

## Section 18 Semiconductor Equations

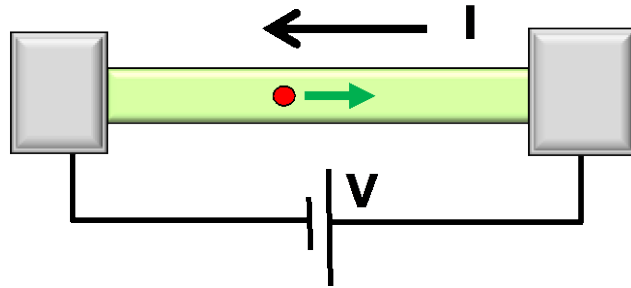
### 18.3 Numerical Solutions

Gerhard Klimeck  
[gekco@purdue.edu](mailto:gekco@purdue.edu)



School of Electrical and  
Computer Engineering

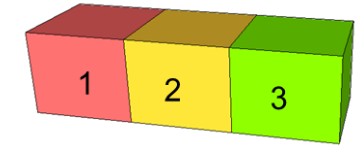
# Section 18 Semiconductor Equations



$$I = G \times V$$

$$= q \times n \times v \times A$$

↑ charge density    ↑ velocity    area



- 18.1 Continuity Equations
- 18.2 Analytical Solutions (Strategy & Examples)
- 18.3 Numerical Solutions
  - » Basic Transport Equations
  - » Gridding and finite differences
  - » Discretizing equations and boundary conditions
  - » Iterative Solution Approach



$$\frac{\partial n}{\partial t} = \frac{1}{q} \nabla \cdot \mathbf{J}_N - r_N + g_N$$

$$\mathbf{J}_N = qn\mu_N E + qD_N \nabla n$$

$$\frac{\partial p}{\partial t} = -\frac{1}{q} \nabla \cdot \mathbf{J}_P + g_P - r_P$$

$$\mathbf{J}_P = qp\mu_P E - qD_P \nabla p$$

$$\nabla \cdot D = q(p - n + N_D^+ - N_A^-)$$

Vid

Vid

Vid

# Preface

- The 5 equations we derived the past few lectures have been used for the longest time in the industry and in academia to understand carrier transport in devices.
- It is useful to know the essentials of how these equations are *implemented* on a modern computer so that one understands some of the finer details involved in creating tools that simulate these phenomena.
- Understanding some of these details helps one become a ‘power user’ of the simulation tools that implement the physics. One also understands the limitations re. numerical issues and applicability ranges of results.

$$\frac{\partial n}{\partial t} = \frac{1}{q} \nabla \cdot \mathbf{J}_N - r_N + g_N$$

$$\mathbf{J}_N = qn\mu_N \mathbf{E} + qD_N \nabla n$$

$$\frac{\partial p}{\partial t} = -\frac{1}{q} \nabla \cdot \mathbf{J}_P + g_P - r_P$$

$$\mathbf{J}_P = qp\mu_P \mathbf{E} - qD_P \nabla p$$

$$\nabla \cdot \mathbf{D} = q(p - n + N_D^+ - N_A^-)$$

# Equations to be solved

Band-diagram  $\longrightarrow$

Diffusion approximation,  
Minority carrier transport,  
Ambipolar transport

$$\nabla \cdot D = q(p - n + N_D^+ - N_A^-)$$

$$\frac{\partial n}{\partial t} = \frac{1}{q} \nabla \cdot \mathbf{J}_N - r_N + g_N$$

$$\mathbf{J}_N = qn\mu_N E + qD_N \nabla n$$

$$\frac{\partial p}{\partial t} = -\frac{1}{q} \nabla \cdot \mathbf{J}_P + g_P - r_P$$

$$\mathbf{J}_P = qp\mu_P E - qD_P \nabla p$$

# 1) The Semiconductor Equations

Conservation Laws: not specific to a particular problem - Universal

$$\nabla \cdot \vec{D} = \rho$$

$$\nabla \cdot \left( \vec{J}_n / -q \right) = \left( g_N - r_N \right)$$

$$\nabla \cdot \left( \vec{J}_p / q \right) = \left( g_P - r_P \right)$$

(steady-state)

Constitutive relations: specific to problem at hand – reflect physics of the problem

$$\vec{D} = \kappa \epsilon_0 \vec{E} = -\kappa \epsilon_0 \vec{\nabla} V$$

$$\rho = q(p - n + N_D^+ - N_A^-)$$

$$\vec{J}_n = nq\mu_n \vec{E} + qD_n \vec{\nabla} n$$

$$\vec{J}_p = pq\mu_p \vec{E} - qD_p \vec{\nabla} p$$

$$g_{N,P} = f(n, p) \text{ etc.}$$

$$\nabla \cdot D = q(p - n + N_D^+ - N_A^-)$$

$$\frac{\partial n}{\partial t} = \frac{1}{q} \nabla \cdot \mathbf{J}_N - r_N + g_N$$

$$\frac{\partial p}{\partial t} = -\frac{1}{q} \nabla \cdot \mathbf{J}_P + g_P - r_P$$

