

## Section 27

# Heterojunction Bipolar Transistor

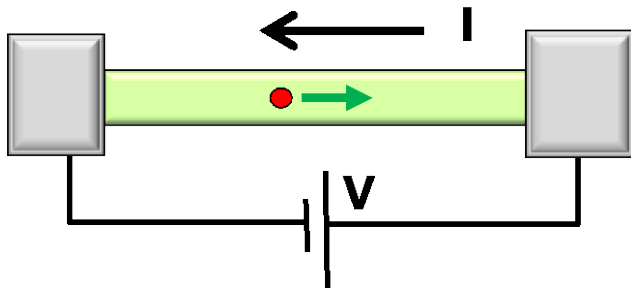
## 27.2 Heterojunction Equilibrium Solution

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School of Electrical and  
Computer Engineering

# Section 27 Heterojunction Bipolar Transistor



$$I = G \times V$$

$$= q \times n \times v \times A$$

↑ charge density   
 ↑ velocity   
 area

$$\frac{n_{i,B}^2}{n_{i,E}^2} = \frac{N_{C,B} N_{V,B} e^{-E_{g,B}\beta}}{N_{C,E} N_{V,E} e^{-E_{g,E}\beta}} \approx e^{(E_{g,E} - E_{g,B})\beta}$$

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$$\nabla \cdot D = q(p - n + N_D^+ - N_A^-) \quad \beta_{DC} = \frac{I_C}{I_B}$$

$$\beta_{poly,ballistic} \rightarrow \frac{n_{i,B}}{n_{i,E}^2} \times \frac{N_E}{N_B} \times \frac{v_{th}}{v_s}$$

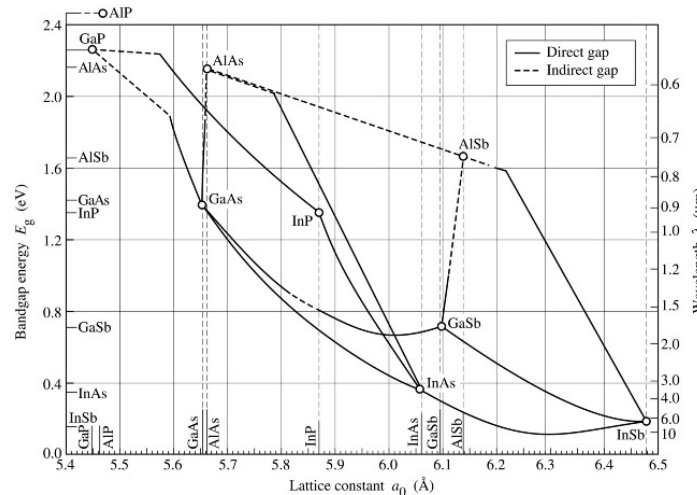
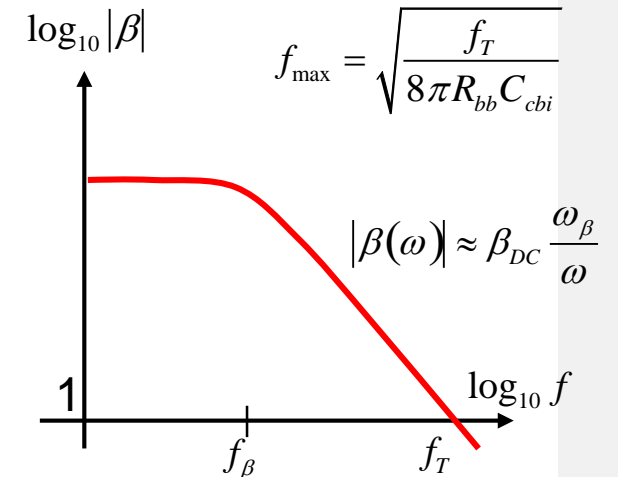


