

Section 18 Semiconductor Equations

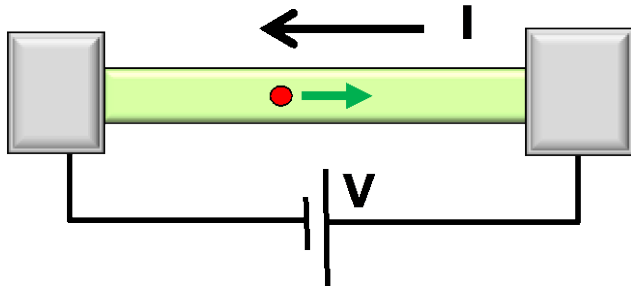
18.2 Analytical Solutions (Strategy & Examples)

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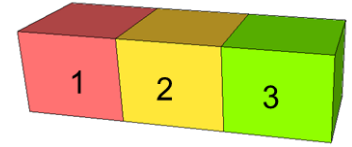
Section 18 Continuity Equations



$$I = G \times V$$

$$= q \times n \times v \times A$$

↑ charge density ↑ density ↑ velocity area



- 18.1 Continuity Equations
- 18.2 Analytical Solutions (Strategy & Examples)
- 18.3



$$\frac{\partial n}{\partial t} = \frac{1}{q} \nabla \cdot \mathbf{J}_N - r_N + g_N$$

$$\mathbf{J}_N = qn\mu_N E + qD_N \nabla n$$

$$\frac{\partial p}{\partial t} = -\frac{1}{q} \nabla \cdot \mathbf{J}_P + g_P - r_P$$

$$\mathbf{J}_P = qp\mu_P E - qD_P \nabla p$$

$$\nabla \cdot D = q(p - n + N_D^+ - N_A^-)$$

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Analytical Solutions

$$\frac{\partial n}{\partial t} = \frac{1}{q} \nabla \cdot \mathbf{J}_N - r_N + g_N$$

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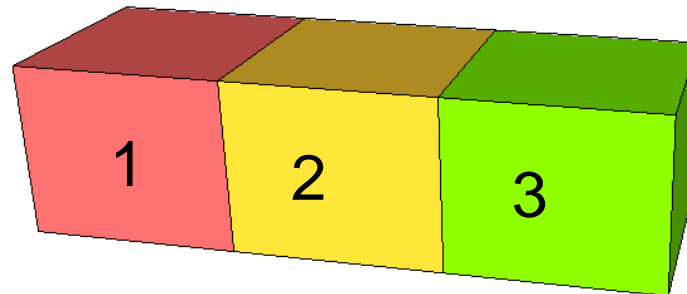
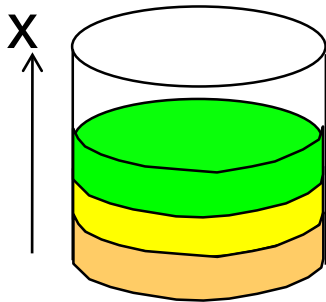
$$\frac{\partial p}{\partial t} = -\frac{1}{q} \nabla \cdot \mathbf{J}_P + g_P - r_P$$

