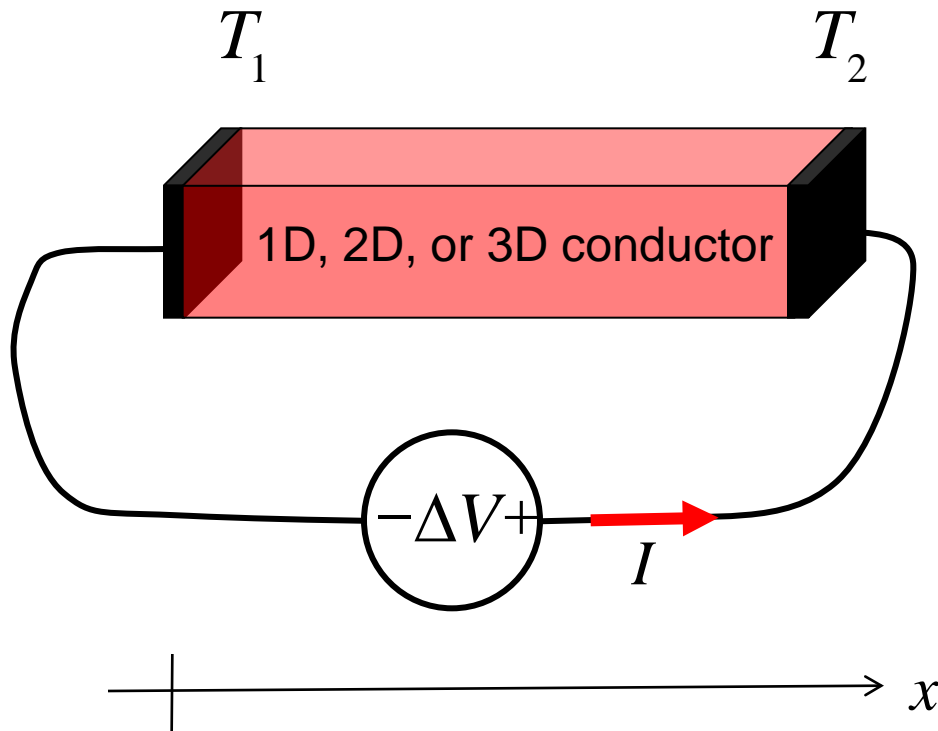


ECE-656: Fall 2009

**Lecture 9:
Coupled Current Equations**

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charge and heat currents



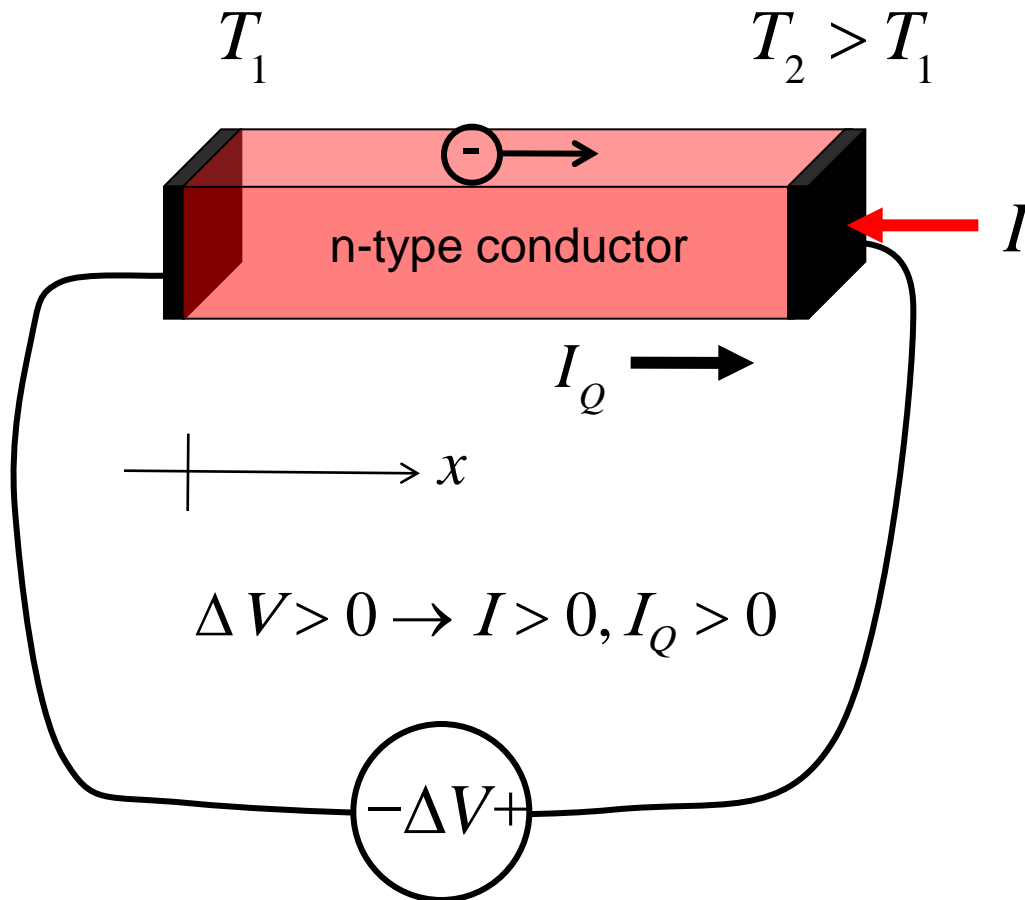
$$I = G\Delta V - [SG]\Delta T$$

$$I_Q = T[SG]\Delta V - K_0\Delta T$$

driving "forces:"

- 1) differences in voltage
(electrochemical potential)
- 2) differences in temperature