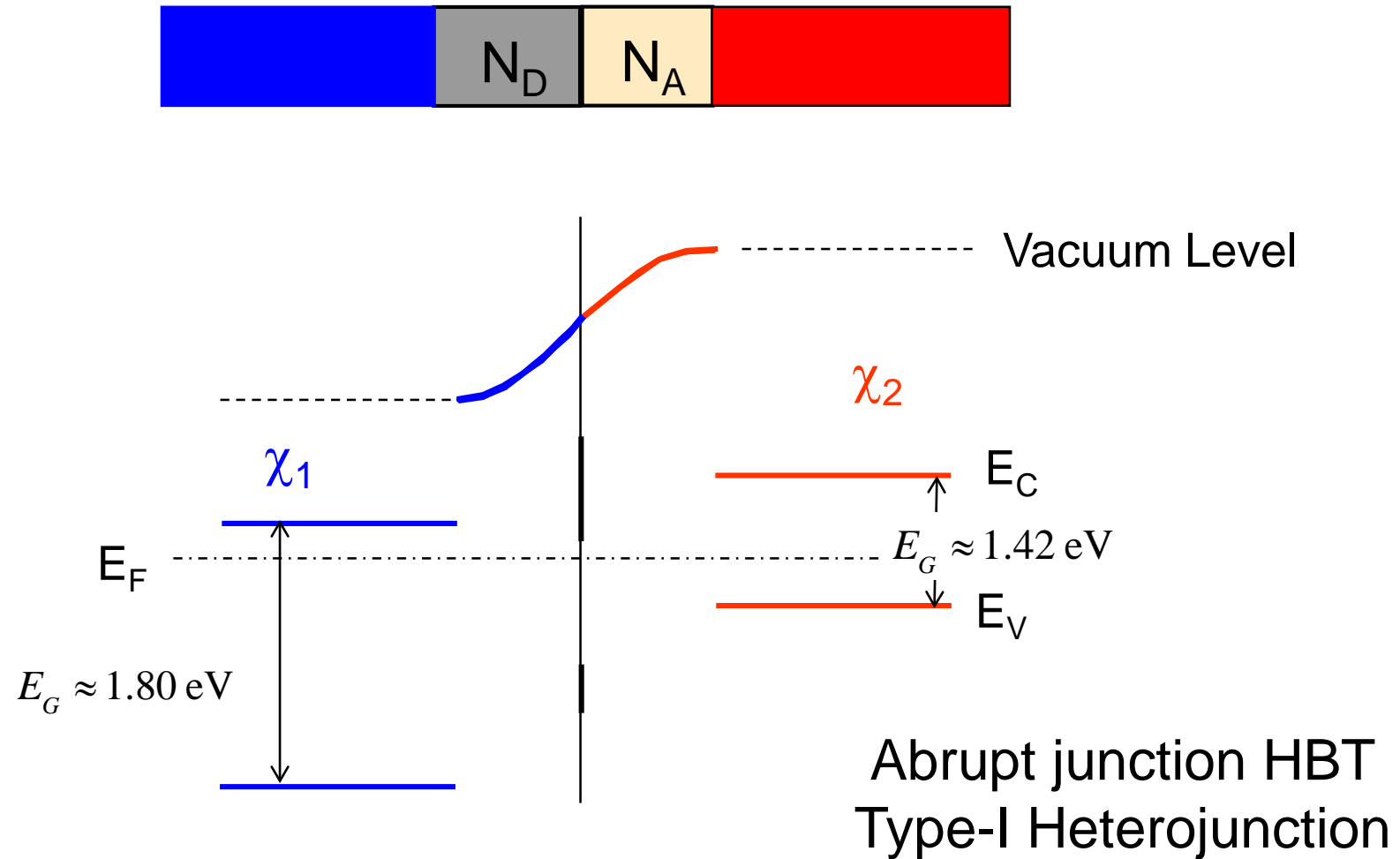
Fig. 7.6. Bandgap energy and lattice constant of various III-V semiconductors at room temperature (adopted from Tien, 1988).



Mark Lundstrom, "Heterostructure Fundamentals," Purdue University, 1995.

