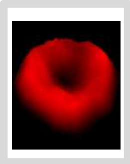

Optical Imaging
Chapter 1 - Introduction

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1.1 Properties of EM Fields

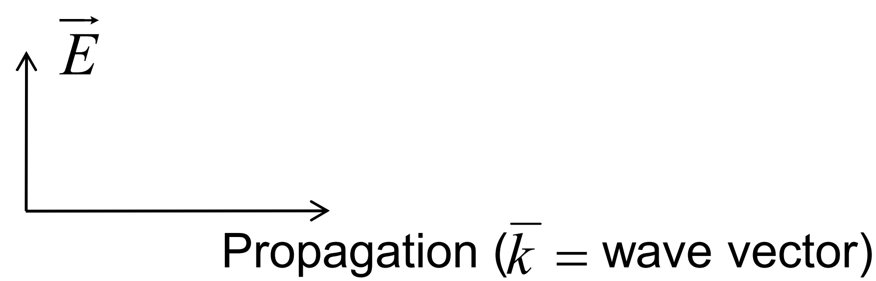
- Amplitude A and phase ϕ are random functions of both time and space:

$$\vec{E}(\vec{r},t) = \vec{A}(\vec{r},t).e^{i\phi(\vec{r},t)} \qquad (1.1)$$

1.1 Properties of EM Fields

a) Polarization:

- Gives the direction of field oscillation
- Generally, light is a transverse wave (unlike sound = longitudinal)



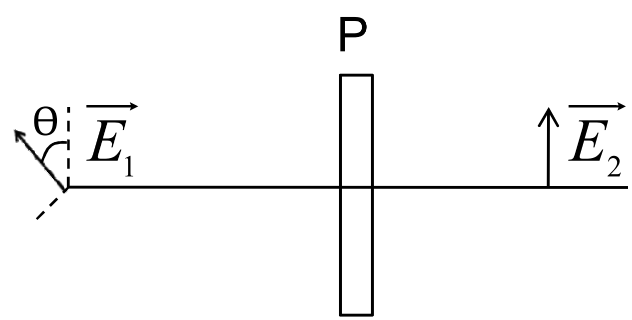
$$|\bar{k}| = \frac{2\pi}{\lambda}$$

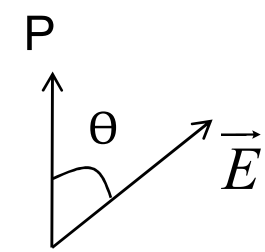
- Anisotropic materials: different optical properties along different axis → useful

1.1 Properties of EM Fields

a) Polarization:

- There is always a basis $\left(\hat{x}, \hat{y}\right)$ for decomposing the field into 2 polarizations (eigen modes); equivalently (right, left) circular polarization is also a basis.
- Dichroism: preferential absorption of one component → one way to create <u>polarizers</u>:





- <u>Malus Law</u>: $\left|E_1\right| = \left|E_2\right|.\cos\theta$ (1.2)

1.1 Properties of EM Fields

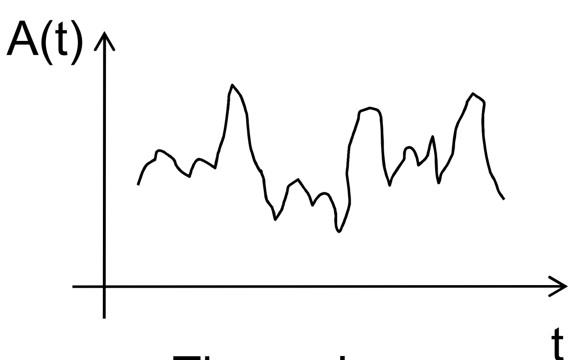
a) Polarization:

- Natural Light → unpolarized → superposition $E_x = E_y$ with no phase relationship between the two
- Circularly polarized → $E_x = E_y$, $\phi_x - \phi_y = \pi/2$!
- Matrix formalism of polarization transformation

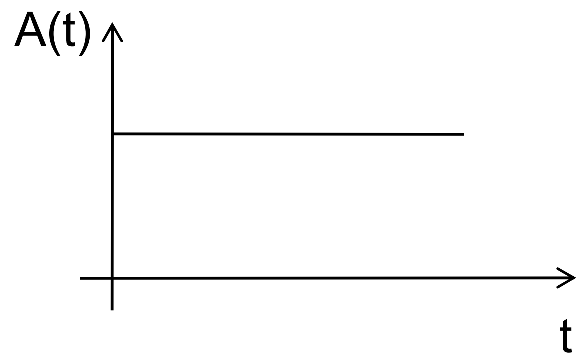
  (Jones – 2x2, complex & Muller – 4x4, real)

