

Section 7

Bandstructure in 1D Periodic Potentials

Gerhard Klimeck

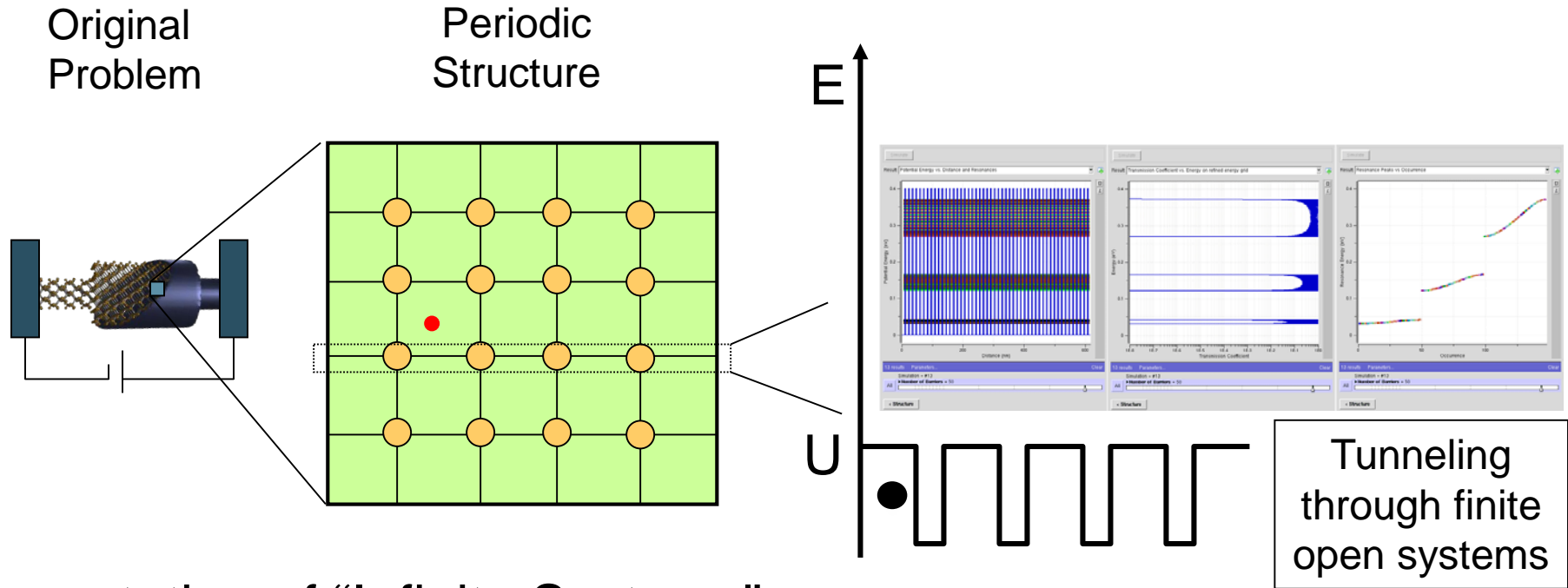
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School of Electrical and
Computer Engineering

Section 7

Bandstructure - in 1D Periodic Potentials



Goal: Representation of “Infinite Systems”

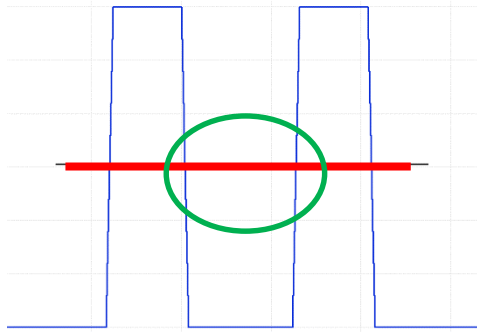
- 7.1 Bandstructure - Problem Formulation
- 7.2 Bandstructure - Solutions
- 7.3 Band Properties

Reference: Vol. 6, Ch. 3

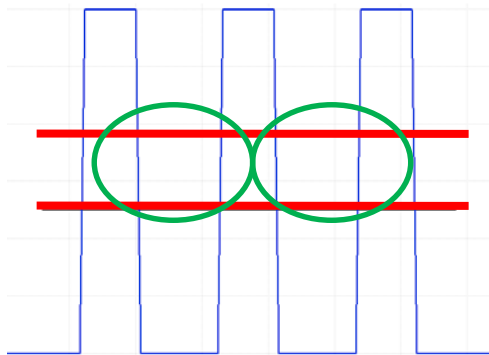
Daniel Mejia, Gerhard Klimeck (2019), "Periodic Potential Lab - Kronig Penney Model," <https://nanohub.org/resources/kronigpenneylab>. (DOI: 10.21981/TT2Y-A185).

