# **Solid State Devices**

# Section 17 Intro to Transport - Drift, Mobility, Diffusion, Einstein Relationship

17.3 Carrier Concentration from Hall Effect

Gerhard Klimeck gekco@purdue.edu

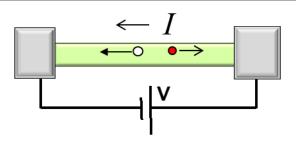


School of Electrical and Computer Engineering





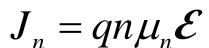
# Intro to Transport - Drift, Mobility, Diffusion, Einstein Relationship



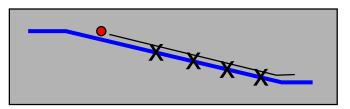
Transport with scattering, non-equilibrium Stat. Mech.

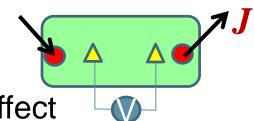
• Drift-diffusion equation with recombination-generation

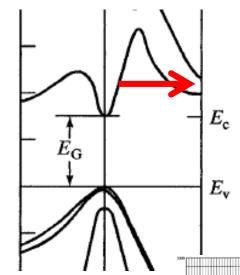
- 17.1 Drift Current
- 17.2 Mobility
  - »Matthiessen Rule
  - »High Field Effects
  - »Mobility Measurement



$$\frac{1}{\mu_n} = \frac{1}{\mu_{ph}} + \frac{1}{\mu_{II}}$$

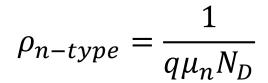


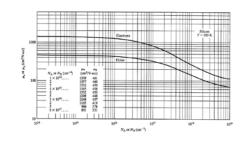


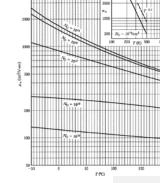


- 17.3 Carrier Concentration from Hall Effect
- 17.4

 $\mu_n$  and  $N_D$  are coupled need the electron density => get the doping





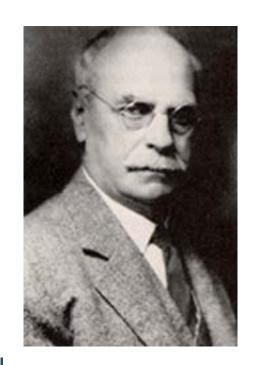


Vid

Consider system to be in local equilibrium

### Measuring Electron Density Hall Effect Trivia - FYI

- Discovered in 1879 by Edwin Herbert Hall while he was working on his doctoral degree at Johns Hopkins University in Baltimore, Maryland.
- Done 18 years before the electron was discovered.
- 4 Nobel Prizes associated directly with it.
- Read the original article:
   "On a New Action of the Magnet on Electric Currents"
   at <a href="http://www.stenomuseet.dk/skoletj/elmag/kilde9.html">http://www.stenomuseet.dk/skoletj/elmag/kilde9.html</a>
   for a fascinating account of the discovery of the effect.



http://en.wikipedia.org/wiki/Hall\_effect and
http://en.wikipedia.org/wiki/Edwin\_Hall

 $\mu_n$  and  $N_D$  are coupled need the electron density => get the doping

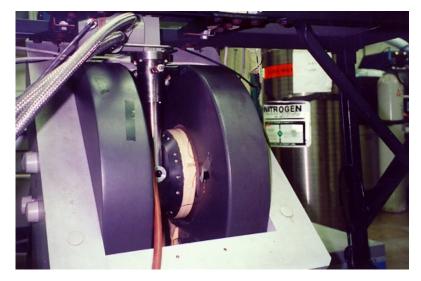




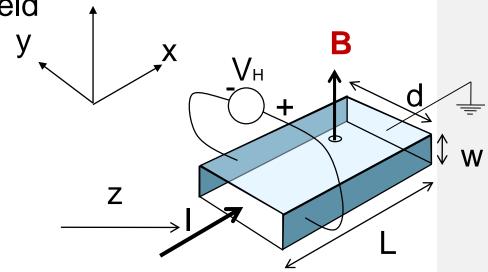
### Set up for Hall Measurement

Force acting on a charged particle in a magnetic field

$$\vec{F} = q\vec{\mathcal{E}} + q\vec{\boldsymbol{v}} \times \vec{\boldsymbol{B}} + \frac{m^*\vec{\boldsymbol{v}}}{\tau}$$



UIC system: 4-300K, 0-1.5 T



- Applying a magnetic field one 'pushes' carriers towards one face
- This induces a voltage across the two faces perpendicular to current flow.
- We will relate this voltage to the current and magnetic field and deduce the density of carriers.

