

# Optical Cavities

Cavity provides the feedback system needed for laser operation!

A cavity mode is a field distribution that reproduces itself in relative shape and in relative phase after a round trip.

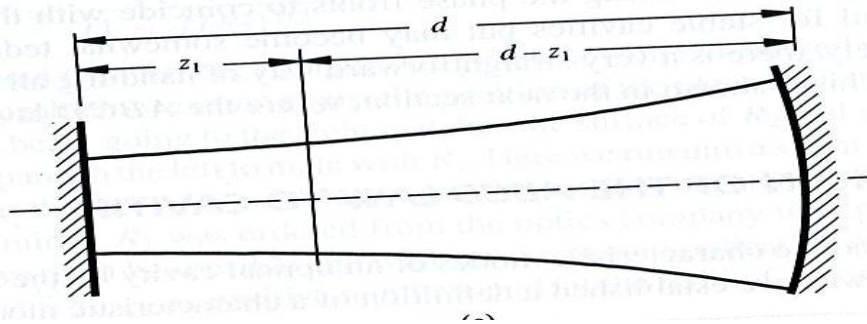
To follow a Gaussian beam through a round trip in a cavity we employ the ABCD law.

1. We assume that the Hermite-Gaussian beams are the characteristic modes of the optical cavity (Mirrors exactly match the surfaces of constant phase of the beams.)

For this require that:

2. The complex beam parameter repeats itself after a round trip!

$$q(z_1 + \text{roundtrip}) = q(z_1)$$



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$$q(z_1) = \frac{Aq(z_1) + B}{Cq(z_1) + D} \quad \begin{pmatrix} A & B \\ C & D \end{pmatrix} \text{ transmission matrix for the unit cell (round trip)}$$

$$Cq^2 + Dq = Aq + B$$

$$B\left(\frac{1}{q}\right)^2 + 2\left(\frac{A-D}{2}\right)\left(\frac{1}{q}\right) - C = 0$$

$$\frac{1}{q} = -\frac{A-D}{2B} \pm \frac{1}{B} \left[ \left(\frac{A-D}{2}\right)^2 + BC \right]^{\frac{1}{2}}$$

$$AD - BC = 1 \quad \frac{A^2}{4} - \frac{AD}{2} + \frac{D^2}{4} + BC \left(1 - \frac{AD}{2} + \frac{AD}{2}\right) = \left(\frac{A+D}{2}\right)^2 - 1$$

$$\frac{1}{q} = -\frac{A-D}{2B} - j \frac{\left[1 - \left(\frac{A+D}{2}\right)^2\right]^{\frac{1}{2}}}{B}$$

$$\frac{1}{q} = \frac{1}{R} - j \frac{\lambda}{\pi w^2}$$

$$R(z_1) = -\frac{2B}{A-D} \quad \frac{\pi w^2(z)}{\lambda} = \frac{B}{\left[1 - \left(\frac{A+D}{2}\right)^2\right]^{\frac{1}{2}}}$$

$$T = \begin{pmatrix} 1 & d+z_1 \\ 0 & 1 \end{pmatrix} \begin{pmatrix} 1 & d-z_1 \\ -\frac{1}{f} & 1 - \frac{d-z_1}{f} \end{pmatrix}$$

If we start the unit cell at the flat mirror ( $z=0$ )

$$T = \begin{pmatrix} 1 & d \\ 0 & 1 \end{pmatrix} \begin{pmatrix} 1 & d \\ -\frac{1}{f} & 1 - \frac{d}{f} \end{pmatrix} = \begin{pmatrix} 1 - \frac{d}{f} & d + d\left(1 - \frac{d}{f}\right) \\ 0 & 1 - \frac{d}{f} \end{pmatrix} \quad (R_L = \infty)$$

$$\frac{\pi w_0^2}{\lambda} = \frac{d + d\left(1 - \frac{d}{f}\right)}{\left[1 - \left(1 - \frac{d}{f}\right)^2\right]^{\frac{1}{2}}} = \frac{2d\left(1 - \frac{d}{R}\right)}{\left[4\left(\frac{d}{R}\right)\left(1 - \frac{d}{R}\right)\right]^{\frac{1}{2}}} \quad f = \frac{R}{2}$$