Reminder

Transmission through Repeated Wells

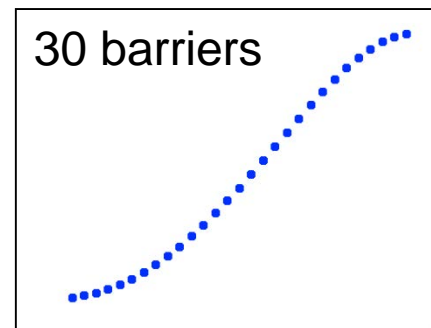
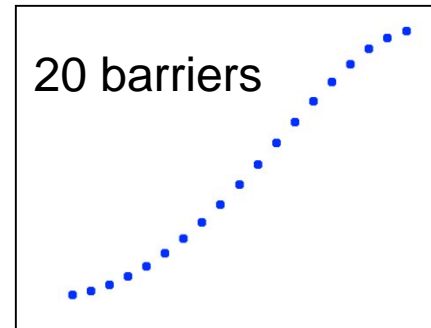
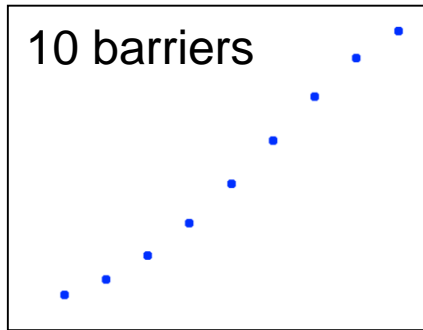


2 barriers => 1 resonance



3 barriers => 2 resonance

n barriers => $n-1$ resonance

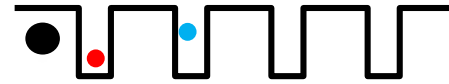


As the number of barriers are increased more and more energy resonances begin to appear and energy bands are formed.

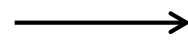
What if the number of wells is infinite?

Reminder: Five Steps for ~~Closed~~ System Analytical Solution

Open



1) $\frac{d^2\psi}{dx^2} + k^2\psi = 0$

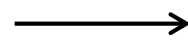


Solution Ansatz
2N unknowns
for N regions

$$\psi(x) = A_+ e^{ikx} + A_- e^{-ikx}$$

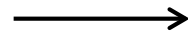
$$\psi(x) = D e^{-\alpha x} + E e^{+\alpha x}$$

2) ~~$\psi(x = -\infty) = 0$~~
 ~~$\psi(x = +\infty) = 0$~~



Boundary Conditions at the **Open** edge
Reduces 2 unknowns

3) $\psi|_{x=x_B^-} = \psi|_{x=x_B^+}$
 $\frac{d\psi}{dx}|_{x=x_B^-} = \frac{d\psi}{dx}|_{x=x_B^+}$



Boundary Condition at each interface:
Set 2N-2 equations for
2N-2 unknowns (for continuous U)

4) Det (coefficient matrix)=0
~~And find E by graphical~~
~~or numerical solution~~

5) ~~$\int_{-\infty}^{\infty} |\psi(x, E)|^2 dx = 1$~~

~~Normalization of unity probability
for wave function~~

Reminder: Five Steps for Closed System Analytical Solution



1) $\frac{d^2\psi}{dx^2} + k^2\psi = 0$

~~Open~~
Periodic Solution Ansatz
2N unknowns
for N regions

$$\psi(x) = A_+ e^{ikx} + A_- e^{-ikx}$$

$$\psi(x) = D e^{-\alpha x} + E e^{+\alpha x}$$

Bloch Theorem

2) ~~$\psi(x = -\infty) = 0$~~
 ~~$\psi(x = +\infty) = 0$~~

Edge => huge number of cells
Boundary Conditions at the ~~Open~~ edge
Reduces 2 unknowns **Periodic**

3) $\psi|_{x=x_B^-} = \psi|_{x=x_B^+}$
 $\frac{d\psi}{dx}|_{x=x_B^-} = \frac{d\psi}{dx}|_{x=x_B^+}$

Boundary Condition at each interface:
Interface => interfaces in periodic cell

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~~And find E by graphical~~
~~or numerical solution~~

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~~Normalization of unity probability~~
~~for wave function~~

Reminder:

Five Steps for ~~Closed~~ System Analytical Solution



1) $\frac{d^2\psi}{dx^2} + k^2\psi = 0$

~~Open~~
Periodic Solution Ansatz
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Boundary Condition at each interface:

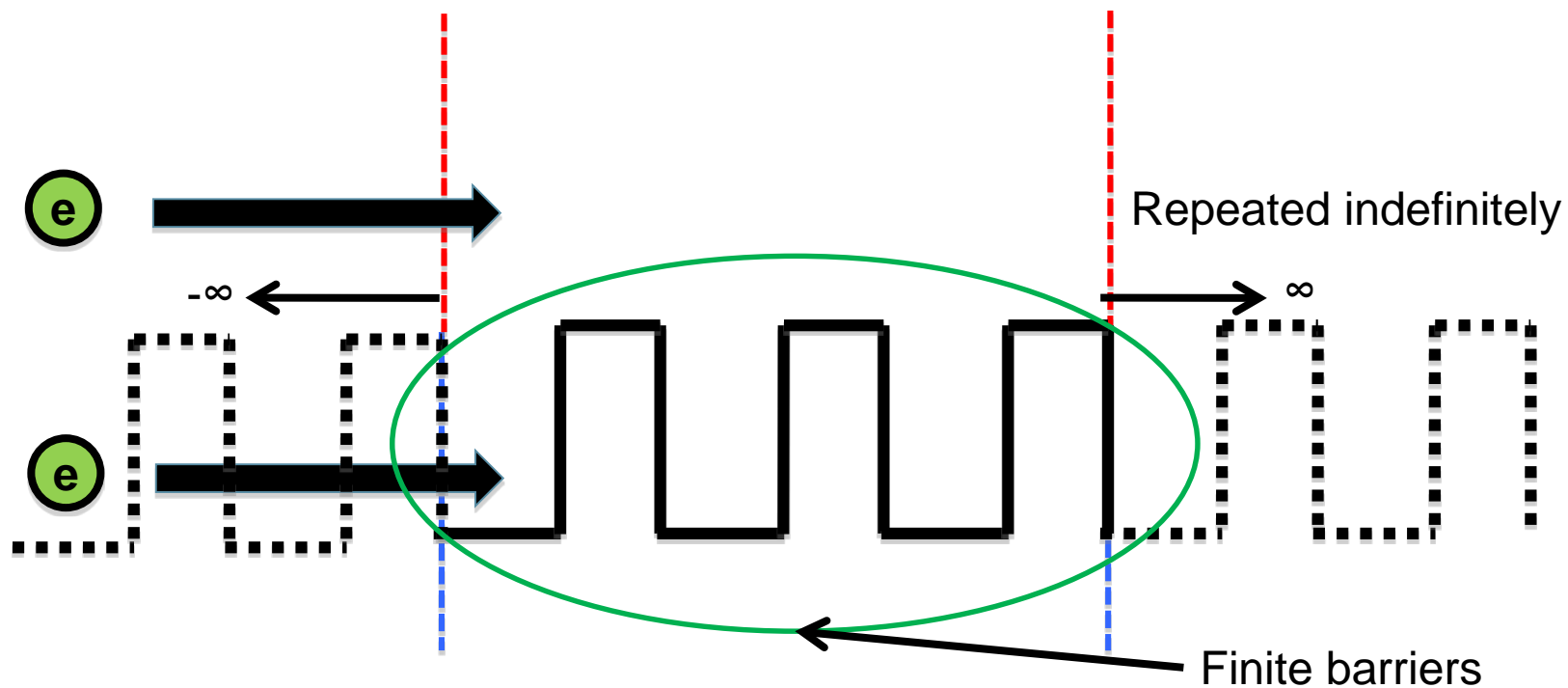
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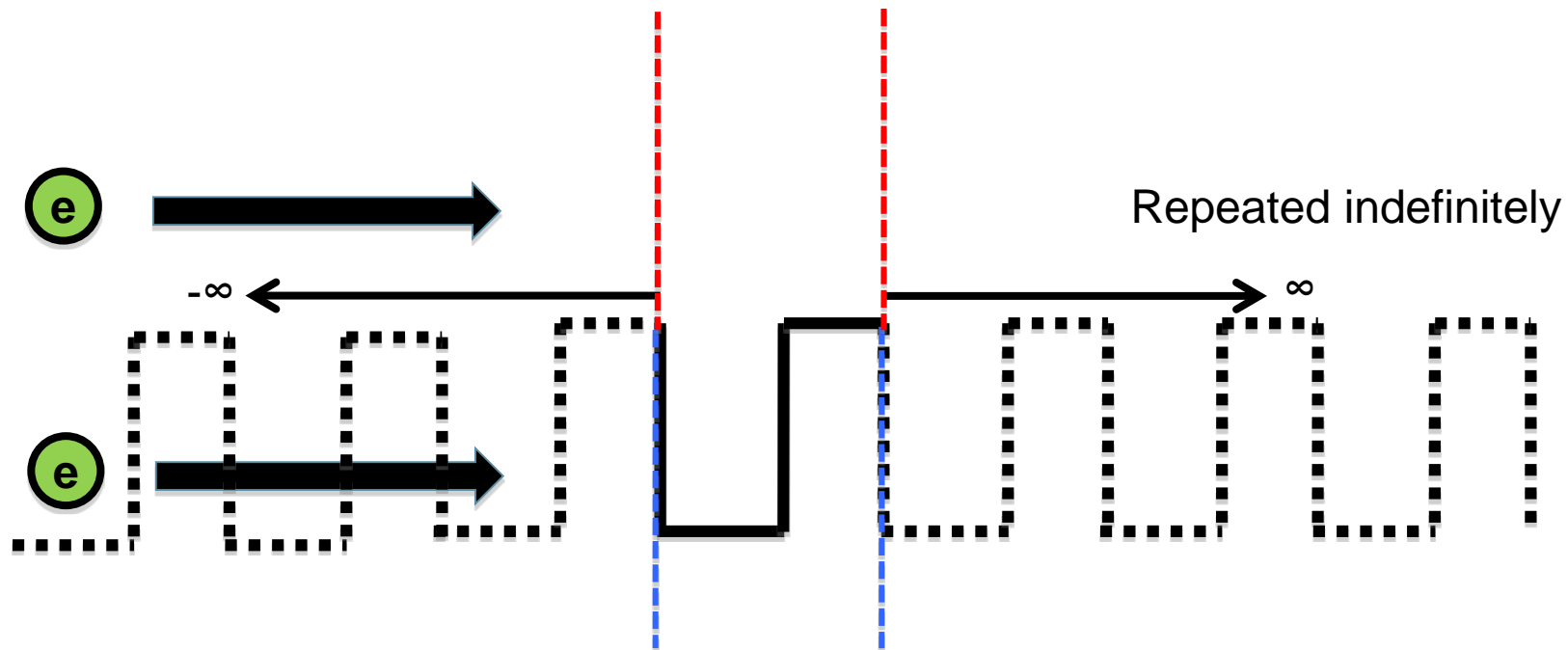
~~Normalization of unity probability
 for wave function~~