signs of the coefficient



$$I = G\Delta V - [SG]\Delta T$$

$$I_Q = T[SG]\Delta V - K_0\Delta T$$

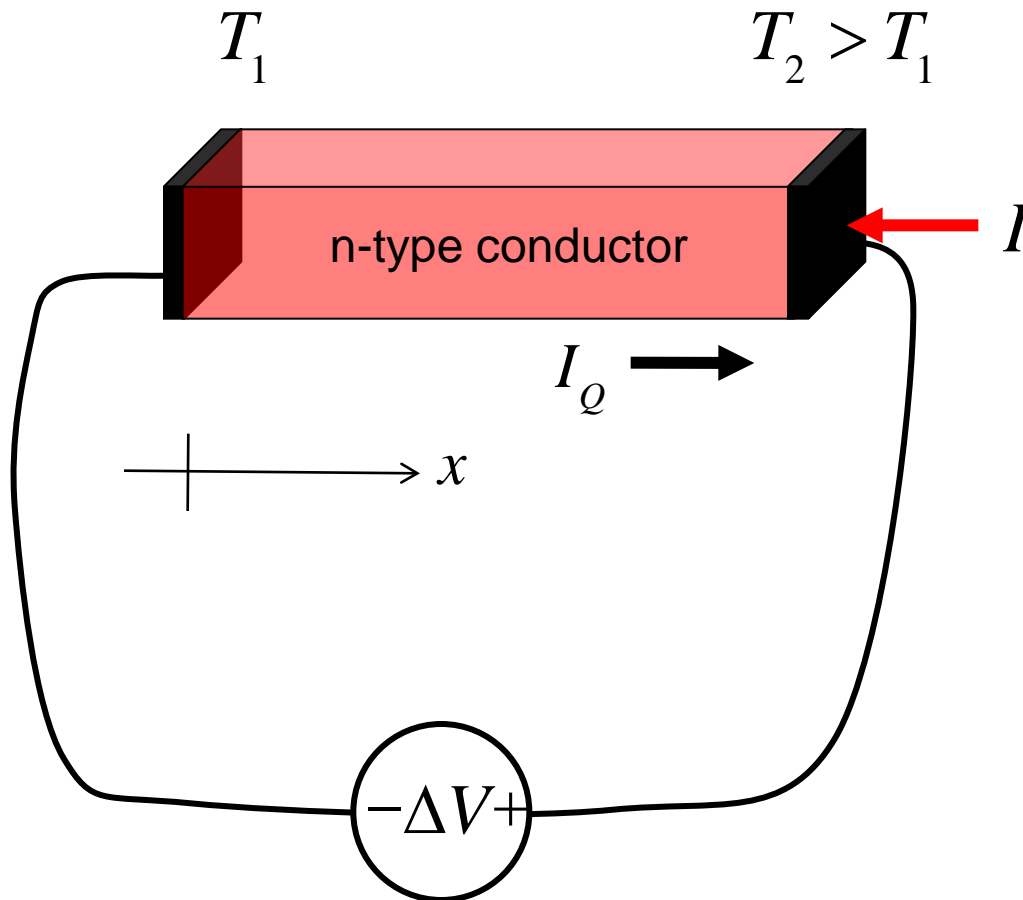
$$G > 0$$

$$[SG] > 0$$

$$K_0 > 0$$

(electronic) thermal
conductance for $\Delta V = 0$

conductance vs. conductivity



$$I = G\Delta V - [SG]\Delta T$$

$$I_Q = T[SG]\Delta V - K_0\Delta T$$

3D conductors

$$G = \sigma \frac{A}{L}$$

$$K_0 = \kappa_0 \frac{A}{L}$$

etc.

temperature gradients

$$I \propto (f_1 - f_2) = (f_1 - f_2)_T + (f_1 - f_2)_{E_F}$$

$$(f_1 - f_2)_T = - \left(\frac{\partial f_0}{\partial E_F} \right) \Delta E_F = \left(- \frac{\partial f_0}{\partial E} \right) (q \Delta V)$$

$$\begin{aligned} (f_1 - f_2)_{E_F} &= \left(- \frac{\partial f_0}{\partial \Theta} \right) \Delta \Theta \\ &= k_B T \left(- \frac{\partial f_0}{\partial E} \right) \times \frac{(E - E_F)}{k_B} \Delta \left(\frac{1}{T} \right) \end{aligned}$$

$$f_0 = \frac{1}{1 + e^{(E - E_F)/k_B T}}$$

$$f_0 = \frac{1}{1 + e^{\Theta}}$$

$$\Theta = (E - E_F) / k_B T$$

The proper driving forces are changes in the electrochemical potential and in the inverse temperature.

coupled flows

$$I = L_{11}\Delta(E_F) + L_{12}\Delta\left(\frac{1}{T}\right)$$

$$I_Q = L_{21}\Delta E_F + L_{22}\Delta\left(\frac{1}{T}\right)$$

$$I_1 = L_{11}(\overset{\cdot}{B})F_1 + L_{12}(\overset{\cdot}{B})F_2$$

$$I_2 = L_{21}(\overset{\cdot}{B})F_1 + L_{22}(\overset{\cdot}{B})F_2$$

I_1, I_2 “generalized fluxes”

F_1, F_2 “generalized forces”

$L_{12} = L_{21}$ Onsager relation

example

temperature differences produce heat currents

pressure differences produce matter currents

heat flow per pressure difference = matter flow per temperature difference

Lars Onsager, Nobel Prize in Chemistry, 1968.

http://en.wikipedia.org/wiki/Onsager_reciprocal_relations

for more about this topic

A.C. Smith, J.F. Janak, and R.B. Adler, *Electronic Conduction in Solids*, McGraw-Hill, New York, 1967.

Irreversible thermodynamics: Chapter 2

Onsager relations: Chapter 3

outline

- 1) Onsager relations
- 2) **Measurement considerations**
- 3) Thermoelectric devices

measurements

$$I = G\Delta V - [SG]\Delta T$$
$$I_Q = T[SG]\Delta V - K_0\Delta T$$

$$\Delta V = RI - S\Delta T$$
$$I_Q = -\pi I - K_e\Delta T$$

two forms of the equations

$$I = G\Delta V - [SG]\Delta T$$

$$I_Q = T[SG]\Delta V - K_0\Delta T$$

$$\Delta V = RI - S\Delta T$$

$$I_Q = -\pi I - K_e\Delta T$$

$$\pi = TS < 0$$

$$K_e = K_0 - T\frac{[SG]^2}{G} > 0$$

electronic thermal
conductance for $I = 0$

$$\Delta V = \frac{1}{G}I + \frac{[SG]}{G}\Delta T$$

$$R = 1/G \quad S = -[SG]/G$$

$$\Delta V = RI - S\Delta T$$

$$I_Q = T[SG](RI - S\Delta T) - K_0\Delta T$$

$$I_Q = T\frac{[SG]}{G}I - (K_0 + T[SG]S)\Delta T$$

$$I_Q = -TSI - \left(K_0 - T\frac{[SG]^2}{G} \right) \Delta T$$

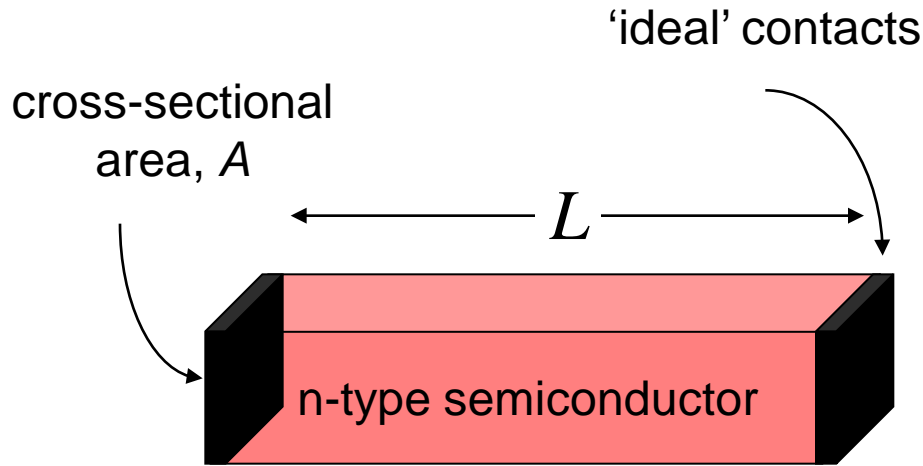
resistivity / conductivity measurements

$$I = G\Delta V - [SG]\Delta T$$
$$I_Q = T[SG]\Delta V - K_0\Delta T$$

$$\Delta V = RI - S\Delta T$$
$$I_Q = -\pi I - K_e\Delta T$$

People generally measure **resistivity** (or **conductivity**) because for bulk materials, these parameters depend on material properties and not on the length of the resistor or its width or cross-sectional area.