$$\mathbf{J}_N = qn\mu_N E + qD_N \nabla n$$

$$\mathbf{J}_P = qp\mu_P E - qD_P \nabla p$$

# 1) The Mathematical Problem

## The “Semiconductor Equations”

$$\nabla \cdot \vec{D} = \rho$$

$$\nabla \cdot \left( \vec{J}_n / -q \right) = \left( S_n - r_n \right)$$

$$\nabla \cdot \left( \vec{J}_p / q \right) = \left( S_p - r_p \right)$$

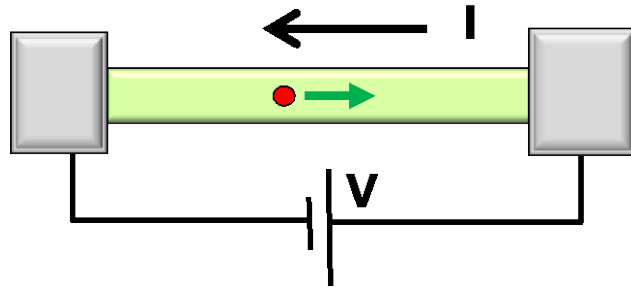
3 coupled, nonlinear,  
second order PDE's  
for the 3 unknowns:

Why are these equations coupled?  
Potential → Field → current → changes  
potential → changes field and so on...

$$p(\vec{r})$$

{  
 Conservations laws: **exact**  
 Transport eqs. (drift-diffusion): **approximate**  
 }

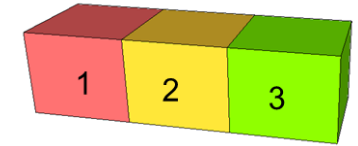
# Section 18 Semiconductor Equations



$$I = G \times V$$

$$= q \times n \times v \times A$$

↑ charge density    ↑ velocity    area



- 18.1 Continuity Equations
- 18.2 Analytical Solutions (Strategy & Examples)
- 18.3 Numerical Solutions
  - » Basic Transport Equations



$$\frac{\partial n}{\partial t} = \frac{1}{q} \nabla \cdot \mathbf{J}_N - r_N + g_N$$

$$\mathbf{J}_N = qn\mu_N E + qD_N \nabla n$$

$$\frac{\partial p}{\partial t} = -\frac{1}{q} \nabla \cdot \mathbf{J}_P + g_P - r_P$$

$$\mathbf{J}_P = qp\mu_P E - qD_P \nabla p$$

$$\nabla \cdot D = q(p - n + N_D^+ - N_A^-)$$

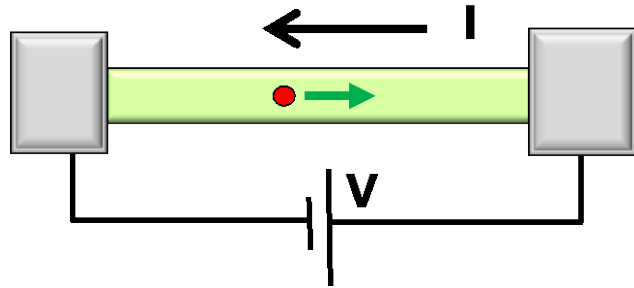
$$\nabla \cdot \vec{D} = \rho$$

$$\nabla \cdot \left( \vec{J}_n / -q \right) = (g_N - r_N)$$

$$\nabla \cdot \left( \vec{J}_p / q \right) = (g_N - r_N)$$

Vid  
Vid  
Vid

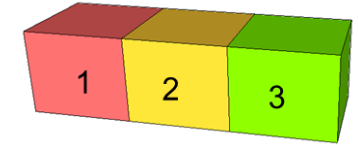
# Section 18 Semiconductor Equations



$$I = G \times V$$

$$= q \times n \times v \times A$$

↑ charge density    ↑ velocity    area



- 18.1 Continuity Equations
- 18.2 Analytical Solutions (Strategy & Examples)
- 18.3 Numerical Solutions
  - » Basic Transport Equations
  - » Gridding and finite differences



$$\frac{\partial n}{\partial t} = \frac{1}{q} \nabla \cdot \mathbf{J}_N - r_N + g_N$$

$$\mathbf{J}_N = qn\mu_N E + qD_N \nabla n$$

$$\frac{\partial p}{\partial t} = -\frac{1}{q} \nabla \cdot \mathbf{J}_P + g_P - r_P$$

