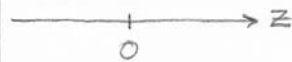
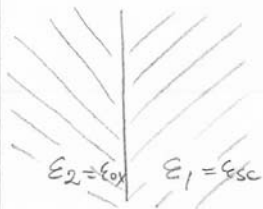


## 2D-System

The potential due to a charge center located at the position  $Q_i(\vec{r}_i, z_i)$  is:

$$\Psi(\vec{r}, z) = \frac{e^2}{4\bar{u}\epsilon_{sc}} \left[ \frac{1}{\sqrt{(\vec{r}-\vec{r}_i)^2 + (z-z_i)^2}} + \frac{\epsilon_{sc} - \epsilon_{ox}}{\epsilon_{sc} + \epsilon_{ox}} \frac{1}{\sqrt{(\vec{r}-\vec{r}_i)^2 + (z+z_i)^2}} \right] \quad \text{for } z > 0$$

$$\Psi(\vec{r}, z) = \frac{e^2}{2\bar{u}(\epsilon_{sc} + \epsilon_{ox})} \frac{1}{\sqrt{(\vec{r}-\vec{r}_i)^2 + (z-z_i)^2}} \quad ; \text{ for } z < 0$$



$$\mathcal{F} \left[ \frac{1}{[(\vec{r}-\vec{r}_i)^2 + (z \pm z_i)^2]^{1/2}} \right] = \frac{2\bar{u}}{q} e^{-q|z-z_i|} e^{-i\vec{k} \cdot \vec{r}_i}$$

If the charge center is located at the interface, then:

$$z_i = 0$$

$$\Psi(q, z, \vec{r}_i) = \frac{2\bar{u}}{q} e^{-q|z|} e^{-i\vec{k} \cdot \vec{r}_i}$$

Then, if we have a single subband case, the probability for scattering from a state  $k_{||}$  in a subband  $n$  to a state  $k'_{||}$  in a subband  $m$  is:

$$S_{nm}(k_{||}, k'_{||}) = \frac{2\bar{u}}{\hbar} |H(k_{||}, k'_{||})|_{nm}^2 \delta(E - E' \pm \hbar\omega)$$

$$\text{where: } \Psi_n(k_{||}) = \frac{1}{\sqrt{A}} e^{i\vec{k}_{||} \cdot \vec{r}} \Psi_n(z) \quad ; \quad E = E_n + \hbar^2 k_{||}^2 / 2m^*$$

$$\Psi_m(k'_{||}) = \frac{1}{\sqrt{A}} e^{i\vec{k}'_{||} \cdot \vec{r}} \Psi_m(z') \quad ; \quad E' = E_m + \hbar^2 k'_{||}{}^2 / 2m^*$$

and:

$$H(k_{||}, k'_{||})_{nm} = \frac{1}{A} \int d\vec{r} e^{i(\vec{k}_{||} - \vec{k}'_{||}) \cdot \vec{r}} \int dz \Psi_m^*(z) H_{qp}(\vec{R}) \Psi_n(z)$$

$$\downarrow$$

$$\Psi(\vec{r}, z) = \frac{e^2}{2\bar{u}(\epsilon_{sc} + \epsilon_{ox})} \frac{1}{\sqrt{r^2 + z^2}}$$

Now, using that  $\mathcal{F} \left[ \frac{1}{\sqrt{r^2 + z^2}} \right] = \frac{2\bar{u}}{q} e^{-q|z|}$ , we get:

$$M(k_{11}, k_{11}')_{nm} = \frac{1}{A} \frac{e^2}{2\bar{\sigma}(\epsilon_{sc} + \epsilon_{ox})} \frac{2\bar{\sigma}}{q_{11}} \int dz \psi_m^*(z) e^{-q_{11}|z|} \psi_n(z)$$

$$= \frac{e^2}{A(\epsilon_{sc} + \epsilon_{ox})} \frac{1}{q_{11}} \int dz \underbrace{\psi_m^*(z) e^{-q_{11}|z|}}_{A_{00}(q_{11})} \psi_n(z)$$

