

NCN Summer School: Nanoscale Transistors: July, 2012

# Lecture 10: Scattering Model

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and

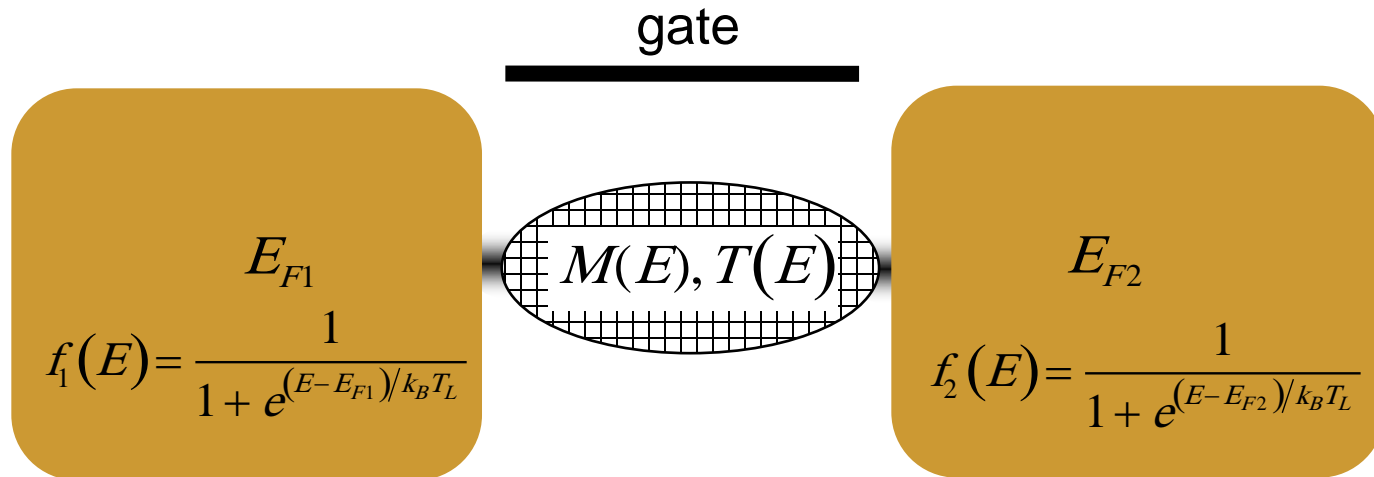
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Lundstrom 7.2012



# Landauer approach to transport



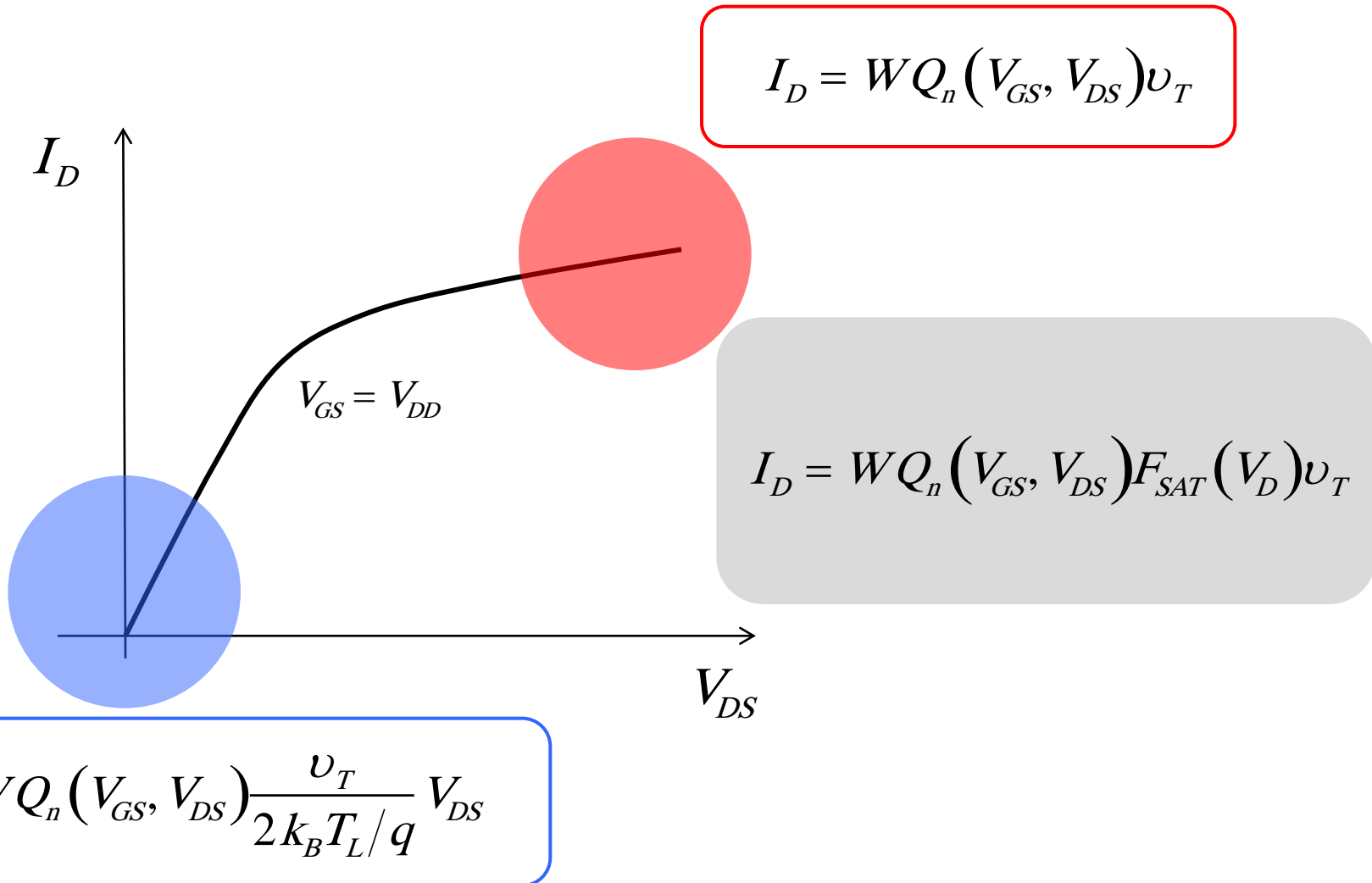
$$I_D = \frac{2q}{h} \int T(E) M(E) (f_1 - f_2) dE$$