$$\mathbf{J}_P = qp\mu_P E - qD_P \nabla p$$

$$\nabla \cdot D = q(p - n + N_D^+ - N_A^-)$$

$$\nabla \cdot \vec{D} = \rho$$

$$\nabla \cdot \left( \vec{J}_n / -q \right) = (g_N - r_N)$$

$$\nabla \cdot \left( \vec{J}_p / q \right) = (g_N - r_N)$$

Vid  
Vid  
Vid

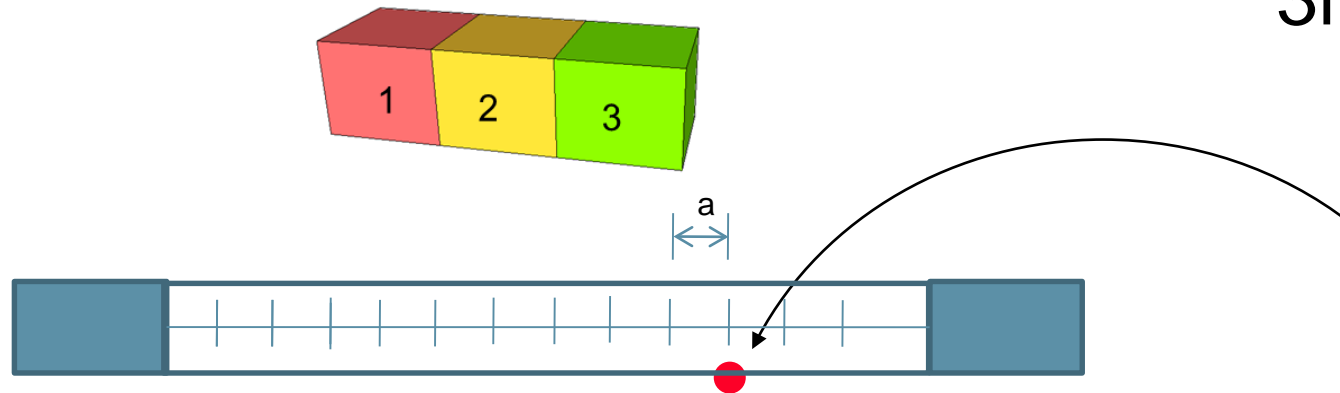


## 2) The Grid

(ii) “exact” numerical solutions

N nodes

3N unknowns



‘Gridding’– total length divided into ‘N’ parts

- equal (uniform gridding) , or
- unequal (adaptive and non-uniform gridding)

Variables described at each point ‘ $i$ ’.

$V_0$  and  $V_{n+1}$  is known because these are voltages at source and drain.

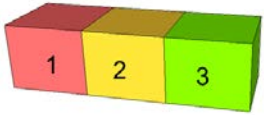
$V_i$

$n_i$

$p_i$

$$\frac{df}{dx} = ?$$

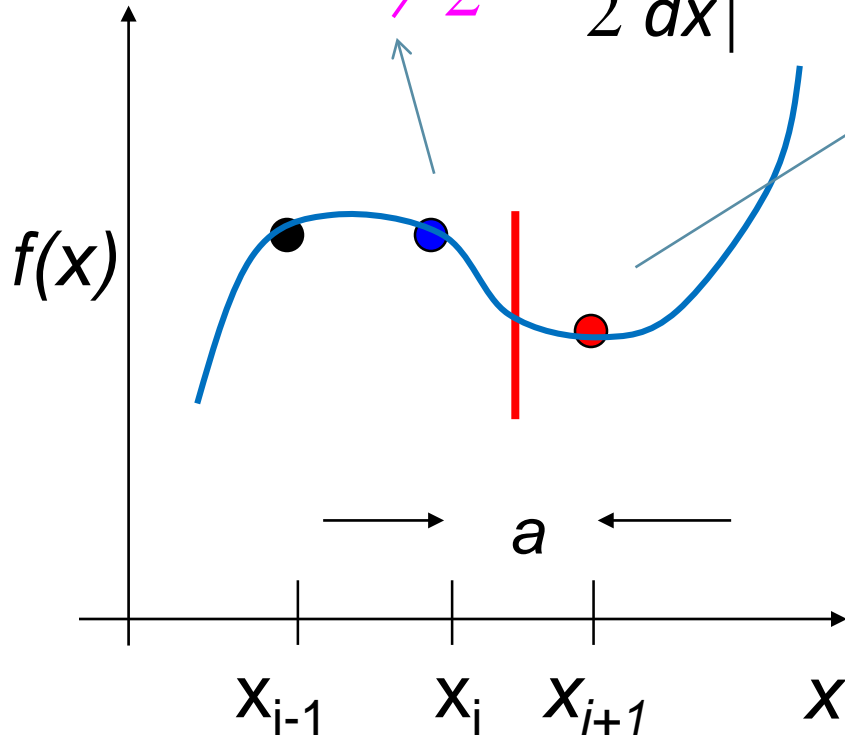
# Finite Difference Expression for Derivative



$$\frac{df}{dx}$$

$$f(x_0) = f(x_0 + a/2) - \frac{a}{2} \left. \frac{df}{dx} \right|_{x_0 + a/2}$$

$$f(x_0 + a) = f(x_0 + a/2) + \frac{a}{2} \left. \frac{df}{dx} \right|_{x_0 + a/2}$$



$$\left. \frac{df}{dx} \right|_{(x_{i+1/2})} = \frac{f_{i+1} - f_i}{a}$$

“centered difference”

$$\frac{d^2 f}{dx^2} = ?$$

# The Second Derivative ...

$$\left. \frac{df}{dx} \right|_{(x_{i+1/2})} = \frac{f_{i+1} - f_i}{a}$$

$$\frac{d^2 f}{dx^2} = ?$$

$$f(x_0 + a) = f(x_0) + a \left. \frac{df}{dx} \right|_{x_0=a} + \frac{a^2}{2} \left. \frac{d^2 f}{dx^2} \right|_{x_0=a} + \dots$$

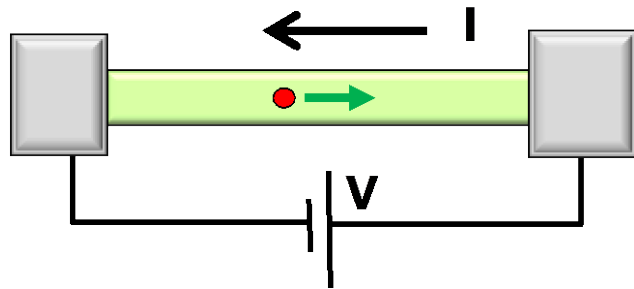
$$f(x_0 - a) = f(x_0) - a \left. \frac{df}{dx} \right|_{x_0=a} + \frac{a^2}{2} \left. \frac{d^2 f}{dx^2} \right|_{x_0=a} - \dots$$

$$f(x_0 + a) + f(x_0 - a) - 2f(x_0) = a^2 \left. \frac{d^2 f}{dx^2} \right|_{x_0=a}$$

$$\left. \frac{d^2 f}{dx^2} \right|_i = \frac{f_{i-1} - 2f_i + f_{i+1}}{a^2}$$

3 point formula, could be extended to N points depending on the number of derivatives we carry in our expansion