 $\mu_n$  and  $N_D$  are coupled need the electron density => get the doping

#### **Drude Model for electrons \***

# Simple classical Newton's law expression m – effective mass v – drift velocity w/ scattering

$$\vec{F} = q\vec{E} + q\vec{v} \times \vec{B} + \frac{m^*\vec{v}}{\tau} = 0$$

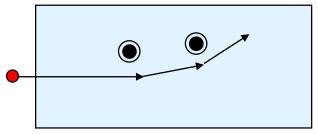
$$m^* \mathbf{v} = -q\tau \mathbf{\mathcal{E}} - q\tau \mathbf{v} \times \mathbf{B}$$

$$\approx -q\tau \mathcal{E} - q\tau \left( -\frac{q\tau \mathcal{E}}{m^*} \right) \times \mathbf{B}$$

$$= -q\tau \mathbf{\mathcal{E}} + \frac{q^2\tau^2}{m^*} \mathbf{\mathcal{E}} \times \mathbf{B}$$

$$\upsilon = -\frac{q\tau \mathbf{\mathcal{E}}}{m^*} + \frac{q^2\tau^2}{m^{*2}}\mathbf{\mathcal{E}} \times \mathbf{B}$$
Must be "small"

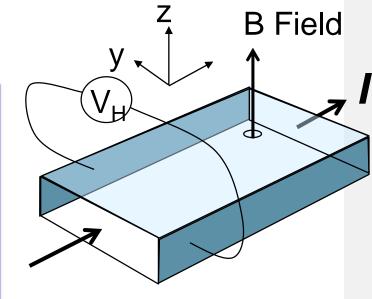
\* Same model works for holes, but with +q instead of -q



Approximation for Weak **B=0** field ...

$$-q\mathcal{E} - \frac{m^* \mathbf{v}}{\tau} \approx 0$$

$$\mathbf{v} = \frac{-q\tau \mathcal{E}}{m^*}$$



Approximation works iff ....

$$\frac{q^2\tau^2B}{m^*q\tau} = \frac{q\tau B}{m^*} \equiv \tau\omega_c << 1$$

Limit of strong scattering

Time for one magnetic orbit  $\frac{1}{\omega_c} \gg \tau$  Relaxation time

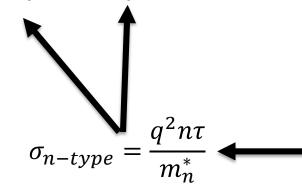
#### Drude Model & Hall Effect ...

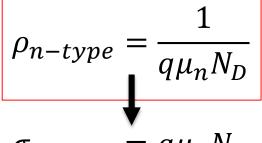
$$\upsilon = -\frac{q\tau \mathbf{\mathcal{E}}}{m^*} + \frac{q^2\tau^2}{m^{*2}} \mathbf{\mathcal{E}} \times \mathbf{B}$$

$$\mathbf{J}_n = -qn\mathbf{v}$$

$$= \frac{q^2 n \tau}{m^*} \mathcal{E} - \frac{q^2 n \tau}{m^*} \frac{q \tau}{m^*} \mathcal{E} \times \mathbf{B}$$

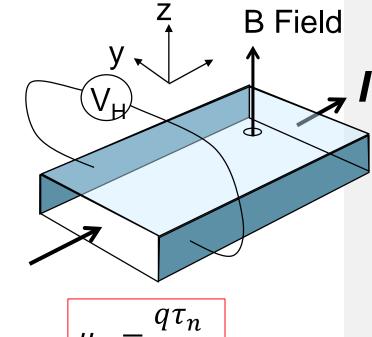
$$= \sigma_0 \mathbf{\mathcal{E}} - \sigma_0 \mu \mathbf{\mathcal{E}} \times \mathbf{B}$$





$$\sigma_{n-type} = q\mu_n N_D$$

$$\sigma_{n-type} = q \frac{q \tau_n}{m_n^*} n$$



$$\mu_n - \overline{m_n^*}$$

$$N_D \approx n$$



#### Drude Model & Hall Effect ...

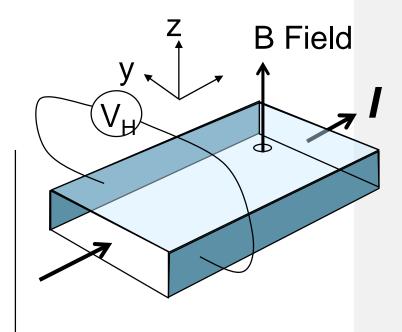
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$$\mathbf{J}_{n} = -qn\mathbf{v}$$

$$= \frac{q^2 n \tau}{m^*} \mathcal{E} - \frac{q^2 n \tau}{m^*} \frac{q \tau}{m^*} \mathcal{E} \times \mathbf{B}$$