any drain bias

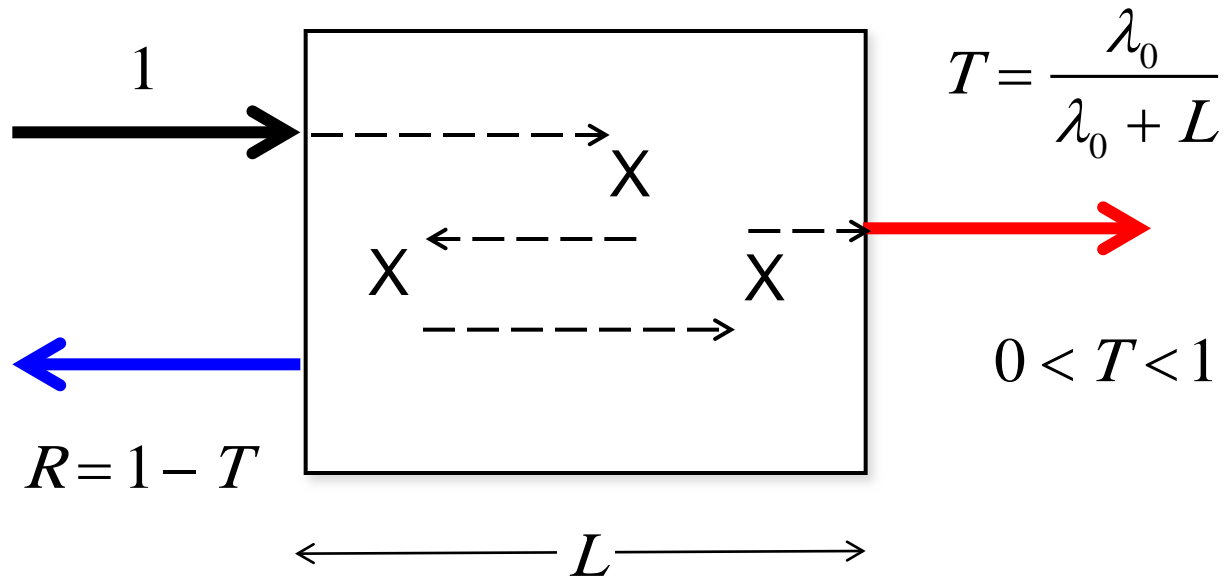
$$I_D = \left[ \frac{2q^2}{h} \int T(E) M(E) \left( -\frac{\partial f_0}{\partial E} \right) dE \right] V$$