Periodic Potential Concept



As the number of barriers is increased the electrons see no difference between the actual structure and a structure that is simply modeled as being repeated indefinitely (Periodic).

Choosing the Smallest Unit Cell



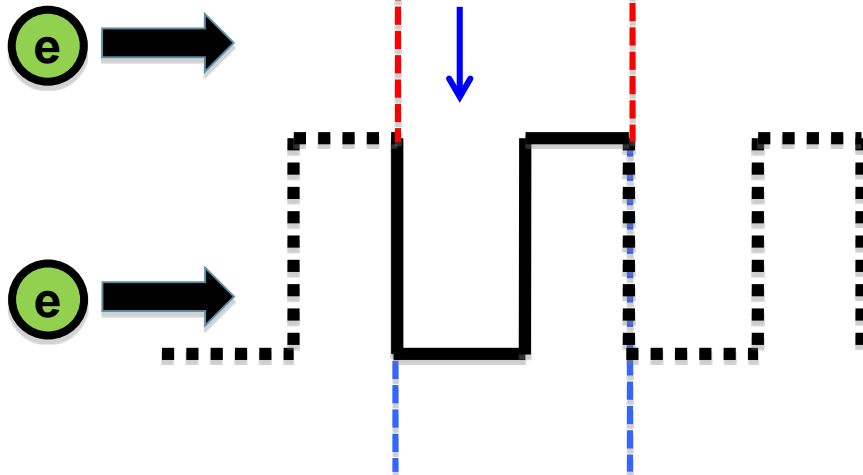
As the number of barriers is increased the electrons see no difference between the actual structure and a structure that is simply modeled as being repeated indefinitely (Periodic).

