



Basics of Transport in Nanostructures

Nick Fang

Course Website: nanoHUB.org
Compass.illinois.edu

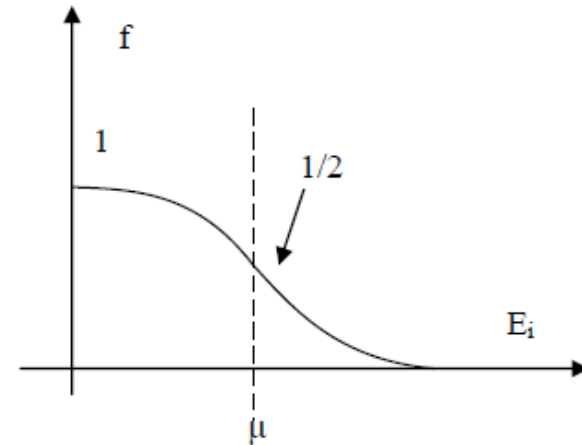


For Quantum Particles



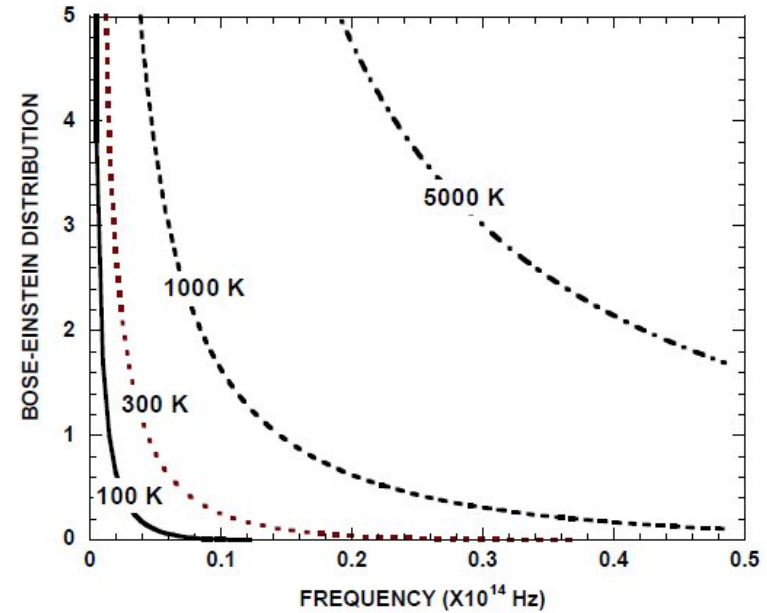
- Electrons: only two states possible (conduction, valence)

$$p(E_i) = \frac{\exp((E_F - E_i)/k_B T)}{1 + \exp((E_F - E_i)/k_B T)}$$



- Photons and Phonons: all possible states of energy $nh\omega$

$$p(\omega) = \frac{1}{\exp(h\omega/k_B T) - 1}$$





How Fast do they move?



- Let's calculate the average kinetic energy

$$\langle E \rangle = \int_{-\infty}^{\infty} dv_x \int_{-\infty}^{\infty} dv_y \int_{-\infty}^{\infty} \frac{m}{2} (v_x^2 + v_y^2 + v_z^2) p(v_x, v_y, v_z) dv_z$$

- For monatomic gas

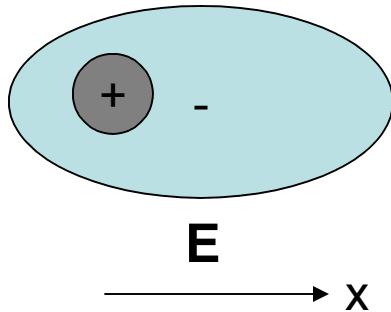
$$\langle E \rangle = \frac{3}{2} k_B T$$

*At room temperature (300 K), this average energy is **39 meV**, or **6.21x10⁻²¹ J**.*

For He gas, $m=6.4 \times 10^{-27}$ kg, $v \sim 1000$ m/s



Photon Excitation in Materials

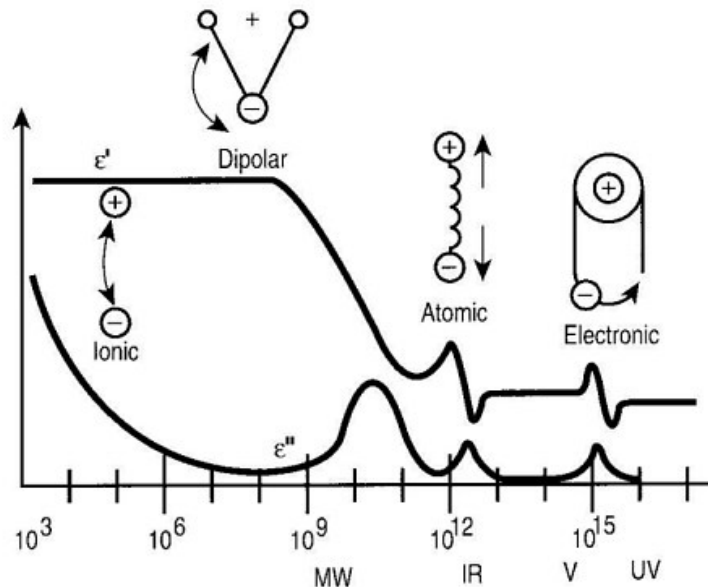


Lorenz Oscillator Model:

$$m \frac{\partial^2 x}{\partial t^2} + \gamma \frac{\partial x}{\partial t} + kx = eE_x$$