$$\mathbf{J}_P = qp\mu_P E - qD_P \nabla p$$

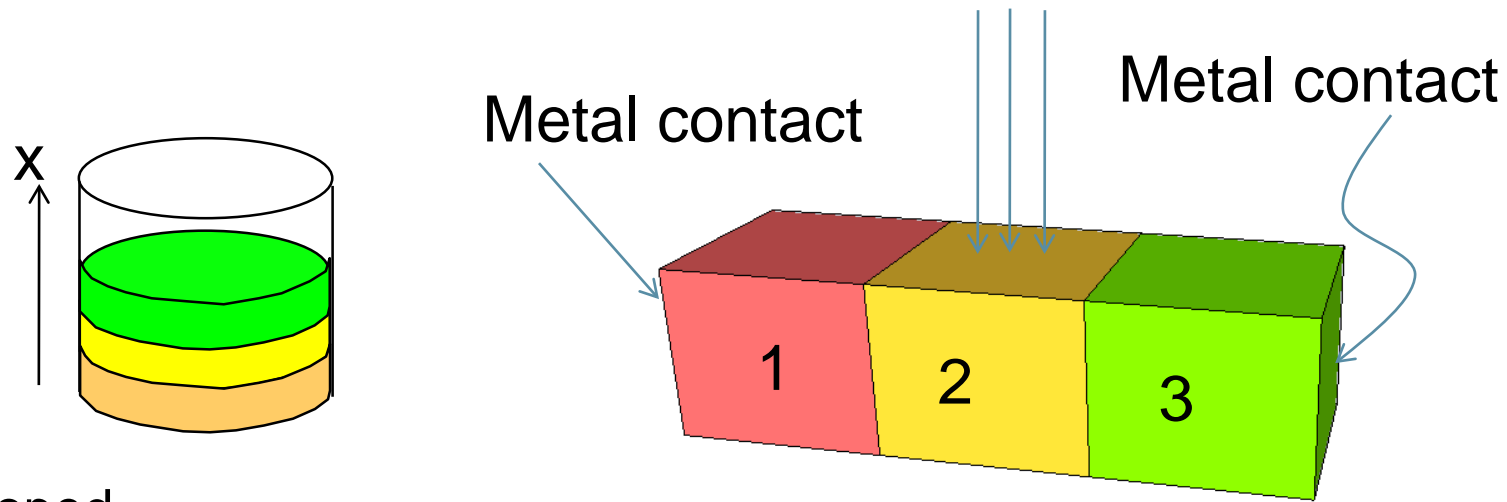
$$\nabla \cdot D = q(p - n + N_D^+ - N_A^-)$$

← Diffusion approximation,
Minority carrier transport,
Ambipolar transport

← Band-diagram



Consider a complicated real device example

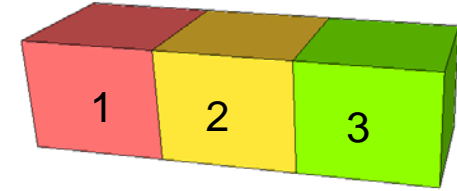


- ❑ Acceptor doped
- ❑ Light turned on in the middle section.
- ❑ The right region is full of mid-gap traps because of dangling bonds due to unpassivated surface.
- ❑ Interface traps at the end of the right region (That's where the dangling bonds are...)
- ❑ The left region is trap free.
- ❑ The left/right regions contacted by metal electrode.

Recall: Analytical Solution of Schrodinger Equation

1) $\frac{d^2\psi}{dx^2} + k^2\psi = 0$

→ 2N unknowns for N regions



2) $\psi(x = -\infty) = 0$
 $\psi(x = +\infty) = 0$

→ Reduces 2 unknowns

3) $\psi|_{x=x_B^-} = \psi|_{x=x_B^+}$
 $\frac{d\psi}{dx}|_{x=x_B^-} = \frac{d\psi}{dx}|_{x=x_B^+}$

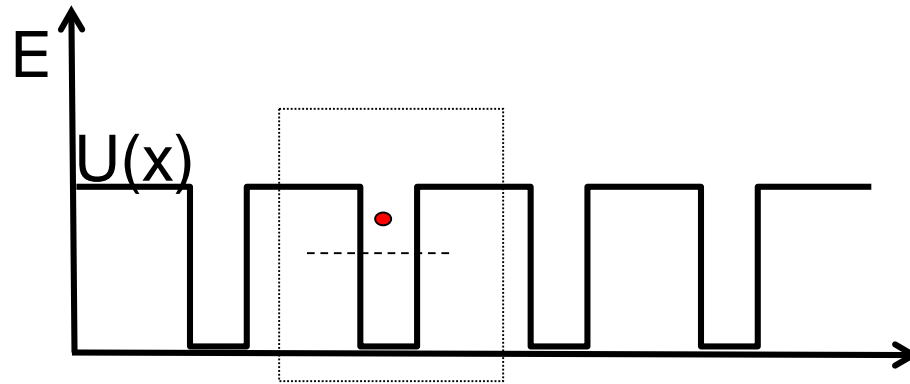
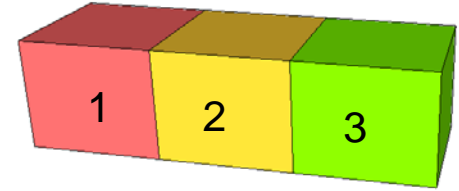
→ Set 2N-2 equations for 2N-2 unknowns (for continuous U)