low drain bias

# ballistic MOSFET (MB)



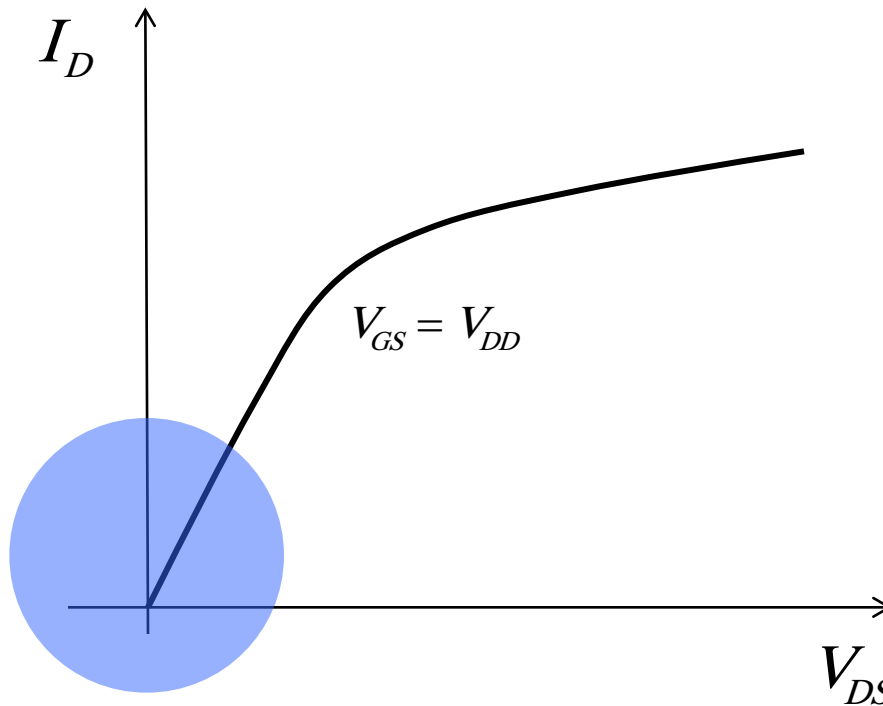
# transmission



$\lambda_0$  is the mean-free-path for backscattering

$$I_D \rightarrow TI_D ?$$

# scattering: linear region



$$T = \frac{\lambda_0}{\lambda_0 + L}$$

$$I_D = WQ_n(V_{GS}, V_{DS}) \frac{v_T}{2k_B T_L / q} V_{DS} \rightarrow I_D = TWQ_n(V_{GS}, V_{DS}) \frac{v_T}{2k_B T_L / q} V_{DS}$$

# example

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For a reasonable carrier mobility and a state-of-the-art channel length of 22 nm (2012), what is  $T$ ?

$$\mu_{\text{eff}} \approx 200 \text{ cm}^2/\text{V}\cdot\text{s}$$

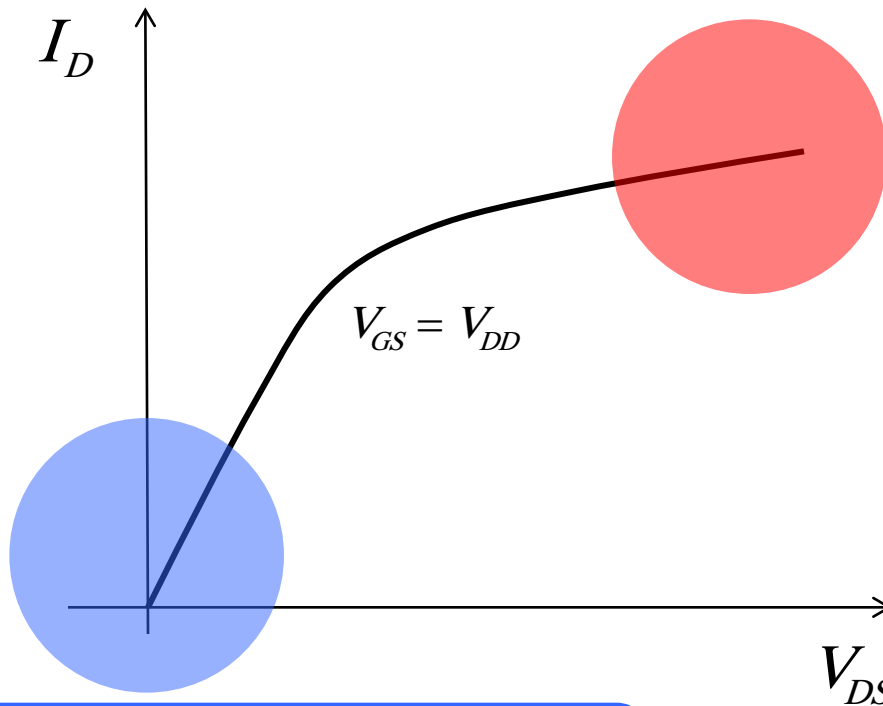
$$\mu_n = \frac{v_T \lambda_0}{2(k_B T/q)}$$

