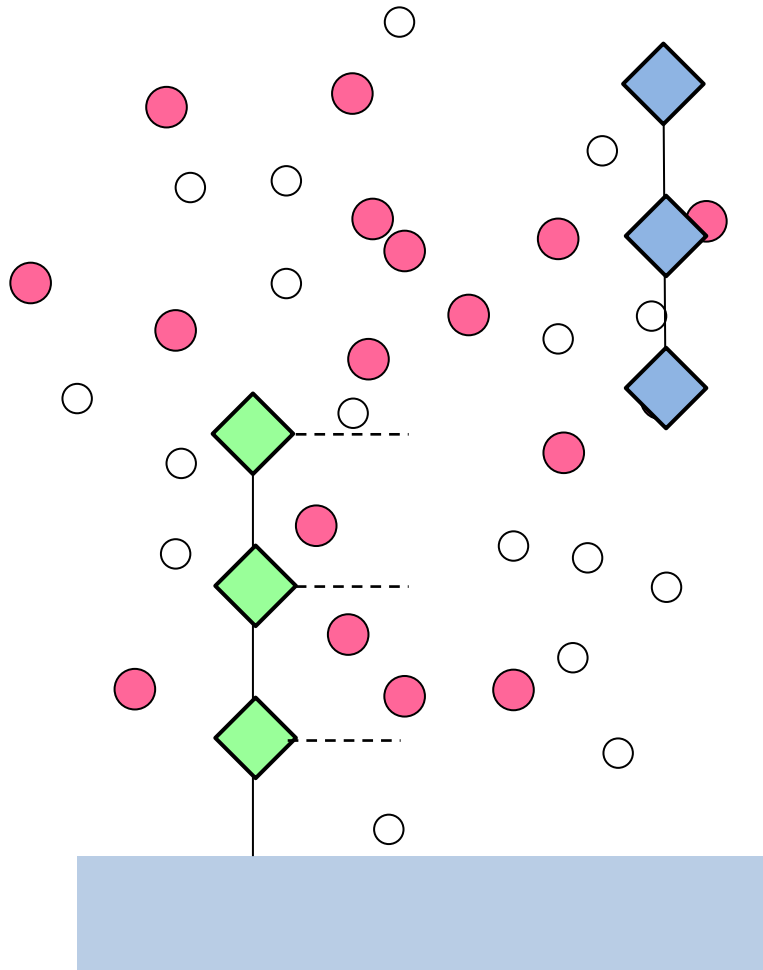
- Background: From a MOSFET to a biosensor
- Biocharges, salt and screening
- Debye theory response of a planar sensor
- Response of planar sensors at high salt concentration: Gouy-Chapman theory
- Conclusions
- Appendix: Derivation of Gouy-Chapman theory

Sensitivity of a MOSFET-based potentiometric sensor



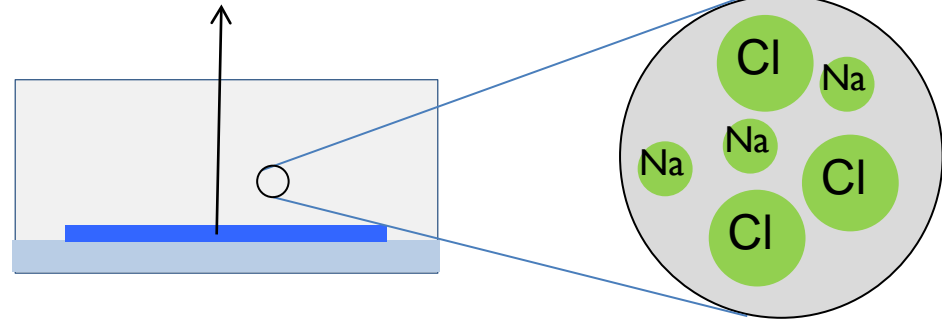
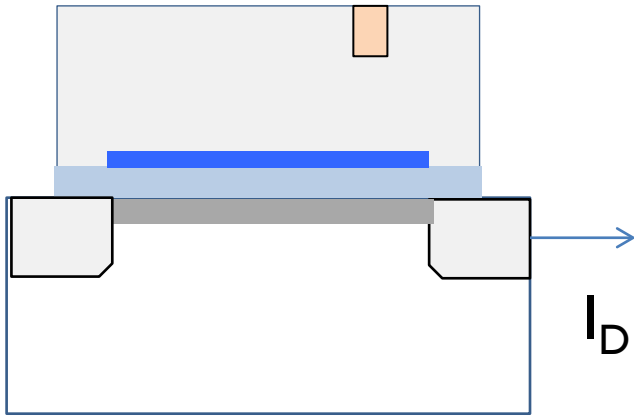
$$S \equiv \frac{I_{D,after} - I_{D,before}}{I_{D,before}} \equiv \frac{\Delta I_D}{I_D} \propto \rho_0$$

DNA binding and Salt screening



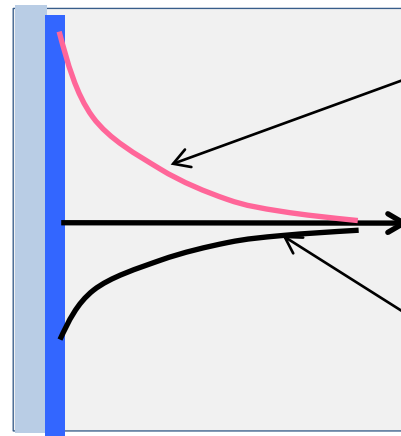
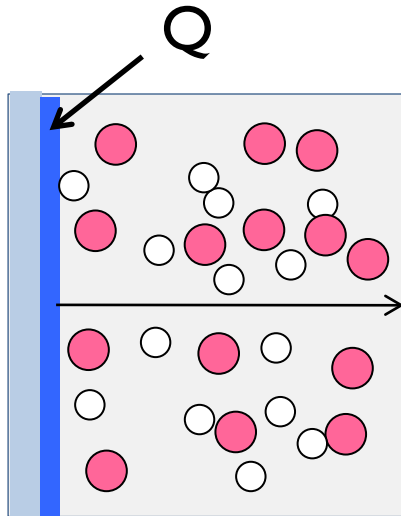
Salt and screening