4) Det(coefficient matrix)=0
And find E by graphical or numerical solution

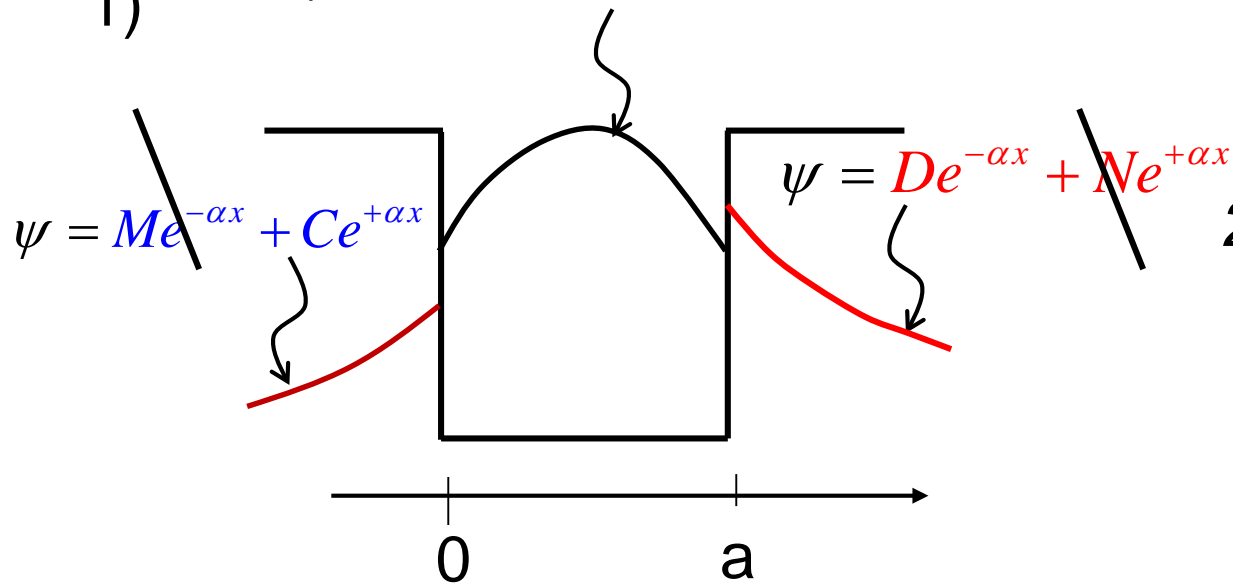
5) $\int_{-\infty}^{\infty} |\psi(x, E)|^2 dx = 1$

for wave function

Recall: Bound-levels in Finite well



1) $\psi = A \sin kx + B \cos kx$



2) *Boundary Conditions ...*

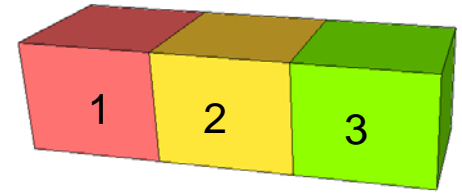
$$\psi(x = -\infty) = 0$$

$$\psi(x = +\infty) = 0$$

Analogously, we solve for our device

Solve the equations in different regions independently.

Bring them together by applying boundary conditions.



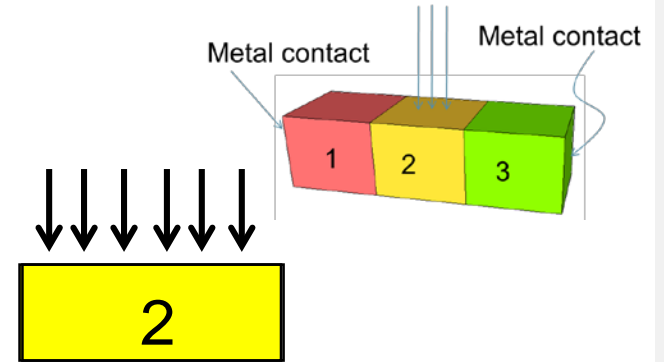
Region 2: Transient, Uniform Illumination, Uniform doping

$$\frac{\partial n}{\partial t} = \frac{1}{q} \nabla \cdot \mathbf{J}_N - r_N + g_N \quad (\text{uniform})$$

