
Theory of high resolution TEM

Lecture 15

Outline

Introductory level theory

- Weak-phase object approximation
- Contrast transfer function
- Envelope functions

Some advanced concepts

- Delocalization
- Using pass bands
- Focal series reconstruction
- Variable C_s Imaging

HRTEM

general idea

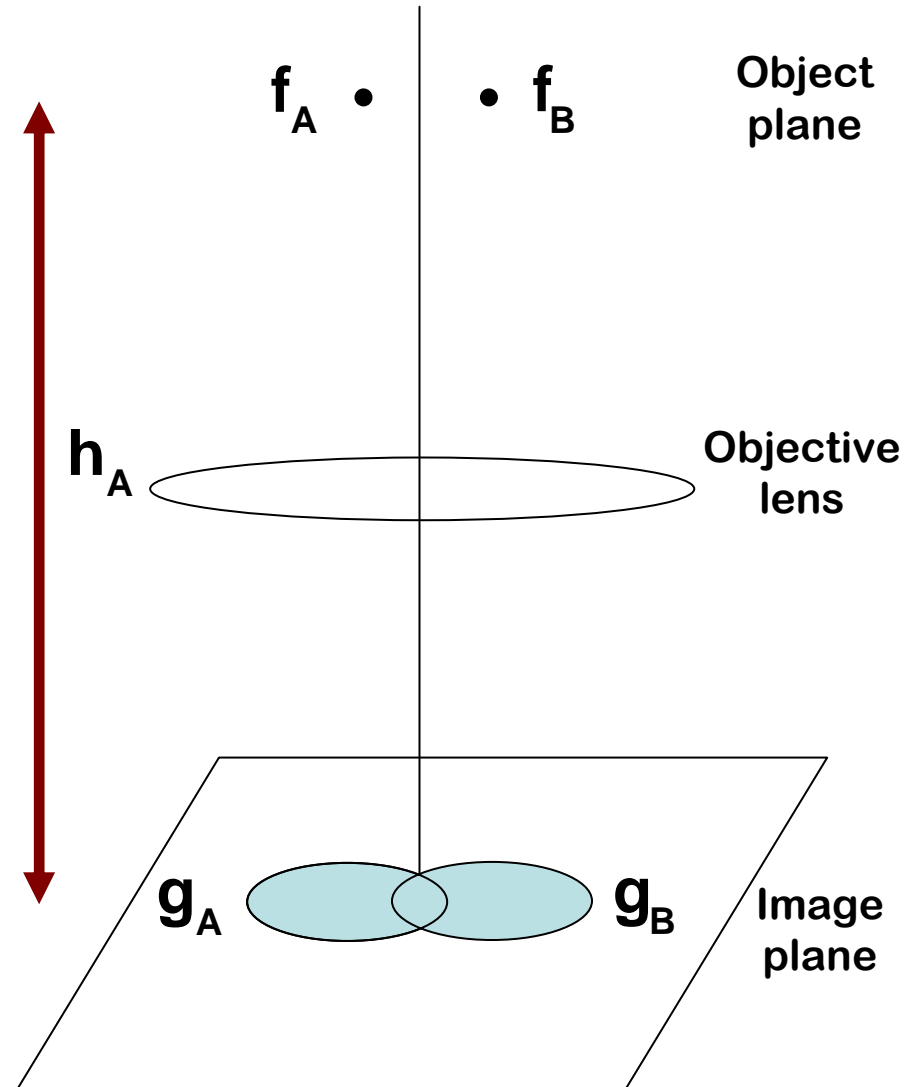
How do we express this mathematically?

$$g(\mathbf{r}) = \int f(\mathbf{r}') \otimes h(\mathbf{r} - \mathbf{r}') d\mathbf{r}'$$

- The \otimes is a mathematical symbol indicating 'convolution'*
- $h(\mathbf{r})$ is thus a 'blurring function'