Choose the smallest cell to reduce computational work load

Solution Ansatz

Choose the Simplest Basis Set

$$\frac{d^2 \psi}{dx^2} + k^2 \psi = 0$$



Propagation above barriers

$$\psi_{n+1} = A_{n+1} \sin kx + B_{n+1} \cos kx$$

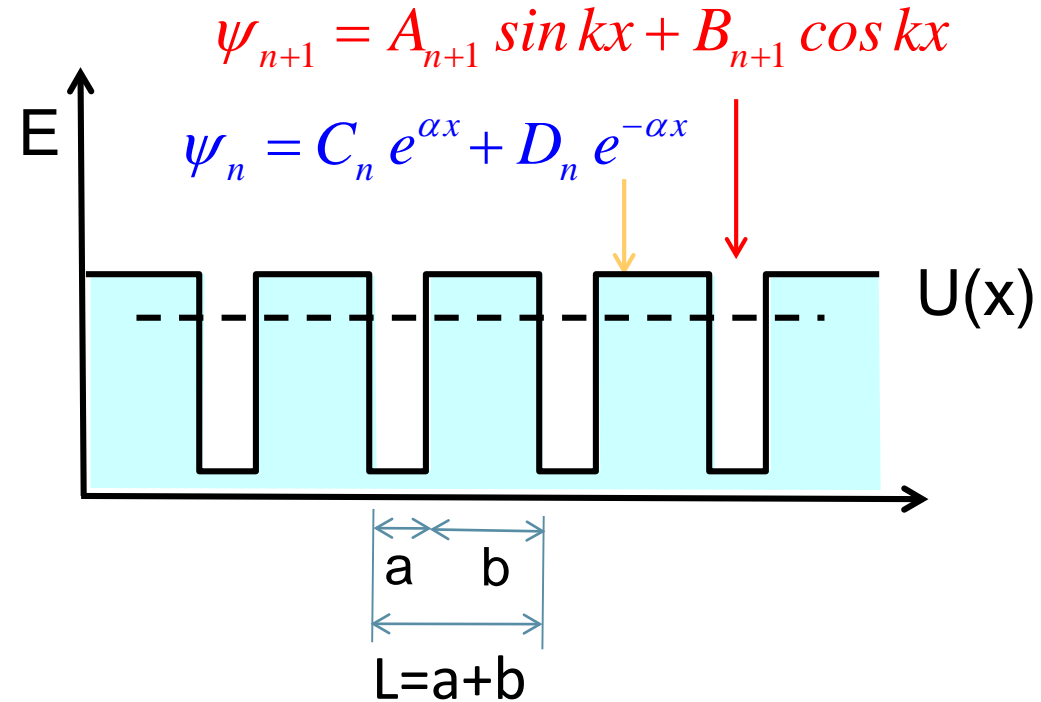
Decays in barriers

$$\psi_n = C_n e^{\alpha x} + D_n e^{-\alpha x}$$

As the number of barriers is increased the electrons see no difference between the actual structure and a structure that is simply modeled as being repeated indefinitely (Periodic).

Choose the smallest cell to reduce computational work load

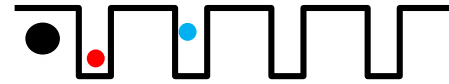
Finally an (almost) Real Problem ...



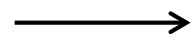
But N atoms have two $2N$ unknown constants to find ...
For large N , isn't there a better way?

Reminder: Five Steps for Closed System Analytical Solution

~~Periodic~~



1) $\frac{d^2\psi}{dx^2} + k^2\psi = 0$



Solution Ansatz
2N unknowns
for N regions

$$\psi(x) = A_+ e^{ikx} + A_- e^{-ikx}$$

$$\psi(x) = D e^{-\alpha x} + E e^{+\alpha x}$$

~~Bloch Theorem~~

2) ~~$\psi(x = -\infty) = 0$~~
 ~~$\psi(x = +\infty) = 0$~~

Edge => huge number of cells

Boundary Conditions at the **Periodic** edge
Reduces 2 unknowns

3) $\psi|_{x=x_B^-} = \psi|_{x=x_B^+}$
 $\frac{d\psi}{dx}|_{x=x_B^-} = \frac{d\psi}{dx}|_{x=x_B^+}$



Boundary Condition at each interface:

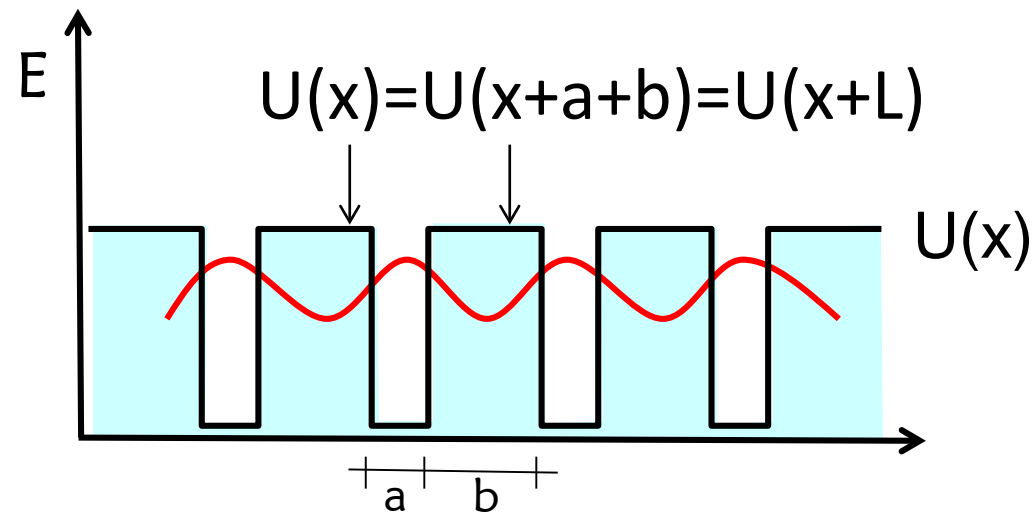
Interface => interfaces in periodic cell