For a single subband, we have:

$$S(k_{11}, k_{11}') = \frac{2\bar{\sigma}}{\hbar} \frac{n_I}{A} \frac{e^4}{(\epsilon_{sc} + \epsilon_{ox})^2} \left( \frac{1}{q_{11} + q_D} \right)^2 A_{00}^2(q_{11})$$

$$\frac{1}{\Gamma_m(k_{11})} = \frac{2\bar{\sigma}}{\hbar} n_I \frac{e^4}{(\epsilon_{sc} + \epsilon_{ox})^2} \iint \frac{1}{(2\bar{\sigma})^2} \frac{\delta(E_{k_{11}'} - E_{k_{11}}) k_{11} dk_{11} d\varphi (1 - \cos\varphi)}{(q_{11} + q_D)^2} A_{00}^2(q_{11})$$

$$E_{11} = \frac{\hbar^2 k_{11}^2}{2m^*} \Rightarrow k_{11} dk_{11} = \frac{m^*}{\hbar^2} dE_{k_{11}}$$

$$\frac{1}{\Gamma_m(k_{11})} = n_I \frac{e^4 m^*}{2\bar{\sigma} \hbar^3 (\epsilon_{sc} + \epsilon_{ox})^2} \int_0^{2\bar{\sigma}} \frac{d\varphi (1 - \cos\varphi)}{(q_{11} + q_D)^2} \underbrace{A_{00}^2(q_{11})}_{\left(\frac{b}{b + q_{11}}\right)^3}$$

$$q_{11} = 2k_{11} \sin(\varphi/2)$$

$$q_D = \frac{e^2 N_s}{2\alpha \hbar \epsilon_0 T} \Rightarrow \text{Debye-Hückel screening parameter}$$

$$\varphi/2 = \left(\frac{b^3}{2}\right)^{1/2} e^{-6z/2}$$

$$\text{where } b = \left(\frac{12 m_0^* N^+ e^2}{K_s \epsilon_0 \hbar^2}\right)^{1/3}$$

$$N^+ = N_{dop} + \frac{11}{32} N_s$$

$$q_{11}^2 = k_{11}^2 - 2k_{11}k_{11}' + k_{11}'^2 = 2k_{11}^2 - 2k_{11}^2 \cos\varphi = 2k_{11}^2 (1 - \cos\varphi)$$

$$q_{11}^2 = 4k_{11}^2 \sin^2(\varphi/2) \Rightarrow q_{11} = 2k_{11} \sin(\varphi/2)$$

$$\frac{1}{\Gamma_m(k_{11})} = \frac{n_I e^4 m^* b^3}{\hbar^3 (\epsilon_{sc} + \epsilon_{ox})^2} \int_0^{2\bar{\sigma}} \frac{k \sin^2(\varphi/2) d\varphi}{(2k_{11} \sin(\varphi/2) + q_D)^2 [b + 2k_{11} \sin(\varphi/2)]^3}$$

$$b \rightarrow \text{large} \rightarrow \frac{1}{\Gamma_m(k_{11})} = \frac{n_I e^4 m^*}{\bar{\sigma} \hbar^3 4\alpha^2} \frac{1}{4k_{11}^2} \cdot 2\bar{\sigma} = \frac{n_I e^4 m^* \cdot \bar{\sigma}}{\bar{\sigma} \hbar \cdot 4\alpha^2 \cdot 4 k_{11}^2 E_k}$$

$$E_k = \frac{\hbar^2 k_{11}^2}{2m^*} \Rightarrow \hbar^2 k_{11}^2 = 2m^* E_k$$

$$\frac{1}{\Gamma_m(k_{11})} = \frac{n_I e^4}{16 \hbar \alpha^2 E_k}$$