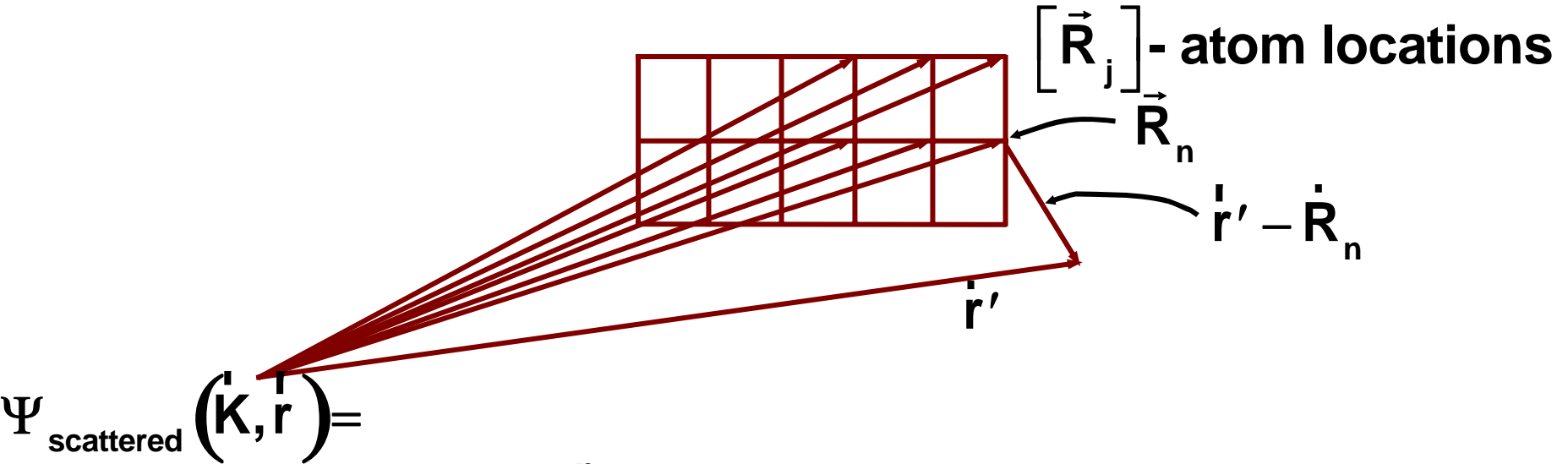
Or: aberrations, etc. limit spatial frequencies in image

↻ Reciprocal space



* <http://mathworld.wolfram.com/Convolution.html>

Scattering from a lattice



$$\Psi_{scattered}(\mathbf{K}, \mathbf{r}) =$$

$$\frac{-m}{2\pi\hbar^2} \frac{\exp\left[2\pi i(\mathbf{k}_d \cdot \mathbf{r})\right]}{|\mathbf{r}|} \int \sum v_{at}(\mathbf{r}' - \mathbf{R}_j) \exp\left[-2\pi i(\mathbf{K} \cdot \mathbf{r}')\right] d^3\mathbf{r}'$$

So scattered wave from an array of N atoms:

$$\Psi_{scatt}(\mathbf{K}) = \sum_j^N f_{el}(\mathbf{R}_j) \exp\left[-2\pi i(\mathbf{K} \cdot \mathbf{R}_j)\right]$$

Remember: ψ is thus the Fourier Transform of the scattering potential

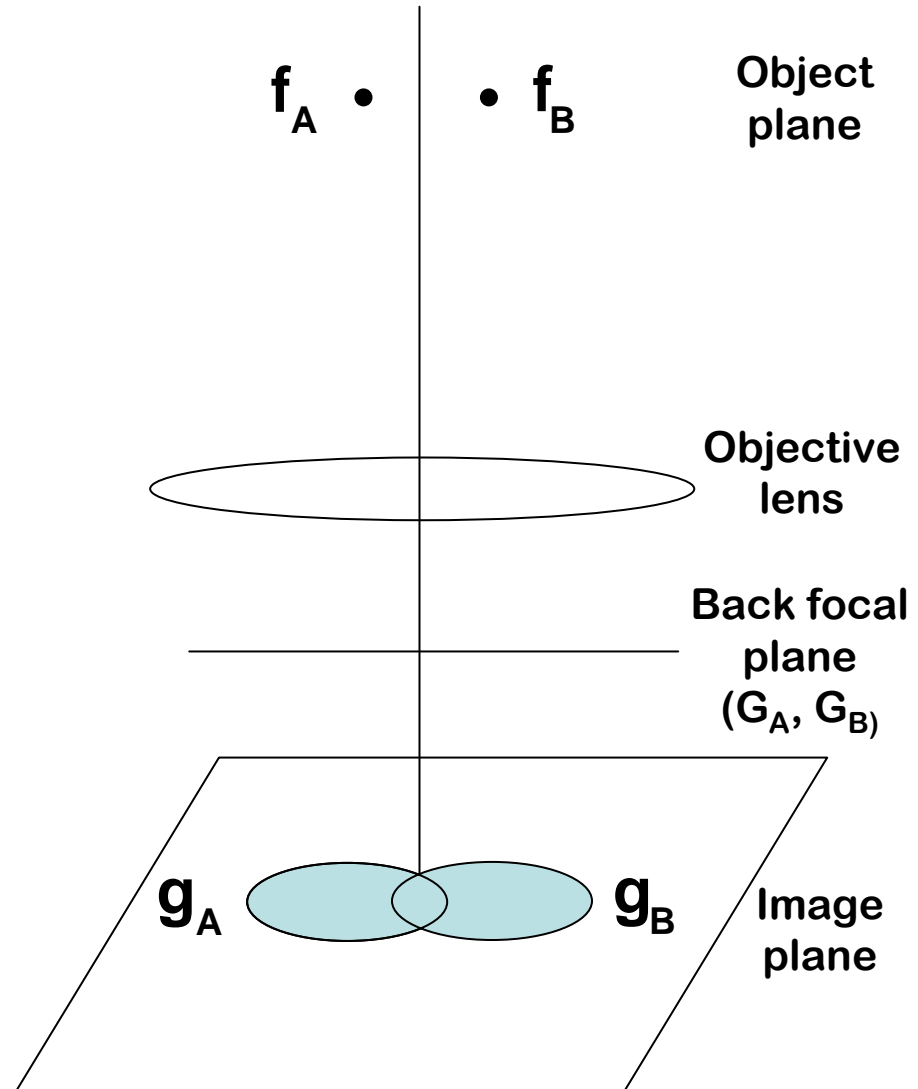
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So, can express $g(\vec{r})$ generally as:

$$g(\vec{r}) = \sum_{\vec{u}} G(\vec{u}) \exp[2\pi i(\vec{u} \cdot \vec{r})]$$

- \vec{u} are reciprocal lattice vectors
- $G(\vec{u})$ is scattering potential at each site
- Thus:
 - $g(\vec{r})$ is FT of $G(\vec{u})$



HRTEM

general idea

Similarly:

- $F(u)$ is FT of $f(r)$
- $H(u)$ is FT of $h(r)$

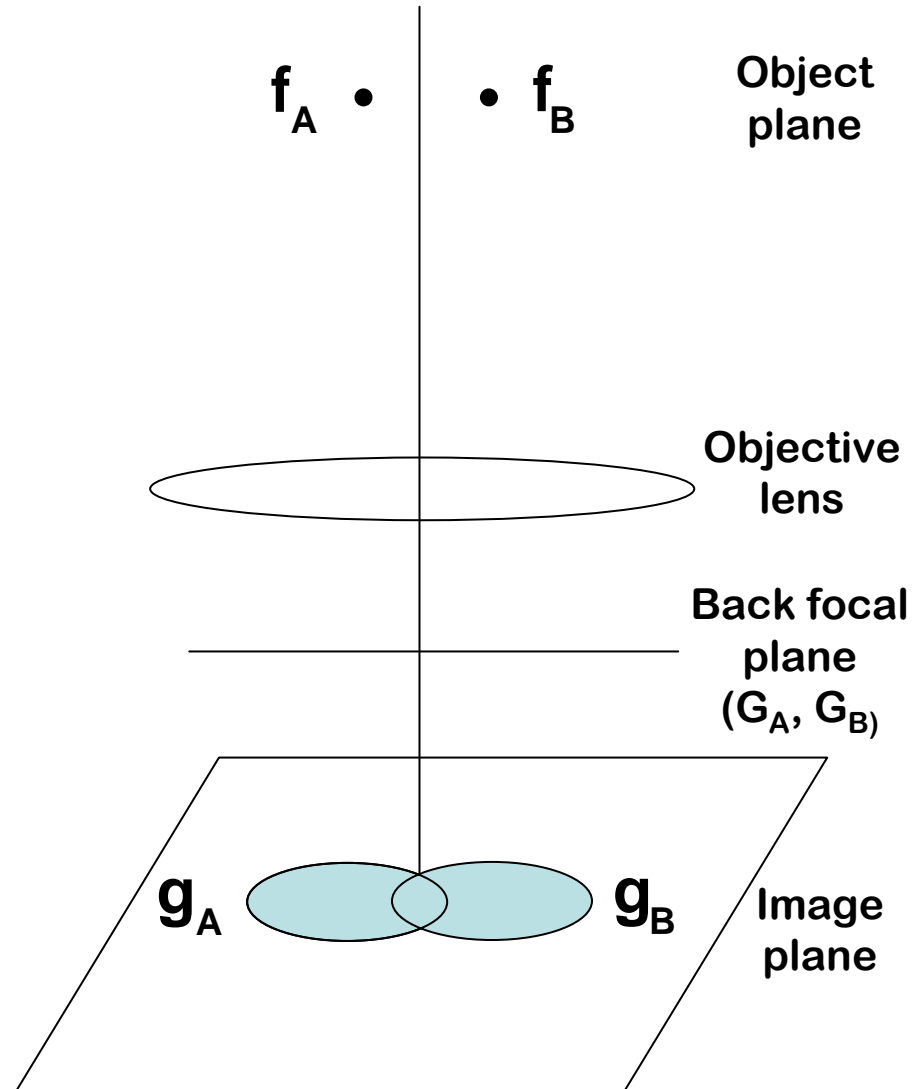
Key relation:

$$G(u) = H(u) \times F(u)$$

Convolution between
functions in **real space**

equals

Multiplication of functions
in **reciprocal space**



“Ideal” Specimen

General expression for specimen:

$$f(x, y) = A(x, y) \exp[-i\phi_t(x, y)]$$

Normalize to $A(x, y) = 1$

$$f(x, y) = \exp[-i\phi_t(x, y)]$$

Phase change depends only on mean inner potential through sample thickness

$$d\phi = \sigma \int_0^t V(x, y, z) dz = \sigma V_t(x, y)$$

Note: σ related to, but not same as cross-sections from earlier

Specimen transfer function

'weak phase object'

Can now have 'specimen transfer function'

- What's the wave look like as it exits sample - 'exit wave'
- This is the 'signal' we want in the end
- Include effect of absorption - $\mu(x,y)$
$$f(x,y) = \exp\left[-i\sigma V_t(x,y) - \mu(x,y)\right]$$
- Called a 'phase object'

"Weak phase object approximation" (WPOA)

- Assume sample very thin - $V_t(x,y) \ll 1$, neglect $\mu(x,y)$, expand above and neglect higher order terms:

$$f(x,y) = 1 - i\sigma V_t(x,y)$$

This is the final "model" of the sample

HRTEM

general idea

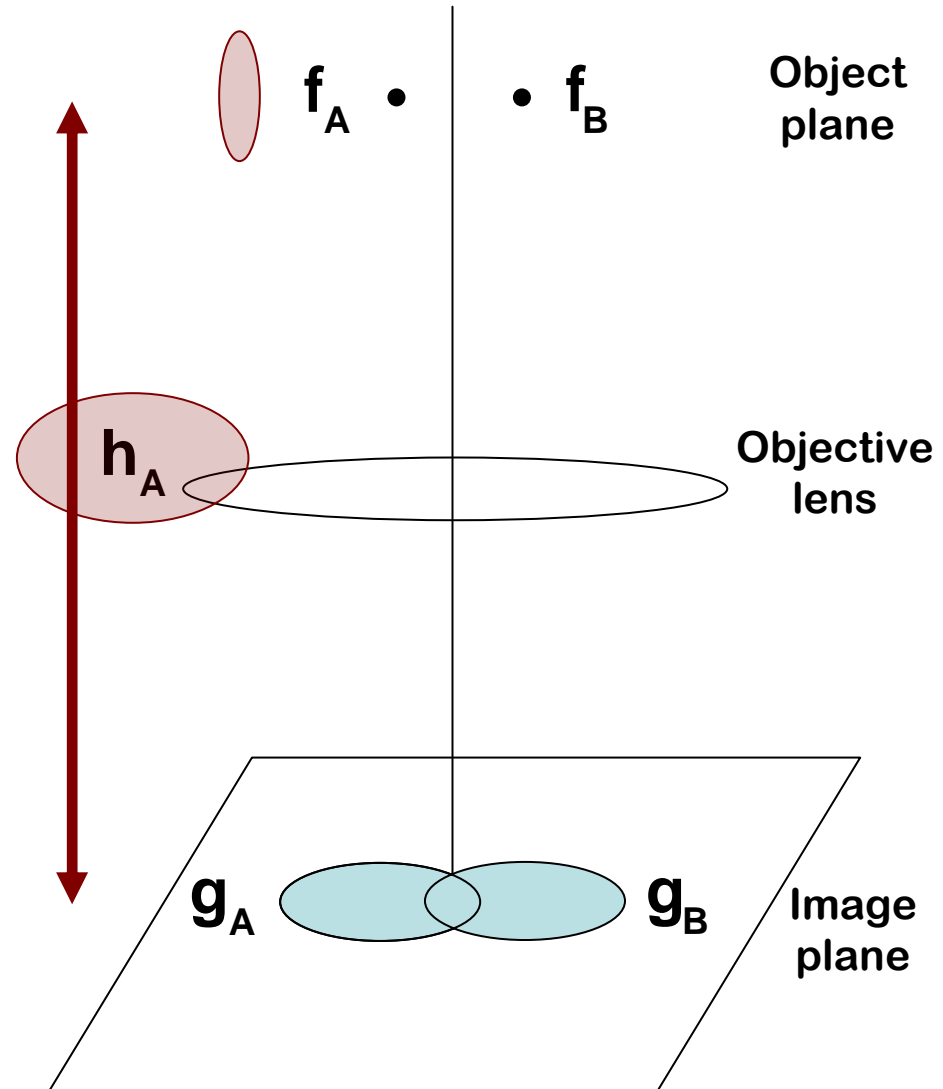
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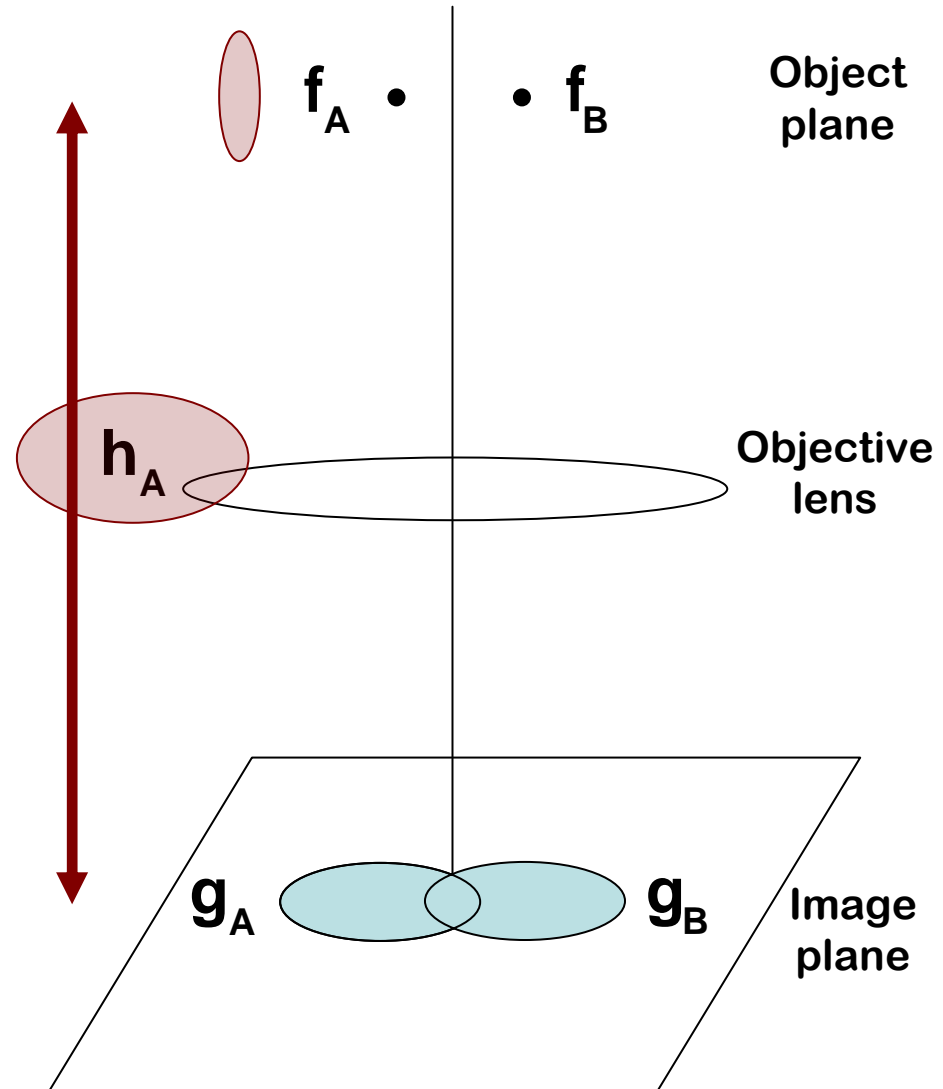
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WPOA & the Blurring Function