resistance and resistivity



$$I = V/R$$

$$R \propto \frac{L}{A} \Omega$$

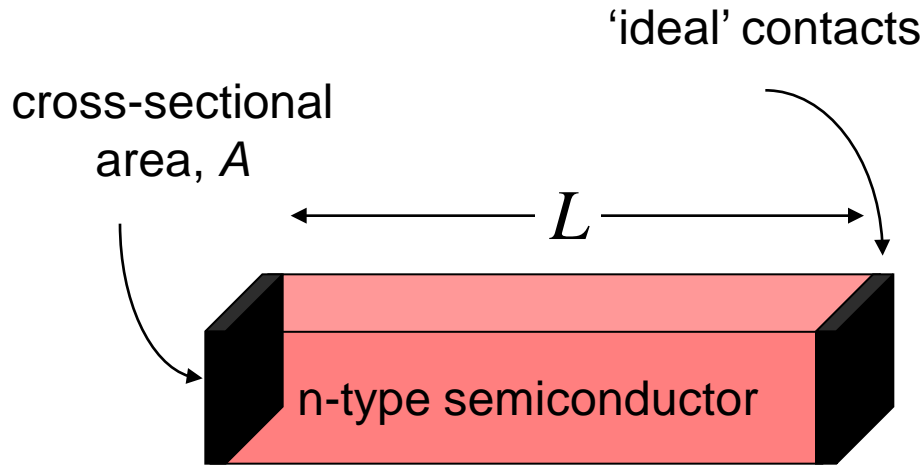
$$R = \rho \frac{L}{A} \Omega$$

resistivity:

$$\rho \text{ } \frac{1}{2}\text{-m}$$

How does the resistivity depend on the parameters of the semiconductor?

conductance and conductivity



$$I = GV = V/R$$

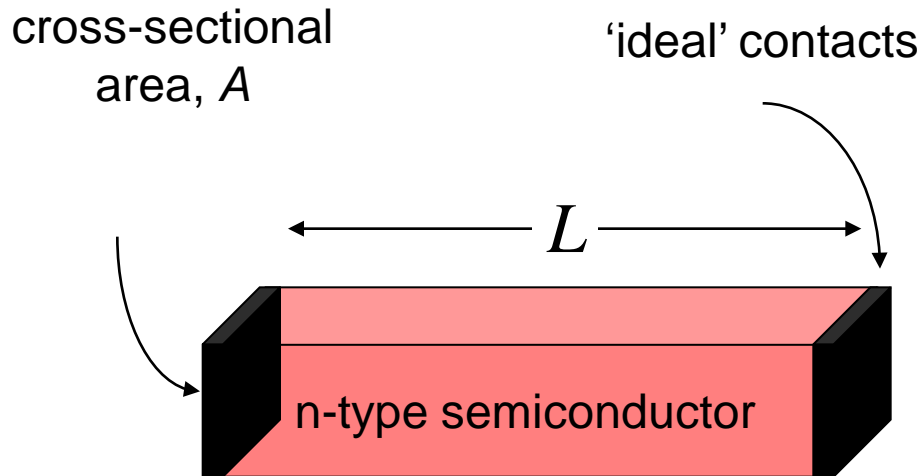
$$G = \frac{1}{R} \propto \frac{A}{L} \text{ S } (\Omega)^{-1}$$

$$G = \sigma \frac{A}{L}$$

conductivity:

$$\sigma = \frac{1}{\rho} \text{ S/m}$$

traditional conductance and conductivity



$$\ominus \rightarrow \langle v \rangle$$

$$Q = (qn)AL$$

$$I = GV \equiv \frac{Q}{\tau} \text{ Amperes}$$

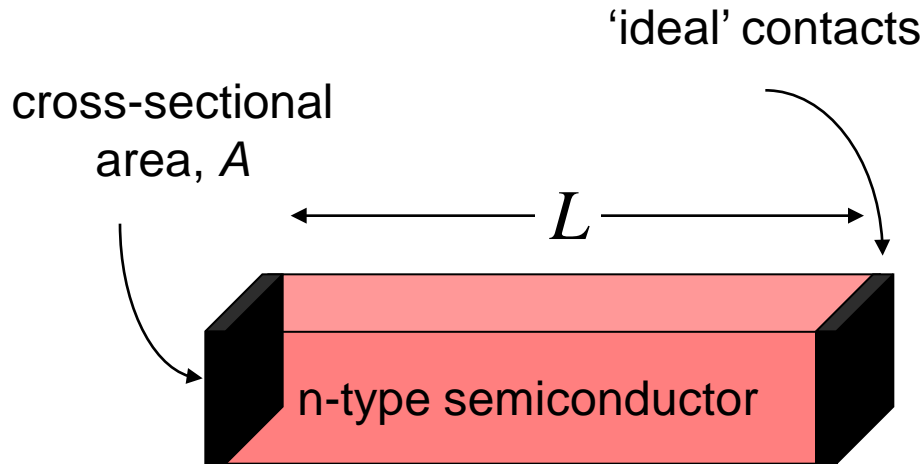
$$\langle v \rangle = \mu_n E = \mu_n \frac{V}{L} \text{ cm/sec}$$

$$\tau = \frac{L}{\langle v \rangle} = \frac{L^2}{\mu_n V}$$

$$I = \left\{ (nq\mu_n) \frac{A}{L} \right\} V = GV$$

$$G = \sigma \frac{A}{L} \quad \sigma = nq\mu_n$$

Landauer conductance and conductivity



$$G = \frac{2q^2}{h} \left(\int T(E) M(E) \left(-\frac{\partial f_0}{\partial E} \right) dE \right)$$

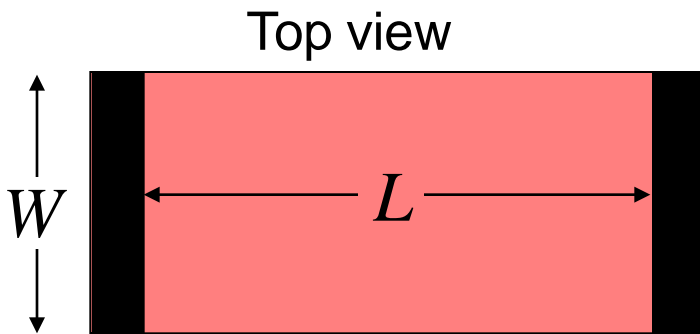
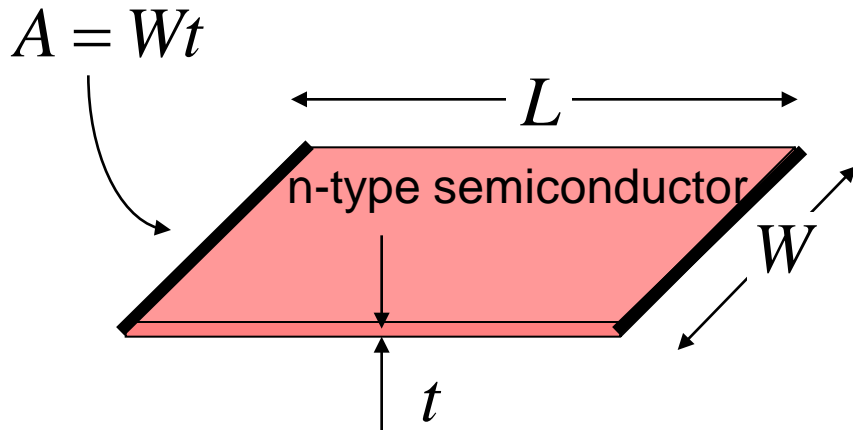
$$\sigma = \frac{G}{A/L} = \frac{2q^2}{h} \left(\int \lambda(E) \frac{M(E)}{A} \left(-\frac{\partial f_0}{\partial E} \right) dE \right) \quad (T(E) \approx \lambda(E)/L)$$