4) Det (coefficient matrix)=0
And find E by graphical
or numerical solution

5) ~~$\int_{-\infty}^{\infty} |\psi(x, E)|^2 dx = 1$~~

~~Normalization of unity probability
for wave function~~

Periodic U(x) and Bloch's Theorem

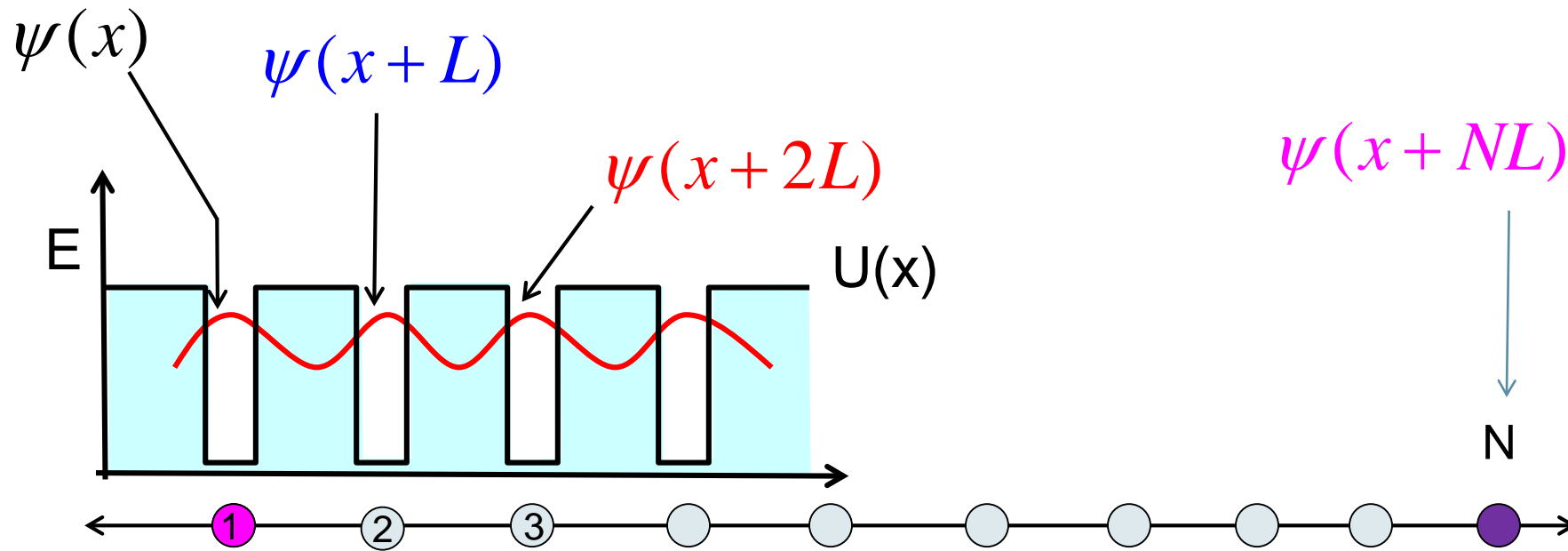


$$|\psi(x)|^2 = |\psi(x+p)|^2 \quad \Rightarrow \quad \psi(x+p) = \psi(x) e^{ikL}$$

p = some period length

not our old (k)