Molecular Polarizability $P_x = ex$ $\epsilon = 1 + n \frac{P}{\epsilon_0 E}$

$$x = \frac{eE_x}{m(\omega_0^2 - \omega^2 + i\gamma\omega/m)}$$

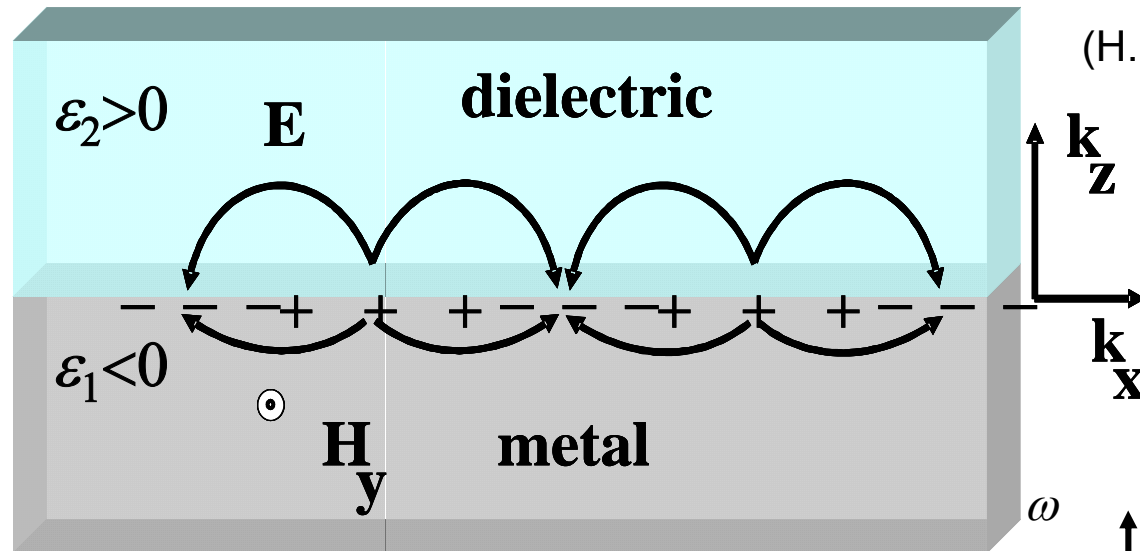


$$\epsilon = 1 + \frac{ne^2}{\epsilon_0 m} \left(\frac{1}{\omega_0^2 - \omega^2 + i\Gamma\omega} \right)$$

<http://en.wikipedia.org/wiki/Permittivity>
#Complex_permittivity

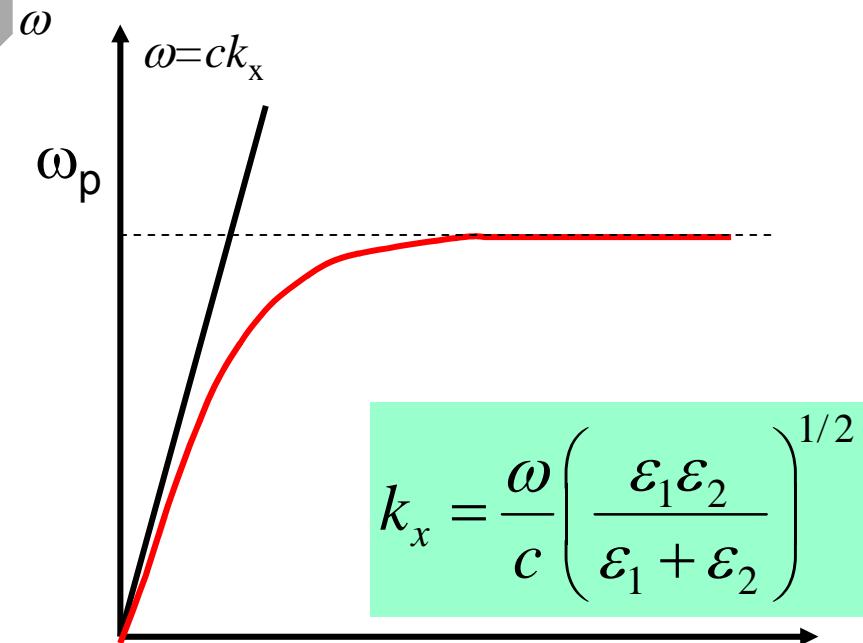


Surface Plasmons



(H. Raether, *Surface Plasmons*, Springer-Verlag, 1988)

- EM waves propagating along the interface between two media with their ϵ of opposite sign.
- Intensity maximum at interface; exponentially decays away from the interface.