$$I = GV$$

$$G \propto \frac{A}{L} S (\Omega)^{-1}$$

$$G = \sigma \frac{A}{L} \quad \sigma = \frac{1}{\rho} \text{ S/m}$$

sheet conductance



$$I = GV$$

$$G = \sigma \frac{A}{L} \quad \sigma = nq\mu_n$$

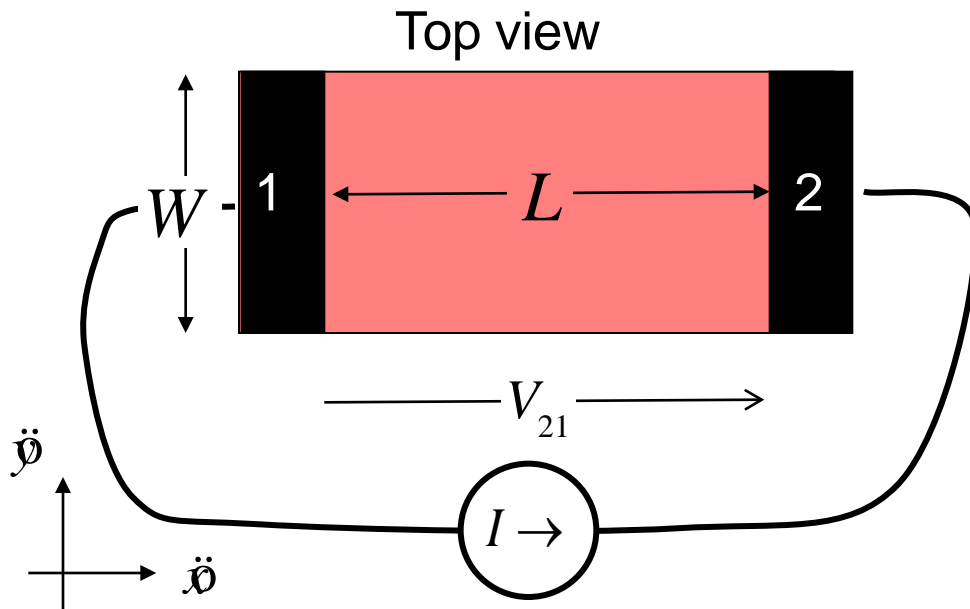
$$G = \sigma \left(\frac{Wt}{L} \right) = nq\mu_n t \left(\frac{W}{L} \right)$$

$$G = \sigma_s \left(\frac{W}{L} \right)$$

$$\sigma_s = n_s q \mu_n$$

'sheet conductance'

2-probe measurements



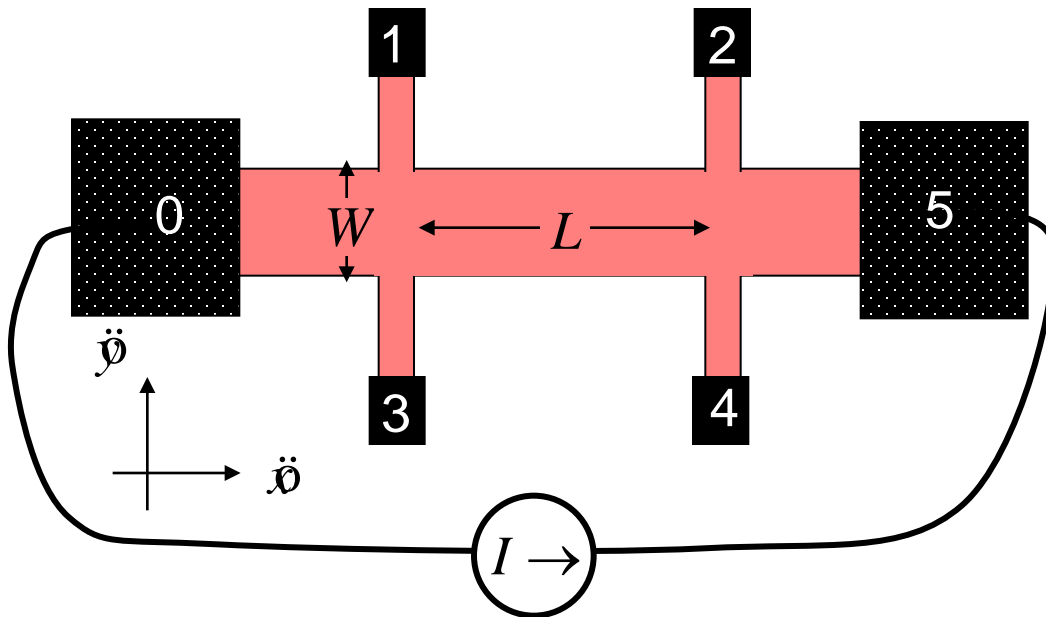
$$R_{DEV} = \rho_S \frac{L}{W}$$

$$V_{21} = I(2R_C + R_{DEV})$$

$$R_{DEV} \neq \frac{V_{21}}{I}$$

4-probe measurements

Top view



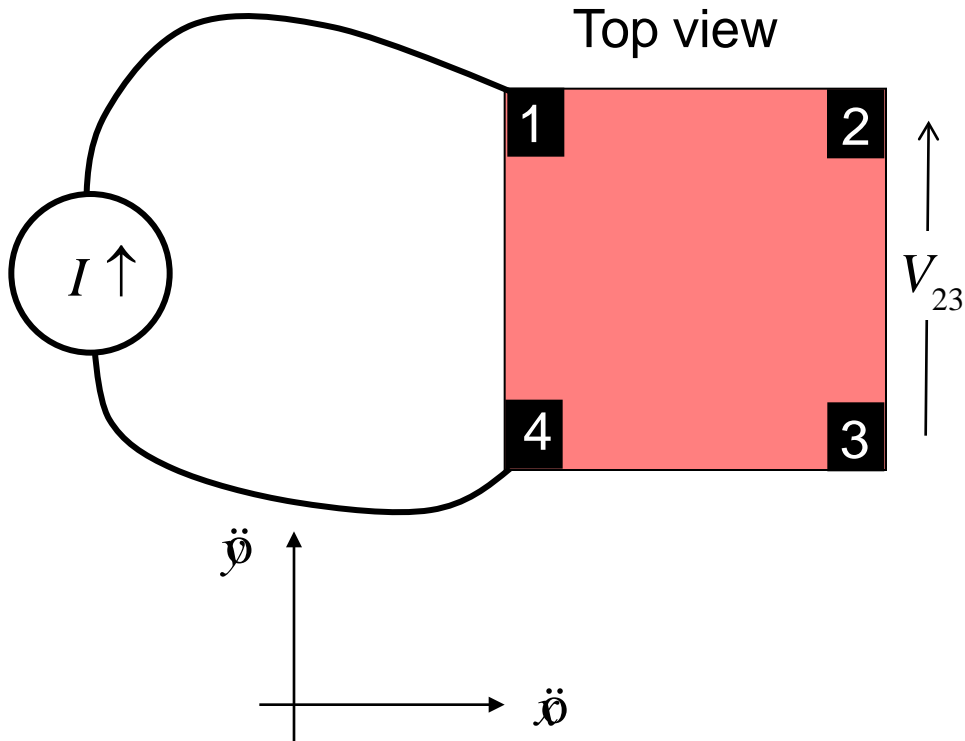
$$R_{DEV} = \rho_S \frac{L}{W}$$

$$V_{21} = I (R_{DEV})$$

(high impedance voltmeter)