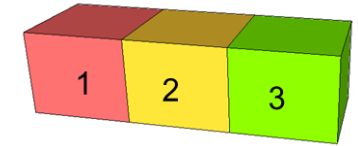
# Section 18 Semiconductor Equations



$$I = G \times V$$

$$= q \times n \times v \times A$$

↑ charge density    ↑ velocity    area



- 18.1 Continuity Equations
- 18.2 Analytical Solutions (Strategy & Examples)
- 18.3 Numerical Solutions
  - » Basic Transport Equations
  - » Gridding and finite differences



$$\frac{\partial n}{\partial t} = \frac{1}{q} \nabla \cdot \mathbf{J}_N - r_N + g_N$$

$$\mathbf{J}_N = qn\mu_N E + qD_N \nabla n$$

$$\frac{\partial p}{\partial t} = -\frac{1}{q} \nabla \cdot \mathbf{J}_P + g_P - r_P$$

$$\mathbf{J}_P = qp\mu_P E - qD_P \nabla p$$

$$\nabla \cdot D = q(p - n + N_D^+ - N_A^-)$$

$$\left. \frac{df}{dx} \right|_{(x_{i+1/2})} = \frac{f_{i+1} - f_i}{a} \quad \left. \frac{d^2 f}{dx^2} \right|_i = \frac{f_{i-1} - 2f_i + f_{i+1}}{a^2}$$

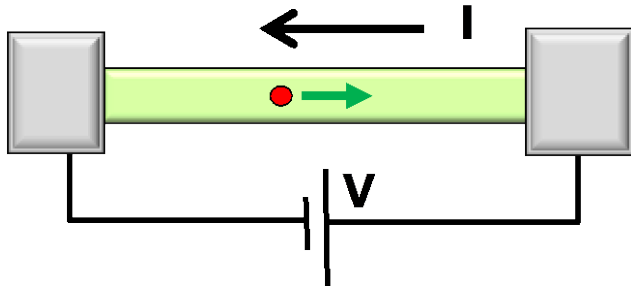
$$\nabla \cdot \vec{D} = \rho$$

$$\nabla \cdot (\vec{J}_n / -q) = (g_N - r_N)$$

$$\nabla \cdot (\vec{J}_p / q) = (g_N - r_N)$$

Vid  
Vid  
Vid

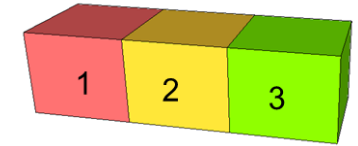
# Section 18 Semiconductor Equations



$$I = G \times V$$

$$= q \times n \times v \times A$$

↑ charge density    ↑ velocity    area



- 18.1 Continuity Equation
- 18.2 Analytical Solutions (Strategy & Examples)
- 18.3 Numerical Solutions
  - » Basic Transport Equations
  - » Gridding and finite differences
  - » Discretizing equations and boundary conditions

$$\frac{\partial n}{\partial t} = \frac{1}{q} \nabla \cdot \mathbf{J}_N - r_N + g_N$$

$$\mathbf{J}_N = qn\mu_N E + qD_N \nabla n$$

$$\frac{\partial p}{\partial t} = -\frac{1}{q} \nabla \cdot \mathbf{J}_P + g_P - r_P$$

$$\mathbf{J}_P = qp\mu_P E - qD_P \nabla p$$

$$\nabla \cdot D = q(p - n + N_D^+ - N_A^-)$$

$$\nabla \cdot \vec{D} = \rho$$

$$\nabla \cdot (\vec{J}_n / -q) = (g_N - r_N)$$

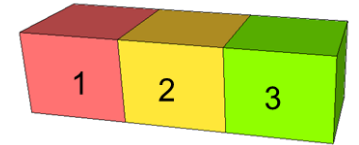
$$\nabla \cdot (\vec{J}_p / q) = (g_N - r_N)$$

status

$$\left. \frac{df}{dx} \right|_{(x_{i+1/2})} = \frac{f_{i+1} - f_i}{a} \quad \left. \frac{d^2 f}{dx^2} \right|_i = \frac{f_{i-1} - 2f_i + f_{i+1}}{a^2}$$

Vid  
Vid  
Vid

## 2) Control Volume



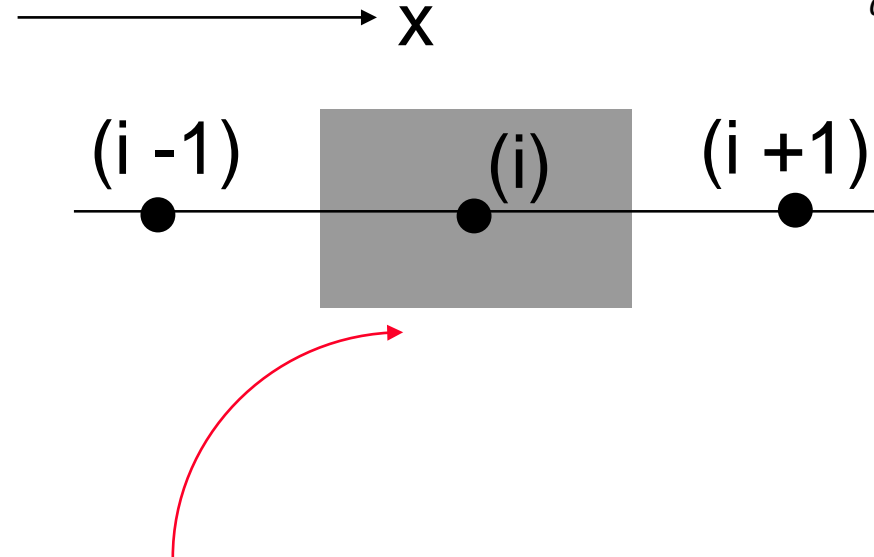
$$\left. \frac{df}{dx} \right|_{(x_{i+1/2})} = \frac{f_{i+1} - f_i}{a}$$