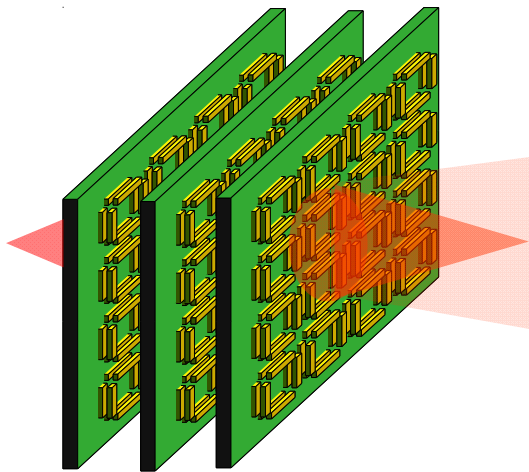




Application: Metamaterials

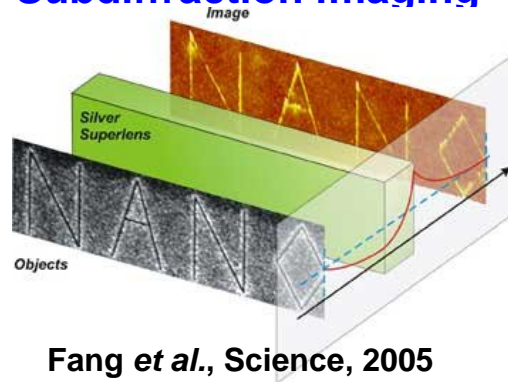


Telecom applications



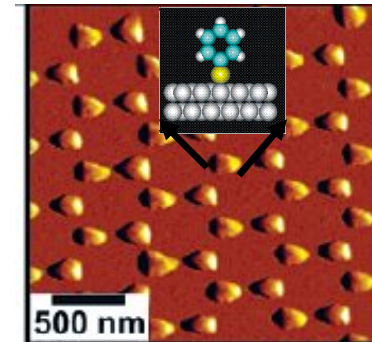
Logeeswaran *et al.*, Appl. Phys. A, 2007

Subdiffraction imaging



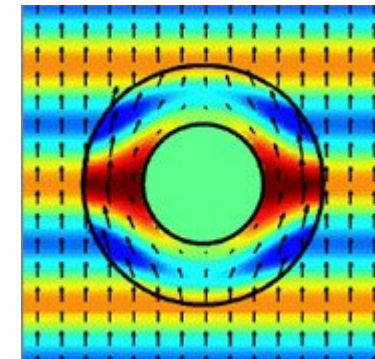
Fang *et al.*, Science, 2005

Sensing



Van Duyne *et al.*, MRS bulletin, 2005

Invisibility cloaks



Chen *et al.*, PRL, 2007

Metamaterials

- **Materials Today's** top 10 advances in material science over the past 50 years
- **Discover** top 100 science stories of the year 2006



