



# **Chapter 2 - Properties of Light**

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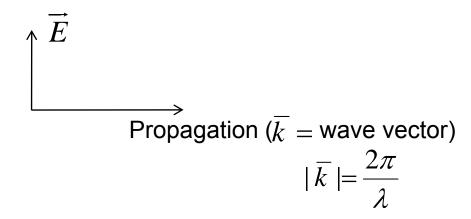


Amplitude A and phase φ are random functions in both <u>time</u> and <u>space</u>:

$$\vec{E}(\vec{r},t) = \vec{e}A(\vec{r},t).e^{i\phi(\vec{r},t)}$$
(1)



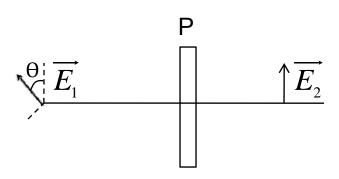
- a) Polarization:
  - Gives the direction of field oscillation
  - Generally, light is a <u>transverse wave</u> (unlike sound = longitudinal)

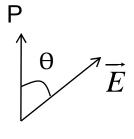


 ■ Anisotropic materials: different optical properties along different axis → useful



- Polarization: a)
  - There is always a basis  $(\hat{x}, \hat{y})$  for decomposing the field into 2 polarizations (eigen modes); equivalently (right, left) circular polarization is also a basis.
  - Dichroism: different absorption for different pol  $\rightarrow$  one way to create polarizers:





Malus Law:  $|E_2| = |E_1| * \cos \theta$ 

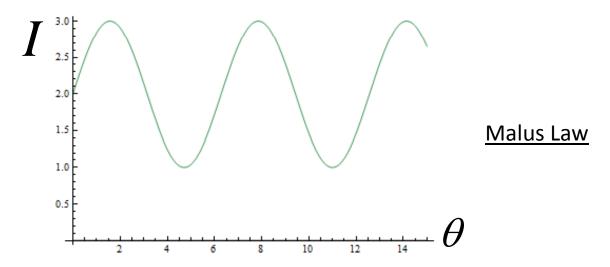
$$|E_2| = |E_1| * \cos \theta$$

(2)



a) Polarization:

$$I = |E|^2$$
;  $I_2 = I_1 * \cos^2 \theta$ 



Birefringence – Different refr. index for different pol.



- a) Polarization:
  - Natural Light  $\rightarrow$  unpolarized  $\rightarrow$  superposition  $E_x = E_y$  with no phase relationship between the two
  - Circularly polarized  $\rightarrow$   $E_x = E_y$ ,  $\phi_x \phi_y = \pi/2$ !
  - Matrix formalism of polarization transformation (Jones 2x2, complex & Muller 4x4, real)

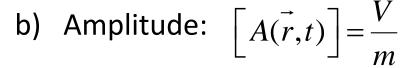
We'll do this later.

$$\begin{pmatrix} E_{x} \\ E_{y} \end{pmatrix} = J \begin{pmatrix} E_{x} \\ E_{y} \end{pmatrix}$$