$$= \sigma_0 \mathcal{E} - \sigma_0 \mu \mathcal{E} \times \mathbf{B}$$

$$\left| \begin{array}{cccc} \overline{x} & \overline{y} & \overline{z} \\ E_x & E_y & E_z \\ B_x & B_y & B_z \end{array} \right|$$



$$N_D \approx n$$

$$\begin{bmatrix} \boldsymbol{J}_{x} \\ \boldsymbol{J}_{y} \end{bmatrix} = \begin{bmatrix} \boldsymbol{\sigma}_{0} & -\boldsymbol{\sigma}_{0} \mu \boldsymbol{B}_{z} \\ \boldsymbol{\sigma}_{0} \mu \boldsymbol{B}_{z} & \boldsymbol{\sigma}_{0} \end{bmatrix} \begin{bmatrix} \boldsymbol{E}_{x} \\ \boldsymbol{E}_{y} \end{bmatrix}$$

$$\boldsymbol{\sigma}_{n-type} = \frac{q^{2}n\tau}{m_{n}^{*}}$$

$$\boldsymbol{\rho}_{n-type} = \frac{1}{q\mu_{n}N_{D}}$$

~0, since voltmeters have a very large internal resistance so very little current flows through

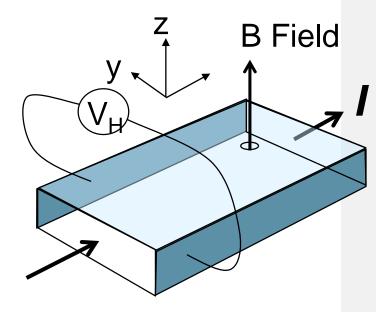
$$\begin{bmatrix} J_x \\ J_y \end{bmatrix} = \begin{bmatrix} \sigma_0 \\ \sigma_0 \mu B_z \end{bmatrix} - \sigma_0 \mu B_z \begin{bmatrix} E_x \\ E_y \end{bmatrix}$$

$$\begin{bmatrix} J_x \\ J_y \end{bmatrix} \approx \begin{bmatrix} \sigma_0 & 0 \\ \sigma_0 \mu B_z & \sigma_0 \end{bmatrix} \begin{bmatrix} E_x \\ E_y \end{bmatrix}$$

$$\begin{bmatrix} J_x \\ 0 \end{bmatrix} = \begin{bmatrix} \sigma_0 & 0 \\ \sigma_0 \mu B_z & \sigma_0 \end{bmatrix} \begin{bmatrix} E_x \\ E_y \end{bmatrix}$$

$$R_{H} = \frac{E_{y}/B_{z}}{J_{x}} = -\frac{1}{qn}$$

In the limit of small B



 $B_z$  is known,  $J_x$  is measured,  $V_h$  is measured =>  $E_v$  is known

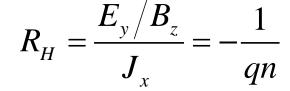
$$N_D pprox n$$
  $\mu_n = rac{q au_n}{m_n^*}$   $\sigma_{n-type} = rac{q^2n au}{q\mu_n N_D}$ 

Why Bz set to zero in one eq and not the other? => analytical simplicity.

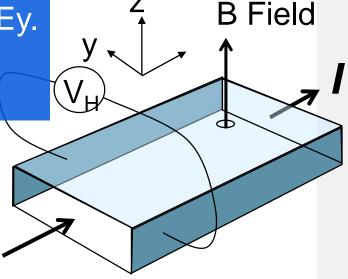
The system is also explicitly solvable in terms of Bz^2 and Ey.

Once Bz^z/Ey is small compared to 1/qn then one can neglect higher orders and come to the same result

$$\begin{bmatrix} J_{x} \\ J_{y} \end{bmatrix} \approx \begin{bmatrix} \sigma_{0} & 0 \\ \sigma_{0} \mu B_{z} & \sigma_{0} \end{bmatrix} \begin{bmatrix} E_{x} \\ E_{y} \end{bmatrix}$$



In the limit of small B



 $B_z$  is known,  $J_x$  is measured,  $V_h$  is measured =>  $E_v$  is known

$$N_D pprox n$$
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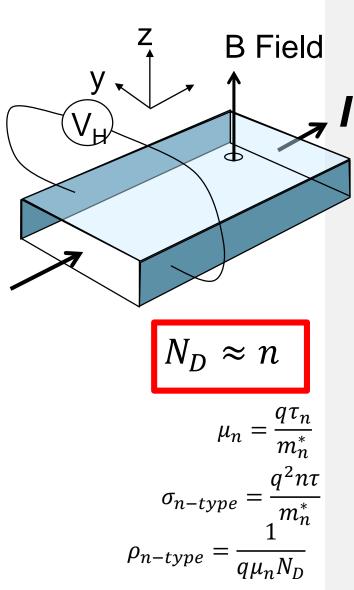
$$R_{H} = \frac{E_{y}/B_{z}}{J_{x}} = -\frac{1}{qn}$$

 $B_z$  is known,  $J_x$  is measured,  $V_h$  is measured =>  $E_y$  is known

By a simple electrical measurement, we now know the concentration of electrons in the sample.