$$\lambda_0 \approx 9 \text{ nm}$$

$$T \approx \frac{\lambda_0}{L + \lambda_0} \approx 0.3$$

# ballistic MOSFET (MB)

$$I_D = TWQ_n(V_{GS}, V_{DS})v_T$$



$$I_D = TWQ_n(V_{GS}, V_{DS}) \frac{v_T}{2k_B T_L / q} V_{DS}$$

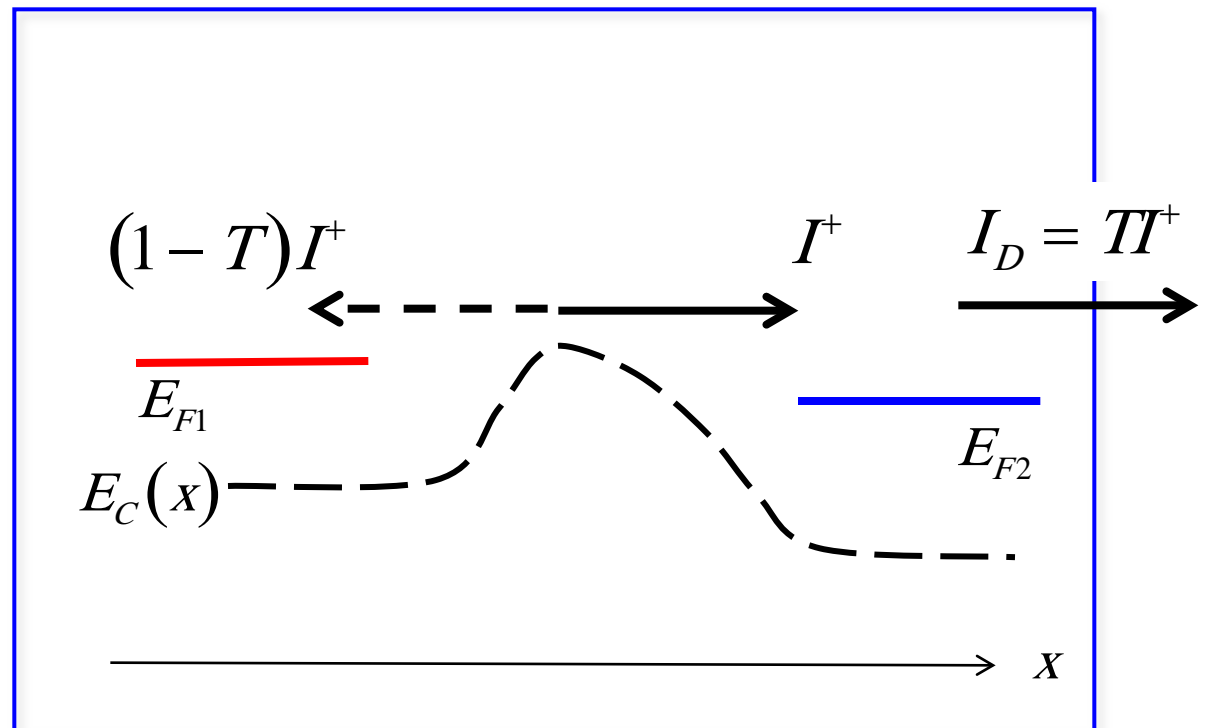
# on-current and transmission

$$Q_n(0) = \frac{I_{BALL}^+}{Wv_T}$$

$$Q_n(0) = \frac{I^+ + (1-T)I^+}{Wv_T}$$

$$I^+ = \frac{I_{BALL}^+}{(2-T)}$$

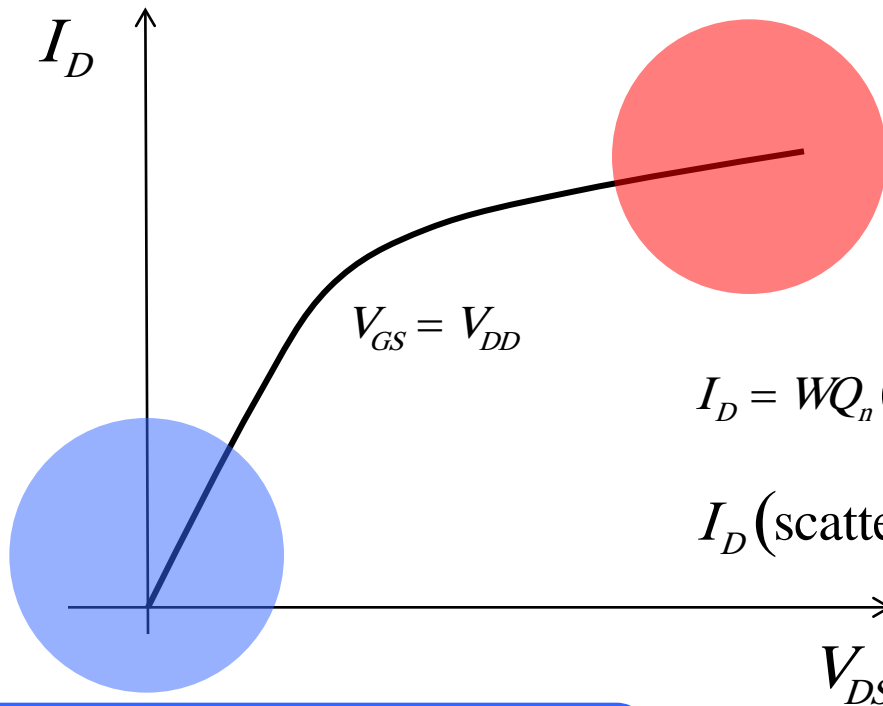
$$I_{ON} = \left( \frac{T}{2-T} \right) I_{BALL}^+$$





