

ECE 659: Quantum Transport Spring 2009
EE 117, MWF 930A-1020P

Course website: <http://cobweb.ecn.purdue.edu/%7Edatta/659.htm>

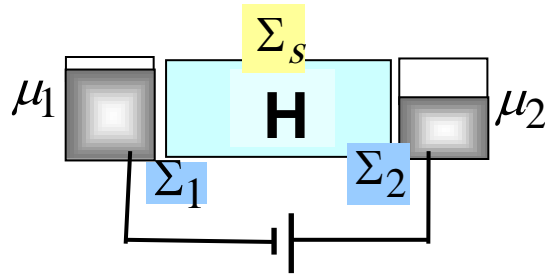
Lecture videos posted at <https://nanohub.org/resources/6172/>

"Input": H-matrix parameters chosen appropriately to match energy levels or dispersion relations. Σ_j for terminal 'j' is in general obtained from $\tau_j g_j \tau_j^+$ where the surface Green function 'g' is calculated from a recursive relation: $g^{-1} = EI - \alpha - \beta g \beta^+$.

$$\Gamma_j = i[\Sigma_j - \Sigma_j^+], \Gamma_s = i[\Sigma_s - \Sigma_s^+]$$

$$\Sigma \equiv \Sigma_s + \sum_j \Sigma_j,$$

and $\Sigma^{in} \equiv \Sigma_s^{in} + \sum_j \Sigma_j^{in}$



NEGF equations:

1. $G(E) = [EI - H - \Sigma_1 - \Sigma_2 - \Sigma_s]^{-1}$

2. $[G^n(E)] = [G \Sigma^{in} G^+]$

3. $A(E) = i[G - G^+] = G \Gamma G^+ = G^+ \Gamma G$

4. $i \hbar_{op} \neq [HG^n - G^n H] + [\Sigma G^n - G^n \Sigma^+] + [\Sigma^{in} G^+ - G \Sigma^{in}]$

- 4a. $I_{a \rightarrow b}(E) = \frac{q}{h} i [H_{ab} G_{ba}^n - G_{ab}^n H_{ba}]$ *a, b: Internal Points*

- 4b. $I_i(E) = \frac{q}{h} ((\text{Trace}[\Sigma_i^{in} A - \Gamma_i G^n])$ *Current/energy at terminal 'i'*

- 4c. $I_i(E) = \frac{q}{h} \sum_j \text{Trace}[\Gamma_i G \Gamma_j G^+] (f_i(E) - f_j(E))$

(used only if Σ_s is zero: coherent transport)

Note: $\Sigma_j^{in} = \Gamma_j f_j$, but Σ_s^{in} cannot in general be written as $\Gamma_s f_s$.

Instead it has to be calculated self-consistently as summarized on next page.

A. For elastic scatterers in equilibrium:

$$[\Sigma_s] = D[G], \quad [\Sigma_s^{in}] = D[G^n] \quad (S)$$

where $D = U_s U_s^*$ describes incoherent processes (no relation to density of states $D(E)$).

B. For inelastic scatterers, with dissipation occurring due to interaction with a reservoir having a spectrum $D(+\varepsilon)$ for absorption and $D(-\varepsilon)$ for emission, See Section 10.3, page 271 of text (QTAT, Quantum Transport, Atom to Transistor):

Replace (S) with

$$\begin{aligned} [\Sigma_s^{in}(E)] &= D(+\varepsilon) [G^n(E - \varepsilon)] \\ [\Gamma_s(E)] &= D(+\varepsilon) [G^n(E - \varepsilon)] + D(+\varepsilon) [G^p(E + \varepsilon)] \end{aligned}$$

(Note that $G^p(E)$ is the "hole density" given by $A(E) - G^n(E)$)

$$[\Sigma_s(E)] = \underbrace{[h(E)]}_{\substack{\text{Hilbert} \\ \text{Transform}}} - \frac{i}{2} [\Gamma_s(E)]$$

For a brief discussion of Hilbert transform see QTAT, top of page 206.

More generally, replace (S) with (summation over repeated indices is implied)

$$(S1) \quad [\Sigma_s^{in}(E)]_{ij} = D_{ik;jl}(+\varepsilon) [G^n(E - \varepsilon)]_{kl}$$

$$(S2) \quad [\Gamma_s(E)]_{ij} = D_{ik;jl}(+\varepsilon) [G^n(E - \varepsilon)]_{kl} + D_{lj;ki}(+\varepsilon) [G^p(E + \varepsilon)]_{kl}$$

$$(S3) \quad [\Sigma_s(E)]_{ij} = \underbrace{[h(E)]_{ij}}_{\substack{\text{Hilbert} \\ \text{Transform}}} - \frac{i}{2} [\Gamma_s(E)]_{ij}$$

$$D_{ik;jl} = [U_s]_{ik} [U_s]_{jl}^*$$

If scatterers are in equilibrium with temperature T , then $\frac{D_{ik;jl}(+\varepsilon)}{D_{lj;ki}(-\varepsilon)} = e^{-\varepsilon/k_B T}$

Can show using (S1) and (S2) that the integrated scattering current is zero, though it

may not be zero at each energy E : $\int dE I_j(E) = \frac{q}{h} \int dE \text{Trace}[\Sigma_s^{in} A - \Gamma_s G^n] = 0$