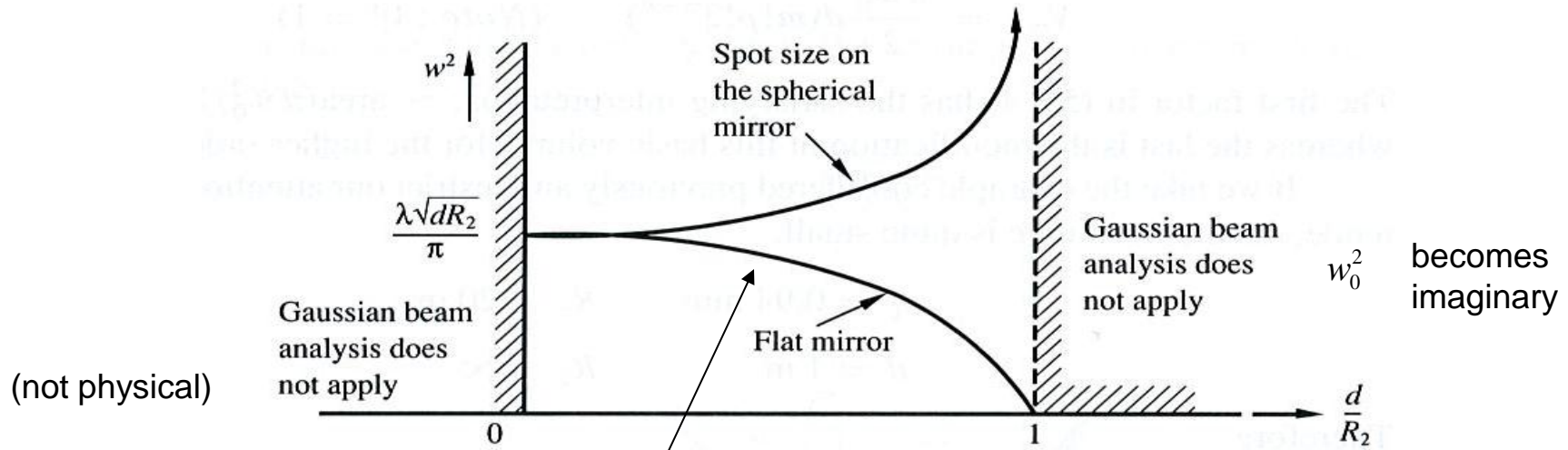
$$\boxed{\frac{\pi w_0^2}{\lambda} = (dR)^{\frac{1}{2}} \left(1 - \frac{d}{R}\right)^{\frac{1}{2}}}$$

$$w^2(z) = w_0^2 \left[ 1 + \left( \frac{z}{z_0} \right)^2 \right] \quad z_0^2 - \text{confocal parameter}$$

$$z_0^2 = \left( \frac{\pi w_0^2}{\lambda} \right)^2$$

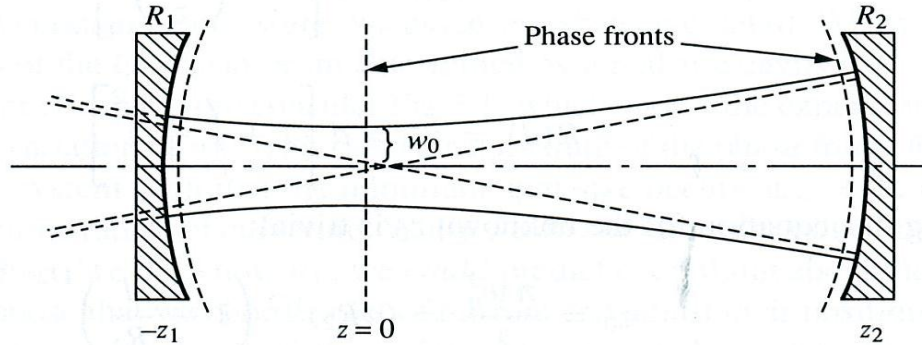
On a spherical mirror

$$\frac{\pi w^2(d)}{\lambda} = (dR)^{\frac{1}{2}} \left( 1 - \frac{d}{R} \right)^{\frac{1}{2}} \left[ 1 + \frac{d^2}{dR \left( 1 - \frac{d}{R} \right)} \right] = \frac{(dR)^{\frac{1}{2}}}{\left( 1 - \frac{d}{R} \right)^{\frac{1}{2}}}$$



$$\frac{\pi w_0^2}{\lambda} = (dR)^{\frac{1}{2}} \left( 1 - \frac{d}{R} \right)^{\frac{1}{2}}$$

The requirement that mirrors must match the surface of the constant phase (so that a cavity mode is excited) allows one to find where the mirrors with  $R_1$  &  $R_2$  should be placed, in general case.