$$\left. \frac{d^2 f}{dx^2} \right|_i = \frac{f_{i-1} - 2f_i + f_{i+1}}{a^2}$$

3 unknowns at each node:

$$V_i, n_i, p_i$$

Need 3 equations  
at each node



“control volume”

# Discretizing Poisson's Equation

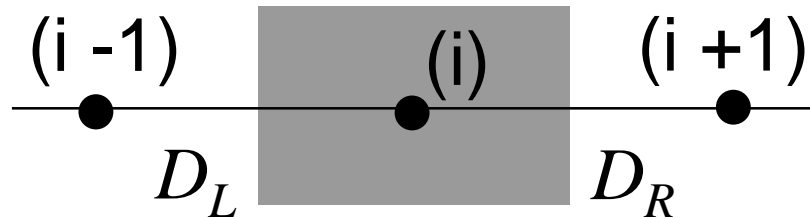
$$\nabla^2 V = -\rho / K_s \epsilon_0$$

$$\nabla \cdot \mathbf{D} = \rho \quad \mathbf{D} = K_s \epsilon_0 \mathbf{E} = -K_s \epsilon_0 \nabla V$$

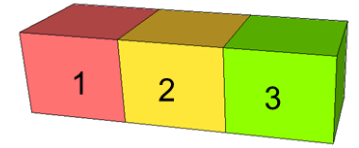
$$\frac{V_{(i-1)} - 2V_{(i)} + V_{(i+1)}}{a^2} = -\frac{q}{K_s \epsilon_0} (p_i - n_i + N_{D,i}^+ - N_{A,i}^-)$$

Since  $V_0$  and  $V_{-1}$  are known, as are carrier concentration on doping (or lack thereof) in contacts, we find  $V_1$  and iterate from this point to solve for potential.

$$F_V^i(V_{i-1}, V_i, V_{i+1}, n_i, p_i) = 0$$



→ Once this potential is found, solve continuity equation to obtain new carrier concentrations

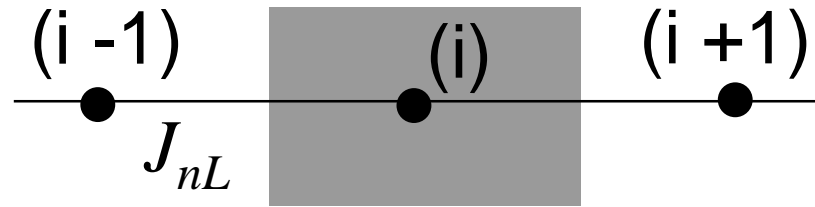


$$\left. \frac{df}{dx} \right|_{(x_{i+1/2})} = \frac{f_{i+1} - f_i}{a}$$

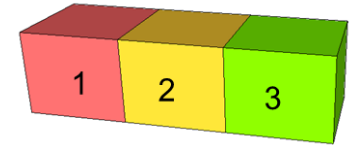
$$\left. \frac{d^2 f}{dx^2} \right|_i = \frac{f_{i-1} - 2f_i + f_{i+1}}{a^2}$$

# Discretizing Continuity Equations

$$\nabla \cdot \vec{J}_n = -q(g_N - r_N)$$



$$J_{nL} = -nq\mu_n \frac{dV}{dx} + kT\mu_n \frac{dn}{dx}$$



$$\left. \frac{df}{dx} \right|_{(x_{i+1/2})} = \frac{f_{i+1} - f_i}{a}$$

$$\left. \frac{d^2 f}{dx^2} \right|_i = \frac{f_{i-1} - 2f_i + f_{i+1}}{a^2}$$

**The simplest approach.....**

$$\frac{J_{nL}}{kT\mu_n} = - \left( \frac{n_{i-1} + n_i}{2} \right) \left( \frac{V_i - V_{i-1}}{a(kT/q)} \right) + \left( \frac{n_i - n_{i-1}}{a} \right)$$

$$F_n^i(V_{i-1}, V_i, n_i, n_{i-1}, p_i, p_{i-1}) = 0$$

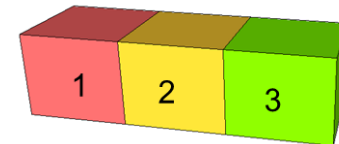
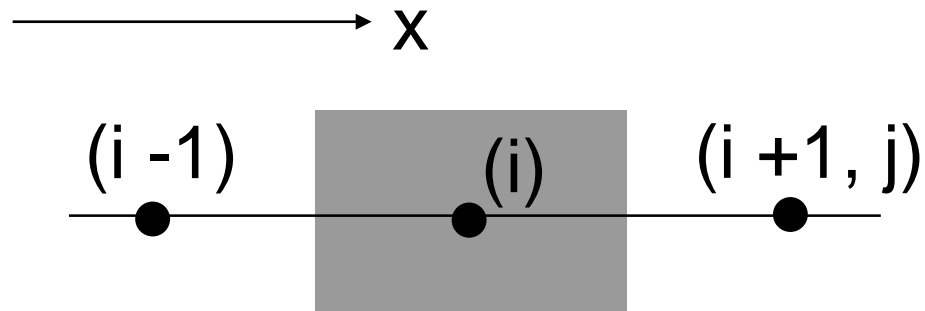


