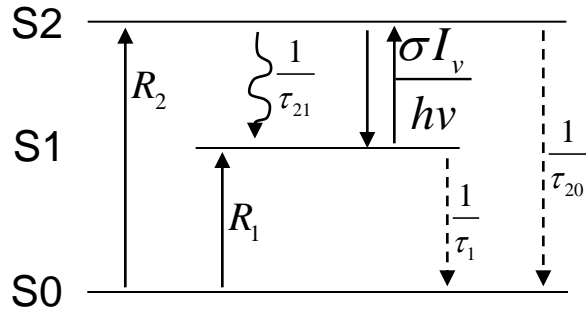


# Power in Laser Oscillation

pump power  $\rightarrow$  laser power

Consider 3-level system and use rate equations (see section 8.3)



$$\frac{g_2}{g_1} = 1$$

use gain coefficient:  $\gamma(\nu) = (N_2 - N_1)\sigma(\nu)$

and solution for  $(N_2 - N_1)$

$$(N_2 - N_1) = \frac{\gamma(\nu)}{\sigma(\nu)} = \frac{R_2 \tau_2 \left(1 - \frac{\tau_1}{\tau_{21}}\right) R_1 \tau_1}{1 + \left(\tau_1 + \tau_2 - \frac{\tau_1 \tau_2}{\tau_{21}}\right) \frac{\sigma(\nu) I_\nu}{h\nu}} \quad (8.3.11c)$$

For  $\tau_{20}^{-1} \ll \tau_{21}^{-1}$ ,  $\tau_2 = \tau_{21}$   $\Rightarrow$

$$(N_2 - N_1) = \frac{\gamma(\nu)}{\sigma(\nu)} = \frac{R_2 \tau_2 \left(1 - \frac{\tau_1}{\tau_2}\right) - R_1 \tau_1}{1 + \tau_2 \frac{\sigma(\nu) I_\nu}{h\nu}}$$

$$(N_2 - N_1) = \frac{\gamma(\nu)}{\sigma(\nu)} = \frac{R_2 \left[1 - \frac{\tau_1}{\tau_2} \left(1 + \frac{R_1}{R_2}\right)\right]}{\frac{1}{\tau_2} + \frac{\sigma(\nu) I_\nu}{h\nu}} = \frac{R}{\frac{1}{\tau_2} + \frac{\sigma(\nu) I_\nu}{h\nu}}$$

effective pump rate  $\swarrow$

$\swarrow$  total decay rate

$$(N_2 - N_1) \left( \frac{1}{\tau_2} + \frac{\sigma(\nu) I_\nu}{h\nu} \right) = R \quad (1)$$

$$\left( \frac{\sigma(\nu)}{h\nu} \right) I_\nu = \left( \frac{R}{N_2 - N_1} - \frac{1}{\tau_2} \right)$$

Below threshold  $I_\nu = 0 \quad \Rightarrow N_2 - N_1 = R\tau_2$

Continue to increase pump until threshold,  $N_t$

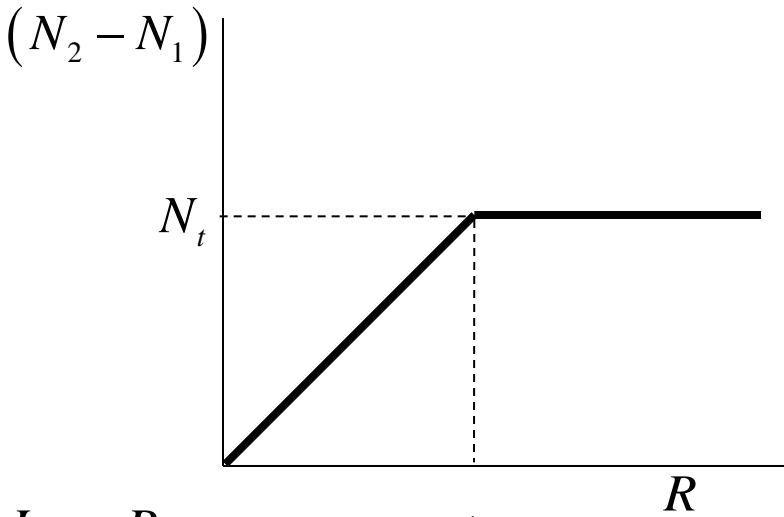
$$N_t = \left( N_2 - \frac{g_2}{g_1} N_1 \right)_{\text{threshold}} = \frac{\gamma_{th}(\nu)}{\sigma(\nu)}$$

$$\gamma_{th} = -\frac{1}{2\ell} \ln R_1 R_2 = \frac{1}{\frac{c}{n} \tau_p} \quad (\text{see section 8.8})$$

↑  
reflectance

$$\sigma = A_{21} \frac{\lambda^2}{8\pi} g(\nu)$$

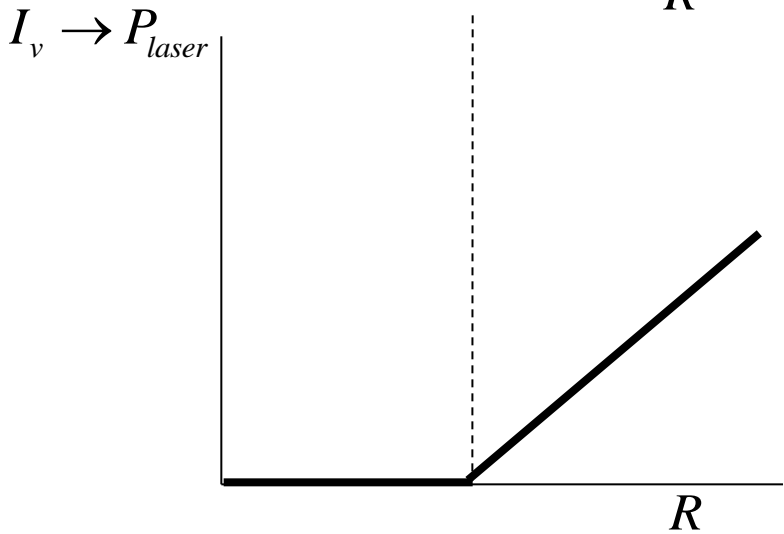
Increasing  $R$  beyond  $N_t/\tau_2$  cannot increase  $(N_2 - N_1)$



$$N_t = \frac{R}{\frac{1}{\tau_2} - \frac{\sigma I_v}{h\nu}} \quad \text{see (1)}$$