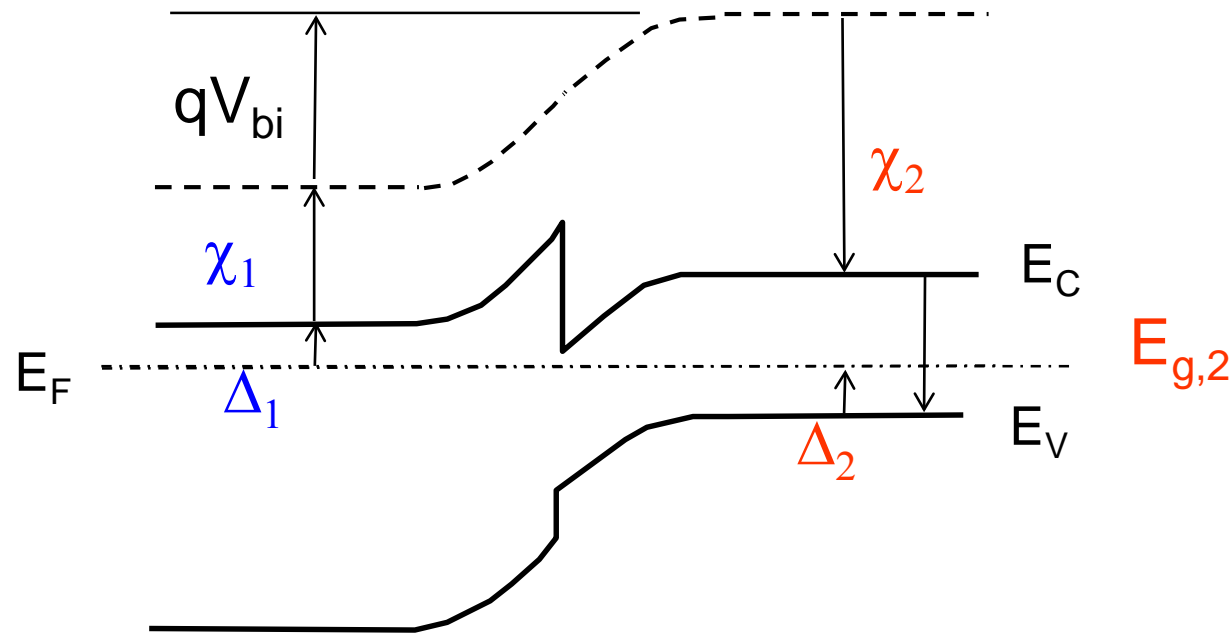
Herbert Kroemer, "Heterostructure bipolar transistors and integrated circuits," Proc. *IEEE*, **70**, pp. 13-25, 1982.

# N-Al<sub>0.3</sub>Ga<sub>0.7</sub>As: p-GaAs (Type-I Heterojunction)



# Built-in Potential: Boundary Condition @Infinity

$$\Delta_1 + \chi_1 + qV_{bi} = E_{g,2} - \Delta_2 + \chi_2$$



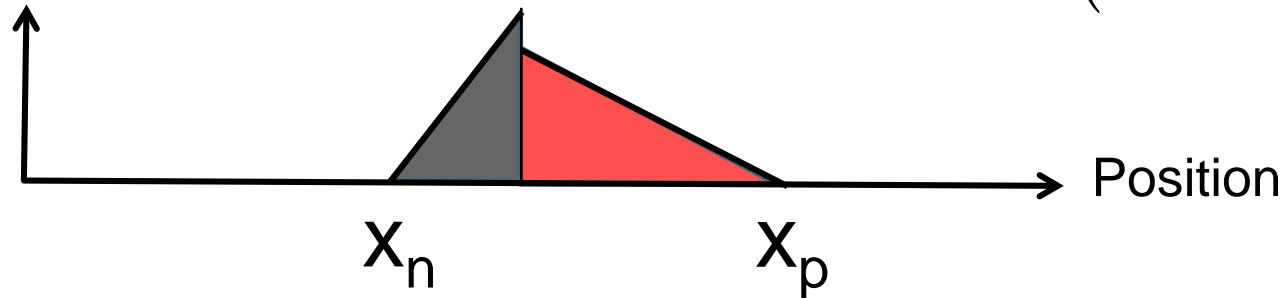
$$qV_{bi} = E_{g,2} - \Delta_2 - \Delta_1 + \chi_2 - \chi_1$$

$$= k_B T \ln \frac{N_A N_D}{N_{V,2} N_{C,1}} e^{-E_{g,2}/k_B T} + (\chi_2 - \chi_1)$$