Salt is ionized in water.



$$U = -\frac{q^2}{4\pi\epsilon_0 r} \times \frac{1}{K_w} \sim k_B T_L$$

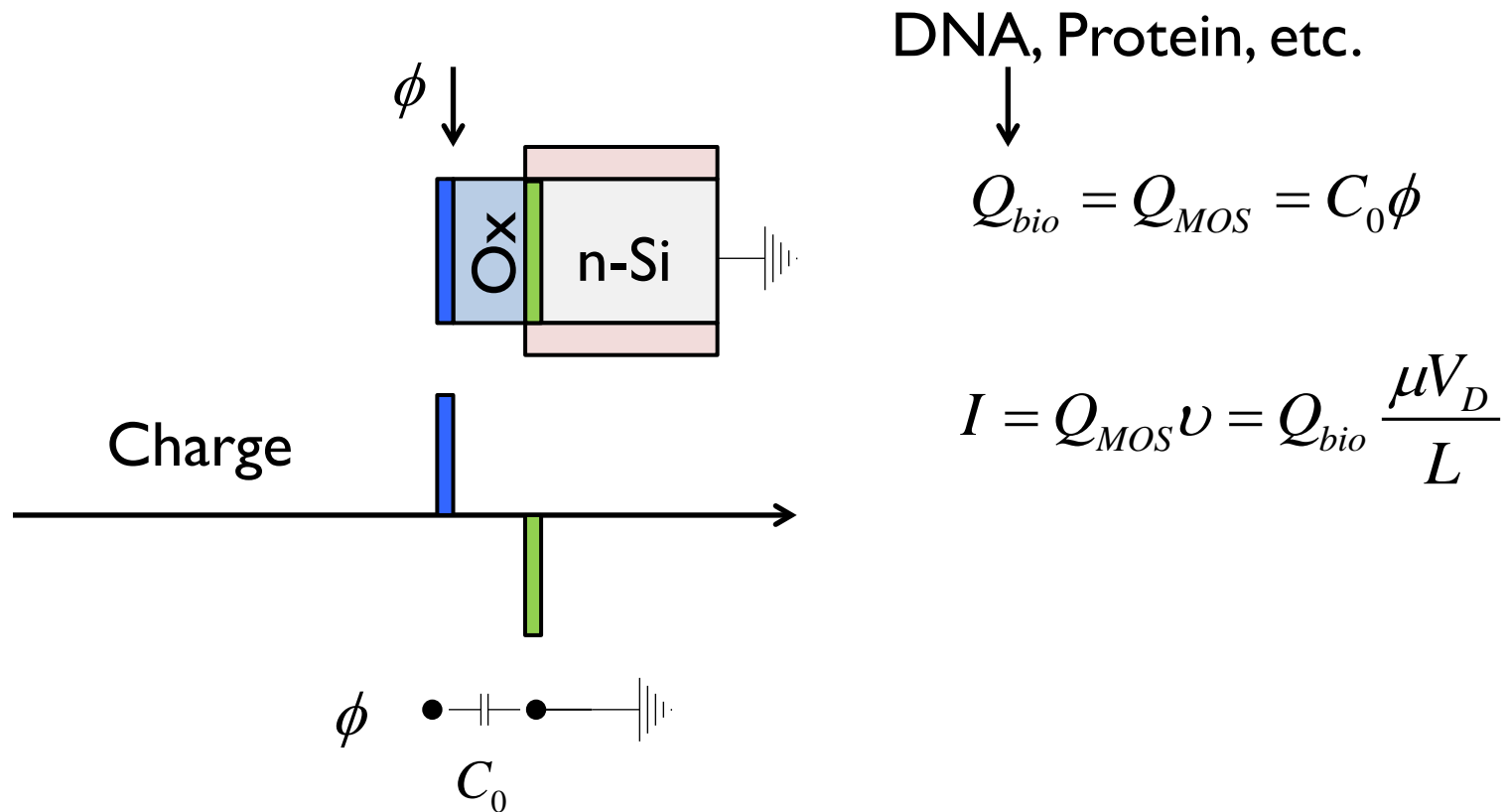
Bjerrum length



$$n^- \sim I_0 N_{av} e^{+q\psi/k_B T_L}$$

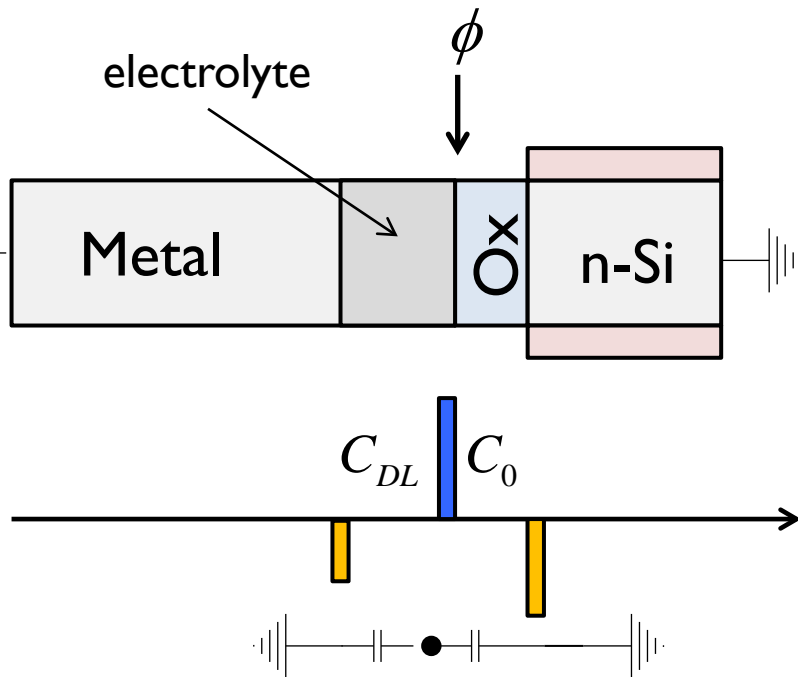
$$n^+ \sim I_0 N_{av} e^{-q\psi/k_B T_L}$$