# ballistic MOSFET (MB)

$$I_D = \left( \frac{T}{2 - T} \right) W Q_n (V_{GS}, V_{DS}) v_T$$



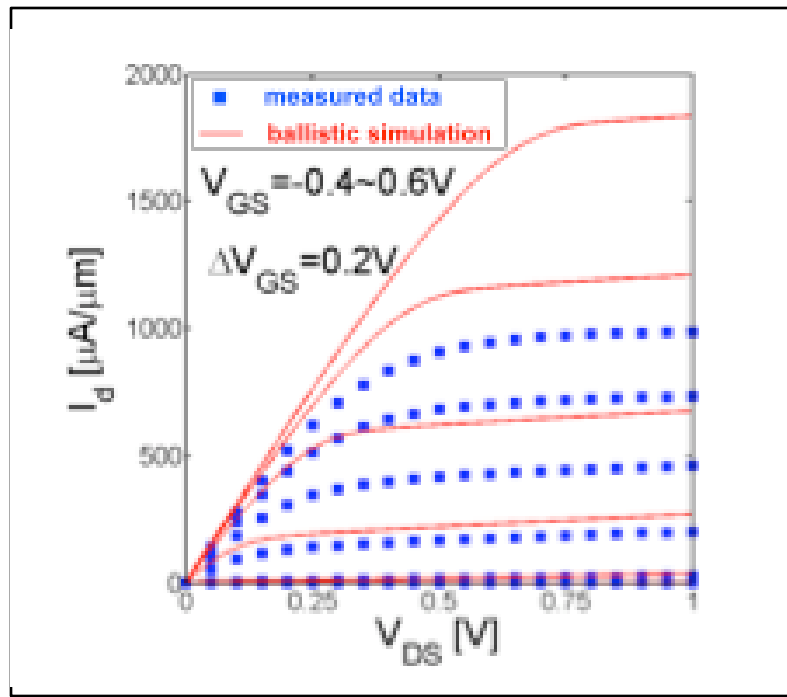
$$I_D = W Q_n (V_{GS}, V_{DS}) T \left( \frac{1 - e^{-qV_{DS}/k_B T}}{(2 - T) + T e^{-qV_{DS}/k_B T}} \right) v_T$$

$$I_D(\text{scattering}) \neq I_D(\text{ballistic})$$

$$I_D = T W Q_n (V_{GS}, V_{DS}) \frac{v_T}{2k_B T_L / q} V_{DS}$$

# recall: ETSOI MOSFET

$$L_G = 40 \text{ nm}$$



$$I_{Dlin} / I_{ballistic} \approx 0.2$$

$$I_{ON} / I_{ballistic} \approx 0.6$$

- Si MOSFETs deliver > one-half of the ballistic on-current. (Similar for the past 15 years.)
- MOSFETs operate closer to the ballistic limit under high  $V_{DS}$ .

A. Majumdar, Z. B. Ren, S. J. Koester, and W. Haensch, "Undoped-Body Extremely Thin SOI MOSFETs With Back Gates," *IEEE Transactions on Electron Devices*, **56**, pp. 2270-2276, 2009.