- Describes how  $I_v$  increases with  $R$

$$\frac{\sigma I_v}{h\nu} N_t = R - \frac{N_t}{\tau_2} \quad \text{for } R \geq \frac{N_t}{\tau_2}$$



$$P = \left( N_t V_g \right) \left( \frac{\sigma I_v}{h\nu} \right) h\nu$$

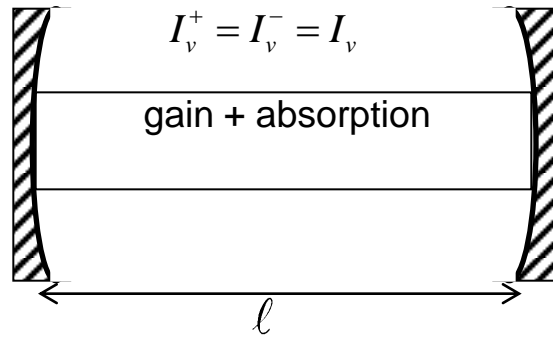
total # of excess Excited atoms      Rate of stimulated emission per atom      energy per photon

$$\frac{P}{V_g h\nu} = R - \frac{N_t}{\tau_2}$$

photon emission rate density      pumping rate density      spontaneous emission rate density

# Optimum Output Coupling for Optical Resonator

Homogeneously broadened. Standing wave.



$$\gamma = \frac{\gamma_0}{1 + (I_v^+ + I_v^-)/I_S}$$

Actually, there is interference between the fields associated with  $I_v^+$  and  $I_v^-$ , which leads to a spatial variation of the saturation. This is complicated, and ignored here.

Optimum coupling: if transmission is too large, losses exceed gain – no generation. If transmission is too small – little goes out.

$$I_v = \frac{I_S}{2} \left( \frac{\gamma_0}{\gamma} - 1 \right) = \frac{I_S}{2} \left( \frac{\gamma_0}{\gamma_t} - 1 \right)$$

clamped at  $\gamma_t$

$$e^{2\gamma_t \ell} \underbrace{R_1 R_2}_{\text{fraction retained after round trip in the cavity}} e^{-2\alpha \ell} = 1 \quad \text{- Threshold condition}$$

fraction retained after round trip in the cavity

$$R_1 R_2 e^{-2\alpha \ell} = 1 - L = S \quad (\text{L-fraction lost in round trip})$$

$$L = 1 - R_1 R_2 e^{-2\alpha \ell}$$

$$(L \ll 1) \Rightarrow 1 - L = e^{-2\gamma_t \ell} \simeq 1 - 2\gamma_t \ell \quad \text{or} \quad 2\gamma_t \ell = -\ln(1 - L) \simeq L$$

$$I_v = \frac{I_s}{2} \left( \frac{\gamma_0}{\gamma_t} - 1 \right)$$

$$2\gamma_t \ell \approx L$$

$$I_v = \frac{I_s}{2} \left( \frac{2\gamma_0 \ell}{L} - 1 \right)$$

$$g_0 \approx \gamma_0 \ell$$

small signal gain/pass

$$\nearrow I_v = \frac{I_s}{2} \left( \underbrace{\frac{2g_0}{L}}_{\text{total losses}} - 1 \right)$$

intensity circulated  
Inside cavity

total  
losses

sum of all  
other losses  
(abs.)

output mirror  
transmission

$$L = L' + T$$

$$L' = L - T \quad T_2 = 1 - R_2$$

$$L = 1 - R_1 R_2 e^{-2\alpha \ell}$$

- transmission through  $M_2$

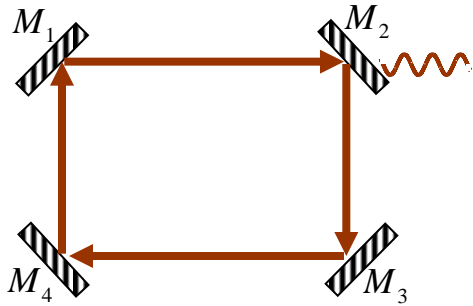
In the text  
 $L' \rightarrow L$

$$I_{out} = I_v T = \frac{I_s}{2} \left( \frac{2g_0}{L' + T} - 1 \right) T$$

- standing wave

$T = T_2$  if  $M_2$  provides  
transmission

for ring lasers:



(for ring cavity)

$$I_{out} = I_S \left( \frac{g_0}{L' + T} - 1 \right) T$$

$T = T_2$  if  $M_2$  provides transmission

Find  $T_{opt}$  :

$$\left[ T \times \left( \frac{g_0}{L' + T} - 1 \right) \right]' = \frac{g_0}{L' + T} - 1 - T \frac{g_0}{(L' + T)^2} = \frac{g_0(L' + T) - g_0 T}{(L' + T)^2} - 1 = 0$$

$$g_0 L' = (L' + T)^2$$

$$T_{opt} = \sqrt{g_0 L'} - L'$$

Useful power output  $P_0$  vs mirror transmission  $T$  for various values of internal loss  $L'$  in He-Ne 632 nm laser.