How does a Ion-sensitive Field Effect Transistor work (accumulation)



The response is proportional to the total charge of the biomolecule

ISFET in electrolyte (accumulation)



$$Q_{bio} = Q_L + Q_R$$

$$= C_{DL}\phi + C_0\phi$$

$$\phi = \frac{Q_{bio}(\text{pH})}{C_{DL}(\text{salt}) + C_0}$$

$$Q_{MOS} = C_0 \frac{Q_{bio}}{C_{DL} + C_0}$$

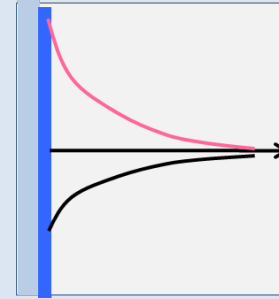
$$I = Q_{MOS}v$$

Electrolyte reduces sensitivity

Calculating C_{DL} : screening in 1D

$$\begin{aligned}\nabla^2\psi &= \frac{d}{dx}\left(\frac{d\psi}{dx}\right) = -\frac{\rho_{DL}}{\kappa_w\epsilon_0} \\ &= -\frac{zq}{\kappa_w\epsilon_0}(n^+ - n^-) \\ &= \frac{2zqI_0N_{av}}{\kappa_w\epsilon_0}\sinh(zq\psi/k_B T) \\ &\approx -\frac{\psi}{L_D^2} \quad (\text{if } \psi \sim k_B T_L)\end{aligned}$$

$$\psi(x) = \psi_0 \times e^{-x/L_D}$$



$$\psi(0) = \psi_0, \quad \psi(\infty) = 0$$

$$n^- = I_0 N_{av} e^{+zq\psi/k_B T_L}$$

$$n^+ = I_0 N_{av} e^{-zq\psi/k_B T_L}$$

$$L_D^2 \equiv \frac{\kappa_w \epsilon_0 k_B T_L}{2z^2 I_0 N_{av} q^2}$$

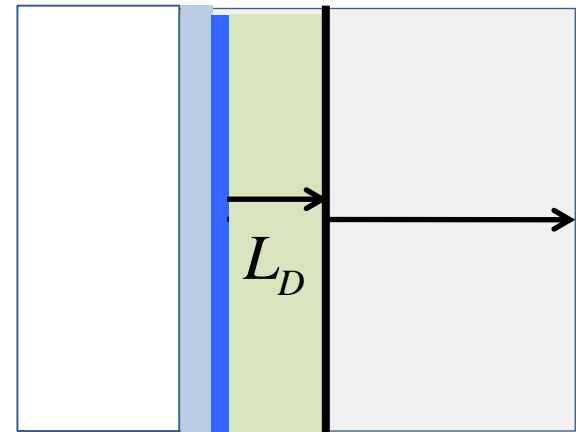
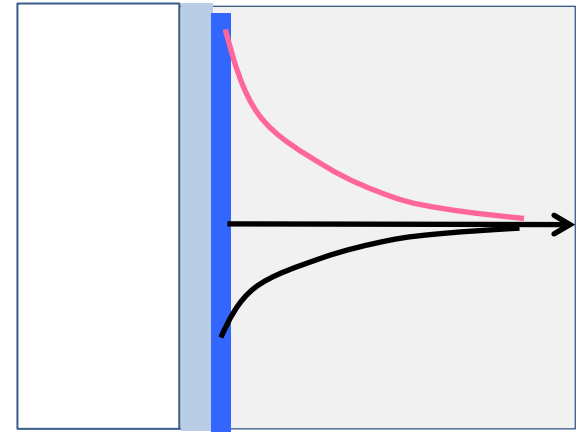
Centroid of DL charge & capacitance

$$\langle x \rangle = \frac{\int_0^{\infty} x \rho_{DL}(x) dx}{\int_0^{\infty} \rho_{DL}(x) dx} = L_D$$

$$\rho_{DL}(x) = -\frac{zq}{\kappa_w \epsilon_0} (n^+ - n^-)$$

$$\approx -\frac{zq}{\kappa_w \epsilon_0} \frac{zq}{k_B T_L} \times \psi$$

$$= -\frac{z^2 q^2}{\kappa_w \epsilon_0 k_B T_L} \times \psi(0) e^{-\frac{x}{L_D}}$$



$$C_{DL} = \frac{\kappa_w \epsilon_0}{\langle x \rangle} = \frac{\kappa_w \epsilon_0}{L_D}$$