# Interface Boundary Conditions

E-field

$$\nabla \cdot D = q(p - n + N_D^+ - N_A^-)$$



$$D(0^-) = D(0^+)$$

Displacement  $D$   
is continuous

$$\kappa_1 \epsilon_0 E(0^-) = \kappa_2 \epsilon_0 E(0^+)$$

$$\kappa_1 \epsilon_0 \left. \frac{dV}{dx} \right|_{0^-} = \kappa_2 \epsilon_0 \left. \frac{dV}{dx} \right|_{0^+}$$

$D$  is not continuous  
if there is a surface charge

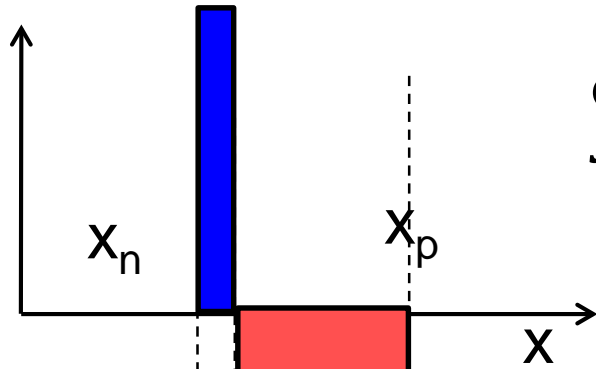
$$\rho_s = \rho_0 \delta(x_0)$$

$$\frac{dD}{dx} = \rho_0 \delta(x_0) + q(p(x) - n(x) + N_D^+(x) - N_A^-(x))$$

$$\oint \vec{D} \cdot d\vec{S} = \rho_{encl} \quad \text{Zero field at edges: } \Rightarrow D(x_n) = D(x_p) \Rightarrow \rho_{encl} = 0$$

# Analytical Solution for Heterojunctions

Charge



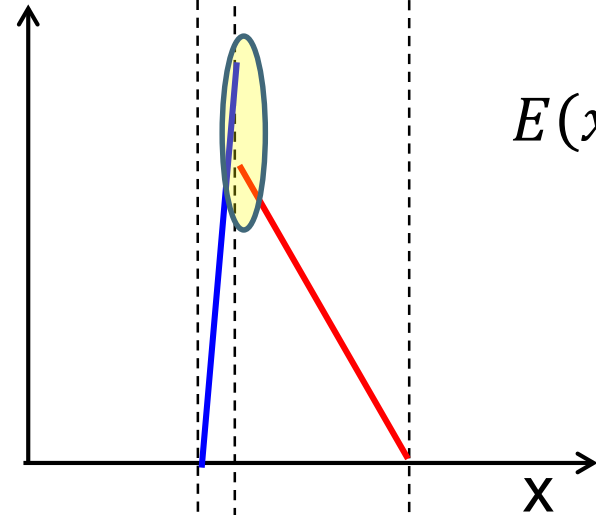
$$\oint \vec{D} dS = \rho_{encl}$$

Zero field at edges:

$$\Rightarrow D(x_n) = D(x_p) \Rightarrow \rho_{encl} = 0$$

$$\Rightarrow N_D x_n = N_A x_p \quad \text{Charge continuity (independent of } k)$$