$R_1, R_2$  are given,  
we look for  $z_1, z_2$

$$1) \quad z_1 + z_2 = d \quad 2) \quad R(z_2) = R_2 = z_2 \left[ 1 + \left( \frac{z_0}{z_2} \right)^2 \right] \quad 3) \quad R(z_1) = -R_1 = -z_1 \left[ 1 + \left( \frac{z_0}{z_1} \right)^2 \right]$$

The wave front on the left at  $z=0$  has a (mathematically) negative radius of curvature, but we know that the mirror  $R_1$  has positive (focusing) properties. We treat all distances  $z_1, z_2$  as positive numbers and let the radii of curvature of the mirrors carry their own sign.

Solving (involving algebra)....

$$z_0^2 = \left( \frac{\pi w_0^2}{\lambda} \right)^2 = \frac{d(R_1 - d)(R_2 - d)(R_1 + R_2 - d)}{(R_1 + R_2 - 2d)^2}$$

$$z_1 = \frac{d(R_2 - d)}{R_1 + R_2 - 2d} \quad z_2 = \frac{d(R_1 - d)}{R_1 + R_2 - 2d}$$

$$R_1 = \infty \quad z_0^2 = d(R_2 - d) \Rightarrow dR_2 \left( 1 - \frac{d}{R_2} \right) \quad z_1 = 0; z_2 = d \quad \frac{\pi w_0^2}{\lambda} = (dR)^{\frac{1}{2}} \left( 1 - \frac{d}{R} \right)^{\frac{1}{2}}$$

Same result

## Summary

- 1) Postulate that Hermite-Gaussian Beams are the normal modes for the cavity.
- 2) Formulate an equivalent transmission system for this cavity showing one round trip. Identify a unit cell.
- 3) Force the complex beam parameter to transform into itself after a round trip by use of the ABCD law.
- 4) Evaluate R and w using:

$$R(z) = -\frac{2B}{A-D} \quad \frac{\pi w^2(z)}{\lambda} = \frac{B}{\left[1 - \left(\frac{A+D}{2}\right)^2\right]^{\frac{1}{2}}}$$

The theory applies for Stable cavity only!

# Mode Volume

$$E_0^2 V = \int_0^d dz \int_{-\infty}^{\infty} dx \int_{-\infty}^{\infty} dy E(x, y, z) E^*(x, y, z)$$

$$E_0^2 V_{m,n} = E_0^2 \int_0^d \frac{w_0^2}{w^2(z)} dz \int_{-\infty}^{\infty} dx \int_{-\infty}^{\infty} dy H_m^2\left(\frac{\sqrt{2}x}{w}\right) e^{-\frac{2x^2}{w^2}} \times H_n^2\left(\frac{\sqrt{2}y}{w}\right) e^{-\frac{2y^2}{w^2}}$$

use  $u = \frac{\sqrt{2}x}{w}$  or  $\frac{\sqrt{2}y}{w}$

$$V_{m,n} = \int_0^d \frac{w_0^2}{2} dz \left[ \int_{-\infty}^{\infty} H_m^2(u) e^{-u^2} du \right] \left[ \int_{-\infty}^{\infty} H_n^2(u) e^{-u^2} du \right]$$

$\underbrace{\hspace{10em}}_{2^n n! \sqrt{\pi}}$

$$V_{m,n} = \underbrace{\frac{\pi w_0^2}{2}}_A \underbrace{d(m!n!2^{m+n})}_B$$

Area x length

Modification for  
high-order modes

Example (textbook)

$$w_0 = 0.94 \text{ mm} \quad R_2 = 20 \text{ m}$$

$$d = 1 \text{ m} \quad R_1 = \infty$$

$$\begin{aligned} \rightarrow V_{0,0} &= 1.38 \text{ cm}^3 \\ &= \frac{\pi w_0^2}{2} \times d \end{aligned}$$

He-Ne laser

$$P = 0.1 \text{ torr (of neon)}$$

Each atom is excited (by the gas discharge) and thus producing a photon at  $632.8 \mu\text{m}$ , say, 10 times per second.

$$\text{Energy per photon } h\nu = \frac{hc}{\lambda} = 3.14 \times 10^{-19} \text{ J} = 1.96 \text{ eV}$$

$$\times \# \text{ of Ne atoms} = 0.1(3.54 \times 10^{16})V_{0,0} = 4.88 \times 10^{15}$$

$$\times (\text{average excitation per atom} = \text{average emission per atom}) = 10 \text{ sec}^{-1}$$

$$= \text{Power} = 15.3 \text{ mW}$$

typical!