Screening in 1D

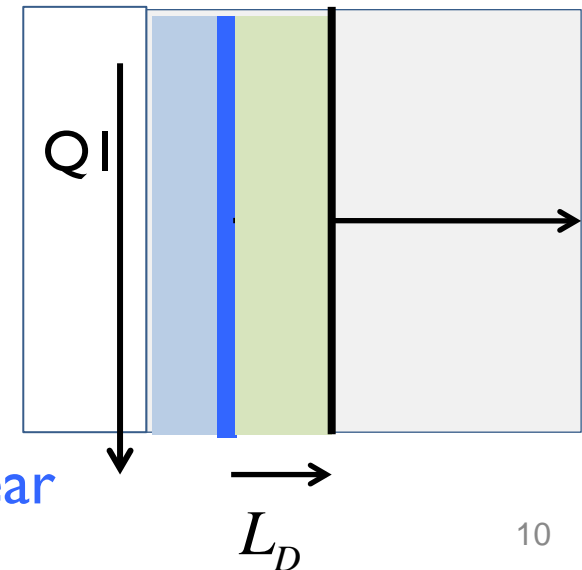
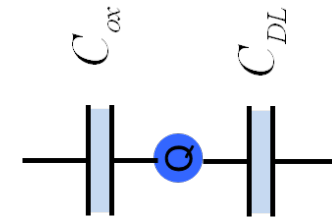
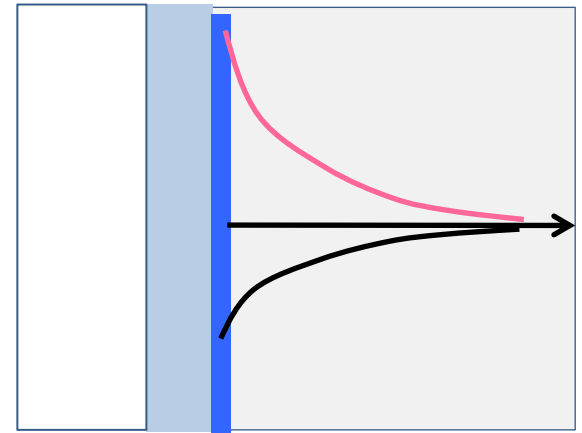
$$Q_{MOS} = \frac{C_{ox}}{C_{ox} + C_{DL}} Q_{bio}$$

$$= \frac{k_{ox} \epsilon_0 / T_{ox}}{k_{ox} \epsilon_0 / T_{ox} + k_w \epsilon_0 / L_D} Q_{bio}$$

$$\rightarrow \frac{k_{ox}}{k_w} \frac{L_D}{T_{ox}} \rho_0 \quad (L_D \sim T_{ox})$$

$$L_D^2 = \kappa_w \epsilon_0 k_B T_L / 2z^2 I_0 N_{av} q^2$$

$$S \propto Q_{MOS} \propto \rho_0 / \sqrt{I_0}$$



Increase in salt density decreases sensitivity
 Response scaled & time-dependent, but still linear

At high salt concentration (Gouy-Chapman theory)

@ low salt
concentration

$$Q_{DL} \sim Q_0 \exp\left(\frac{zq\psi_0}{2k_B T_L}\right)$$

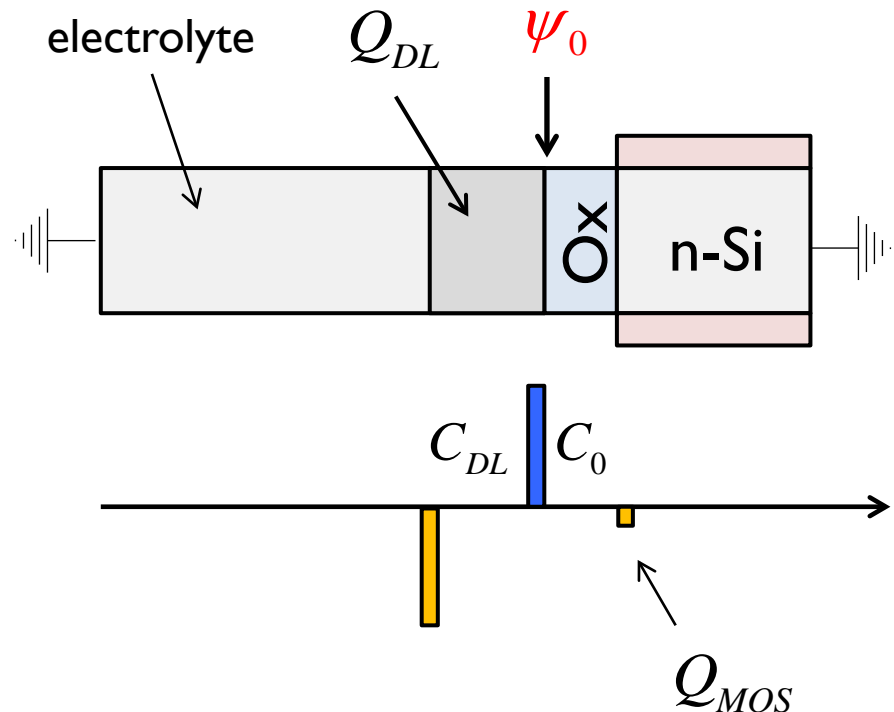
$$Q_{DL} \propto C_{DL}\psi_0 = \frac{k_w \epsilon_0}{L_D} \psi_0$$

