

LECTURE #10

WE HAVE SHOWN THAT THE STATES OF THE HYDROGEN ATOM CAN BE SPECIFIED WITH QUANTUM NUMBERS.

(1) n : PRINCIPAL QUANTUM NUMBER = $n=1, 2, 3, \dots$
 \hookrightarrow THIS SPECIFIES THE ENERGY LEVEL

(2) l : AZIMUTHAL QUANTUM NUMBER, $l \leq n-1$

NOMENCLATURE:

$l=0 \rightarrow$ "s" \rightarrow SHARP

$l=1 \rightarrow$ "p" \rightarrow PRINCIPAL

$l=2 \rightarrow$ "d" \rightarrow DIFFUSE

$l=3 \rightarrow$ "f" \rightarrow FUNDAMENTAL

$l=4 \rightarrow$ "g"

$l=5 \rightarrow$ "h"

(3) m_l : MAGNETIC QUANTUM NUMBER, $|m_l| \leq l$

WHAT DO l AND m_l REPRESENT PHYSICALLY?

RECALL: CLASSICAL NOTION OF ANGULAR MOMENTUM

$$\vec{L} = \vec{r} \times \vec{p} = \begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ x & y & z \\ p_x & p_y & p_z \end{vmatrix}$$

$$\therefore L_x = y p_z - z p_y, \quad L_y = z p_x - x p_z, \quad L_z = x p_y - y p_x$$

\Rightarrow QUANTUM MECHANICALLY, WE CAN DEFINE ANGULAR MOMENTUM OPERATORS.

$$L_x = \frac{\hbar}{i} \left(y \frac{\partial}{\partial z} - z \frac{\partial}{\partial y} \right), \quad L_y = \frac{\hbar}{i} \left(z \frac{\partial}{\partial x} - x \frac{\partial}{\partial z} \right), \quad L_z = \frac{\hbar}{i} \left(x \frac{\partial}{\partial y} - y \frac{\partial}{\partial x} \right)$$

CONSIDER THE COMMUTATION RELATION FOR L_x, L_y :

$$\begin{aligned} [L_x, L_y] &= [y p_z - z p_y, z p_x - x p_z] \\ &= [y p_z, z p_x] - [y p_z, x p_z] - [z p_y, z p_x] + [z p_y, x p_z] \end{aligned}$$

FROM THE CANONICAL COMMUTATION RELATIONS

$$\left[\text{i.e., } [r_i, p_j] = -[p_i, r_j] = i\hbar \delta_{ij}, \quad [r_i, r_j] = [p_i, p_j] = 0 \right]$$

WE KNOW THAT THE ONLY OPERATORS THAT FAIL TO COMMUTE ARE x WITH p_x , y WITH p_y , AND z WITH p_z .

$$\therefore [L_x, L_y] = y p_x [p_z, z] + x p_y [z, p_z] = i\hbar (x p_y - y p_x) = i\hbar L_z$$

\Rightarrow FUNDAMENTAL COMMUTATION RELATIONS FOR ANGULAR MOMENTUM:

$$[L_x, L_y] = i\hbar L_z; \quad [L_y, L_z] = i\hbar L_x; \quad [L_z, L_x] = i\hbar L_y$$

NOTE: L_x, L_y , AND L_z ARE INCOMPATIBLE OBSERVABLES

HOWEVER, CONSIDER $L^2 = L_x^2 + L_y^2 + L_z^2$

$$\begin{aligned} [L^2, L_x] &= [L_x^2, L_x] + [L_y^2, L_x] + [L_z^2, L_x] \\ &= L_y [L_y, L_x] + [L_y, L_x] L_y + L_z [L_z, L_x] + [L_z, L_x] L_z \\ &= L_y (-i\hbar L_z) + (-i\hbar L_z) L_y + L_z (i\hbar L_y) + (i\hbar L_y) L_z \\ &= 0 \end{aligned}$$

SIMILARLY, $[L^2, L_y] = 0$ AND $[L^2, L_z] = 0$

OR MORE COMPACTLY, $[L^2, \vec{L}] = 0$

i.e., L^2 IS COMPATIBLE WITH EACH COMPONENT OF \vec{L}

\Rightarrow WE CAN FIND SIMULTANEOUS EIGENSTATES OF L^2 AND L_z :

$$L^2 f = \lambda f \quad \text{AND} \quad L_z f = \mu f$$

INTRODUCE ANGULAR MOMENTUM LADDER OPERATOR:

$$L_{\pm} = L_x \pm i L_y$$

$$\begin{aligned} [L_z, L_{\pm}] &= [L_z, L_x] \pm i [L_z, L_y] \\ &= i\hbar L_y \pm i(-i\hbar L_x) = \pm\hbar (L_x \pm i L_y) = \pm\hbar L_{\pm} \end{aligned}$$

$$\therefore [L_z, L_{\pm}] = \pm\hbar L_{\pm}$$

OF COURSE, $[L^2, L_{\pm}] = 0$

$$L^2(L_{\pm} f) = L_{\pm}(L^2 f) = L_{\pm}(\lambda f) = \lambda(L_{\pm} f)$$

\Rightarrow IF f IS AN EIGENFUNCTION OF L^2 AND L_z , THEN SO IS $L_{\pm} f$ WITH THE SAME EIGENVALUE λ .

$$\begin{aligned} \therefore L_z(L_{\pm} f) &= (L_z L_{\pm} - L_{\pm} L_z) f + L_{\pm} L_z f = \pm\hbar L_{\pm} f + L_{\pm}(\mu f) \\ &= (\mu \pm \hbar)(L_{\pm} f) \end{aligned}$$