“Hall bar geometry”

van der Pauw technique



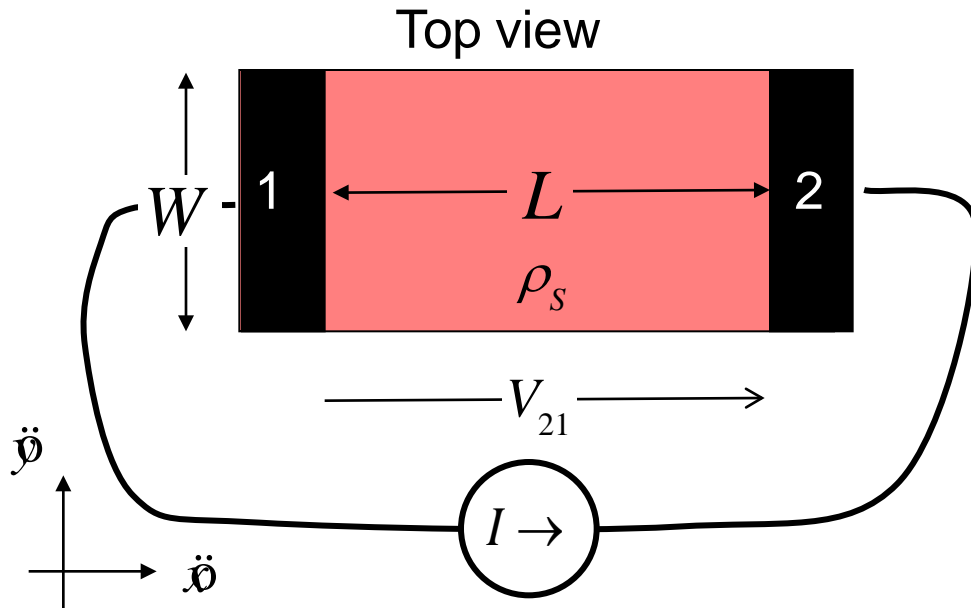
$$R_{1,4/2,3} \equiv \frac{V_{23}}{I_{14}} \quad R_{2,1/3,4} \equiv \frac{V_{34}}{I_{21}}$$

$$e^{-\frac{\pi}{\rho_S} R_{1,4/2,3}} + e^{-\frac{\pi}{\rho_S} R_{2,1/3,4}} = 1$$

$$\Rightarrow \rho_S$$

Force I through two contacts, measure V between the other two contacts.

isothermal vs. adiabatic



$$R_{DEV} = \rho_S \frac{L}{W}$$

$$V_{21} = R_{DEV} I - S \Delta T$$

$$I_Q = -\pi I - K_e \Delta T$$

i) $\Delta T = 0$

$$V_{21} = R_{DEV} I \rightarrow R_{DEV} = V_{21} / I$$

ii) $I_Q = 0$

$$\Delta T = (-\pi / K_e) I$$

$$V_{21} = R_{DEV} I + (S^2 T / K_e) I$$

$$V_{21} / I = R_{DEV} + (S^2 T / K_e)$$

for more about low-field measurements

D.C. Look, *Electrical Characterization of GaAs Materials and Devices*, John Wiley and Sons, New York, 1989.

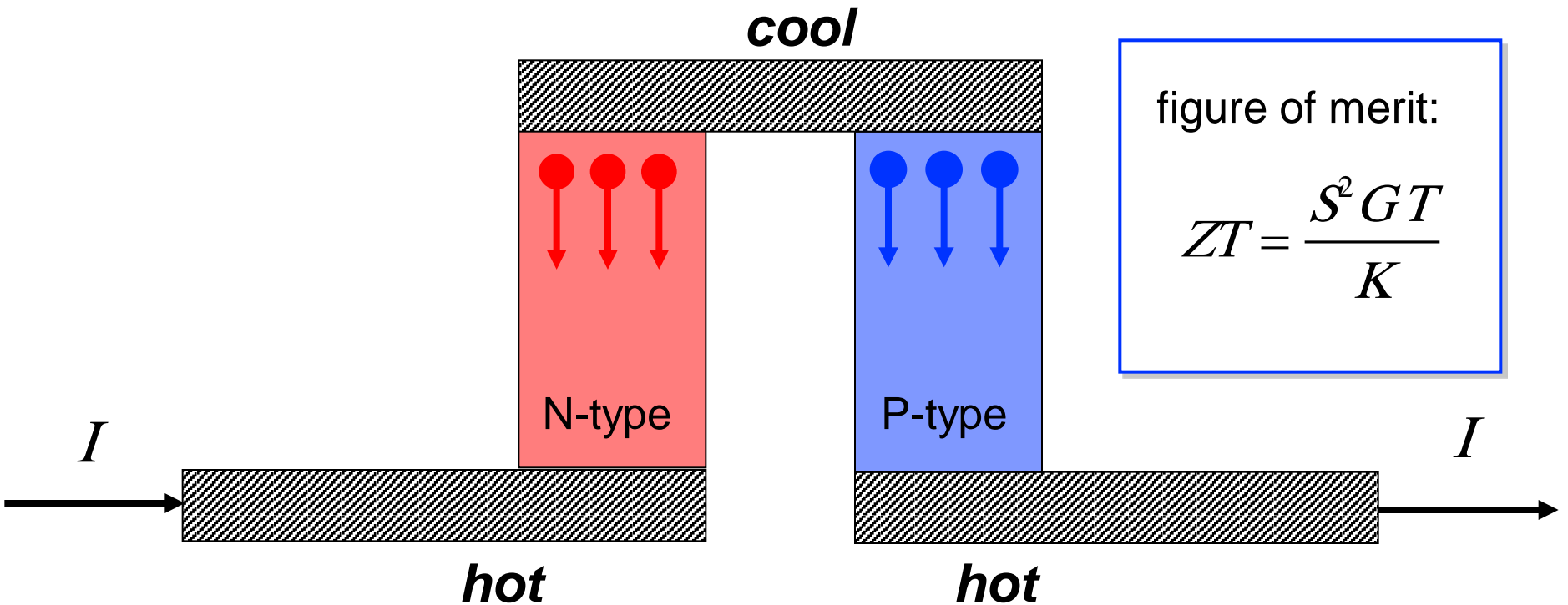
L.J. van der Pauw, “A method of measuring specific resistivity and Hall effect of discs of arbitrary shape,” *Phillips Research Reports*, vol. 13, pp. 1-9, 1958.