$$Q_0 \equiv \sqrt{2k_B T_L k_w \epsilon_0 I_0 N_{avg}}$$

$$Q_{bio} = Q_{MOS} + Q_{DL}$$

$$Q_{bio} = C_{ox}\psi_0 + Q_{DL}$$

$$= C_{ox}\psi_0 + Q_0 \exp\left(\frac{zq\psi_0}{2k_B T_L}\right)$$

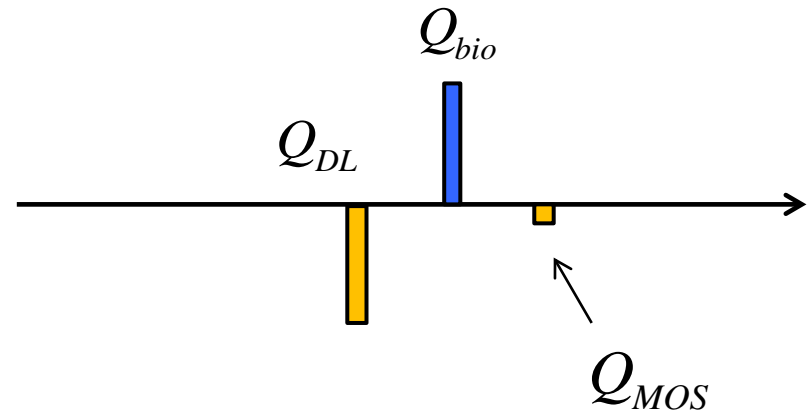


Strong screening (perturbative solution)

If MOS charge is small

$$Q_{bio} \approx Q_0 \exp\left(\frac{zq\psi_0}{2k_B T_L}\right)$$

$$\psi_0 = \frac{2k_B T_L}{zq} \ln \frac{Q_{bio}}{Q_0}$$



$$Q_{MOS} \approx C_{ox} \psi_0 \propto \frac{2k_B T_L}{zq} \ln \frac{Q_{bio}}{Q_0}$$

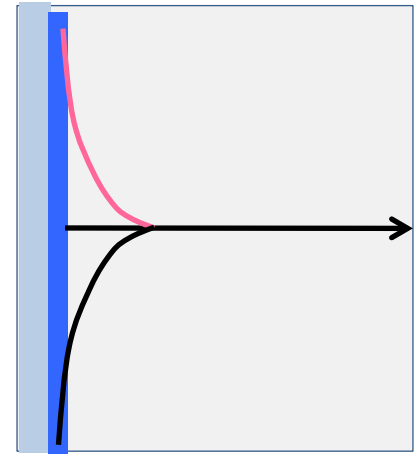
$$= \frac{k_B T_L}{q} \ln \frac{\sigma_T N(t)}{Q_0} = \frac{k_B T_L}{q} \ln \left(\frac{\sigma_T \times k \rho_0 \times t^{(3-D_F/2)}}{Q_0} \right) \leftarrow$$

Strong screening (perturbative solution)

$$Q_{MOS} \approx \frac{k_B T_L}{q} \ln \left(\frac{\sigma_T \times k \rho_0 \times t^{(3-D_F/2)}}{Q_0} \right)$$

$$S(t) \equiv \frac{\Delta I_D}{I_D} \propto Q_{MOS}$$

$$= c_1 \ln(\rho_0) + c_2 \frac{(3-D_F) \ln(t)}{2} - c_3 \ln(I_0) + c_4$$



Response logarithmic
in density, as observed
experimentally

Increases with time,
but logarithmically

Increasing salt
decreases response

Conclusions

- Potentiometric sensors rely on charges of biomolecules for detection
- Many biomolecules can only survive at high –salt concentration
- The screening associated with the salt reduces charge reflected in the channel.
- The screening depends on a complex interplay of multiple variables.

Review questions

- If we bring a pair of DNA in water vs. in air, which configuration would have stronger repulsion?
- What is the role of salt in water in stabilizing the DNA pair? Can it play the same role in air?
- Why does increasing salt concentration reduce sensitivity? Explain physically.
- If we reduce the oxide thickness of a biosensor , would it increase or decrease the sensitivity? What about increasing the dielectric constant?

References

- Nair, Pradeep R., and Muhammad A. Alam. "Screening-limited response of nanobiosensors." *Nano letters* 8.5 (2008): 1281-1285.
- Nair, Pradeep R., and Muhammad A. Alam. "Design considerations of silicon nanowire biosensors." *Electron Devices, IEEE Transactions on* 54.12 (2007): 3400-3408.
- Vacic, Aleksandar, et al. "Determination of molecular configuration by debye length modulation." *Journal of the American Chemical Society* 133.35 (2011): 13886-13889.
- Baumgartner, Stefan, et al. "Predictive simulations and optimization of nanowire field-effect PSA sensors including screening." *Nanotechnology* 24.22 (2013): 225503.

Appendix

- Derivation of the Gouy-Chapman theory for high salt concentration

Strong screening: Gouy-Chapman model

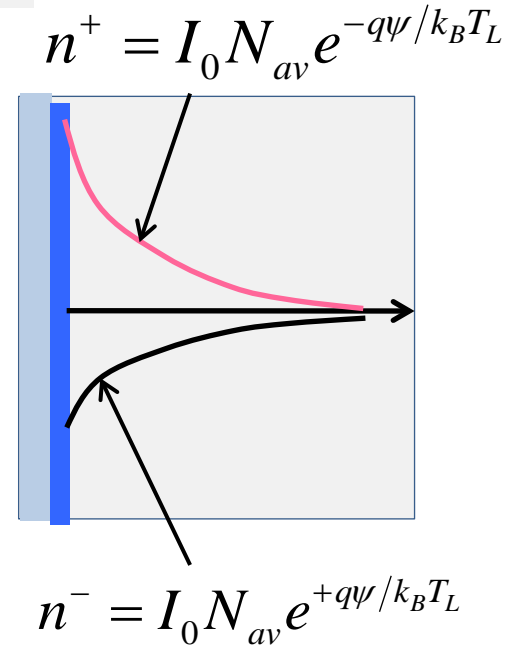
Prove that:
$$Q_{DL} = \sqrt{8k_B T_L \kappa_w \epsilon_0 I_0 N_{av}} \times \sinh\left(\frac{zq\psi_0}{k_B T_L}\right)$$