Microscopic Transport Theory

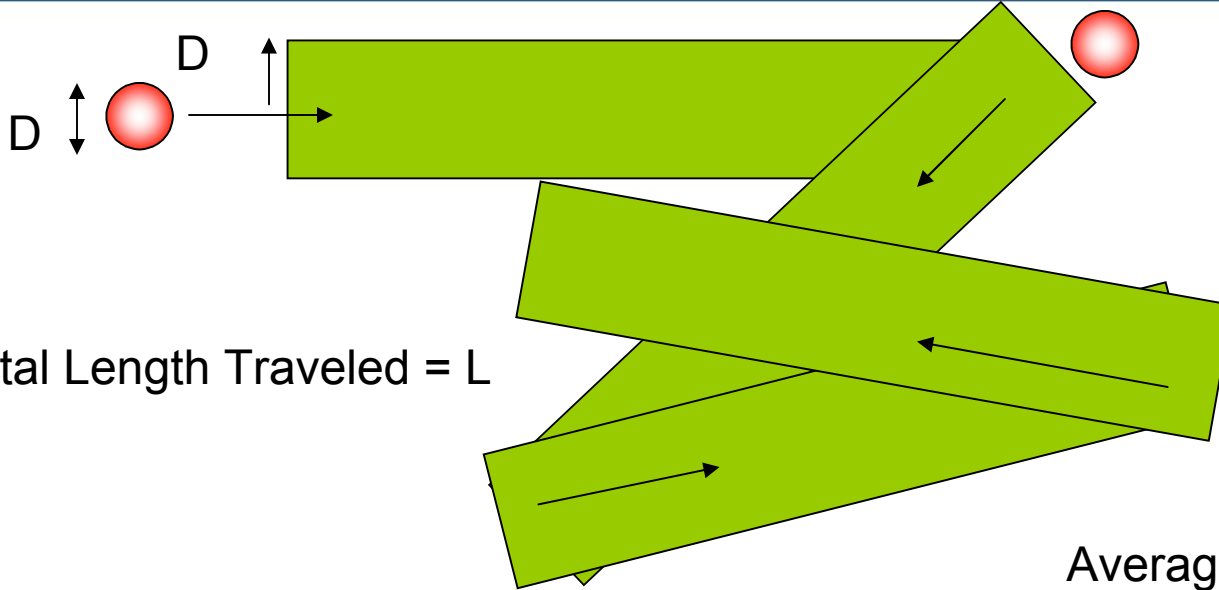


To understand nanoscale transport and energy conversion, we need to know:

- **How much energy/momentum can a particle have?**
- **How many particles have the specified energy E ?**
- **How fast do they move?**
- **How do they interact with each other?**
- **How far can they travel?**



How Far Can They Travel?



Total Length Traveled = L

Average Distance between Collisions, $\lambda_{mc} = L/(\text{\#of collisions})$

E.G. Ideal Gas:

Total Collision Volume

$$\text{Swept} = \pi D^2 L$$

Number Density of Molecules = n

Total number of molecules encountered in swept collision volume $\sim n\pi D^2 L$

Mean Free Path

$$\lambda_{mc} = \frac{L}{n\pi D^2 L} = \frac{1}{n\sigma}$$

σ : collision cross-sectional area
 $\sim \text{nm}^2$



Mean Free Path for Gas Molecules



Number Density of
Molecules from Ideal

Gas Law:

$$n = P/k_B T$$

k_B : Boltzmann constant

$$1.38 \times 10^{-23} \text{ J/K}$$

Mean Free Path:

$$\lambda_{mc} = \frac{1}{n\sigma} = \frac{k_B T}{P\sigma}$$

Typical Numbers:

Diameter of Molecules, $D \approx 2 \text{ \AA} = 2 \times 10^{-10} \text{ m}$

Collision Cross-section: $\sigma \approx 1.3 \times 10^{-19} \text{ m}^2$

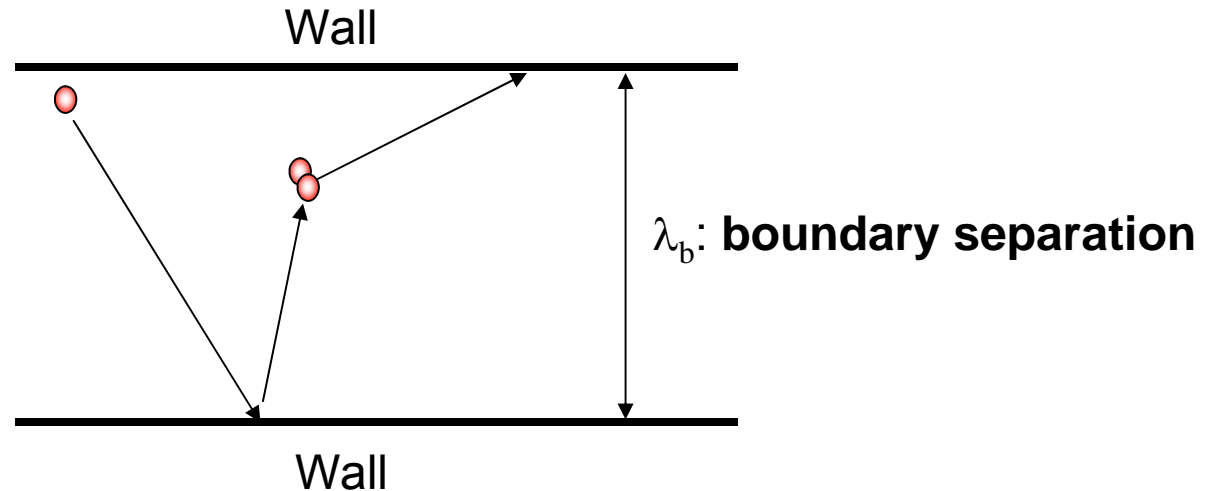
Mean Free Path at Atmospheric Pressure:

$$\lambda_{mc} \approx \frac{1.38 \times 10^{-23} \times 300}{10^5 \times 1.3 \times 10^{-19}} \approx 3 \times 10^{-7} \text{ m or } 0.3 \mu\text{m}$$

At 1 Torr pressure, $\lambda_{mc} \approx 200 \mu\text{m}$; at 1 mTorr, $\lambda_{mc} \approx 20 \text{ cm}$



Effect of Nanoscale confinement



Effective Mean Free Path:

$$\frac{1}{\lambda} = \frac{1}{\lambda_{mc}} + \frac{1}{\lambda_b}$$

The smaller dimension governs collision time!