$$\mathbf{J}_N = qn\mu_N E + qD_N \nabla n$$

$$\frac{\partial(n'_0 + \Delta n)}{\partial t} = -\frac{\Delta n}{\tau_n} + G$$

Recall Shockley-Read-Hall



Acceptor doped

$$\frac{\partial p}{\partial t} = \frac{-1}{q} \nabla \cdot \mathbf{J}_p - r_p + g_p \quad (\text{uniform})$$

$$\mathbf{J}_p = qp\mu_p E - qD_p \nabla p$$

Electric field still zero because new carriers balance

$$\nabla \cdot D = q(p - n + N_D^+ - N_A^-) = q(p_0 + \Delta p - n_0 - \Delta n + N_D^+ - N_A^-) = 0$$

Example: Transient, Uniform Illumination, Uniform doping, No applied electric field

$$\frac{\partial(n_0 + \Delta n)}{\partial t} = -\frac{\Delta n}{\tau_n} + G$$

$$\frac{\partial(\Delta n)}{\partial t} = -\frac{\Delta n}{\tau_n} + G$$

$$\Delta n(x, t) = A + Be^{-t/\tau_n}$$

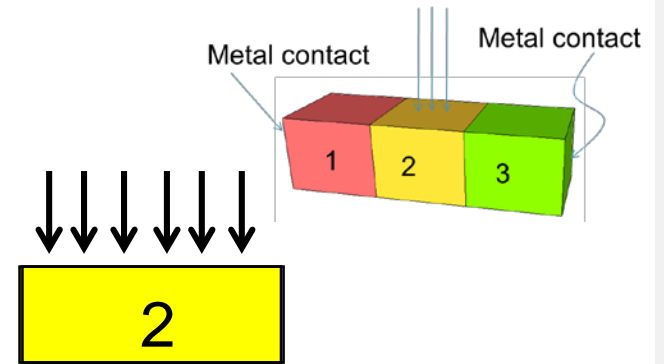
No carriers yet generated...

$$t = 0, \quad \Delta n(x, 0) = 0 \Rightarrow A = -B$$

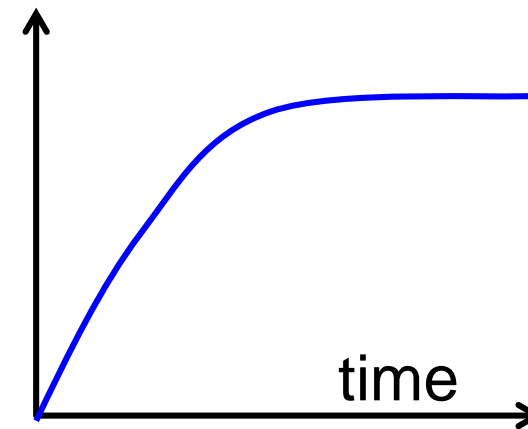
$$t \rightarrow \infty, \quad \Delta n(x, \infty) = G\tau_n = A$$

Steady state, no change
in carriers with time...

$$\Delta n(x, t) = G\tau_n \left(1 - e^{-t/\tau_n}\right)$$



Acceptor doped



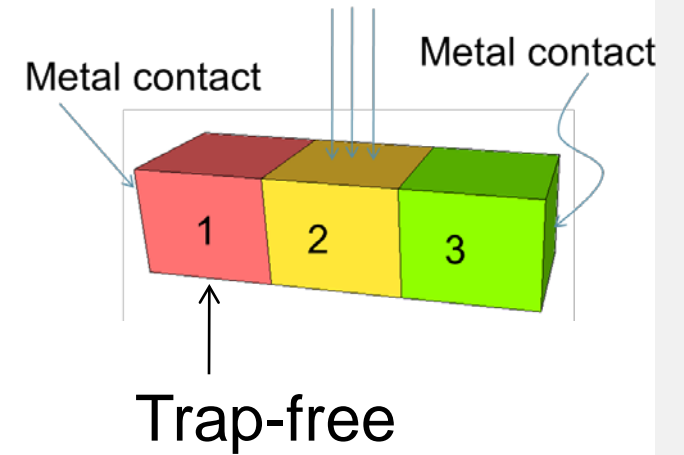
Region 1: One sided Minority Diffusion at steady state