$\Rightarrow L_{\pm} f$ IS AN EIGENFUNCTION OF L_z WITH NEW EIGENVALUE $m \pm \hbar$

L_+ IS A "RAISING OPERATOR" THAT INCREASES THE EIGENVALUE OF L_z BY \hbar

L_- IS A "LOWERING OPERATOR" THAT DECREASES THE EIGENVALUE OF L_z BY \hbar

NOTE: WE CANNOT RAISE L_z INDEFINITELY SINCE EVENTUALLY THE z -COMPONENT WOULD EXCEED THE TOTAL

\therefore THERE MUST EXIST A "TOP RUNG" SUCH THAT:

$$L_+ f_+ = 0$$

LET $\hbar l$ BE THE EIGENVALUE OF L_z AT THE TOP RUNG:

$$\therefore L_z f_+ = \hbar l f_+ ; L^2 f_+ = \lambda f_+$$

$$\begin{aligned} \text{NOTE: } L_{\pm} L_{\mp} &= (L_x \pm i L_y)(L_x \mp i L_y) = L_x^2 + L_y^2 \mp i(L_x L_y - L_y L_x) \\ &= L^2 - L_z^2 \mp i(i\hbar L_z) \end{aligned}$$

$$\Rightarrow L^2 = L_{\pm} L_{\mp} + L_z^2 \mp \hbar L_z$$

$$\therefore L^2 f_+ = (L_- L_+ + L_z^2 + \hbar L_z) f_+ = (0 + \hbar^2 l^2 + \hbar^2 l) f_+ = \lambda f_+$$

$$\Rightarrow \lambda = \hbar^2 l(l+1)$$

SIMILARLY, THERE MUST EXIST A "BOTTOM RUNG" SUCH THAT:

$$L_- f_b = 0$$

LET $\hbar \bar{l}$ BE THE EIGENVALUE OF L_z AT THE BOTTOM RUNG.

$$L_z f_b = \hbar \bar{l} f_b ; L^2 f_b = \lambda f_b$$

$$\therefore L^2 f_b = (L_+ L_- + L_z^2 - \hbar L_z) f_b = (0 + \hbar^2 \bar{l}^2 - \hbar^2 \bar{l}) f_b = \lambda f_b$$

$$\Rightarrow \lambda = \hbar^2 \bar{l}(\bar{l} - 1)$$

$$\therefore l(l+1) = \bar{l}(\bar{l} - 1)$$

$\Rightarrow \bar{l} = l + 1$ (WHICH IS ABSURD SINCE THE BOTTOM RUNG
WOULD BE HIGHER THAN THE TOP RUNG)

OR $\bar{l} = -l$

\Rightarrow THE EIGENVALUES OF L_z ARE $m\hbar$ WHERE m RUNS
FROM $-l$ TO $+l$ IN N INTEGER STEPS

$$\therefore l = -l + N \Rightarrow l = \frac{N}{2} \Rightarrow l \text{ MUST BE AN INTEGER
OR HALF INTEGER}$$

IN SUMMARY,

$$L^2 f_l^m = \hbar^2 l(l+1) f_l^m ; L_z f_l^m = \hbar m f_l^m$$

WHERE $l = 0, \frac{1}{2}, 1, \frac{3}{2}, \dots$; $m = -l, -l+1, \dots, l-1, l$

$\underbrace{\hspace{10em}}$
2l+1 RUNGS ON THE LADDER

TO DETERMINE THE EIGENFUNCTIONS, WE NEED TO EXPRESS
 L_z AND L^2 IN SPHERICAL COORDINATES:

$$L_z = \frac{\hbar}{i} \frac{\partial}{\partial \phi}, \quad L^2 = -\hbar^2 \left[\frac{1}{\sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial}{\partial \theta} \right) + \frac{1}{\sin^2 \theta} \frac{\partial^2}{\partial \phi^2} \right]$$

$$\therefore L^2 f_\ell^m = -\hbar^2 \left[\frac{1}{\sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial}{\partial \theta} \right) + \frac{1}{\sin^2 \theta} \frac{\partial^2}{\partial \phi^2} \right] f_\ell^m = \hbar^2 \ell(\ell+1) f_\ell^m$$

* THIS IS PRECISELY THE ANGULAR EQUATION!

$$L_z f_\ell^m = \frac{\hbar}{i} \frac{\partial}{\partial \phi} f_\ell^m = \hbar m f_\ell^m$$

* THIS IS EQUIVALENT TO THE Φ EQUATION!

$\Rightarrow f_\ell^m = Y_\ell^m(\theta, \phi) \rightarrow$ THE SPHERICAL HARMONICS ARE THE EIGENFUNCTIONS OF L^2 AND L_z

WHEN WE SOLVED THE SCH. EQ. FOR SPHERICALLY SYMMETRIC POTENTIALS BY SEPARATION OF VARIABLES, WE INADVERTENTLY CONSTRUCTED SIMULTANEOUS EIGENFUNCTIONS OF THREE COMMUTING OPERATORS H , L^2 , AND L_z :

$$H\psi = E\psi, \quad L^2\psi = \hbar^2 \ell(\ell+1)\psi, \quad L_z\psi = \hbar m\psi$$

NOTE: FROM SEPARATION OF VARIABLES, WE FOUND THAT ℓ (AND ALSO m) COULD ONLY BE INTEGER VALUED. HOWEVER, THE LADDER OPERATORS YIELDED HALF-INTEGERS VALUES, WHAT IS THE PHYSICAL SIGNIFICANCE OF THE HALF-INTEGERS VALUES?