# Three Discretized Equations

$$F_V^i = 0$$

$$F_n^i = 0$$

$$F_p^i = 0$$



3 unknowns at each node

N nodes

3N unknowns and 3N equations (coupled to each other)

$$F_n^i (V_{i-1}, V_i, n_i, n_{i-1}, p_i, p_{i-1}) = 0$$

# Numerical Solution - Poisson Equation Only

3 unknowns at each node

N nodes

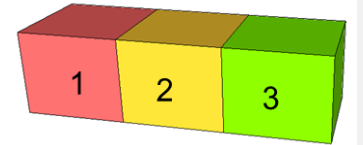
3N unknowns and 3N equations (coupled to each other)

$$F_n^i(V_{i-1}, V_i, n_i, n_{i-1}, p_i, p_{i-1}) = 0$$

$$F_V^i(V_{i-1}, V_i, V_{i+1}, n_i, p_i) = 0$$

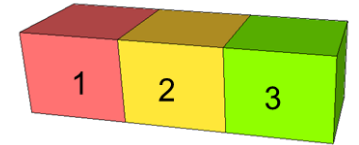
linear if  $n_i$  and  $p_i$  are known  $[A] \vec{V} = \vec{b}$

$$[A]: \left( \begin{array}{c} \diagdown \\ \diagdown \\ \diagdown \\ \diagdown \end{array} \right) \quad \vec{V} = \begin{bmatrix} V_1 \\ V_2 \\ \vdots \\ V_N \end{bmatrix}$$



# Boundary conditions

Contacts are assumed large and in equilibrium  
→ detailed balance and law of mass-action apply!!



$$n_0 p_0 = n_i^2$$

$$n_{N+1} p_{N+1} = n_i^2$$

Dopant density

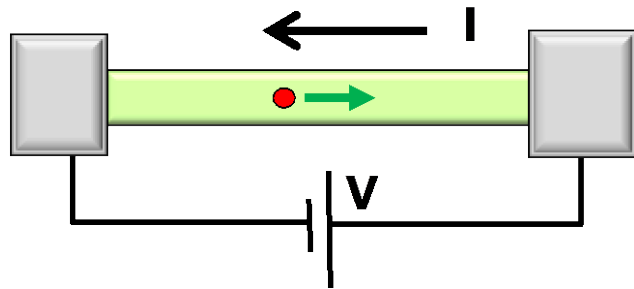


$$V = V_A$$

One could have unequal materials on the two contact sides, one must be careful to use the right intrinsic concentration <-> material.

$$V = 0$$

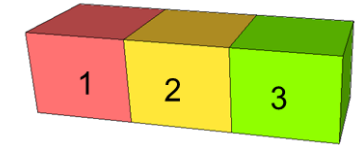
# Section 18 Semiconductor Equations



$$I = G \times V$$

$$= q \times n \times v \times A$$

↑ charge density    ↑ velocity    area



- 18.1 Continuity Equations
- 18.2 Analytical Solutions (Strategy & Examples)
- 18.3 Numerical Solutions
  - » Basic Transport Equations
  - » Gridding and finite differences
  - » Discretizing equations and boundary conditions

status

$$\frac{\partial n}{\partial t} = \frac{1}{q} \nabla \cdot \mathbf{J}_N - r_N + g_N$$

$$\mathbf{J}_N = qn\mu_N E + qD_N \nabla n$$

$$\frac{\partial p}{\partial t} = -\frac{1}{q} \nabla \cdot \mathbf{J}_P + g_P - r_P$$

$$\mathbf{J}_P = qp\mu_P E - qD_P \nabla p$$

$$\nabla \cdot D = q(p - n + N_D^+ - N_A^-)$$

$$\left. \frac{df}{dx} \right|_{(x_{i+1/2})} = \frac{f_{i+1} - f_i}{a} \quad \left. \frac{d^2 f}{dx^2} \right|_i = \frac{f_{i-1} - 2f_i + f_{i+1}}{a^2}$$

$$\nabla \cdot \vec{D} = \rho$$

$$\nabla \cdot (\vec{J}_n / -q) = (g_N - r_N)$$

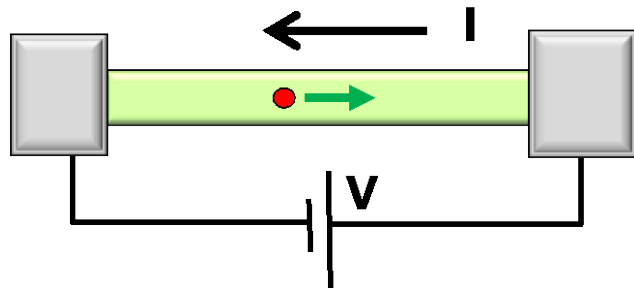
$$\nabla \cdot (\vec{J}_p / q) = (g_N - r_N)$$