$$\frac{\partial n}{\partial t} = 0 \text{ (steady-state)}$$

$$r_N = 0 \text{ (trap free)}$$

$$g_N = 0 \text{ (no generation)}$$

Steady state
Acceptor doped



$$\cancel{\frac{\partial n}{\partial t}} = \frac{1}{q} \frac{dJ_n}{dx} - \cancel{r_N} + \cancel{g_N}$$

$$E = 0 \quad \mathbf{J}_N = qn\cancel{\mu_N}E + qD_N \frac{dn}{dx}$$

$$D_N \frac{dn}{dx} \neq 0 \text{ (due to insertion of electrons from central region)}$$

$$0 = D_N \frac{d^2 n}{dx^2}$$

Example: One sided Minority Diffusion

$$0 = D_N \frac{d^2 n}{dx^2}$$

$$\Delta n(x, t) = C + Dx'$$

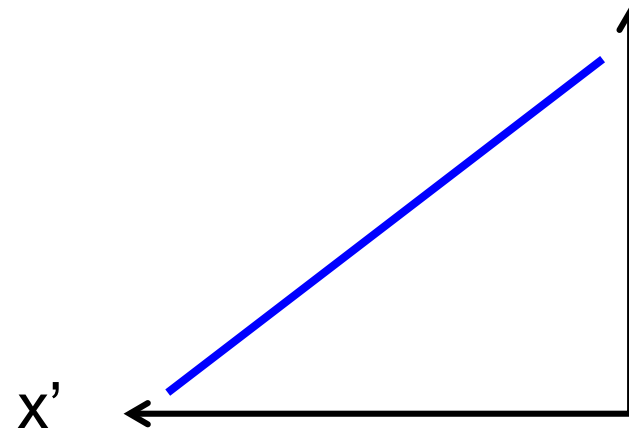
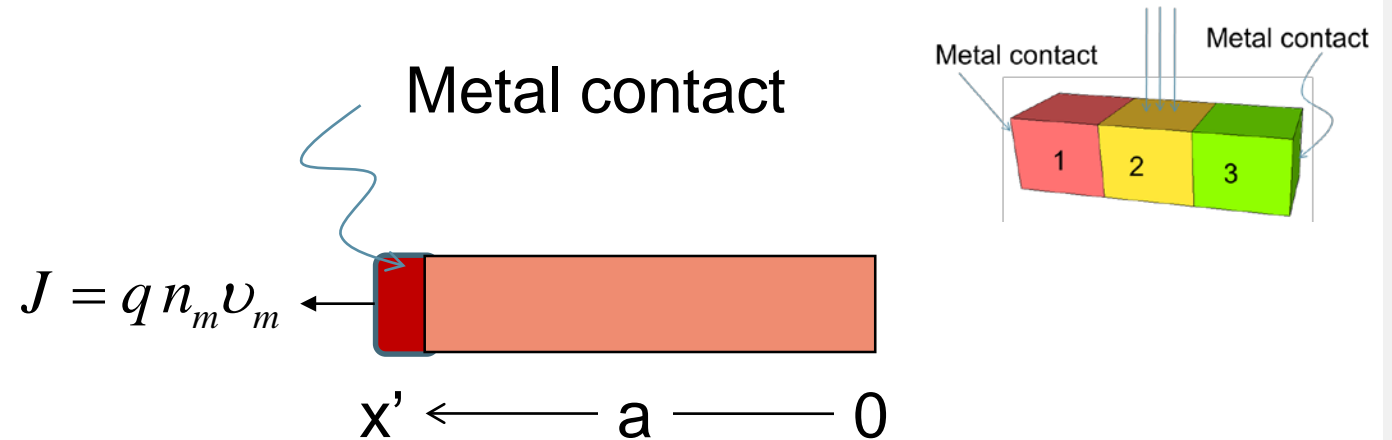
$$x = a, \quad \Delta n(x' = a) = 0 \Rightarrow C = -Da$$

(Metal has high electron density
as compared to semiconductor)

$$x = 0', \quad \Delta n(x' = 0') = C$$

Just substitute $x=0$ in above eqn.