1.1 Properties of EM Fields

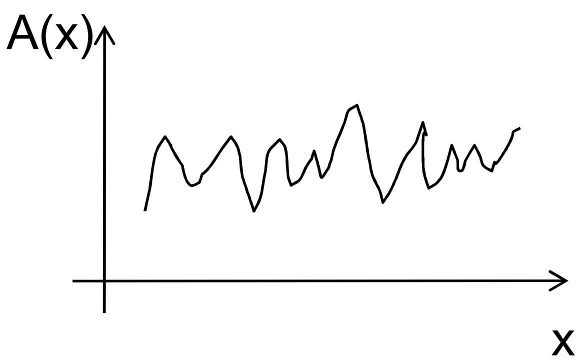
b) Amplitude: $\left[A(\vec{r},t)\right]=\frac{V}{m}$



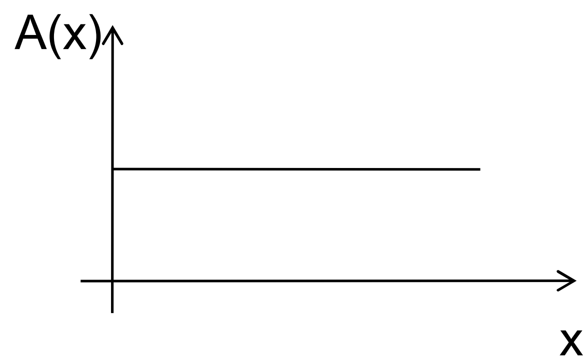
•Thermal source



•Stabilized laser



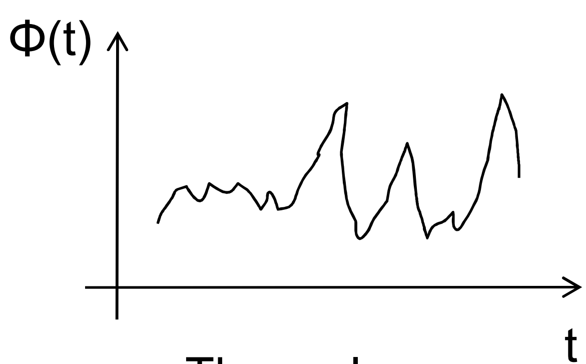
•Arbitrary field



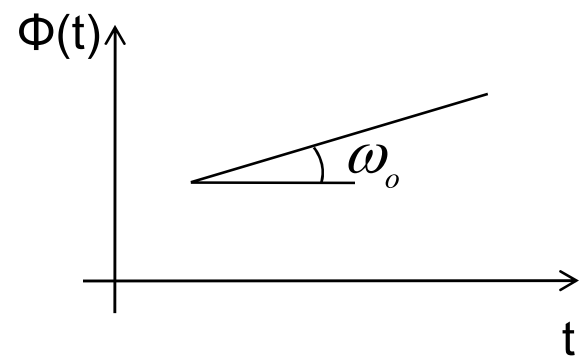
•Plane Wave

1.1 Properties of EM Fields

c) Phase: $[\Phi] = \text{rad}$

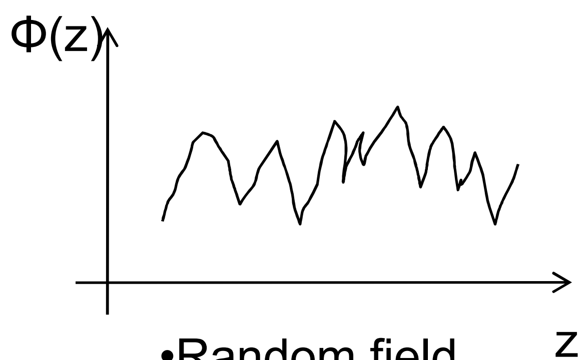


•Thermal source

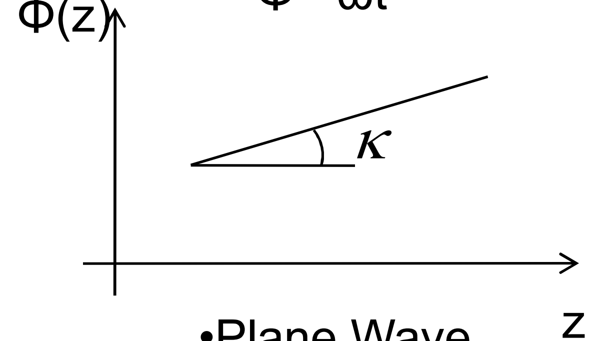


•Laser at freq ω_o

$\Phi = \omega t$



•Random field



•Plane Wave

$\Phi = kz$

1.1 Properties of EM Fields

c) Phase: [Φ] = rad

- For quasi-monochromatic fields, plane wave

$$\phi = \omega t - \vec{k} \cdot \vec{r}$$

- $k = \dfrac{\omega}{c} = \dfrac{2\pi\upsilon}{c} = \dfrac{2\pi}{Tc} = \dfrac{2\pi}{\lambda} =$ wave number (1.3)

1.2 The frequency domain representation

- Random variable E(t) has a frequency-domain counterpart:

$$E(\omega) = A(\omega)e^{i\phi(\omega)} \quad (1.4)$$

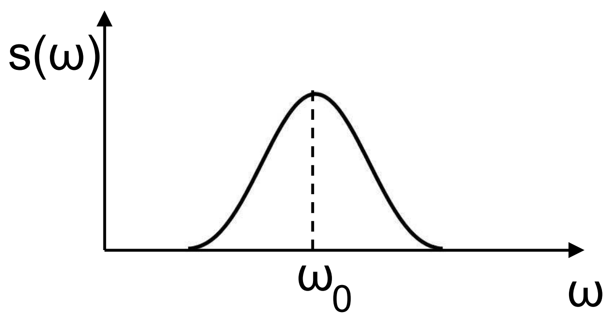