Vid

Vid

Vid

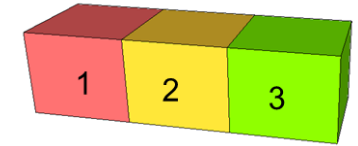
# Section 18 Semiconductor Equations



$$I = G \times V$$

$$= q \times n \times v \times A$$

↑ charge density    ↑ velocity    area



- 18.1 Continuity Equations
- 18.2 Analytical Solutions (Strategy & Examples)
- 18.3 Numerical Solutions
  - » Basic Transport Equations
  - » Gridding and finite differences
  - » Discretizing equations and boundary conditions
  - » Iterative Solution Approach

$$\frac{\partial n}{\partial t} = \frac{1}{q} \nabla \cdot \mathbf{J}_N - r_N + g_N$$

$$\mathbf{J}_N = qn\mu_N E + qD_N \nabla n$$

$$\frac{\partial p}{\partial t} = -\frac{1}{q} \nabla \cdot \mathbf{J}_P + g_P - r_P$$

$$\mathbf{J}_P = qp\mu_P E - qD_P \nabla p$$

$$\nabla \cdot D = q(p - n + N_D^+ - N_A^-)$$

$$\left. \frac{df}{dx} \right|_{(x_{i+1/2})} = \frac{f_{i+1} - f_i}{a} \quad \left. \frac{d^2 f}{dx^2} \right|_i = \frac{f_{i-1} - 2f_i + f_{i+1}}{a^2}$$

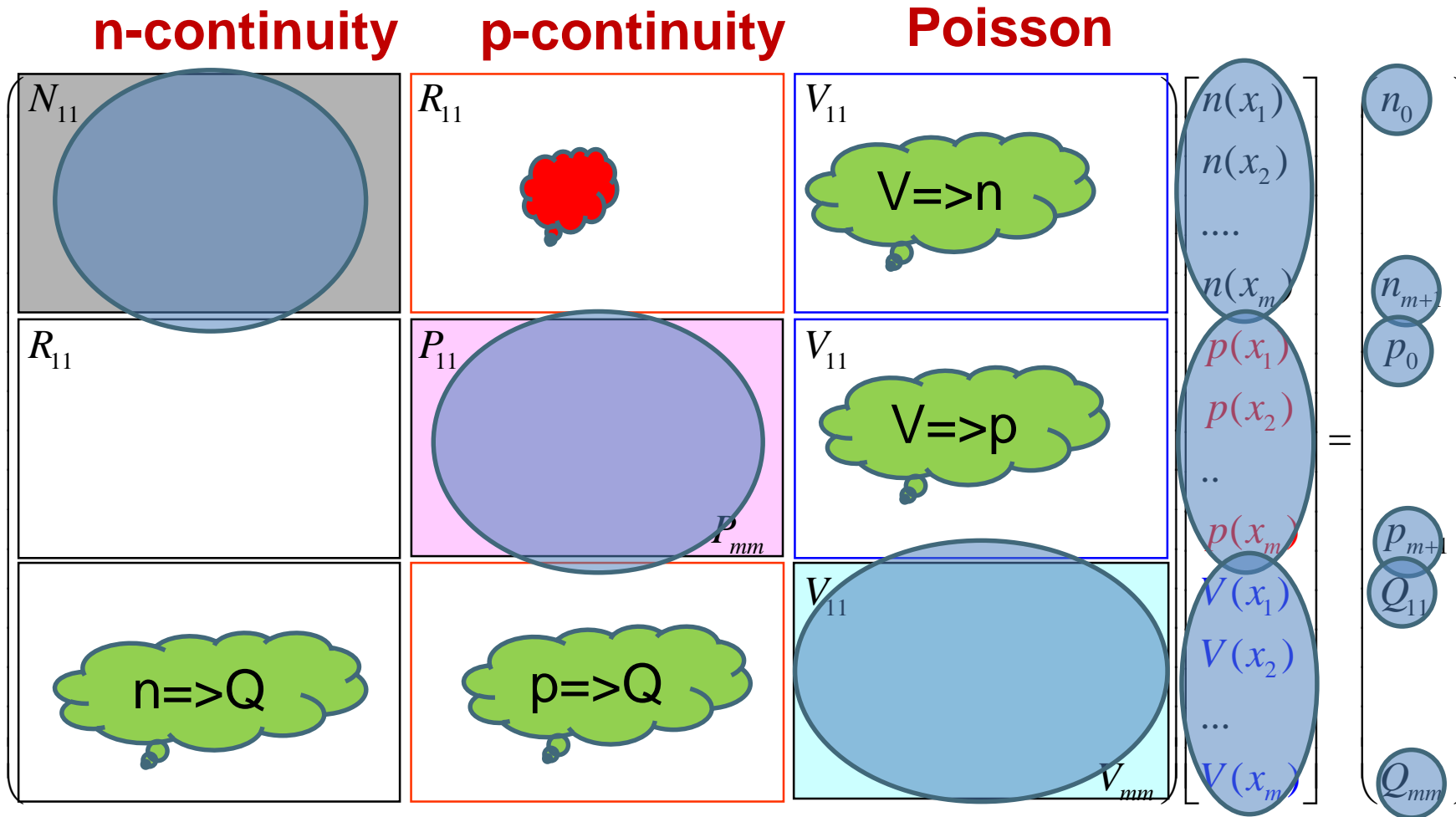
← status =  $\rho$

$$\nabla \cdot (\vec{J}_n / -q) = (g_N - r_N)$$

$$\nabla \cdot (\vec{J}_p / q) = (g_N - r_N)$$

Vid  
Vid  
Vid

# Numerical Solution...



Off-diagonal terms are Poisson-Continuity equations talking to each other.  
 Recombination-generation terms also feed into continuity equations.