# transmission under low and high $V_{DS}$

$$I_{Dlin} / I_{ballistic} \approx 0.2$$

$$I_{Dlin} / I_{ballistic} = T_{lin}$$

$$I_{ON} / I_{ballistic} \approx 0.6$$

$$I_{ON} / I_{ballistic} \approx T_{sat} / (2 - T_{sat})$$

$$I_{Dlin} = T W Q_n (V_{GS}, V_{DS}) \frac{v_T}{2 k_B T_L / q} V_{DS}$$

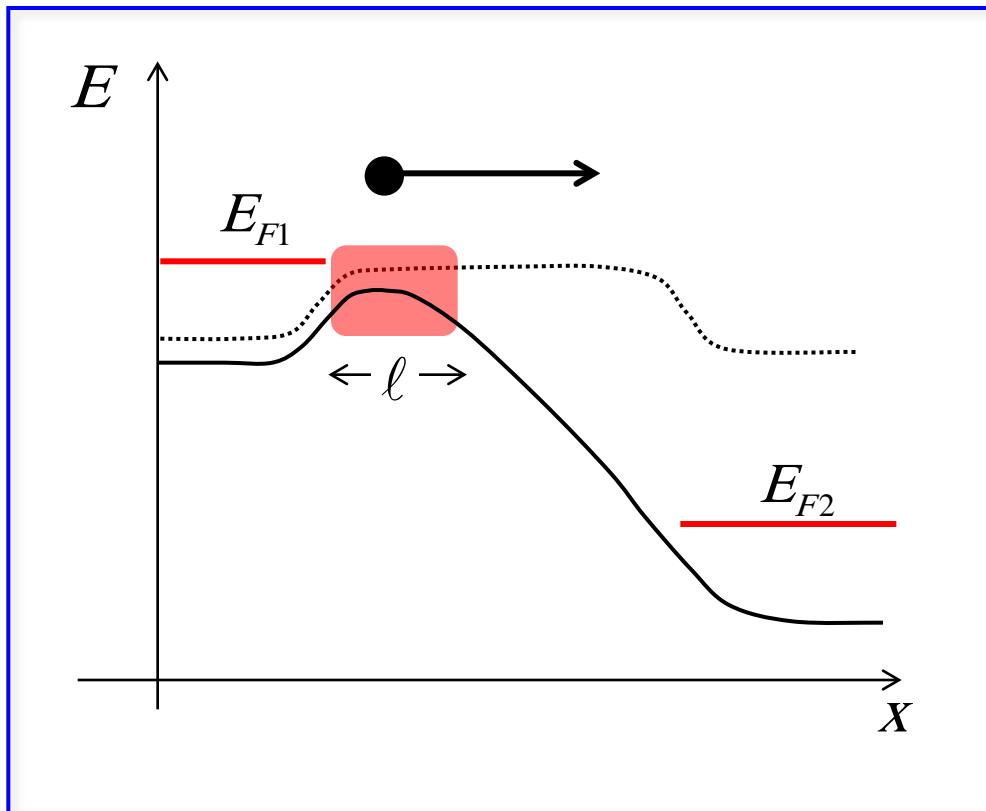
$$I_{ON} = \left( \frac{T}{2 - T} \right) W Q_n (V_{GS}, V_{DS}) v_T$$