E-field

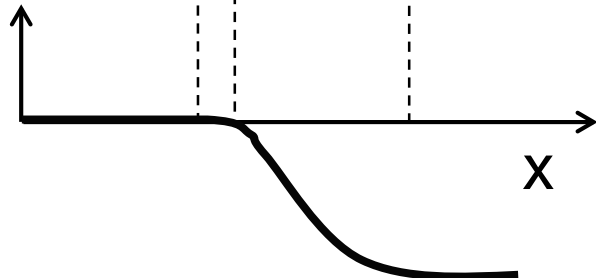


$$E(x) = \int_{-x_n}^x \frac{qN_D}{k_{S,E}\epsilon_0} dx'$$

$$E(0^-) = \frac{qN_D x_n}{k_{S,E}\epsilon_0}$$

$$E(0^+) = \frac{qN_A x_p}{k_{S,B}\epsilon_0}$$

Potential

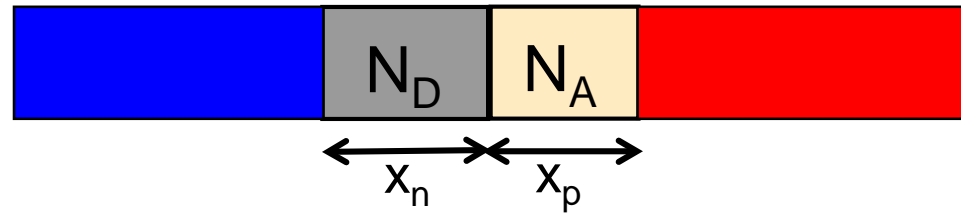


$$V(x) = - \int_{-x_n}^x E(x) dx'$$

$$V_{bi} = \frac{E(0^-)x_n}{2} + \frac{E(0^+)x_p}{2}$$

$$= \frac{qN_D x_n^2}{2k_{S,E}\epsilon_0} + \frac{qN_A x_p^2}{2k_{S,B}\epsilon_0}$$

# Base Emitter Depletion Region



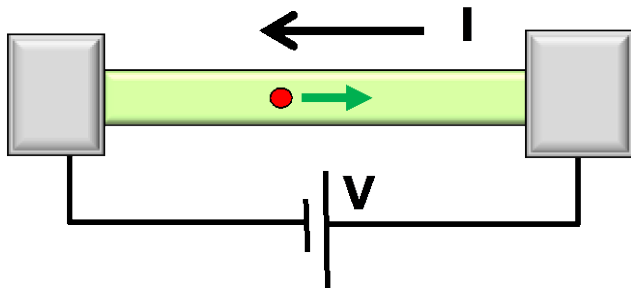
$$N_E x_{n, BE} = N_B x_{p, BE}$$

$$V_{bi} = \frac{q N_E x_{n, BE}^2}{2 \kappa_{s, E} \epsilon_0} + \frac{q N_B x_{p, BE}^2}{2 \kappa_{s, B} \epsilon_0}$$

$$x_n = \sqrt{\frac{2 \epsilon_0}{q} \frac{\kappa_{s, E} \kappa_{s, B} N_B}{N_E (\kappa_{s, E} N_B + \kappa_{s, B} N_E)}} V_{bi}$$

$$x_p = \sqrt{\frac{2 \epsilon_0}{q} \frac{\kappa_{s, E} \kappa_{s, B} N_E}{N_B (\kappa_{s, E} N_B + \kappa_{s, B} N_E)}} V_{bi}$$

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$$\nabla \cdot D = q(p - n + N_D^+ - N_A^-)$$

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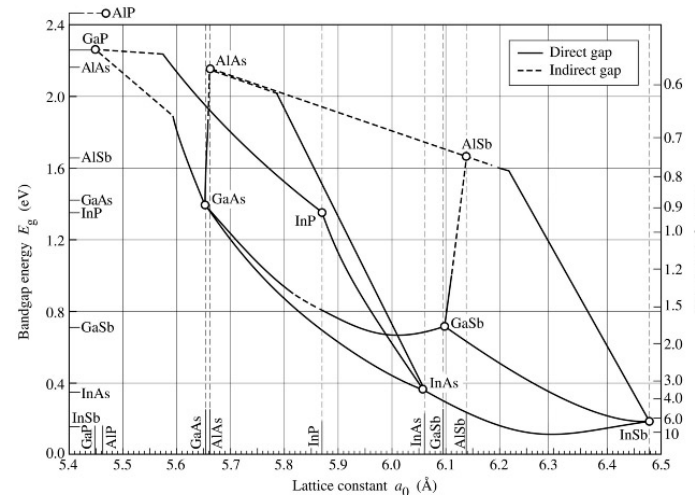
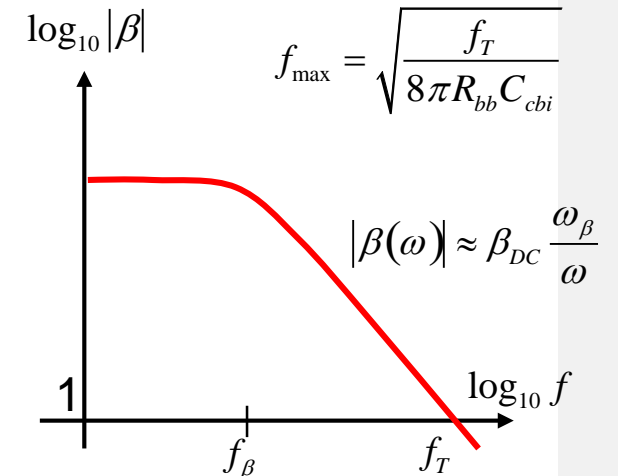