Electron exit wave is now modified by blurring function

$$\Psi(x, y) = [1 - i\sigma V_t(x, y)] \otimes h(x, y)$$

Represent $h(x, y)$ as general wave:

$$h(x, y) = \cos(x, y) + i \sin(x, y)$$

$$\Psi(x, y) = [1 - i\sigma V_t(x, y)] \otimes \sin(x, y) + i \cos(x, y)$$

$$\Psi(x, y) = 1 + \sigma V_t(x, y) \otimes \sin(x, y) - i \cos(x, y)$$

Intensity:

$$I = \Psi \Psi^* = 1 + 2\sigma V_t(x, y) \otimes \sin(x, y)$$

HRTEM

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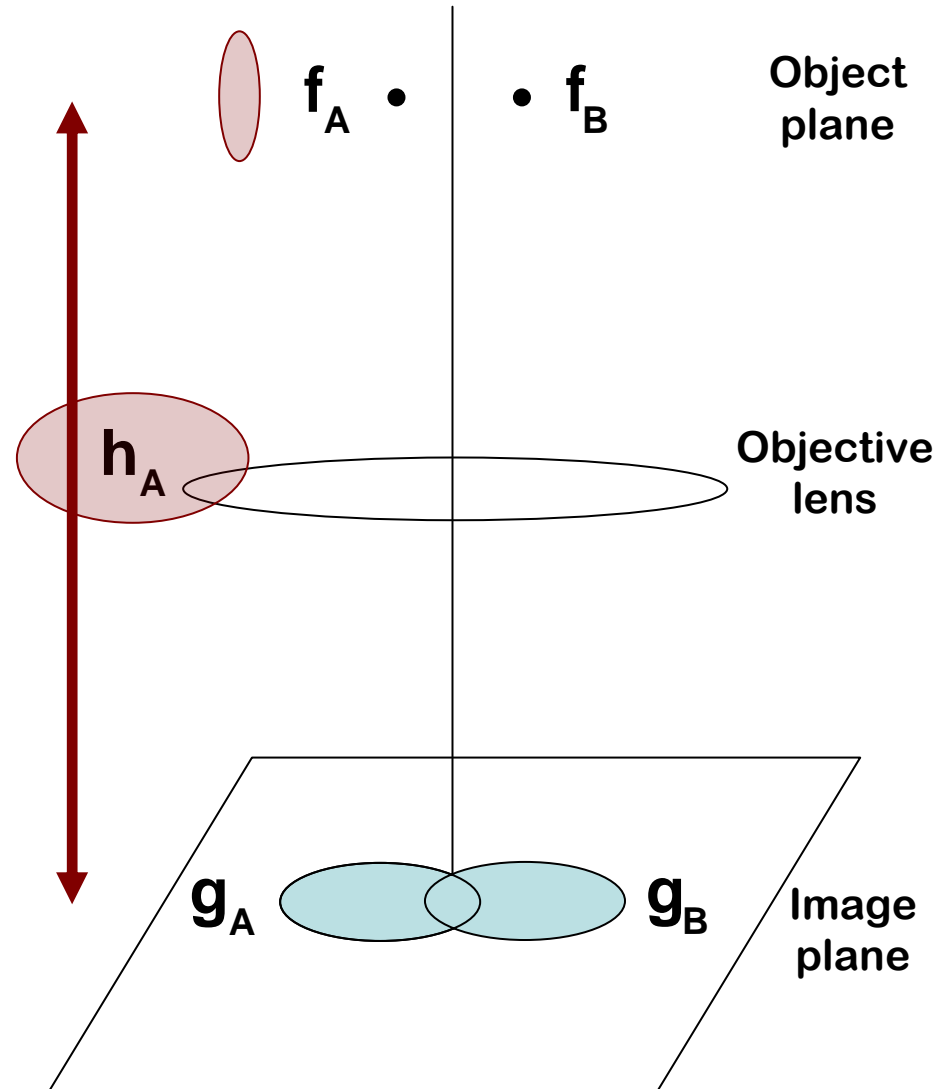
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Objective lens transfer function

What is internal to $H(u)$?

- Aperture function: $A(u)$
- Coherence function: $E(u)$
- Aberration function: $B(u)$

$$H(u) = A(u)E(u)B(u)$$

Let's dissect these terms ...

Aperture function: $A(u)$

- Determined by size of objective aperture
- Objective aperture (or column liner if no OA used) limits spatial frequencies

Objective lens transfer function

Aberration function: $B(u)$

- Consider only spherical aberration & defocus
- Point in object becomes disc of diameter $\delta(\theta)$ in image:

$$\delta(\theta) = C_s \theta^3 + \Delta f \theta$$

- Integrate over all θ :

$$D(\theta) = \int_0^\theta \delta(\theta) d\theta = \frac{C_s \theta^4}{4} + \frac{\Delta f \theta^2}{2}$$