Phase-factor for N-cells



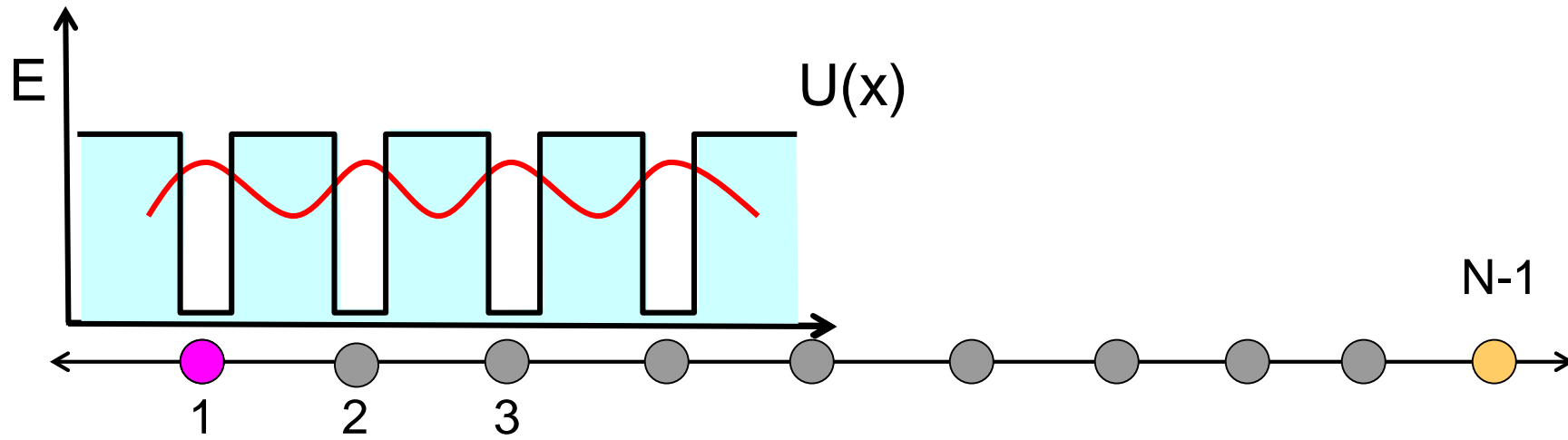
$$\psi[x+L] = \psi(x)e^{ikL}$$

$$\psi[x+2L] = \psi(x+L)e^{ikL}$$

$$= \psi(x)e^{ikL \times 2}$$

$$\psi[x+NL] = \psi(x)e^{ikLN}$$

Step 2: Periodic Boundary Condition

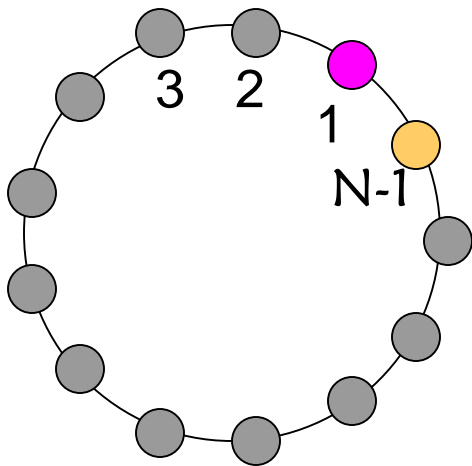


$$\psi[x + NL] = \psi(x)e^{ikLN}$$

$$e^{ikLN} = 1 \equiv e^{\pm i2\pi n}$$

$$k = \pm \frac{2\pi n}{NL} \quad n = -\frac{N}{2}, \dots, -1, 0, 1, \dots, \frac{N}{2}$$

$$k_{\max} = \frac{\pi}{L}, \quad k_{\min} = -\frac{\pi}{L}$$



Step 3: Boundary Conditions

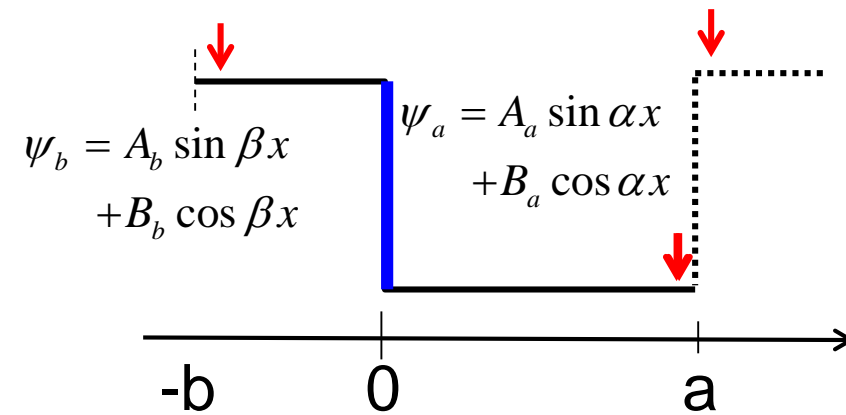
$$\psi|_{x=0^-} = \psi|_{x=0^+}$$

$$\frac{d\psi}{dx}\bigg|_{x=0^-} = \frac{d\psi}{dx}\bigg|_{x=0^+}$$

$$B_a = B_b$$

$$\alpha A_a = \beta A_b$$

$$\beta \equiv i\sqrt{2m(U_0 - E)/\hbar^2} \quad \alpha \equiv \sqrt{2mE/\hbar^2}$$



$$\psi_a|_{x=a} = \psi_b|_{x=-b} e^{ikL}$$

$$\frac{d\psi_a}{dx}\bigg|_{x=a} = \frac{d\psi_b}{dx}\bigg|_{x=-b} e^{ikL}$$

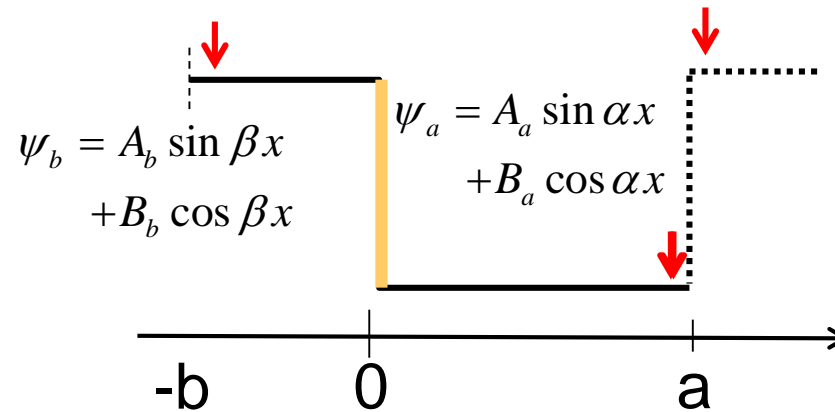
$$A_a \sin \alpha a + B_a \cos \alpha a = e^{ik(a+b)} [-A_b \sin \beta b + B_b \cos \beta b]$$

$$\alpha A_a \sin \alpha a - \alpha B_a \cos \alpha a = e^{ik(a+b)} [\beta A_b \sin \beta b + \beta B_b \cos \beta b]$$

