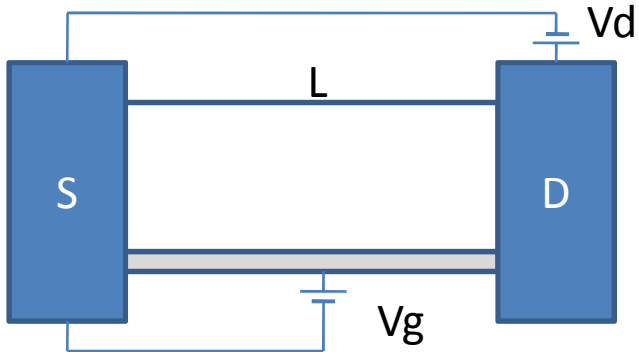


ECE 659 Quantum Transport: Atom to Transistor

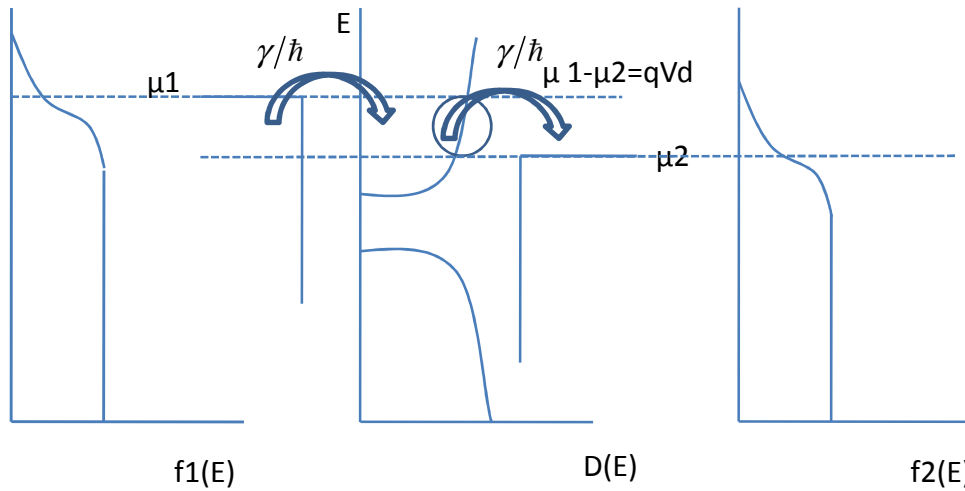
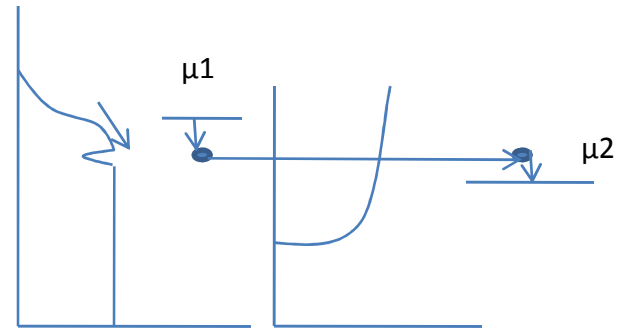
Lecture 4: Landauer Model

Supriyo Datta

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$$I = \frac{q}{h} \int dE \bar{T}(E) (f_1 - f_2)$$



$$\text{Energy lost} = \mu_1 - \mu_2 = qVd$$

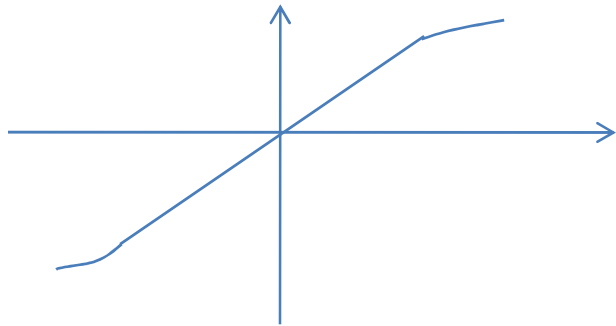
$$G = \frac{q^2}{h} \int dE \left(-\frac{\partial f}{\partial E} \right) \bar{T}(E)$$

Usual expression $G = \frac{\sigma A}{L}$, $\sigma = qn \frac{\overbrace{q\tau}^{\mu_n}}{m}$

gives the idea that all electrons conduct

$$T = \underbrace{M}_{Dv_z} \frac{\overbrace{\lambda}^{2v_z\tau}}{\lambda + L}$$

Limits of linearity of I-V characteristics



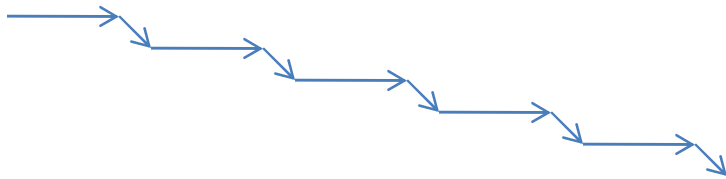
Limit of non linearity:

$E - \mu < kT$

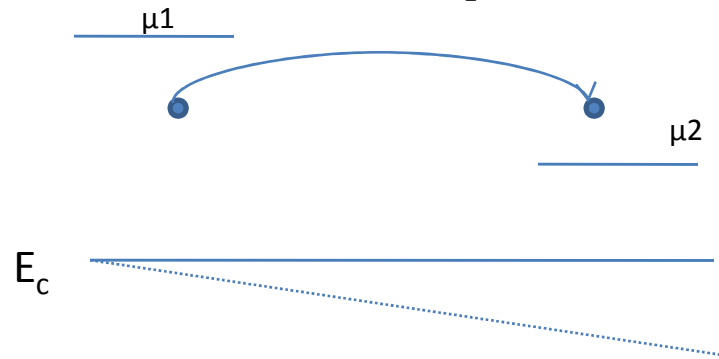
$$\frac{25mV}{0.1\mu m} = 2.5KV / cm$$

Breakdown fields: $\frac{1V}{0.1\mu m} = 10^5 V / cm$

Series of Landauer devices:



Instead of $J = \sigma E$
we should use $J = \sigma \nabla \frac{\mu}{q}$



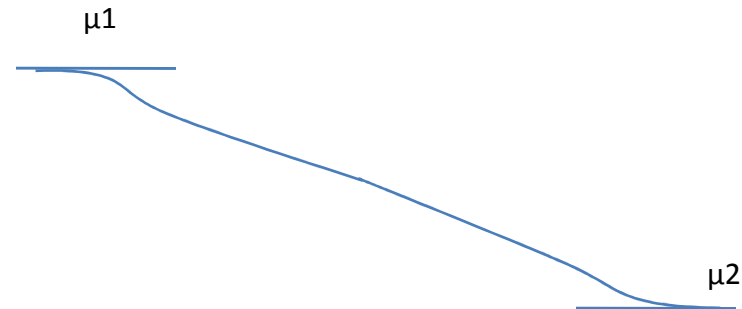
In presence of electric field

Poisson's Equation:

$$\nabla^2 U = \frac{q}{\epsilon} (n - n_0)$$

$$\nabla^2 U = \frac{q^2}{\epsilon} (\mu - U) D$$

$$\left[\nabla^2 + \frac{1}{\lambda_D^2} \right] U = \frac{q^2 \mu D}{\epsilon}$$



Two resistances in series

$$G = \sigma \frac{A}{L + \lambda}$$

$$R = \frac{\rho}{A} 2(L + \lambda) ?$$

$$R = \frac{\rho}{A} (L + \lambda)$$

or

$$R = \frac{\rho}{A} (2L + \lambda) ?$$