- Replace θ w/ λu (Bragg's law, small angles $\Rightarrow 2\theta = \lambda u$)

$$D(u) = C_s \frac{\lambda^4 u^4}{4} + \Delta f \frac{\lambda^2 u^2}{2}$$

Objective lens transfer function

Interested in phase only: $\chi(\mathbf{u})$

$$\begin{aligned}\chi(\mathbf{u}) &= \frac{2\pi}{\lambda} \mathbf{D}(\mathbf{u}) \\ &= \frac{2\pi}{\lambda} \left[\mathbf{C}_s \frac{\lambda^4 \mathbf{u}^4}{4} + \Delta \mathbf{f} \frac{\lambda^2 \mathbf{u}^2}{2} \right] \\ \chi(\mathbf{u}) &= \frac{1}{2} \pi \mathbf{C}_s \lambda^3 \mathbf{u}^4 + \pi \Delta \mathbf{f} \lambda \mathbf{u}^2\end{aligned}$$

Aberration function $\chi(\mathbf{u})$ is a complicated function of \mathbf{C}_s and $\Delta \mathbf{f}$

Equivalence between real and reciprocal space

So - why did we do all of this?

Well, recall that in real space:

$$I = \Psi\Psi^* = 1 + 2\sigma V_t(x, y) \otimes \sin(x, y)$$

We make an equivalence in reciprocal space:

$$H(u) = 2A(u)E(u)\sin\chi(u)$$

Which means that our image - $G(u)$ - is:

$$G(u) = H(u)F(u) = \boxed{2A(u)E(u)\sin\chi(u)} F(u)$$

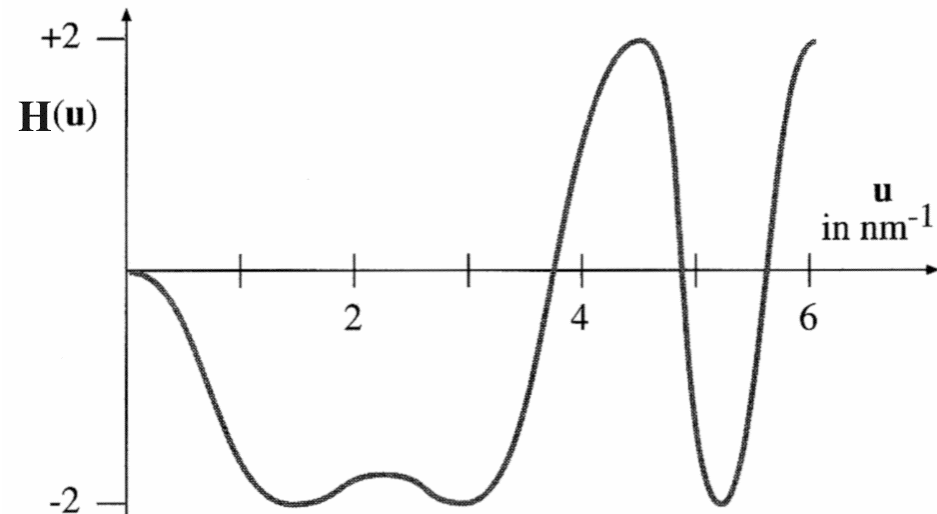
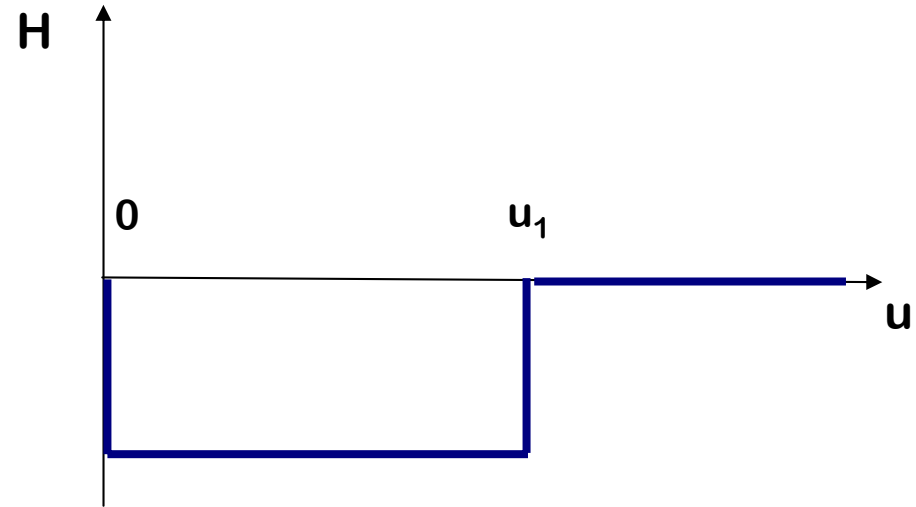
$H(u)$ is thus a spatial frequency filter