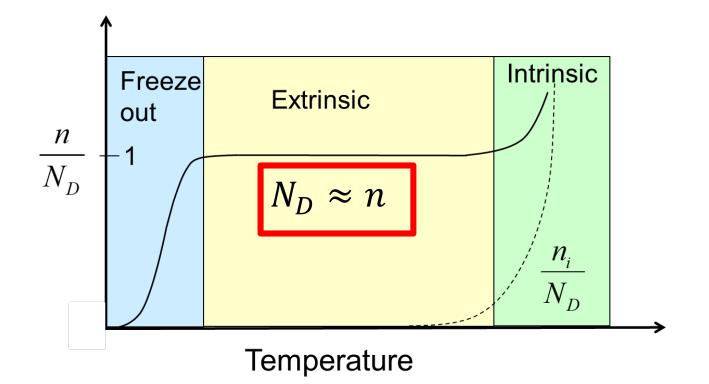
From the intrinsic concentration of carriers in the semiconductor and temperature we can deduce the doping concentration in the sample

#### In the limit of small B



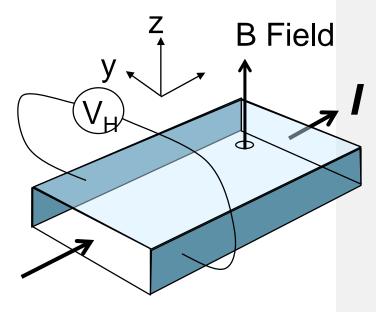
$$R_{H} = \frac{E_{y}/B_{z}}{J_{x}} = -\frac{1}{qn}$$

 $B_z$  is known,  $J_x$  is measured,  $V_h$  is measured =>  $E_y$  is known



Need multiple temperature dependent measurements to determine doping

#### In the limit of small B

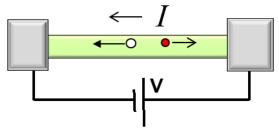


$$\mu_n = \frac{q\tau_n}{m_n^*}$$

$$\sigma_{n-type} = \frac{q^2n\tau}{m_n^*}$$

$$\rho_{n-type} = \frac{1}{q\mu_n N_D}$$

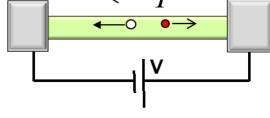
# Intro to Transport - Drift, Mobility, Diffusion, Einstein Relationship



 $J_n = qn\mu_n \mathcal{E}$ 

Transport with scattering, non-equilibrium Stat. Mech.

Drift-diffusion equation with recombination-generation



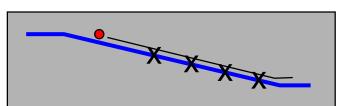
- 17.1 Drift Current
- 17.2 Mobility »Matthiessen Rule

  - »Mobility Measurement



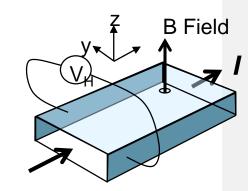


• 17.4

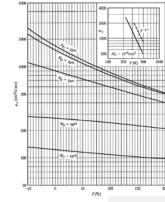




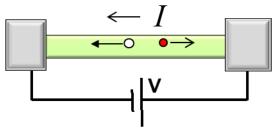
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$$R_H = \frac{E_y / B_z}{J_x} = -\frac{1}{qn}$$

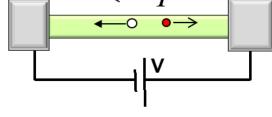


## Intro to Transport - Drift, Mobility, Diffusion, Einstein Relationship



Transport with scattering, non-equilibrium Stat. Mech.

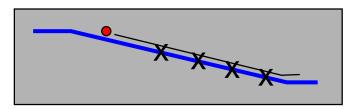
Drift-diffusion equation with recombination-generation



- 17.1 Drift Current
- 17.2 Mobility »Matthiessen Rule
  - »High Field Effects
  - »Mobility Measurement

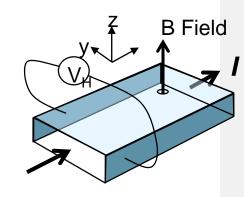
 $J_n = qn\mu_n \mathcal{E}$ 

$$\frac{1}{\mu_n} = \frac{1}{\mu_{ph}} + \frac{1}{\mu_{II}}$$





- 17.3 Carrier Concentration from Hall Effect
- 17.4 Physics of diffusion Einstein Relationship



$$R_H = \frac{E_y/B_z}{J_x} = -\frac{1}{qn}$$

