

Quantum Transport:

ATOM TO TRANSISTOR

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Lecture 32: Non-equilibrium Density Matrix
Ref. Chapter 9.2



Network for Computational Nanotechnology



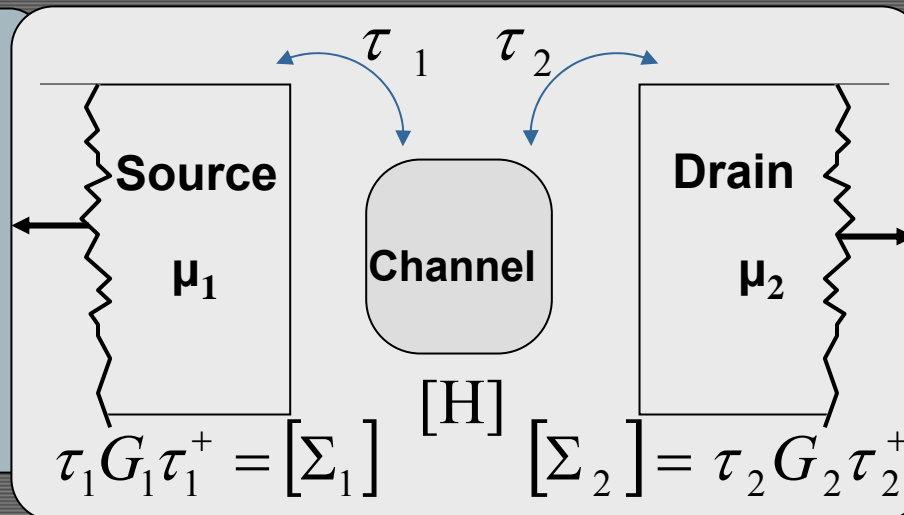
Current Flow

00:00

- **Reminder: We're discussing the current flow through the channel connected to the source and drain contacts with two different chemical potentials.**

- **Towards the beginning of the course we discussed this problem for a single level device and we derived the equations with common sense argument without much use of quantum mechanics.**

- **Huge size; however the effect on the channel can be written using Σ_1 which has the same size as that of the channel.**



- **Huge size; however the effect on the channel can be written using Σ_2 which has the same size as that of the channel.**

$$\left[EI - H_1 + i0^- \right]^{-1} = G_1$$

$$G_2 = \left[EI - H_2 + i0^+ \right]^{-1}$$

Density Matrix; Current

$$[\rho] = \int \frac{dE}{2\pi} [G^n(E)] \text{ want to derive}$$

$$G^n = ([A_1]f_1 + [A_2]f_2)$$

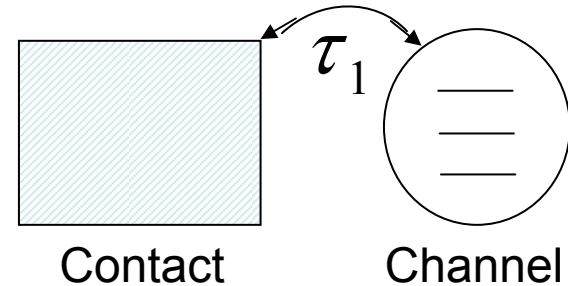
$$A_1 = G\Gamma_1G^+ \quad A_2 = G\Gamma_2G^+$$

$$G = (EI - H - \Sigma_1 - \Sigma_2)^{-1}$$

$$I_i = \left(-\frac{q}{\hbar} \right) \int \frac{dE}{2\pi} \tilde{I}_i(E) \text{ want to derive}$$

$$\tilde{I}_i(E) = \text{Trace}(\Gamma_i A) f_1 - \text{Trace}(\Gamma_i G^n)$$

- Next, we want to derive G^n .



- Schrödinger Equation

- General Form $E\Psi = \bar{H}\Psi$

- Partitioned $E \begin{Bmatrix} \psi \\ \Phi \end{Bmatrix} = \begin{bmatrix} H & \tau \\ \tau^+ & H_1 + i0^+ \end{bmatrix} \begin{Bmatrix} \psi \\ \Phi \end{Bmatrix}$

- Contact before being connected to channel: $(EI - H_1 + i0^+) \Phi_1 = 0$

- What happens after connection?