Internal Energy and Specific Heat



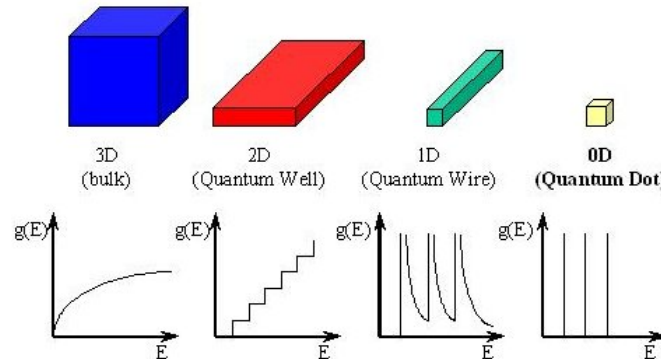
- Now we know the energy and momentum of particles/carriers in the material, we can start counting the properties
- E.G. Internal energy

Boltzmann Distribution

Density of States

Energy of Carrier at Given States

$$p_i = \frac{1}{Z} e^{-E_i/k_B T}$$



Translation
Vibration
Rotation



E.G. Internal Energy of Photons



Photon Energy
at Given States

$$\hbar\omega = \hbar ck \quad \text{In vacuum, 3D}$$

Bose-Einstein
Distribution

$$p(\hbar\omega) = \frac{1}{\exp(\hbar\omega/k_B T) - 1}$$

Density of
States

$$D(\omega)d\omega = \frac{4\pi\omega^2}{Vc^3}d\omega$$

Total Internal
Energy:

$$U = N \int_0^\infty \hbar\omega p(\omega) D(\omega) d\omega$$



Thermal Radiation (Planck's Law)



- Total Internal Energy of Vacuum Photons:

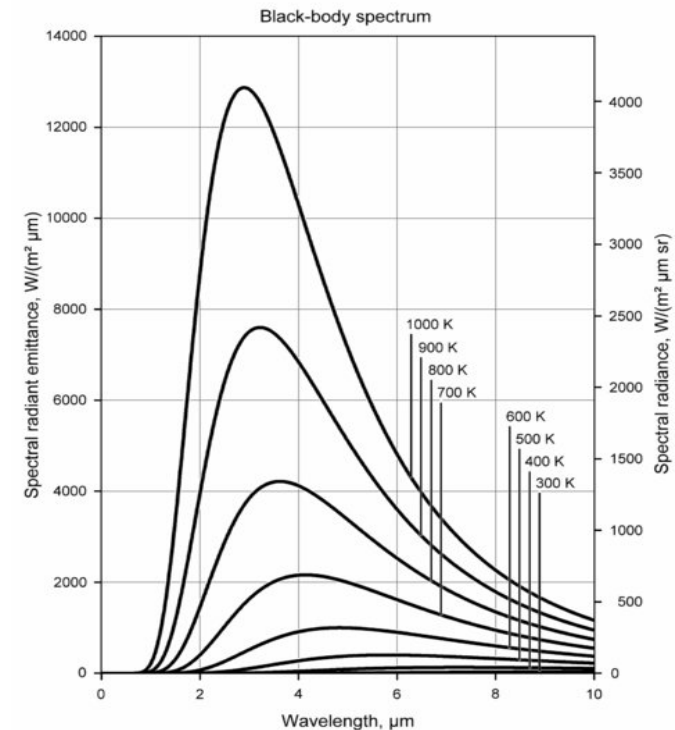
$$U = \frac{4\pi N}{Vc^3} \int_0^{\infty} \frac{h\omega^3}{\exp(h\omega/k_B T) - 1} d\omega$$

Converting the distribution to wavelength,

$$\omega = \frac{2\pi c}{\lambda}$$

$$P(\lambda) = \frac{1}{\lambda^5} \frac{(2\pi hc)^4}{\exp(2\pi hc / \lambda k_B T) - 1} d\lambda$$

describes the spectral radiance of electromagnetic radiation at temperature T.



http://upload.wikimedia.org/wikipedia/commons/8/85/BlackbodySpectrum_lin_150dpi_en.png



Thermal Radiation (Stefan-Boltzmann)



Define: $x = h\omega / k_B T$

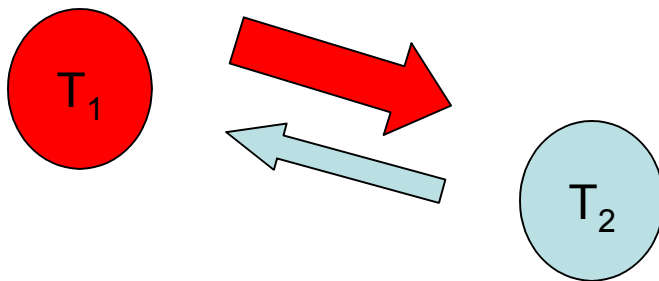
$$U(T) = \frac{4\pi N(k_B T)^4}{V(hc)^3} \int_0^\infty \frac{x^3}{\exp(x) - 1} dx$$