$$|E|^2 = I \rightarrow \text{Stokes Vect. Dim 4}$$

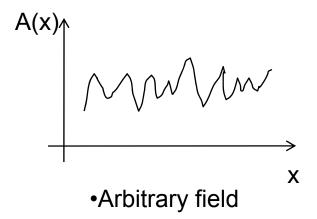
$$J_{ij}\in\mathbb{C}$$

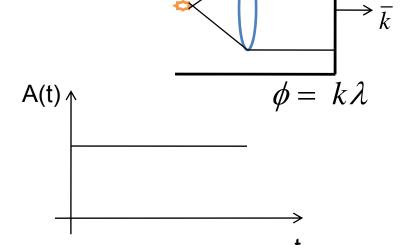


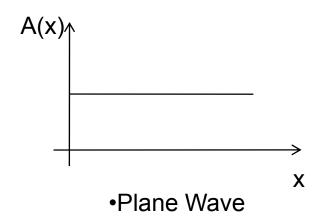




Thermal source

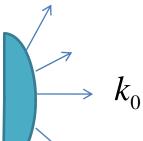






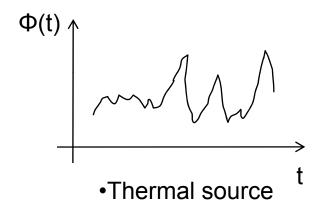
Stabilized laser

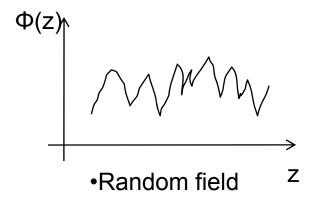


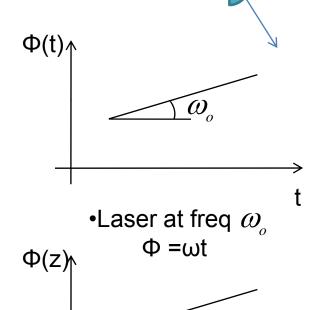


Ζ

c) Phase:  $[\Phi] = rad$ 







Plane Wave

Φ=kz



- c) Phase:  $[\Phi] = rad$ 
  - For quasi-monochromatic fields, plane wave

$$\phi = \omega t - \vec{k} \cdot \vec{r}$$

• 
$$k = \frac{\omega}{c} = \frac{2\pi\upsilon}{c} = \frac{2\pi}{Tc} = \frac{2\pi}{\lambda} = \text{wave number}$$
 (3)

$$\lambda = cT$$
;  $T = \frac{1}{N}$ ;  $\omega = 2\pi N$ 



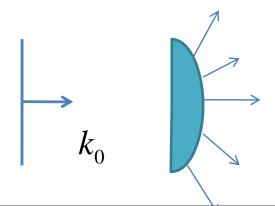
# 2.2 The frequency domain representation

Random variable E(t) has a frequency-domain counterpart:

$$E(\omega) = A(\omega)e^{i\phi(\omega)} \tag{4}$$

Similarly E(x) has a frequency-domain pair:

$$E(\xi) = A(\xi)e^{i\phi(\xi)} \tag{5}$$



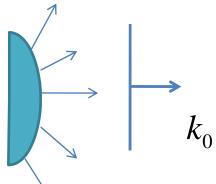
$$\phi(\omega) = k \cdot z = n(\omega) \cdot k_0 \cdot z$$

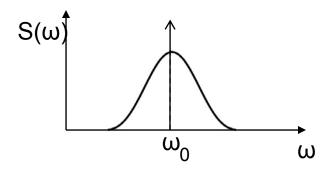
$$\uparrow_{k_0} = \frac{2\pi}{\lambda}$$

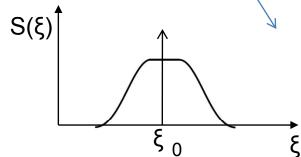


# 2.2 The frequency domain representation

- a) Spectral amplitude:
  - Optical Spectrum:  $S(\omega) = |A(\omega)|^2$
  - Angular Spectrum:  $S(\xi) = |A(\xi)|^2$







- $[\xi] = m^{-1} =$ Spatial Frequency (connects to angular spectrum)
- Tipically:  $t \leftrightarrow \omega$  Will follow similar equations  $x \xi$
- The information contained is the same  $(t, \omega)$  and  $(x, \xi)$

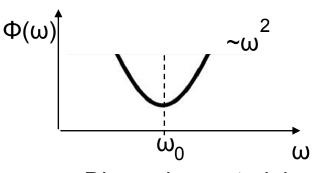


# 2.2 The frequency domain representation

- b) Spectral phase:
  - Phase delay of each spectral component

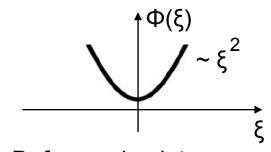
Optical Frequency

$$\alpha = chirp$$



•Dispersive material (linear chirp)

**Spatial Frequency** 



 Defocused point source (1<sup>st</sup> order aberration) A point is mapped to a blur



• Full similarity between  $(t,\omega)$  and  $(x,\xi)$ 

$$\frac{d}{d\omega}\omega^2 \sim \omega$$



#### 2.3 Measurable Quantities

The information about the system under investigation may be contained in <u>polarization</u> and:

■ A(t), 
$$\phi(t)$$
  
■ A( $\omega$ ),  $\phi(\omega)$   
■ A(x),  $\phi(x)$   
■ A( $\xi$ ),  $\phi(\xi)$   
(x, $\xi$ )  
8 quantities

Experimentally, we have access only to:

$$I = \left\langle \left| A(t) \right|^2 \right\rangle = \text{time average}$$

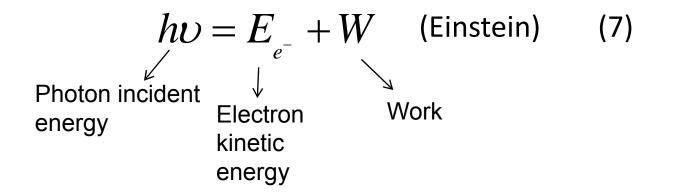


#### 2.3 Measurable Quantities

Experimentally, we have access only to:

$$I = \langle |A(t)|^2 \rangle = \text{time average}$$
 (6)

• i.e the phtodetectors (photodiode, CCD, retina, etc) produce photoelectrons:





#### 2.3 Measurable Quantities

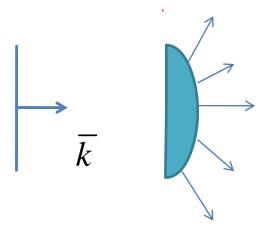
- All detectors sensitive to power/energy
- However, all 8 quantities can be accessed via various tricks
- Eg1: Want  $I(\lambda)$  → measure  $I(\theta)$  and use a device with  $\theta(\lambda)$
- Eg2: Want  $\phi \rightarrow$  use interferometry  $\rightarrow I(\phi)\alpha |E_1||E_2|\cos(\phi_1 \phi_2)$



 Space - momentum or energy-time cannot be measured simmultaneously with infinite accuracy

For photons:

$$\begin{cases} E = \hbar \omega \\ -\frac{1}{p} = \hbar \bar{k}; \ p = \frac{h}{\lambda} \end{cases}$$





a) 
$$t-\omega$$

$$\hbar\Delta\omega\Delta t = \text{constant}$$

$$\rightarrow |\Delta\omega\Delta t \simeq 2\pi$$

- Implications:
  - 1- short pulses require broad spectrum
- 2-high spectral resolution requires long time of measurement



$$\frac{\overline{k}_{i}}{\overline{q}}$$

$$\underline{\overline{k}_{i}} \underbrace{\overline{k}_{s}}_{\Theta} \underbrace{\overline{\Delta p}} = h(\overline{k}_{s} - \overline{k}_{i}) = h\overline{q}$$

$$\Rightarrow \left| \Delta x | \overline{q} | \approx \pi \right|$$

$$\Delta x |q| \approx \pi$$
 ;  $|q| = 2k \sin(\frac{\theta}{2})$ 

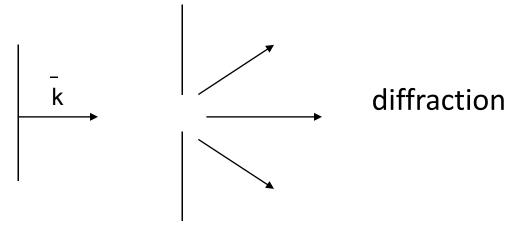
$$\rightarrow \Delta x \frac{2\sin(\theta/2)}{2} \approx 1$$

$$\theta \sim \frac{\lambda}{\Delta x}$$

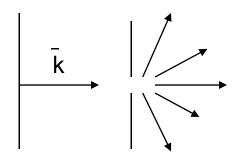
$$\rightarrow \Delta x_{\min} \approx$$

 $\Delta x_{\min} \approx \frac{\lambda}{2}$  - meaning of resolution





■ Smaller aperture → Higher angles



- If aperture  $<\frac{\lambda}{2}$ , light doesn't go through (easily)
- Eg: Microwave door



- We will encounter these relationships many times later
- Fourier may have understood this uncertainity principle way before <u>Heisenberg!</u>