- Similarly E(x) has a frequency-domain pair:

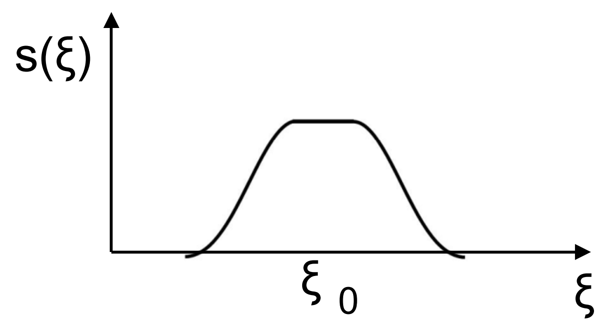
$$E(\xi) = A(\xi)e^{i\phi(\xi)} \quad (1.5)$$

1.2 The frequency domain representation

a) Spectral amplitude:

- Optical Spectrum: $s(\omega) = |A(\omega)|^2$
- Angular Spectrum: $s(\xi) = |A(\xi)|^2$





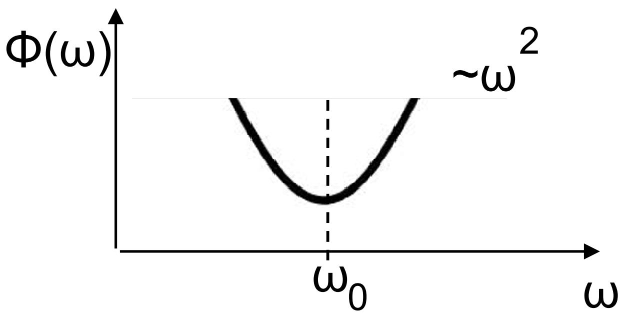
- $[\xi] = m^{-1}$ = Spatial Frequency (connects to angular spectrum)
- Tipically: t → ω, x → ξ } Will follow similar equations
- The information contained is the same (t, ω) and (x, ξ)

1.2 The frequency domain representation

b) Spectral phase:

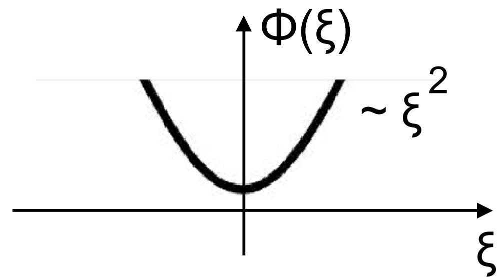
- Phase delay of each spectral component

Optical Frequency



•Dispersive material
(linear chirp)

Spatial Frequency



•Defocused point source
(1st order aberration)

- Full similarity between (t,ω) and (x,ξ)

1.3 Measurable Quantities

- The information about the system under investigation may be contained in polarization and:
  - A(t), ϕ(t)
  - A(ω), ϕ(ω)