Fig. 7.6. Bandgap energy and lattice constant of various III-V semiconductors at room temperature (adopted from Tien, 1988).

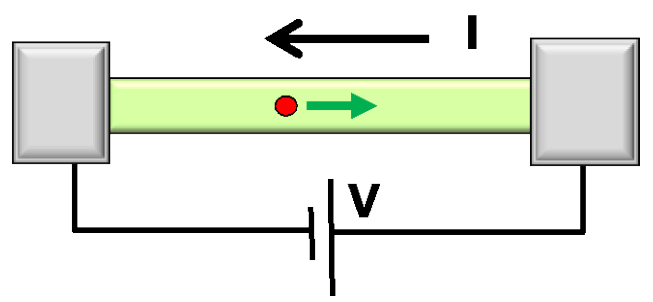


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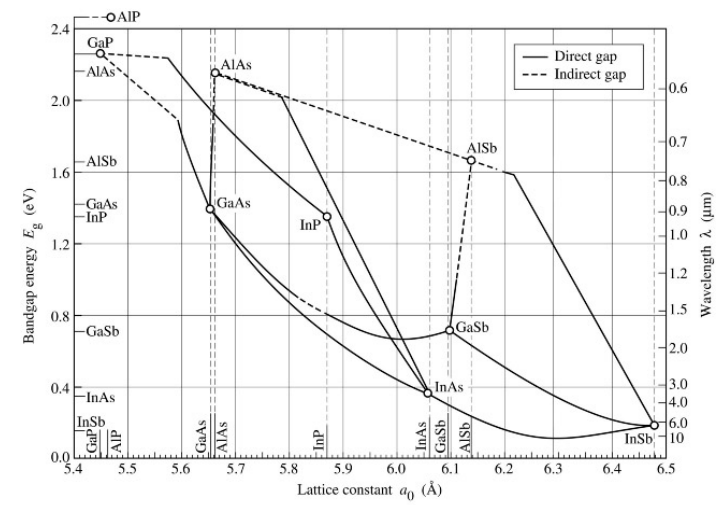
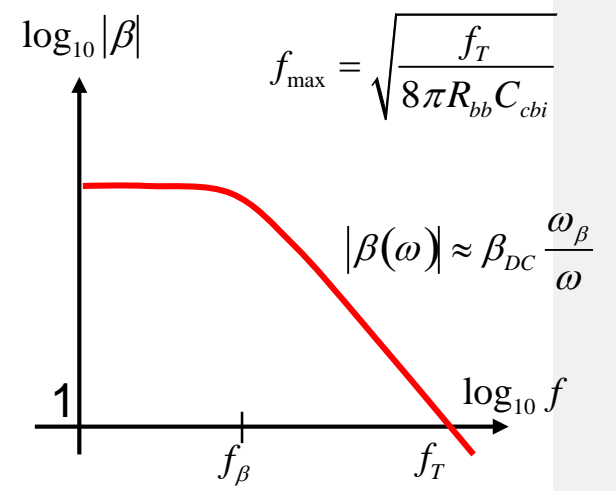


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