$$\Delta n(x, t) = \Delta n(x = 0') \left(1 - \frac{x'}{a}\right)$$



Region 3: Steady state Minority Diffusion with recombination

$$\frac{\partial n}{\partial t} = 0 \text{ (steady-state)}$$

$$r_N \neq 0 \text{ (not trap free)}$$

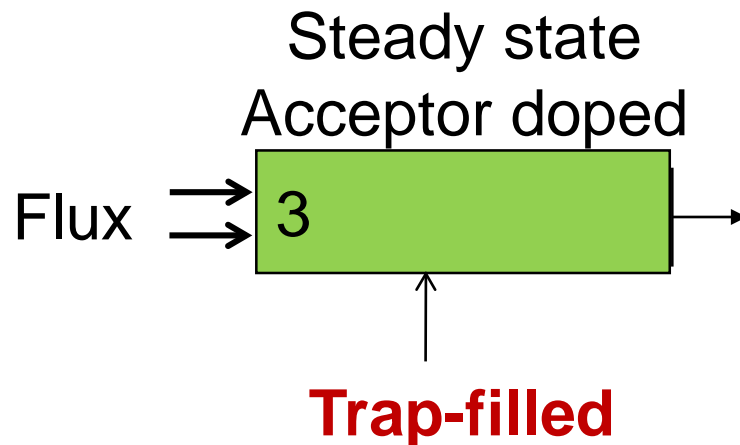
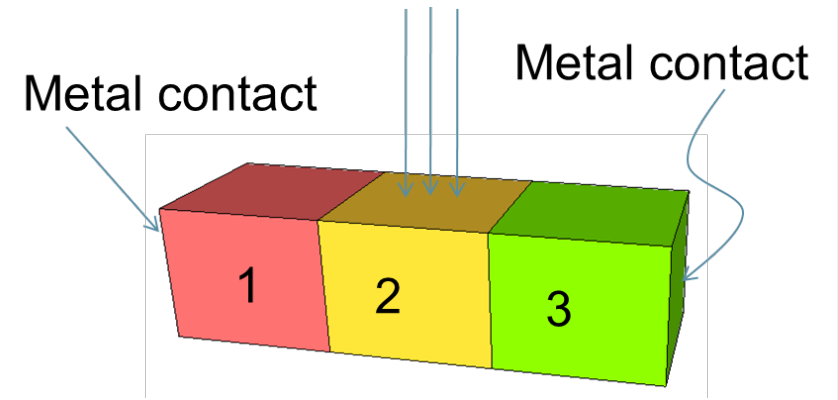
$$g_N = 0 \text{ (no generation)}$$

$$E = 0$$

$$D_N \frac{dn}{dx} \neq 0 \text{ (due to insertion of electrons from central region)}$$

$$0 = D_N \frac{d^2(n_0 + \Delta n)}{dx^2} - \frac{\Delta n}{\tau_n}$$

$$0 = D_N \frac{d^2 \Delta n}{dx^2} - \frac{\Delta n}{\tau_n}$$



Diffusion with Recombination ...