- Schrödinger Equation becomes:

$$E \begin{Bmatrix} \psi \\ \Phi_1 + \chi \end{Bmatrix} = \begin{bmatrix} H & \tau_1 \\ \tau_1^+ & H_1 - i0^+ \end{bmatrix} \begin{Bmatrix} \psi \\ \Phi_1 + \chi \end{Bmatrix}$$

- So we have:

$$E\psi = [H]\psi + [\tau_1]\chi + [\tau]\Phi_1$$

$$E\Phi_1 + E\chi = [\tau_1^+]\psi + [H_1 - i0^+]\Phi_1 + [H_1 + i0^+]\chi$$

- Since Φ_1 is an eigenstate of the reservoir,

$$E\Phi_1 + E\chi = [\tau_1^+]\psi + \cancel{[H_1 - i0^+]\Phi_1} + [H_1 + i0^+]\chi$$

- Grouping the similar terms,

$$[E - H_1 + i0^+]\chi = [\tau_1^+]\psi$$

- The solution is:

$$\chi = [G_1][\tau_1^+]\psi$$

- Substituting back into the first equation,

$$[EI - H - \Sigma_1]\psi = [\tau_1]\Phi_1$$

$$([\Sigma_1] = \tau_1 G_1 \tau_1^+)$$

- We are used to dealing with the Schrödinger equation like

$$[EI - H]\psi = 0$$

- But here we have a driven system that we solve for:

$$[EI - H - \Sigma_1]\psi = [\tau_1]\Phi_1$$

Towards Electron Density

- In case of two contacts, we would have:

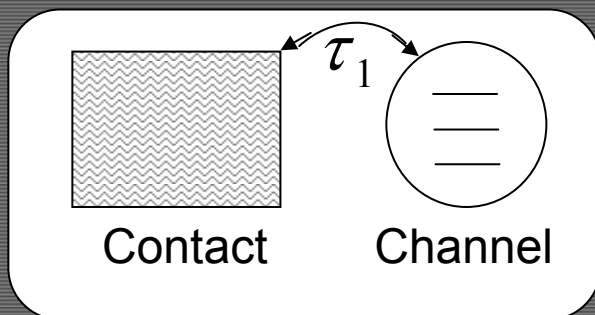
$$[EI - H - \Sigma_1 - \Sigma_2]\psi = [\tau_1]\Phi_1 + [\tau_2]\Phi_2$$

For now, we keep things simple and continue with one contact; solving for ψ , we have: $\psi = [G]S_1$

Where,

$$(S_1 = [\tau_1]\Phi_1) \text{ and } G = (EI - H - \Sigma_1)^{-1}$$

- Next step is to find the electron density inside the channel.



- There are millions of eigenstates inside the reservoir which we can denote by α and each of them has its own eigenvalue. i.e.

$$(E - \varepsilon_\alpha + i0^+) \Phi_\alpha = 0$$

and each solution will give rise to a certain ψ .

A Very Important Physical Concept

- We know that the net electron density is $\psi^*\psi$. Now the question is:

$$\left| \sum_\alpha \psi_\alpha \right|^2 \quad \text{or} \quad \sum_\alpha |\psi_\alpha|^2$$

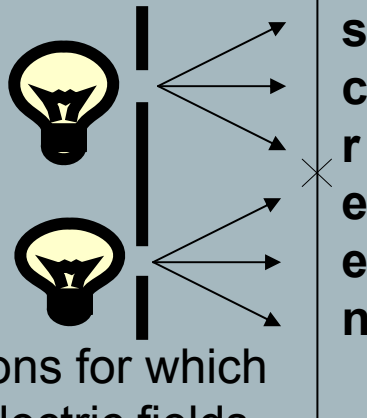
Wrong Correct

- Note that this is only true as long as α represents an eigenstate of the contact.
- The reason relies on coherence discussed on the next page.

Young's Double-Slit Experiment Intensity vs. Electric Field

29:50

- There are lots of eigenstates in the contact, however there is NO coherence between them and in that sense they all act independently.
- For instance, consider Young's double-slit experiment:
 - For thermal sources without coherence we should add intensities whereas for coherent sources like laser we might have had situations for which we'd want to add the electric fields.
 - Electrons are more like thermal sources and they act incoherently.
 - We should find ψ due to each excitation in the contact and then add the intensities ($\psi\psi^*$).



$$\psi_{\alpha} = [G] S_{\alpha}$$

$$(S_1 = [\tau_1] \Phi_1)$$

$$\psi_{\alpha} = [G] [[\tau_1] \Phi_{\alpha}]$$

- Electron Density

$$\{\psi_{\alpha}\} \{\psi_{\alpha}\}^+ = \phi_2 \begin{Bmatrix} \times \\ \otimes \\ \times \\ \times \end{Bmatrix} \left\{ \begin{matrix} \times \\ \times \\ \otimes \\ \times \end{matrix} \right\}^* \phi_3 \text{ Row 2} = \begin{pmatrix} \text{Column 3} \\ \phi_2 \phi_3 \end{pmatrix}$$

- The significant ones from the point of view of electron density are the diagonal ones.