$$T_{lin} \approx 0.2$$

$$T_{sat} \approx 0.7$$

**Why?**

# scattering under high $V_{DS}$



$$T_{lin} = \frac{\lambda_0}{\lambda_0 + L}$$

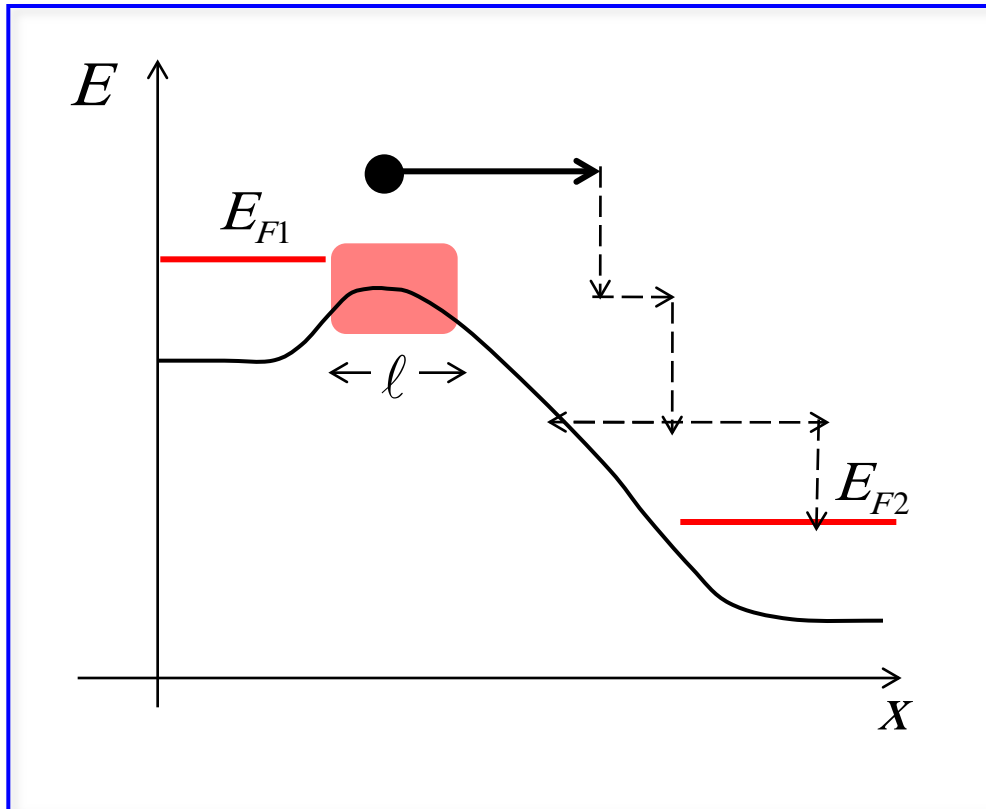
$$L \rightarrow l$$

$$T_{sat} = \frac{\lambda_0}{\lambda_0 + l}$$

$$l \ll L$$

$$T_{sat} > T_{lin}$$

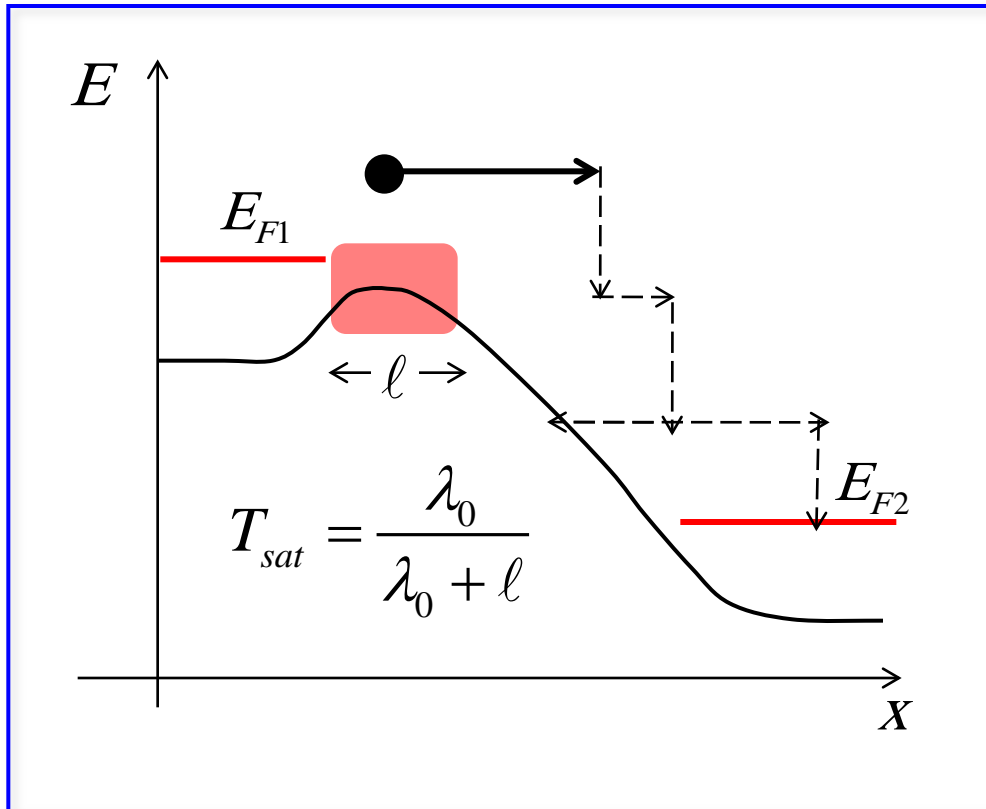
# operation near the “ballistic limit”



$$T_{sat} = \frac{\lambda_0}{\lambda_0 + l}$$

Operation near the ballistic limit current just means that  $T_{sat} \rightarrow 1$ , it does not imply that there is little scattering.

# is mobility relevant at the nanoscale?



- mobility is related to the near-eq. MFP
- backscattering in the critical region is also controlled by the near-eq. MFP.
- mobility determines the on-current
- but the MFP near the drain is very short.

# connection to traditional model: linear

$$I_D = \frac{W}{L} \mu_{eff} C_{inv} (V_{GS} - V_T) V_{DS}$$

$$I_D = T \left( WC_{inv} (V_{GS} - V_T) \frac{v_T}{(2k_B T_L / q)} \right) V_{DS} \quad T = \frac{\lambda_0}{\lambda_0 + L}$$

$$I_D = \frac{W}{L + \lambda_0} \mu_{eff} C_{inv} (V_{GS} - V_T) V_{DS}$$

$$I_D = \frac{W}{L} \mu_{app} C_{inv} (V_{GS} - V_T) V_{DS}$$

$$1/\mu_{app} = 1/\mu_{eff} + 1/\mu_B$$

# connection to traditional model: linear region

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$$I_D = \frac{W}{L} \mu_{eff} C_{inv} (V_{GS} - V_T) V_{DS}$$

$$I_D = \frac{W}{L} \mu_{app} C_{inv} (V_{GS} - V_T) V_{DS}$$

“Mathiessen’s Rule”

$$\frac{1}{\mu_{app}} = \frac{1}{\mu_{eff}} + \frac{1}{\mu_B}$$

$\mu_{eff}$  : measured

$$\mu_B = \frac{v_T L}{2 k_B T_L / q}$$



## connection to traditional model: saturation

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$$I_D = WC_{inv} \nu_{sat} (V_{GS} - V_T)$$