The emissive power of
Black body radiation:

$$E(T) = \sigma T^4$$

Stefan-Boltzmann's Law

$$\sigma = 5.67 \times 10^{-8} \text{ W/m}^2 \text{ K}^4 .$$



$$q = \sigma (T_1^4 - T_2^4)$$



Internal Energy of Phonons



- E.G. in a bulk solid

Phonon Energy
at Given States

$$\hbar\omega = \hbar\sqrt{\frac{K}{m}} \sin(ka)$$

Boltzmann
Distribution

$$p(\omega) = \frac{1}{\exp(\hbar\omega/k_B T) - 1}$$

Density of
States

$$D(\omega)d\omega \approx \frac{4\pi\omega^2}{V(a\omega_D)^{3/2}} d\omega$$

Debye
Approximation

$$\omega < \omega_D = \sqrt{K/m}$$

Total Internal
Energy:

$$U = 3N \int_0^{\omega_D} \hbar\omega p(\omega) D(\omega) d\omega$$



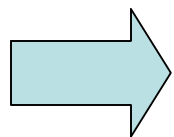
Specific Heat Capacity



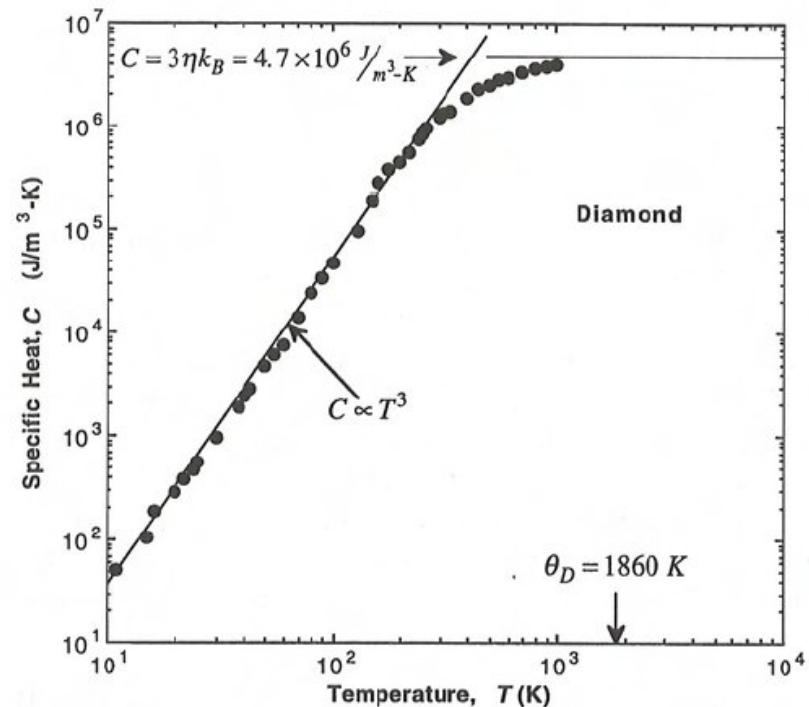
- The specific heat capacity is defined by change of internal energy per unit temperature change:

$$C_V = \frac{\partial U}{\partial T}$$

$$U \propto T^4$$



$$C_V \propto T^3 \quad \text{At low temperature}$$



Specific heat of diamond

(Touloukian and Buyco, 1970).

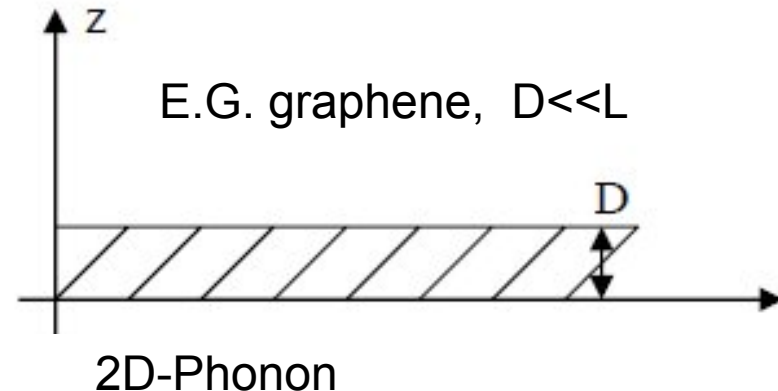


Size Effect on Heat Capacity



Modified Density of States

$$D(\omega)d\omega \approx \frac{2\pi\omega(n)}{V(a\omega_D)}d\omega$$



$$U \propto T^3 \quad C_V = \frac{\partial U}{\partial T} \propto T^2 \text{ (film)}$$

Nanoparticles:

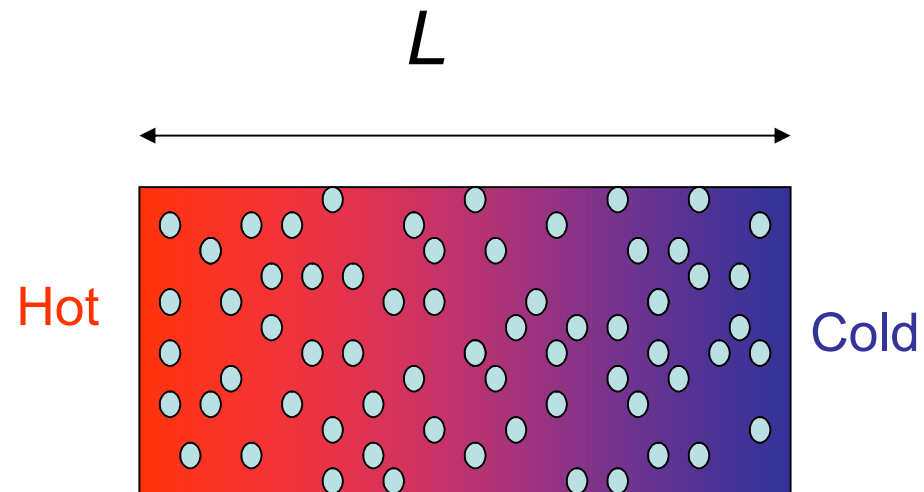
$$U \propto T^{-1} \quad C_V = \frac{\partial U}{\partial T} \propto T^{-2} \frac{\exp(T_E/T)}{(\exp(T_E/T) - 1)^2}$$



Transport Properties



- E.G. Heat Conduction



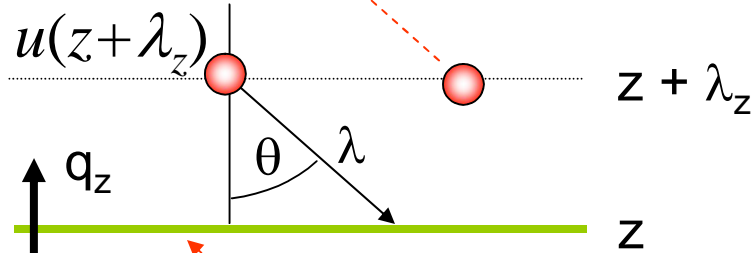
In micro-nano scale thermal and fluid systems, often $L < \text{mean free path of collision of energy carriers}$ & Fourier's law breaks down
→ Particle transport theories or molecular dynamics methods