Step 4: Det(matrix)=0 for Energy-levels

$$B_a = B_b$$

$$\alpha A_a = \beta A_b$$

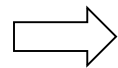


$$A_a \sin \alpha a + B_a \cos \alpha a =$$

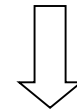
$$e^{ik(a+b)} [-A_b \sin \beta b + B_b \cos \beta b]$$

$$\alpha A_a \sin \alpha a - \alpha B_a \cos \alpha a =$$

$$e^{ik(a+b)} [\beta A_b \sin \beta b + \beta B_b \cos \beta b]$$



$$4) \begin{pmatrix} 0 & 1 & 0 & -1 \\ \alpha & 0 & \beta & 0 \\ * & * & & \\ * & & & \end{pmatrix} \begin{bmatrix} A_a \\ B_a \\ A_b \\ B_b \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \\ 0 \end{bmatrix}$$



$$\frac{1-2\xi}{2\xi\sqrt{1-\xi}} \times \dots = \cos kL \quad \xi \equiv \frac{E}{U_0} \quad \alpha_0 \equiv \sqrt{\frac{2mU_0}{\hbar^2}}$$

Reminder: Five Steps for Closed System Analytical Solution

1) $\frac{d^2\psi}{dx^2} + k^2\psi = 0$

~~Periodic~~

Solution Ansatz
2N unknowns
for N regions

$$\psi(x) = A_+ e^{ikx} + A_- e^{-ikx}$$

$$\psi(x) = D e^{-\alpha x} + E e^{+\alpha x}$$

Bloch Theorem

2) ~~$\psi(x = -\infty) = 0$~~
 ~~$\psi(x = +\infty) = 0$~~

Edge => huge number of cells

Boundary Conditions at the **Periodic** edge
Reduces 2 unknowns

3) $\psi|_{x=x_B^-} = \psi|_{x=x_B^+}$
 $\frac{d\psi}{dx}|_{x=x_B^-} = \frac{d\psi}{dx}|_{x=x_B^+}$

Boundary Condition at each interface:

Interface => interfaces in periodic cell

4) Det (coefficient matrix)=0
And find E by graphical
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5) ~~$\int_{-\infty}^{\infty} |\psi(x, E)|^2 dx = 1$~~

~~Normalization of unity probability
for wave function~~

$$\frac{1-2\xi}{2\xi\sqrt{1-\xi}} \times \dots = \cos kL \quad \xi \equiv \frac{E}{U_0} \quad \alpha_0 \equiv \sqrt{\frac{2mU_0}{\hbar^2}}$$

Five Steps for **Periodic** System Analytical Solution

Bloch Theorem

1) $\frac{d^2\psi}{dx^2} + k^2\psi = 0$ \longrightarrow

Solution Ansatz
4 unknowns
for 1 periodic cell

$$\begin{aligned}\psi_a &= A_a \sin \alpha x \\ &+ B_a \cos \alpha x \\ \psi_b &= A_b \sin \beta x \\ &+ B_b \cos \beta x\end{aligned}$$

2) $\psi[x + NL] = \psi(x)e^{ikLN}$ \longleftarrow

Periodic Boundary Conditions
Imposes symmetry

3) $\psi_a|_{x=a} = \psi_b|_{x=-b} e^{ikL}$
 $\frac{d\psi_a}{dx}|_{x=a} = \frac{d\psi_b}{dx}|_{x=-b} e^{ikL}$ \longleftarrow

Boundary Condition in periodic cell:
Set 4 equations for 4 unknowns

4) Det (coefficient matrix)=0
And find E by graphical
or numerical solution

$$\frac{1-2\xi}{2\xi\sqrt{1-\xi}} \times \dots = \cos kL \quad \xi \equiv \frac{E}{U_0} \quad \alpha_0 \equiv \sqrt{\frac{2mU_0}{\hbar^2}}$$

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Bandstructure - in 1D Periodic Potentials

• 7.1 Bandstructure - Problem Formulation

- » Kronig-Penney Model setup
- » Bloch theorem
- » Analytical solution process



$$\psi[x + NL] = \psi(x)e^{ikLN}$$

$$\frac{1-2\xi}{2\xi\sqrt{1-\xi}} \times \dots = \cos kL \quad \xi \equiv \frac{E}{U_0} \quad \alpha_0 \equiv \sqrt{\frac{2mU_0}{\hbar^2}}$$

• 7.2 Bandstructure - Solutions

• 7.3 Band Properties

Reference: Vol. 6, Ch. 3

Daniel Mejia, Gerhard Klimeck (2019), "Periodic Potential Lab - Kronig Penney Model,

" <https://nanohub.org/resources/kronigpenneylab>. (DOI: 10.21981/TT2Y-A185).

One Video
Segment

One Video
Segment

One Video
Segment

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Bandstructure - in 1D Periodic Potentials

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- » Kronig-Penney Model setup
- » Bloch theorem
- » Analytical solution process

$$\psi[x + NL] = \psi(x)e^{ikLN}$$

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• 7.2 Bandstructure - Solutions

- » Bandgaps
- » Comparison to finite system model



• 7.3 Band Properties

Reference: Vol. 6, Ch. 3

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One Video Segment

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