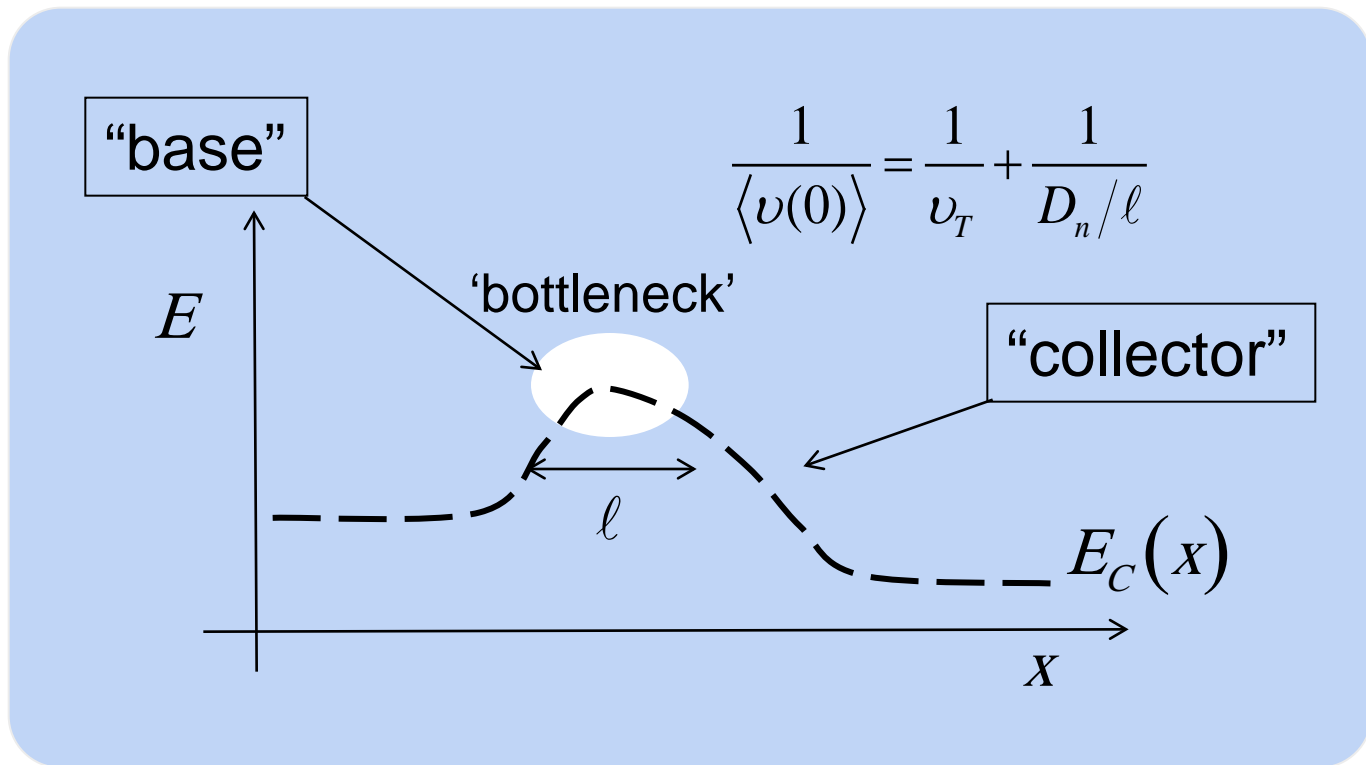
$$I_D = WC_{inv} \nu_T \left( \frac{T}{2 - T} \right) (V_{GS} - V_T) \quad T = \frac{\lambda_0}{\lambda_0 + \ell}$$

$$I_D = W \left[ \frac{1}{\nu_T} + \frac{1}{(D_n/\ell)} \right]^{-1} C_{inv} (V_{GS} - V_T)$$

*how do we interpret this result?*

# current transmission in a MOSFET

$$I_D = W \langle v(0) \rangle C_{inv} (V_{GS} - V_T)$$



# VS model for the MOSFET

$$I_D = W Q_n (V_{GS}, V_{DS}) \langle v(0) \rangle$$

$$\langle v(0) \rangle = \left\{ \frac{V_D / V_{DSAT}}{\left[ 1 + (V_D / V_{DSAT})^\beta \right]^{1/\beta}} \right\} v_{SAT} = F_{SAT}(V_{DS}) v_{SAT}$$

$$V_{DSAT} = \frac{v_{SAT} L}{\mu_{eff}}$$

$$v_{SAT}$$
$$\mu_{eff}$$

# a more physical VS model for a MOSFET

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$$1) \quad \mu_{eff} \rightarrow \mu_{app} \quad \frac{1}{\mu_{app}} = \frac{1}{\mu_{eff}} + \frac{1}{\mu_B}$$

$$2) \quad v_{sat} \rightarrow \left( \frac{T_{sat}}{2 - T_{sat}} \right) v_{inj}^{ball} = \left[ \frac{1}{v_{inj}^{ball}} + \frac{1}{D_n / \ell} \right]^{-1}$$

## wrap-up

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- 1) Scattering reduces the MOSFET current.
- 2) For small  $V_{DS}$ , the effective of scattering is stronger than for high  $V_{DS}$ .
- 3) In the end, the expressions for  $I_D$  look remarkably similar to the traditional expressions, but with parameters that are more physical at the nanoscale.

# references

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For more discussion of scattering in nanoscale MOSFETs, see:

Mark Lundstrom and Jing Guo, *Nanoscale Transistors: Physics, Modeling, and Simulation*, Springer, New York, 2006.