Total Number of Electrons

$$\{\psi_{\alpha}\}^+ \{\psi_{\alpha}\} = \text{Trace} \left[\{\psi_{\alpha}\} \{\psi_{\alpha}\}^+ \right]$$

- Based on $\psi_\alpha = [G][\tau_1]\Phi_\alpha$

we have:

$$\{\psi_\alpha\}\{\psi_\alpha\}^+ = [G][\tau_1]\Phi_\alpha\Phi_\alpha^+[\tau_1]^+[G]^+$$

- Summing over all states that are occupied,

$$\sum_\alpha f_1(\varepsilon_\alpha)\{\psi_\alpha\}\{\psi_\alpha\}^+ = \sum_\alpha f_1(\varepsilon_\alpha)[G][\tau_1]\Phi_\alpha\Phi_\alpha^+[\tau_1]^+[G]^+$$

Making use of the definition of delta function:

$$\left(\sum_\alpha f_1(\varepsilon_\alpha) = \int dE f_1(E) \sum_\alpha \delta(E - \varepsilon_\alpha) \right)$$

- We have: $\sum_\alpha f_1(\varepsilon_\alpha)\{\psi_\alpha\}\{\psi_\alpha\}^+ = \int dE f_1(E) G \tau_1 \underbrace{\sum_\alpha \Phi_\alpha \Phi_\alpha^+ \delta(E - \varepsilon_\alpha) \tau_1^+}_{[a_1]} G^+$

Spectral Function $[a_1]$

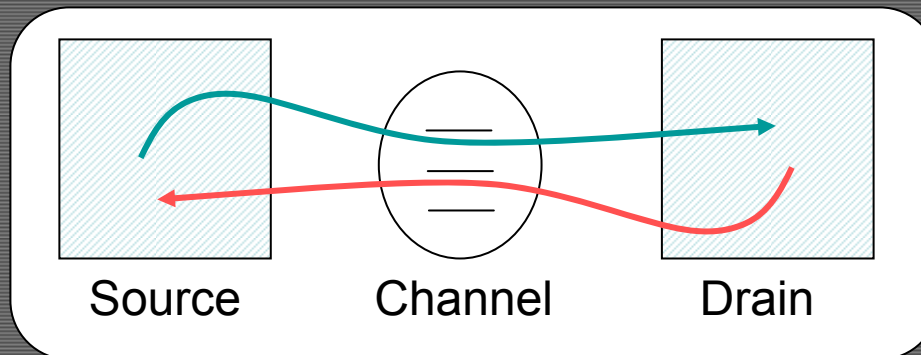
$$\sum_{\alpha} f_1(\varepsilon_{\alpha}) \{ \psi_{\alpha} \} \{ \psi_{\alpha} \}^{\dagger} = \int dE f_1(E) G \tau_1 \sum_{\alpha} \Phi_{\alpha} \Phi_{\alpha}^{\dagger} \delta(E - \varepsilon_{\alpha}) \tau_1^{\dagger} G^{\dagger}$$

$$\Rightarrow \int dE G^n = \int dE f_1(E) G \underbrace{\tau_1 a_1 \tau_1^{\dagger}}_{\Gamma_1} G^{\dagger}$$

- We wanted to derive the electron density inside the channel:

$$G^n = \left(\begin{array}{l} [A_1] f_1 + [A_2] f_2 \\ (G \Gamma_1 G^{\dagger} = A_1) \end{array} \right) \quad \left(\begin{array}{l} \\ (A_2 = G \Gamma_2 G^{\dagger}) \end{array} \right)$$

- Note: $\Gamma_1 = i(\Sigma_1 - \Sigma_1^{\dagger}) = i(\tau_1 G_1 \tau_1^{\dagger} - \tau_1 G_1^{\dagger} \tau_1^{\dagger}) = \tau_1 a_1 \tau_1^{\dagger}$
 $\Sigma_1 = \tau_1 G_1 \tau_1^{\dagger} \quad \tau_1 G_1 \tau_1^{\dagger} = \Sigma_1^{\dagger}$



General Philosophy

- Start from something that describes the device and contact together.
- Contact is full of eigenstates which are trying to fill up the device; find the resulting wave function.
- Each eigenstate acts on the device independently.
- Do the states crossing from right to left and from left to right block one another because of Exclusion Principle?
- No. As long as you start from two states that are orthogonal, even if they overlap, they can't block each other.