Contrast Transfer Function

Ideal transfer function passes all frequencies up to a given u

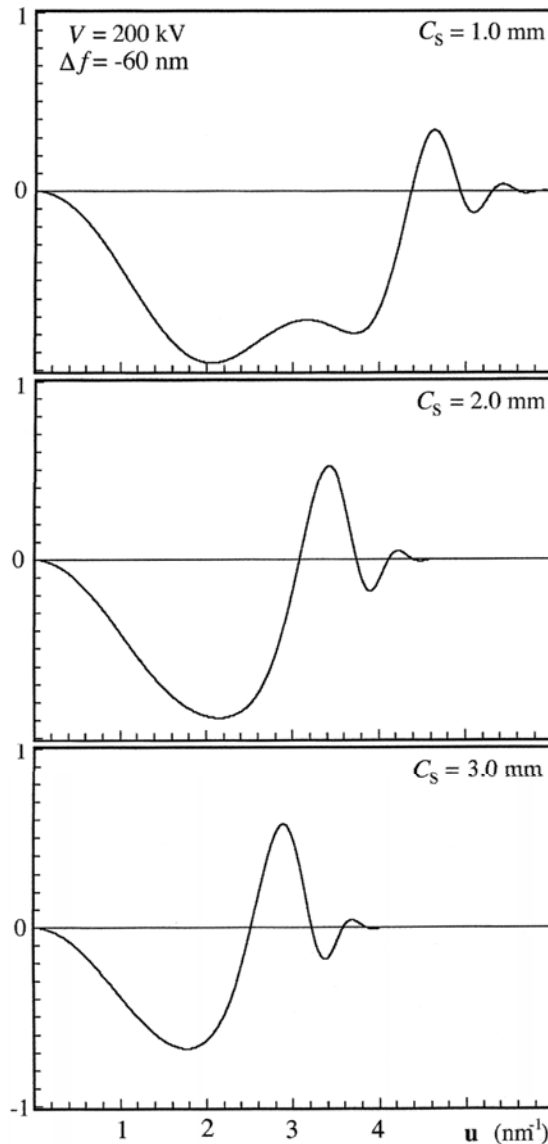
Real transfer functions modify frequencies

Example to right is close to ideal (Scherzer defocus)

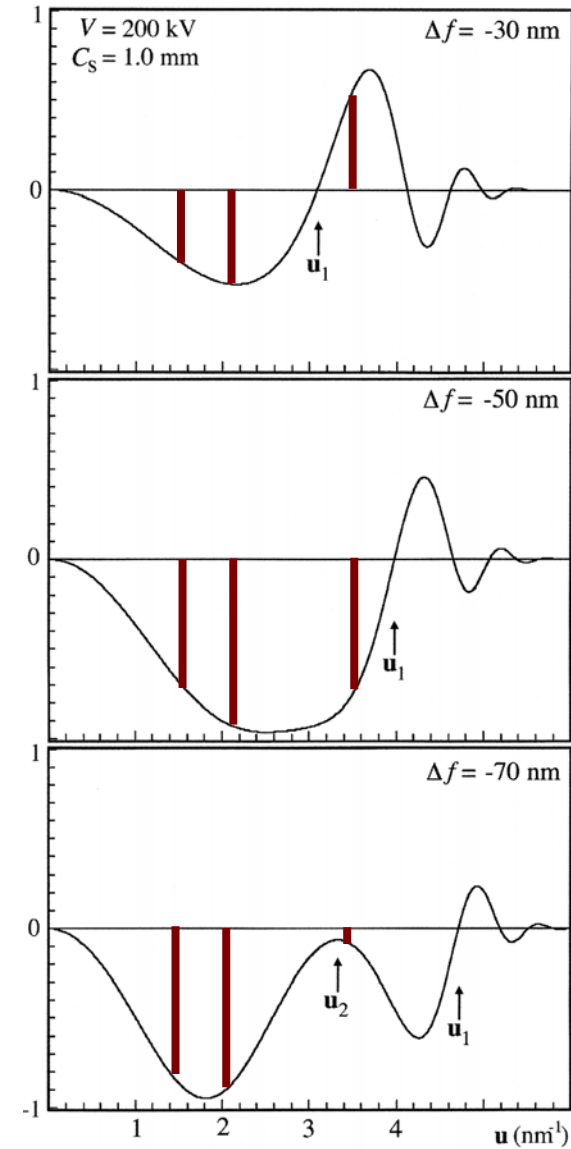


Contrast Transfer Function

Vary C_s



Vary defocus



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