Lundstrom, Chapter 4, Sec. 7

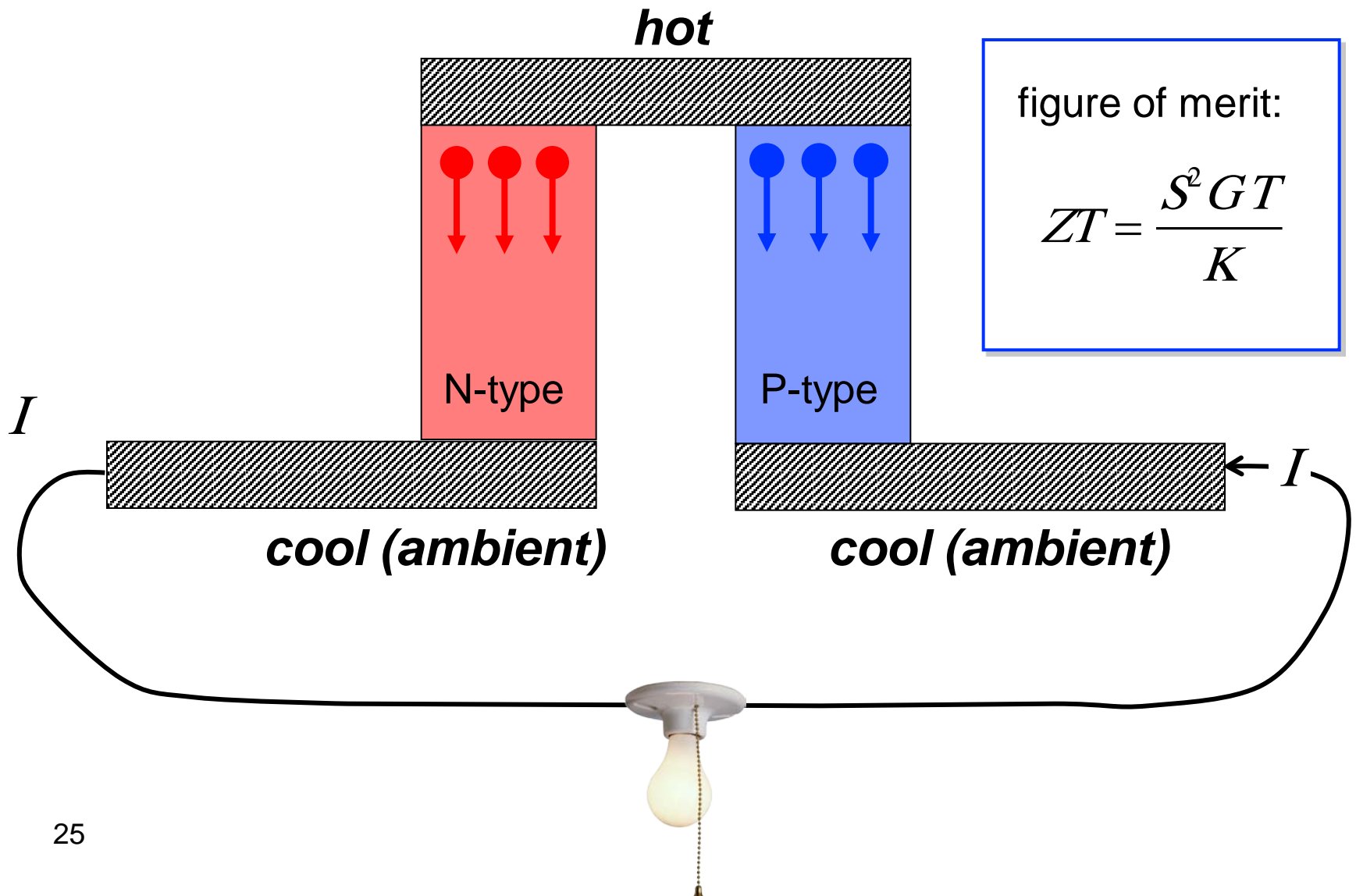
outline

- 1) Onsager relations
- 2) Measurement considerations
- 3) Thermoelectric devices**

thermoelectric devices: cooling



thermoelectric devices: power generation

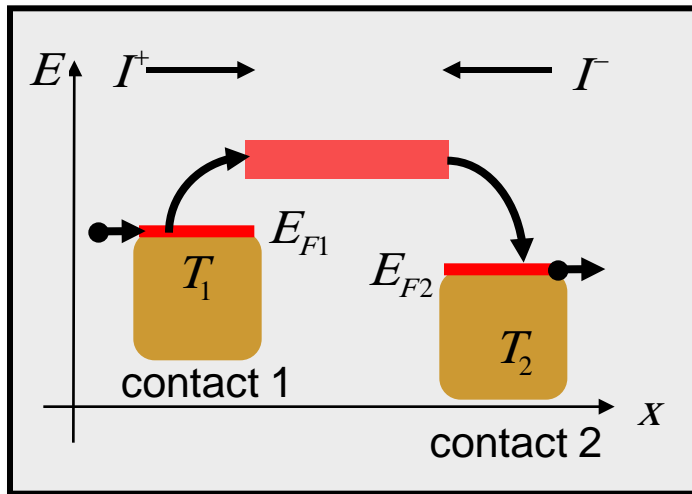


the TE industry: 2009

- 1) NASA: Radioisotope Thermoelectric Generators for deep-space missions.
- 2) DOD: cooling for lasers, night vision, sensors, guidance systems, etc.
- 1) World market for commercial TE power generation: \$30-50M/yr
- 2) World market for commercial TE cooling: \$200-250M/yr

Automobile market (seat cooler/heater) is growing. Vehicle waste heat recovery to improve fuel efficiency is a potential 'killer-app.'

TE heat engine



If we heat contact 1, and keep contact 2 cool, what is the maximum efficiency with which we can heat into useful work?

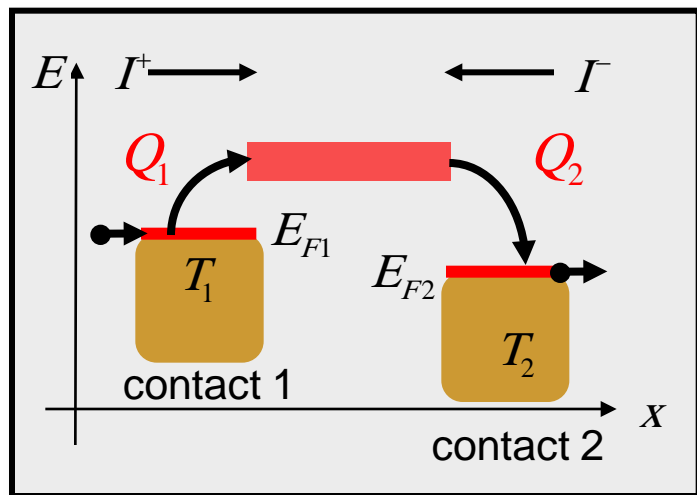
Answer:

Maximum efficiency:

$$\eta_{\max} = 1 - \frac{T_2}{T_1} = 1 - \frac{T_{\text{COLD}}}{T_{\text{HOT}}}$$

Carnot's theorem

what's necessary for current to flow?



$$\frac{Q_1}{T_1} < \frac{Q_2}{T_2}$$

for current to flow

$$I(E) = \frac{2q}{h} T(E) M(E) (f_1 - f_2)$$

Current flows when there is a difference in Fermi levels.

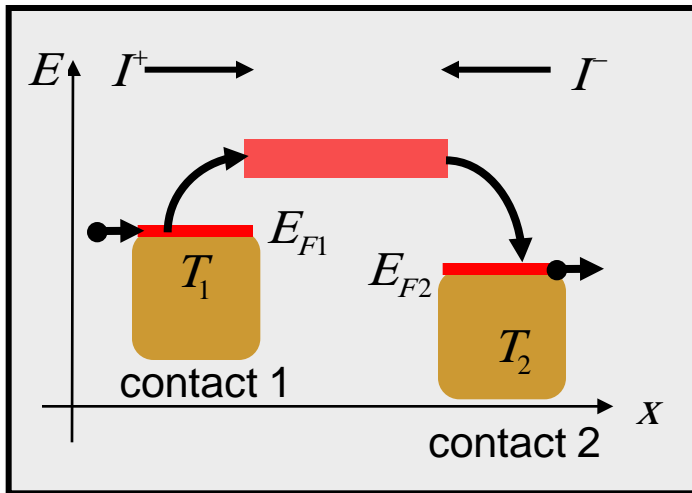
$$f_1(E) = \frac{1}{1 + e^{(E - E_{F1})/k_B T_1}}$$

$$f_2(E) = \frac{1}{1 + e^{(E - E_{F2})/k_B T_2}}$$

$$I(E) > 0 \Rightarrow f_1 > f_2$$

$$\Rightarrow \frac{(E - E_{F1})}{k_B T_1} < \frac{(E - E_{F2})}{k_B T_2}$$

maximum efficiency of a TE heat engine



$$Q_1/T_1 < Q_2/T_2$$

Maximum efficiency:

$$\eta_{\max} = 1 - \frac{T_2}{T_1} = 1 - \frac{T_{\text{COLD}}}{T_{\text{HOT}}}$$