$$D_N \frac{d^2 \Delta n}{dx^2} - \frac{\Delta n}{\tau_n} = 0$$

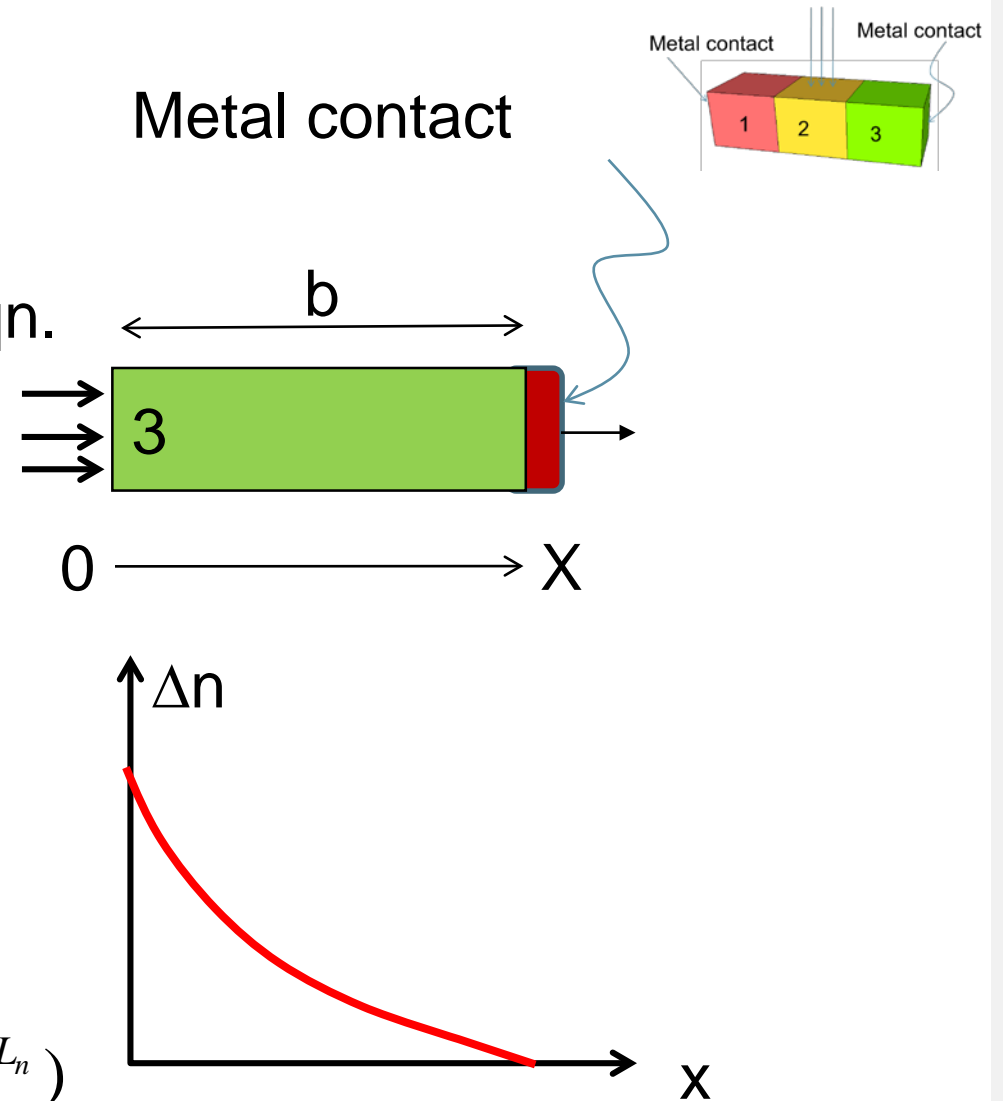
Functionally similar to Schrodinger eqn.

$$\Delta n(x, t) = E e^{x/L_n} + F e^{-x/L_n}$$

$$x = b, \quad \Delta n(x = b) = 0 \Rightarrow F = -E e^{2b/L_n}$$

$$x = 0, \quad \Delta n(x = 0) = E + F = \Delta n(x = 0)$$

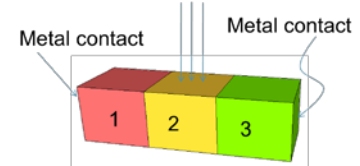
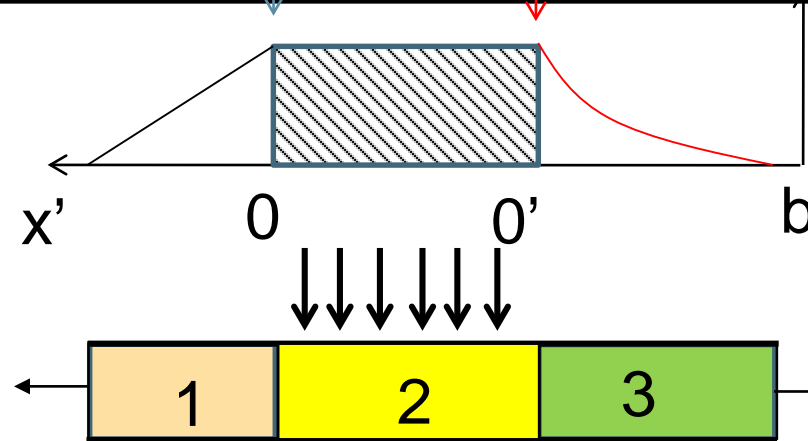
$$\Delta n(x, t) = \frac{\Delta n(0)}{(1 - e^{2b/L_n})} (e^{x/L_n} - e^{2b/L_n} e^{-x/L_n})$$