    } (t,ω)
  - A(x), ϕ(x)
  - A(ξ), ϕ(ξ)

    } (x,ξ)

  } 8 quantities

- Experimentally, we have access only to:

$$I = \left\langle \left| A(t) \right|^2 \right\rangle = \text{time average}$$

1.3 Measurable Quantities

- Experimentally, we have access only to:

$$I = \left\langle \left| A(t) \right|^2 \right\rangle = \text{time average} \qquad (1.6)$$

- i.e the phtodetectors (photodiode, CCD, retina, etc) produce photoelectrons:

$$h\upsilon = E_{e^-} + W \qquad \text{(Einstein)} \qquad (1.7)$$

$h\upsilon$ → Photon incident energy

E_{e^-} → Electron kinetic energy

W → Work

1.3 Measurable Quantities

- All detectors sensitive to power/energy
- However, all 8 quantities can be accessed via various tricks
- Eg1: Want I(λ) → measure I(θ) and use a device with $\theta(\lambda)$
- Eg2: Want ϕ → use interferometry → I(ϕ) $\alpha \left|E_1\right|\left|E_2\right|\cos(\phi_1 - \phi_2)$

1.4 Uncertainty Principle

- Space - momentum or energy-time cannot be measured simmultaneously with infinite accuracy

$$\begin{cases} \Delta\overline{x} \cdot \Delta\overline{p} = \text{constant} \simeq h \\ \Delta E \Delta t = \text{constant} \end{cases}$$

- For photons:

$$\begin{cases} E = \hbar\omega \\ \overline{p} = \hbar\overline{k} \end{cases}$$

1.4 Uncertainty Principle

a) t $-\omega$

$$\hbar \Delta\omega \Delta t = \text{constant}$$

$$\rightarrow \boxed{\Delta\omega \Delta t \simeq 2\pi}$$

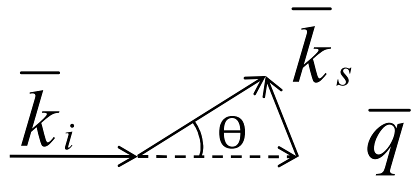
- Implications:

  1- short pulses require broad spectrum

  2-high spectral resolutioon requires long time of measurement

1.4 Uncertainty Principle

b) x $-\xi$



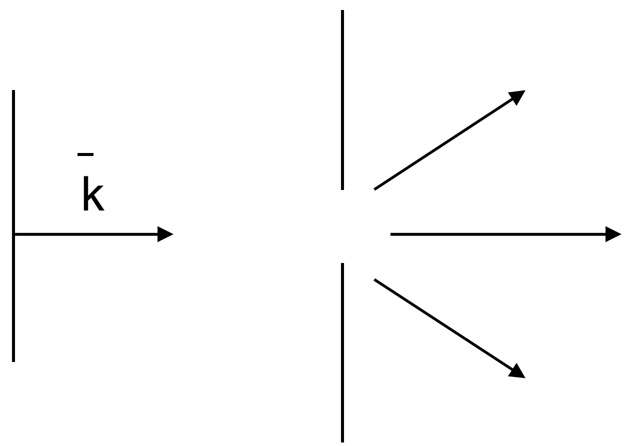
$$\overline{\Delta p} = h(\bar{k}_s - \bar{k}_i) = h\bar{q}$$

→ $\boxed{\Delta x \left|\bar{q}\right| \simeq \lambda\pi}$; $\left|\bar{q}\right| = 2k\sin(\frac{\theta}{2})$

→ $\Delta x \frac{2\sin(\theta/2)}{\lambda} \simeq 1$; $\boxed{\theta \sim \frac{1}{\Delta x}}$

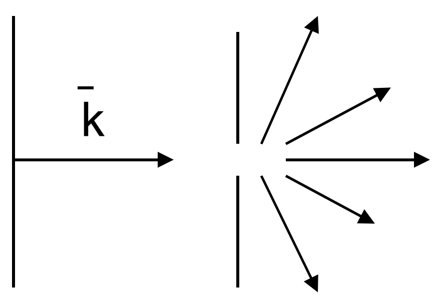
→ $\boxed{\Delta x_{\min} = \frac{\lambda}{2}}$ - meaning of resolution

1.4 Uncertainty Principle



diffraction

- Smaller aperture → Higher angle



- If aperture < $\frac{\lambda}{2}$, light doesn't go through
- Eg: Microwave door

1.4 Uncertainty Principle

- We will encounter these relationships many times later
- Fourier seems to have understood this uncertainity principle way before Heisenberg!