Carnot's theorem

Thermal energy absorbed from contact 1:

$$Q_1 = (E - E_{F1})$$

Thermal energy dissipated into contact 2:

$$Q_2 = (E - E_{F2})$$

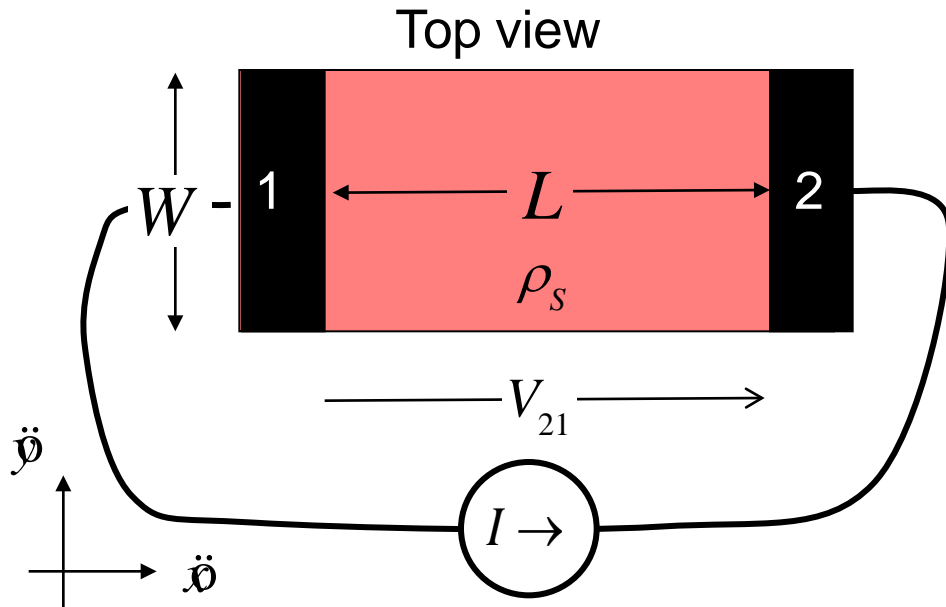
Net thermal energy absorbed from contact 1 that can be converted into useful work:

$$Q_1 - Q_2$$

Efficiency:

$$\eta = \frac{Q_1 - Q_2}{Q_1} = 1 - \frac{Q_2}{Q_1} = 1 - \frac{Q_2/T_2}{Q_1/T_1} \left(\frac{T_2}{T_1} \right)$$

maximum temperature difference

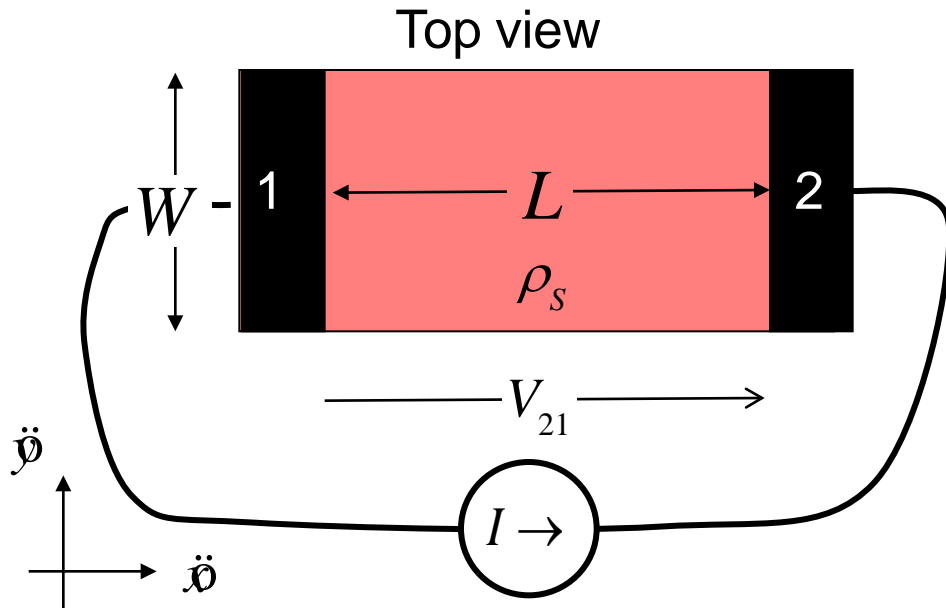


$$\Delta V = RI - S\Delta T$$

$$I_Q = -\pi I - K_e \Delta T$$

If we force a current through this thermoelectric device, what is the maximum temperature difference, $\Delta T = T_2 - T_1$, that we can produce?

maximum temperature difference



$$\Delta V = RI - S\Delta T$$

$$I_Q = -\pi I - K_e \Delta T$$

$$I_{Q1} = -TSI - K_e \Delta T - \frac{1}{2} I^2 R$$

For a fixed current, the temperature will rise until $I_Q = 0$

$$I_{opt} = -TS/R$$

$$\Delta T|_{max} = \frac{1}{2} \frac{T^2 S^2 G}{K}$$

$$\Delta T = -\frac{(TSI + I^2 R/2)}{K_e}$$

$$\frac{\partial \Delta T}{\partial I} = -\frac{(TS + IR)}{K_e} = 0$$

maximum temperature difference

$$\Delta T \Big|_{\max} = \frac{1}{2} \frac{T^2 S^2 G}{K}$$

$$\frac{\Delta T}{T} \Big|_{\max} = \frac{1}{2} \frac{S^2 GT}{K} = \frac{1}{2} ZT$$

$$ZT = \frac{S^2 GT}{K_e + K_L}$$

dimensionless figure of merit

K_e : electronic heat conductance

K_L : lattice heat conductance

figure of merit

Power factor:

- large Seebeck coefficient
- high conductance
- depends on material parameters

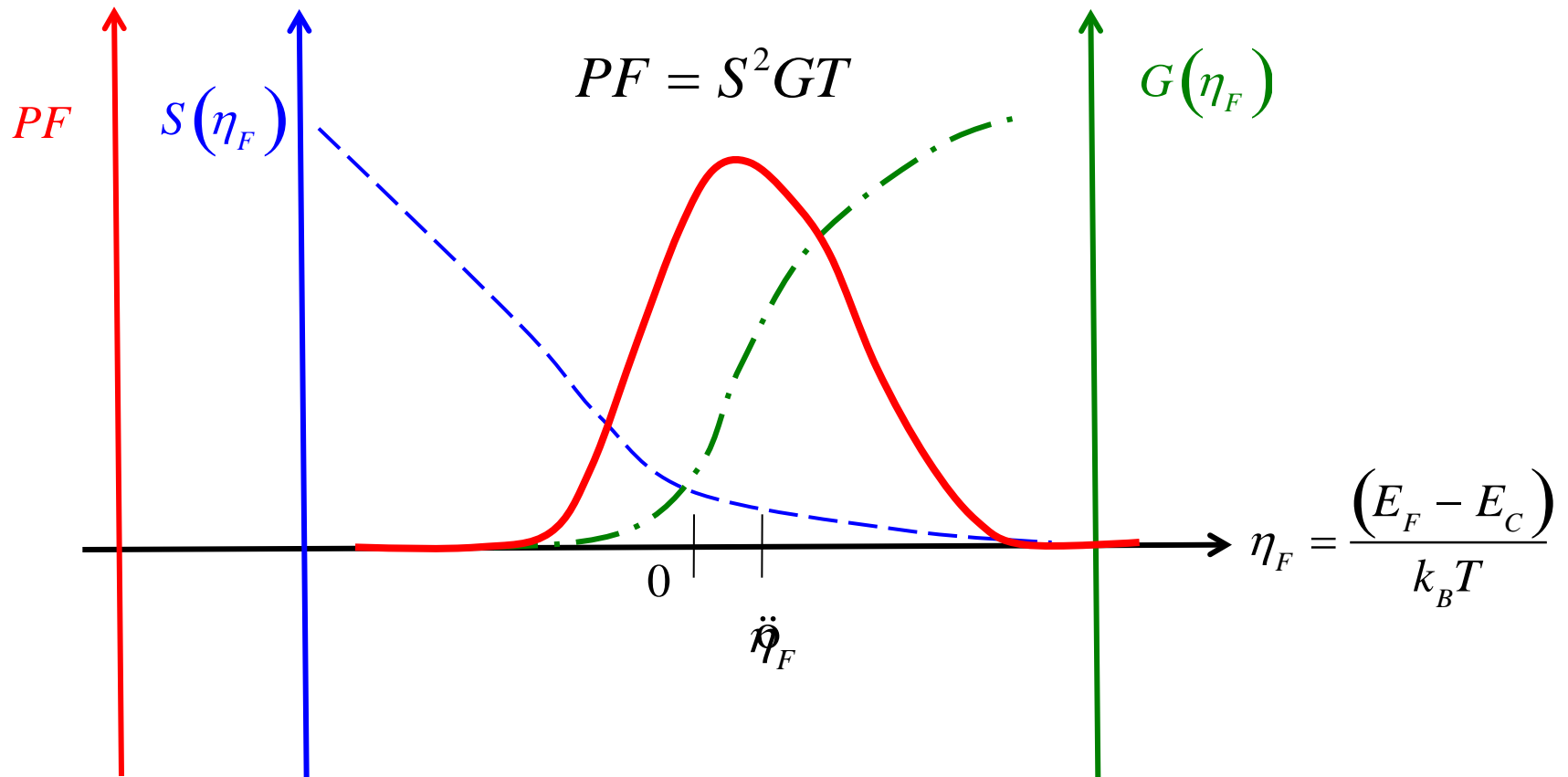
$$ZT = \frac{S^2 GT}{K_e + K_L}$$

thermal conductance:

- $K_L > K_e$
- need short mfp for phonons
(alloys with large mass difference –
e.g. Bi_2Te_3 , SiGe)
- nanostructuring for reducing mfp
- “phonon glass electron crystal”

G.D. Mahan and J.O. Sofo, “The Best Thermoelectric,” *Proc. Natl. Acad. Sci.*, **93**, 7436, 1996

optimal power factor

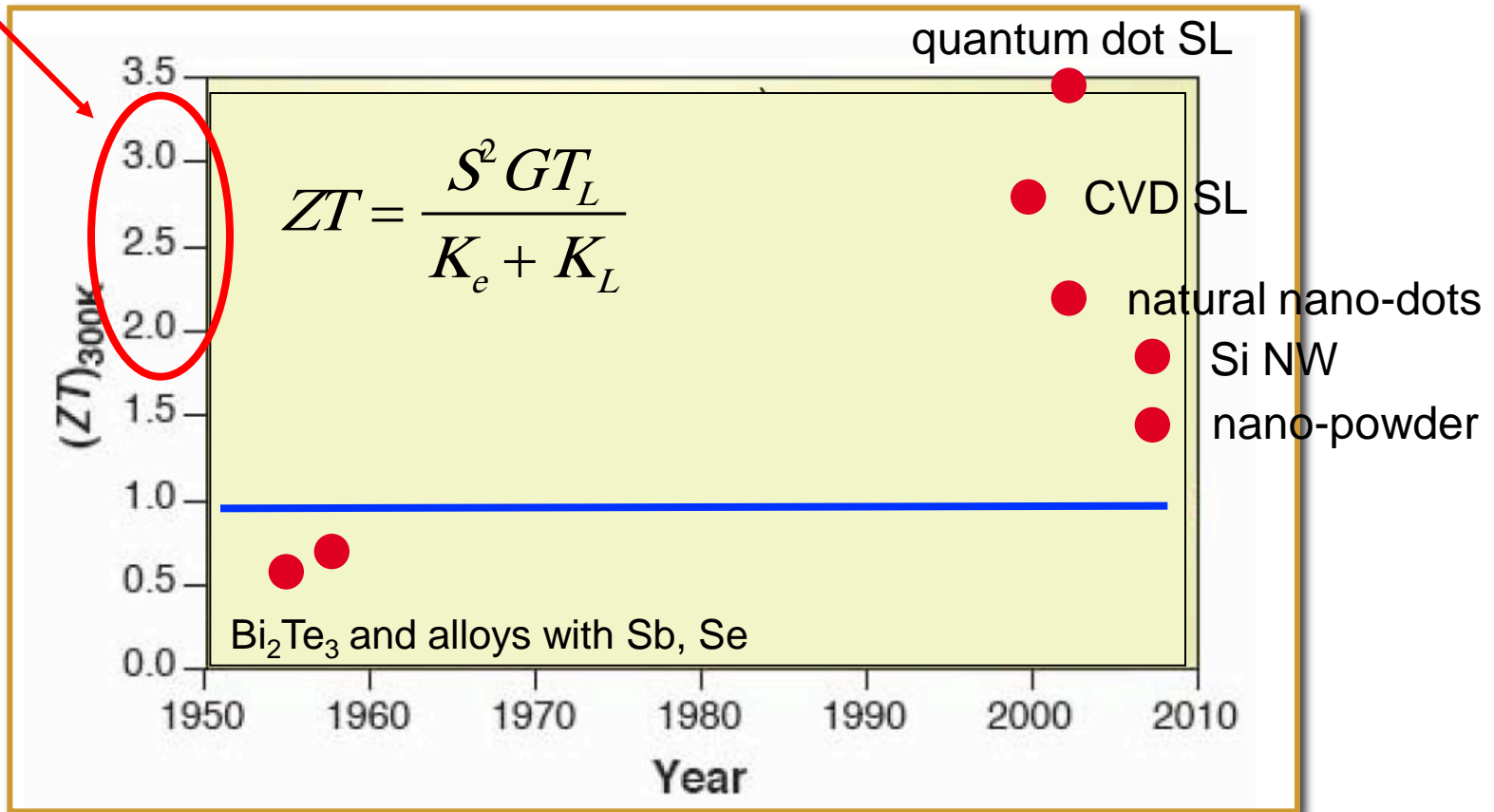


The optimum power factor occurs when E_F is near E_C .

history and prospects

holy grail

nanostructured thermoelectrics



questions

- 1) Onsager relations
- 2) Measurement considerations
- 3) Thermoelectric devices