Combining them all ...

$$\Delta n_2(x) = G\tau_n =$$

$$\Delta n_2(0) = \Delta n_2(0')$$



Match boundary condition

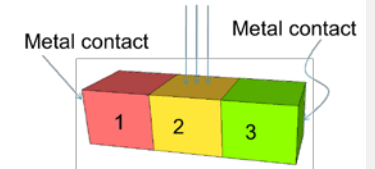
$$\Delta n_1(x') = \Delta n(x=0) \left(1 - \frac{x'}{a}\right) = G\tau_n \left(1 - \frac{x'}{a}\right)$$

$$\Delta n(x) = \frac{\Delta n(0')}{(1 - e^{2b/L_n})} (e^{x/L_n} - e^{2b/L_n} e^{-x/L_n}) = \frac{G\tau_n (e^{x/L_n} - e^{2b/L_n} e^{-x/L_n})}{(1 - e^{2b/L_n})}$$

Calculating current

$$\mathbf{J}_N = qn\mu_N E + qD_N \frac{dn}{dx}$$

Analytical Solutions Summary



- 1) Continuity Equations form the basis of analysis of all the devices we will study in this course.
- 2) Full numerical solution of the equations are possible and many commercial software are available to do so.
- 3) Analytical solutions however provide a great deal of insight into the key physical mechanism involved in the operation of a device.

$$\frac{\partial n}{\partial t} = \frac{1}{q} \nabla \cdot \mathbf{J}_N - r_N + g_N$$

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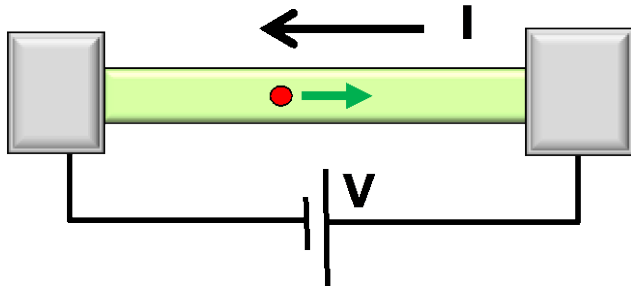
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Section 18

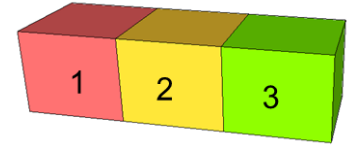
Continuity Equations



$$I = G \times V$$

$$= q \times n \times v \times A$$

↑ charge density ↑ density ↑ velocity area



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$$\frac{\partial n}{\partial t} = \frac{1}{q} \nabla \cdot \mathbf{J}_N - r_N + g_N$$

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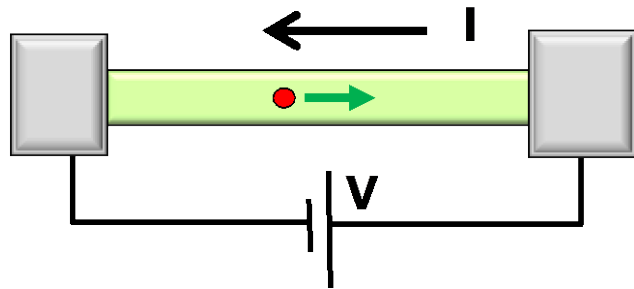
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Vid
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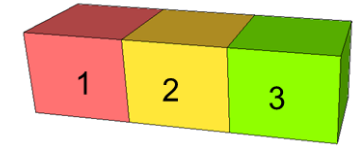
Section 18 Continuity Equations



$$I = G \times V$$

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↑ charge density ↑ velocity area



- 18.1 Continuity Equations
- 18.2 Analytical Solutions (Strategy & Examples)
- 18.3 Numerical Solutions
 - » Basic Transport Equations
 - » Gridding and finite differences
 - » Discretizing equations and boundary conditions
 - » Iterative Solution Approach



$$\frac{\partial n}{\partial t} = \frac{1}{q} \nabla \cdot \mathbf{J}_N - r_N + g_N$$

$$\mathbf{J}_N = qn\mu_N E + qD_N \nabla n$$

$$\frac{\partial p}{\partial t} = -\frac{1}{q} \nabla \cdot \mathbf{J}_P + g_P - r_P$$

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