Boundary conditions: $\psi(0) = \psi_0, \quad \psi(\infty) = 0$

$$\frac{d^2\psi}{dx^2} = -\frac{\rho_{DL}}{\kappa_w \epsilon_0} = -\frac{zq}{\kappa_w \epsilon_0} (n^+ - n^-)$$

Recall that

$$n^+ = I_0 N_{av} e^{-q\psi/k_B T_L} \quad n^- = I_0 N_{av} e^{+q\psi/k_B T_L}$$



$$\frac{d^2\psi}{dx^2} = -\frac{\rho_{DL}}{\kappa_w \epsilon_0} = -\frac{zq}{\kappa_w \epsilon_0} (n^+ - n^-) = \frac{2zqI_0 N_{av}}{\kappa_w \epsilon_0} \sinh(zq\psi/k_B T)$$

Gouy-Chapman model (2)

$$\frac{d^2\psi}{dx^2} = \frac{2zqI_0N_{av}}{\kappa_w\epsilon_0} \sinh(zq\psi/k_B T)$$

$$2 \frac{d\psi}{dx} \times \frac{d^2\psi}{dx^2} = \frac{d}{dx} \left(\frac{d\psi}{dx} \right)^2 = 2 \frac{d\psi}{dx} \times \frac{2zqI_0N_{av}}{\kappa_w\epsilon_0} \sinh\left(\frac{zq\psi}{k_B T}\right)$$

$$\int_0^\infty d \left(\frac{d\psi}{dx} \right)^2 = \int_{\psi_0}^0 \frac{4zqI_0N_{av}}{\kappa_w\epsilon_0} \sinh\left(\frac{zq\psi}{k_B T}\right) d\psi$$

$$\left(\frac{d\psi}{dx} \right)^2 \Big|_{x=\infty} - \left(\frac{d\psi}{dx} \right)^2 \Big|_{x=0} = \frac{4I_0N_{av}k_B T}{\kappa_w\epsilon_0} \left(\cosh\left(\frac{zq\psi_0}{k_B T}\right) - 1 \right)$$

Gouy-Chapman model (3)

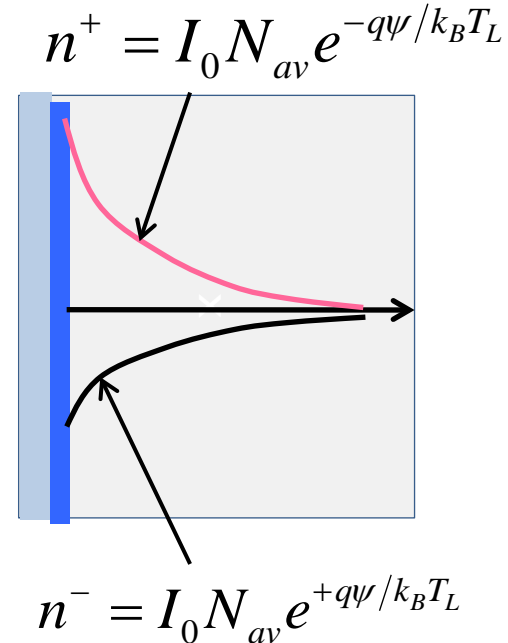
$$\left(\frac{d\psi}{dx}\right)^2 \Big|_{x=0} = \frac{4I_0 N_{av} k_B T}{\kappa_w \epsilon_0} \times 2 \sinh^2 \frac{zq\psi_0}{2k_B T}$$

$$\left(\frac{d\psi}{dx}\right) \Big|_{x=0} = \sqrt{\frac{8I_0 N_{av} k_B T}{\kappa_w \epsilon_0}} \sinh \frac{zq\psi_0}{2k_B T}$$

$$\frac{Q_{DL}}{k_w \epsilon_0} = \mathcal{E} = \frac{d\psi}{dx} \Big|_{x=0} \quad \text{Gauss's law @x=0}$$

$$Q_{DL} = k_w \epsilon_0 \sqrt{\frac{8I_0 N_{av} k_B T}{\kappa_w \epsilon_0}} \sinh \frac{zq\psi_0}{2k_B T}$$

$$= \sqrt{8k_B T_L k_w \epsilon_0 I_0 N_{av}} \times \sinh \left(\frac{zq\psi_0}{2k_B T_L} \right)$$



Gouy-Chapman model (4)