### 3) Uncoupled Numerical Solution

The semiconductor equations are nonlinear!  
(but they are linear individually)

**Uncoupled** solution proced



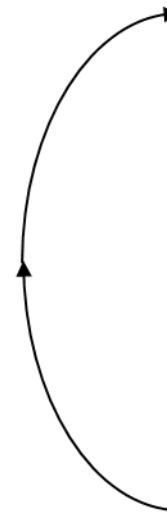
repeat  
until  
satisfied

Guess  $V, n, p$

Solve Poisson  
for new  $V$

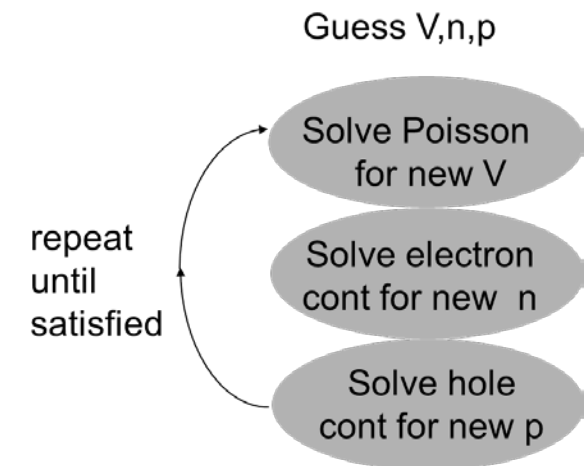
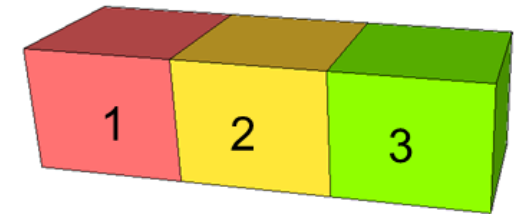
Solve electron  
cont for new  $n$

Solve hole  
cont for new  $p$



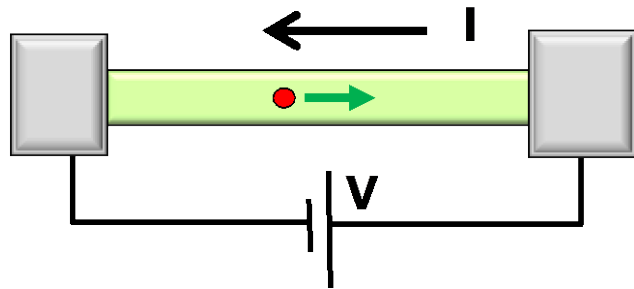
# Summary

- 1) Two methods to solve drift-diffusion equation consistently – analytical and numerical.
- 2) Analytical solution provides great insight and the solution methodology is similar to that of Schrodinger equations.
- 3) Numerical solution is more versatile.
  - 1) One begins with a set of equations and boundary conditions,
  - 2) discretize the equations on a grid with  $N$  nodes to obtain  $3N$  nonlinear equations in  $3N$  unknowns, and
  - 3) solve the system of nonlinear equations by iteration.





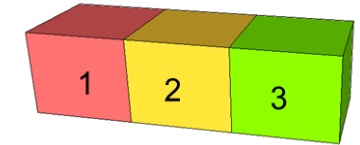
# Section 18 Semiconductor Equations



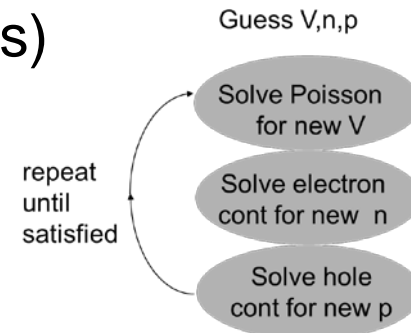
$$I = G \times V$$

$$= q \times n \times v \times A$$

↑ charge density    ↑ velocity    area



- 18.1 Continuity Equations
- 18.2 Analytical Solutions (Strategy & Examples)
- 18.3 Numerical Solutions
  - » Basic Transport Equations
  - » Gridding and finite differences
  - » Discretizing equations and boundary conditions
  - » Iterative Solution Approach



$$\frac{\partial n}{\partial t} = \frac{1}{q} \nabla \cdot \mathbf{J}_N - r_N + g_N$$

$$\mathbf{J}_N = qn\mu_N E + qD_N \nabla n$$

$$\frac{\partial p}{\partial t} = -\frac{1}{q} \nabla \cdot \mathbf{J}_P + g_P - r_P$$

$$\mathbf{J}_P = qp\mu_P E - qD_P \nabla p$$

$$\nabla \cdot D = q(p - n + N_D^+ - N_A^-)$$

$$\left. \frac{df}{dx} \right|_{(x_{i+1/2})} = \frac{f_{i+1} - f_i}{a} \quad \left. \frac{d^2 f}{dx^2} \right|_i = \frac{f_{i-1} - 2f_i + f_{i+1}}{a^2}$$

$$\nabla \cdot \vec{D} = \rho$$

$$\nabla \cdot \left( \vec{J}_n / -q \right) = (g_N - r_N)$$

$$\nabla \cdot \left( \vec{J}_p / q \right) = (g_N - r_N)$$

Vid  
Vid  
Vid