Kinetic Theory of Energy Transport



Cold



Hot

Hot

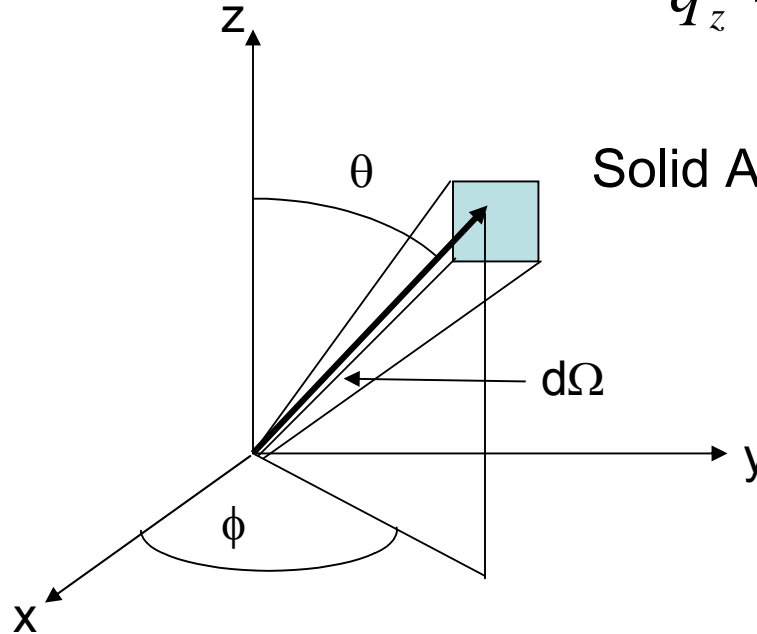
Net Energy Flux

$$q_z = \frac{1}{2} v_z [u(z - \lambda_z) - u(z + \lambda_z)]$$

through Taylor expansion of u

$$q_z = -v_z \lambda_z \frac{du}{dz}$$

Solid Angle, $d\Omega = \sin\theta d\theta d\phi$



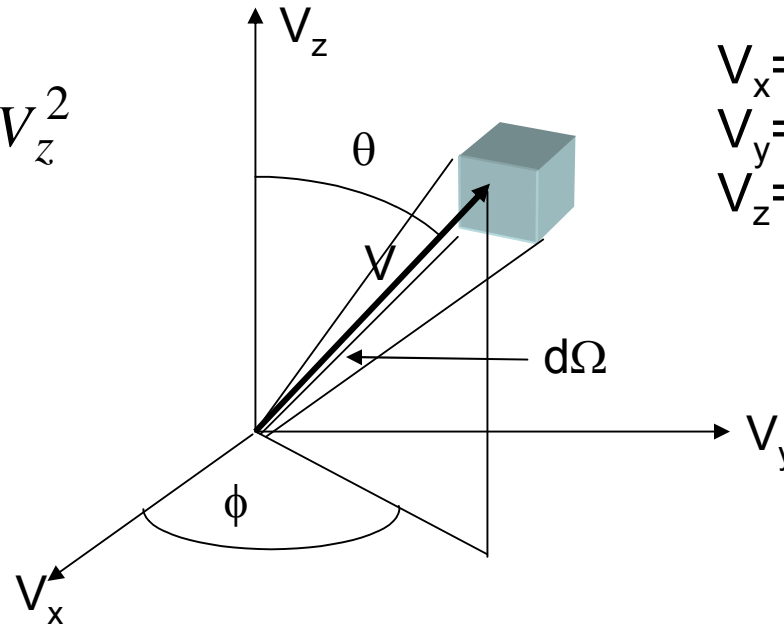


With a bit More Geometrics



Velocity:

$$V^2 = V_x^2 + V_y^2 + V_z^2$$



$$\begin{aligned} V_x &= V \sin \theta \cos \phi \\ V_y &= V \sin \theta \sin \phi \\ V_z &= V \cos \theta \end{aligned}$$

$$q = -(\cos^2 \theta) v \lambda \frac{du}{dz}$$



Averaging over all the solid angles



$$q_z = -v\lambda \frac{du}{dz} \left[\frac{\int_{\varphi=0}^{2\pi} \int_{\theta=0}^{\pi/2} \cos^2 \theta \sin \theta d\theta d\varphi}{\int_{\varphi=0}^{2\pi} \int_{\theta=0}^{\pi/2} \sin \theta d\theta d\varphi} \right] = -v\lambda \frac{du}{dz} \left[\frac{\int_{\varphi=0}^{2\pi} \int_{\theta=0}^{\pi/2} \cos^2 \theta \sin \theta d\theta d\varphi}{2\pi} \right] = -\frac{1}{3} v\lambda \frac{du}{dz}$$

Assuming **local thermodynamic equilibrium**: $u = u(T)$

$$q_z = -\frac{1}{3} v\lambda \frac{du}{dT} \frac{dT}{dz} = -\frac{1}{3} C v\lambda \frac{dT}{dz}$$

Thermal
Conductivity

$$k = \frac{1}{3} C v\lambda$$

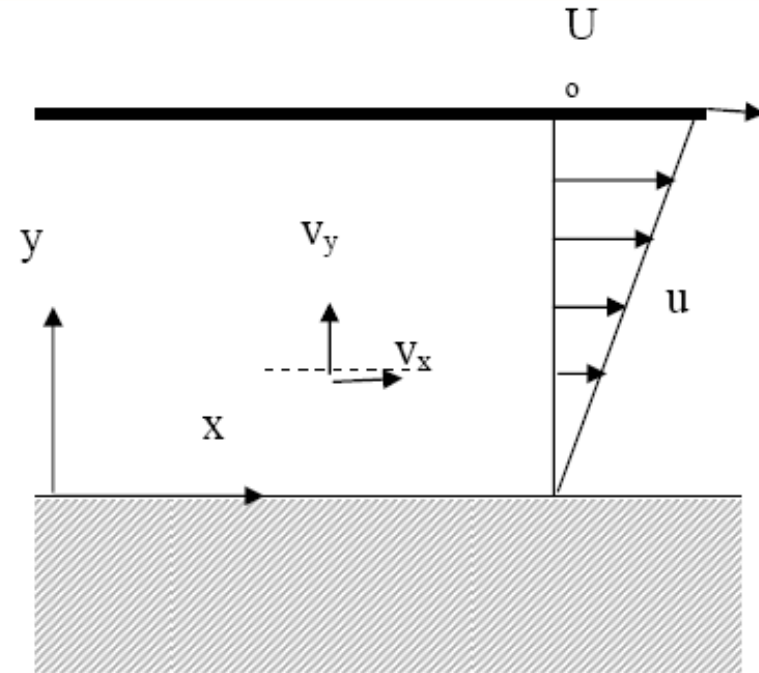


Likewise...



- Newton's shear stress Law

$$\mu = l_{mc} \frac{nk_B T}{\langle v \rangle}$$



$$P(v_x, v_y, v_z) = \left(\frac{m}{2\pi k_B T} \right)^{3/2} e^{-m \left[(v_x - u)^2 + v_y^2 + v_z^2 \right] / 2k_B T}$$



Additional Reading



- Tien, Majumdar, Gerner, “Microscale Energy Transport”, Chapter 1, Taylor&Francis (pdf online)
- ECE 598EP: Hot Chips: Atoms to Heat Sinks
<http://poplab.ece.illinois.edu/teaching.html>