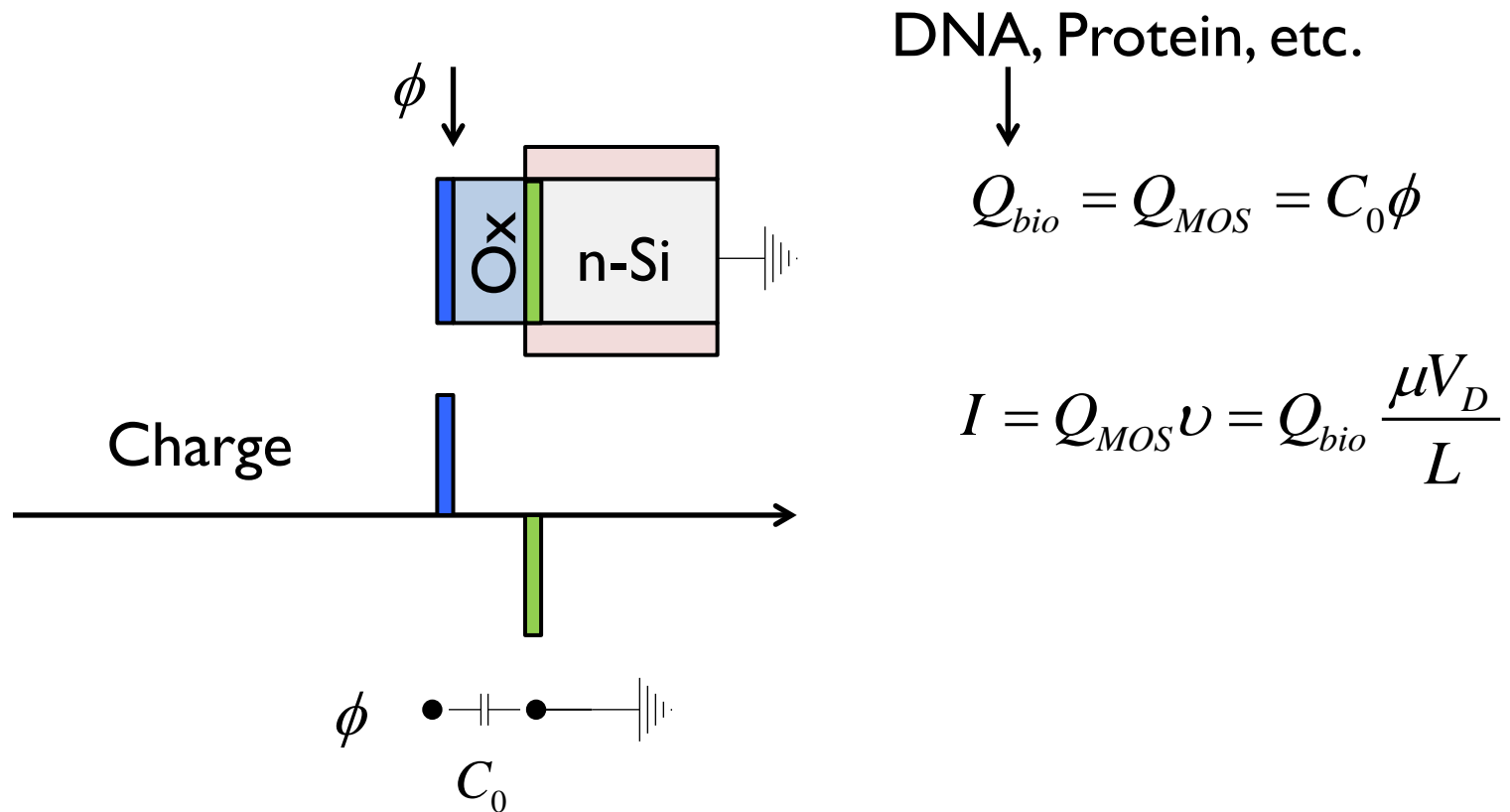
$$\begin{aligned} Q_{DL} &= k_w \varepsilon_0 \sqrt{\frac{8I_0 N_{av} k_B T}{\kappa_w \varepsilon_0}} \sinh \frac{zq\psi_0}{2k_B T} \\ &= \sqrt{8k_B T_L k_w \varepsilon_0 I_0 N_{av}} \times \sinh \left(\frac{zq\psi_0}{2k_B T_L} \right) \\ &\approx Q_0 \times \exp \frac{zq\psi_0}{2k_B T_L} \end{aligned}$$

$$Q_0 \equiv \sqrt{2k_B T_L k_w \varepsilon_0 I_0 N_{avg}}$$

Appendix

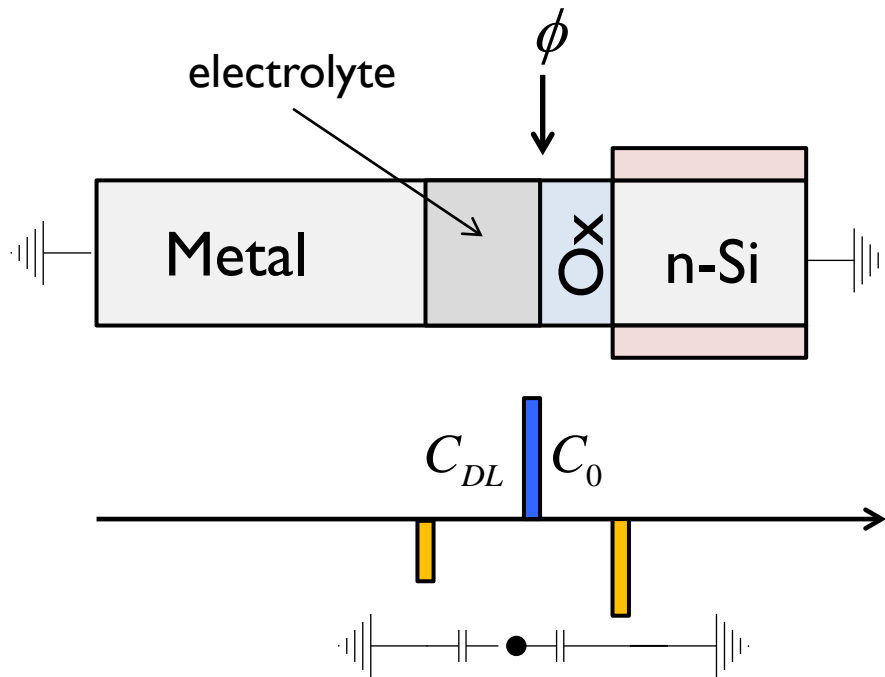
- Modes of operation of a field-effect transistor

How does a Ion-sensitive Field Effect Transistor work (accumulation)



The role of charge can be viewed as shifting the flat-band voltage

ISFET in electrolyte (accumulation)



$$Q_{bio} = Q_L + Q_R$$

$$= C_{DL}\phi + C_0\phi$$

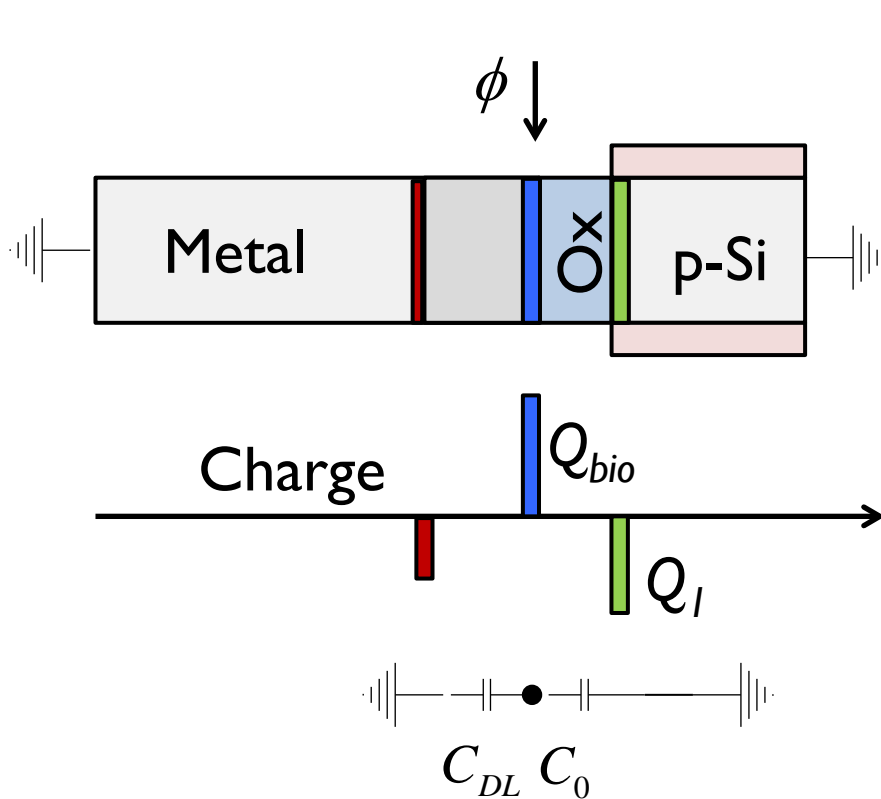
$$\phi = \frac{Q_{bio}(\text{pH})}{C_{DL}(\text{salt}) + C_0}$$

$$Q_{MOS} = C_0 \frac{Q_{bio}}{C_{DL} + C_0}$$

$$I = Q_{MOS}v$$

Electrolyte reduces sensitivity

How does an ISFET work (inversion)



$$\phi = \frac{Q_{bio}}{C_{DL} + C_0}$$

$$Q_i = C_0 \frac{Q}{C_{DL} + C_0}$$

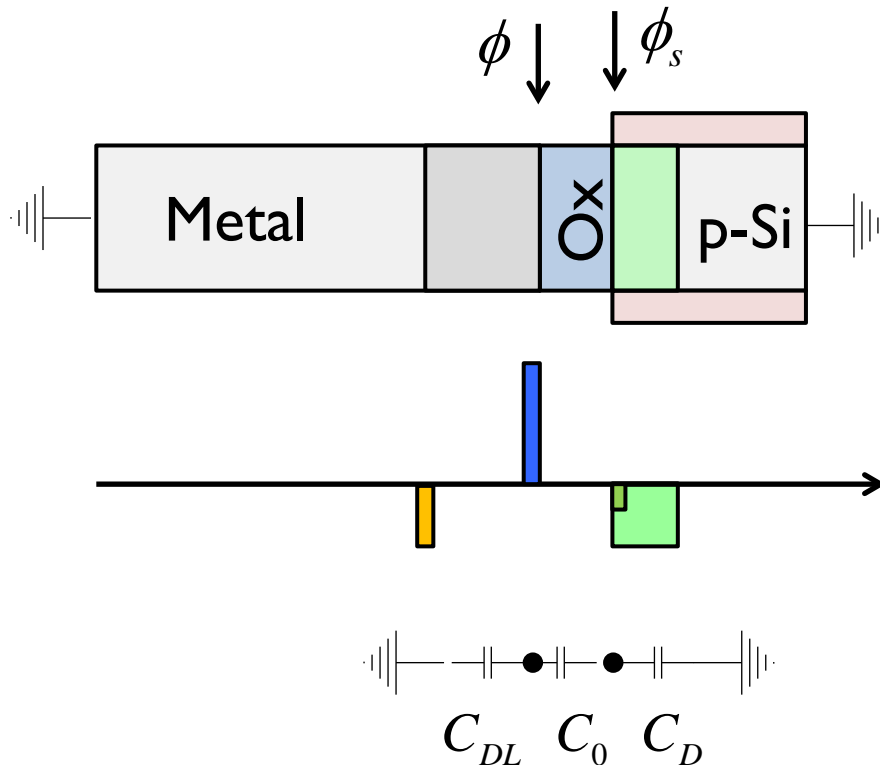
$$Q_T = C_O (V_G - V_T) + Q_{MOS}$$

$$= C_O (V_G + \phi - V_T)$$

$$I_D = Q_T \nu$$

$$\Delta I_D = Q_{MOS} \nu$$

How does an ISFET work (subthreshold)



$$m = \frac{C_0}{C_D + C_0}$$

$$\phi_s = \frac{V_G + \phi}{m}$$

$$\phi = \frac{Q_{bio}}{C_{DL} + C_0}$$

$$I = q \frac{n_i^2}{N_A} e^{\frac{q\phi_s}{k_B T}} \times v_D$$

$$= q \frac{n_i^2}{N_A} e^{\frac{q(V_G + \phi)}{mkT}} \times v_D$$

Although I have used ϕ as an effective gate voltage, it should be really be viewed as a change in flat-band voltage