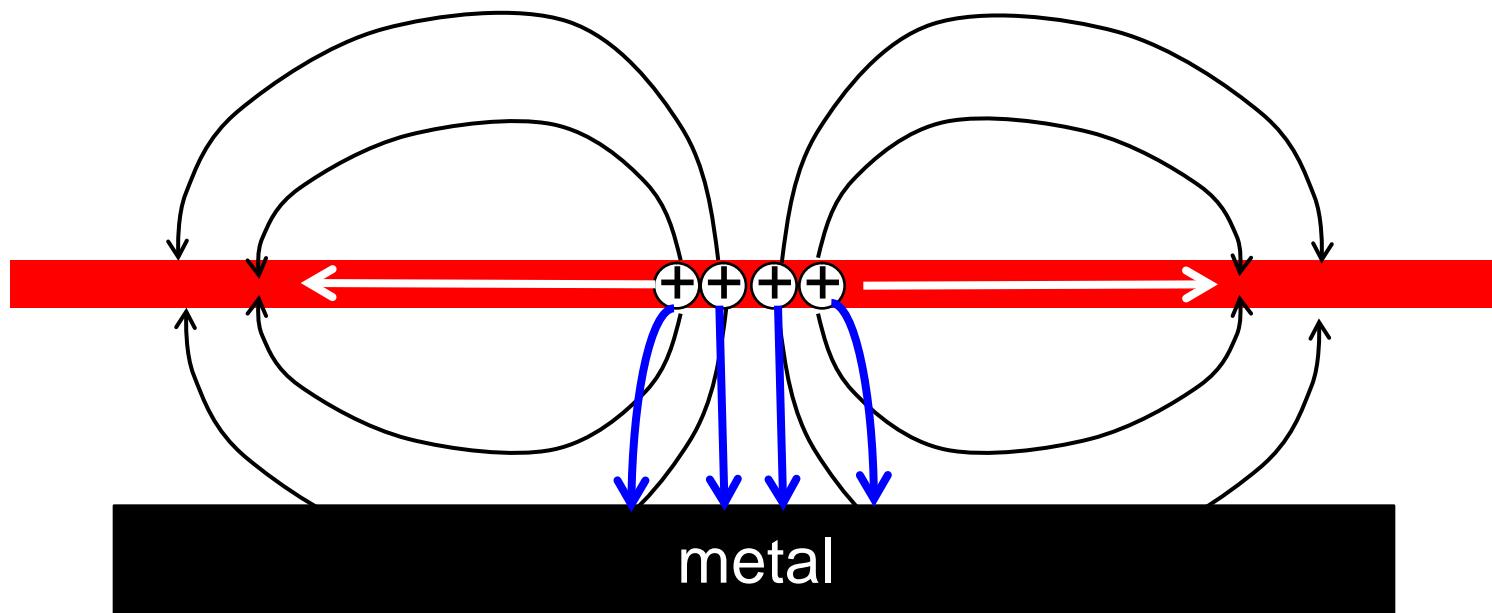


ECE-656: Fall 2009

Lecture 23: Phonon Scattering I

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screening in 2D / 1D

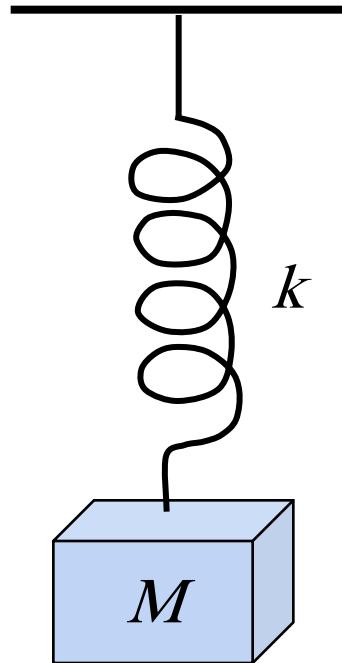


geometric screening

outline

- 1) About phonons**
- 2) Electron-phonon coupling**
- 3) Summary**

springs



$$E = \left(n + \frac{1}{2} \right) \hbar \omega$$

$$U = \frac{1}{2} k (x - x_0)^2$$

$$F = -\frac{dU}{dx} = -k(x - x_0)$$

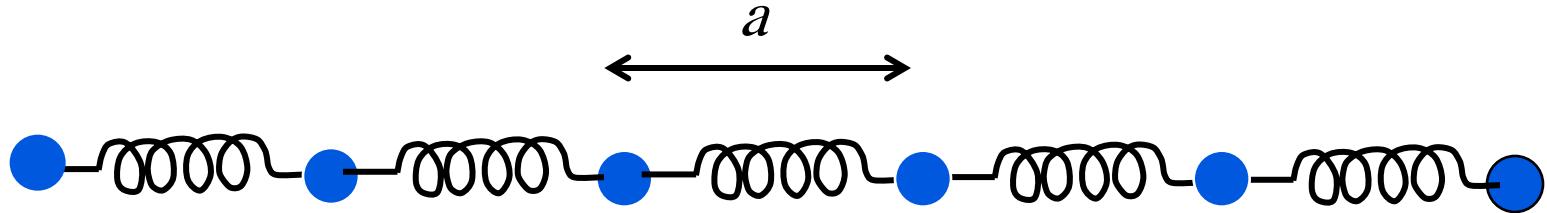
$$M \frac{d^2 x}{dt^2} = -k(x - x_0)$$

$$x(t) - x_0 = A e^{i \omega t}$$

$$\omega = \sqrt{k/M}$$

$$E \sim A^2$$

lattice vibrations



$$\vec{u}(x) = A \hat{e}_\nu e^{i(\beta x - \omega t)}$$

$\longrightarrow X$

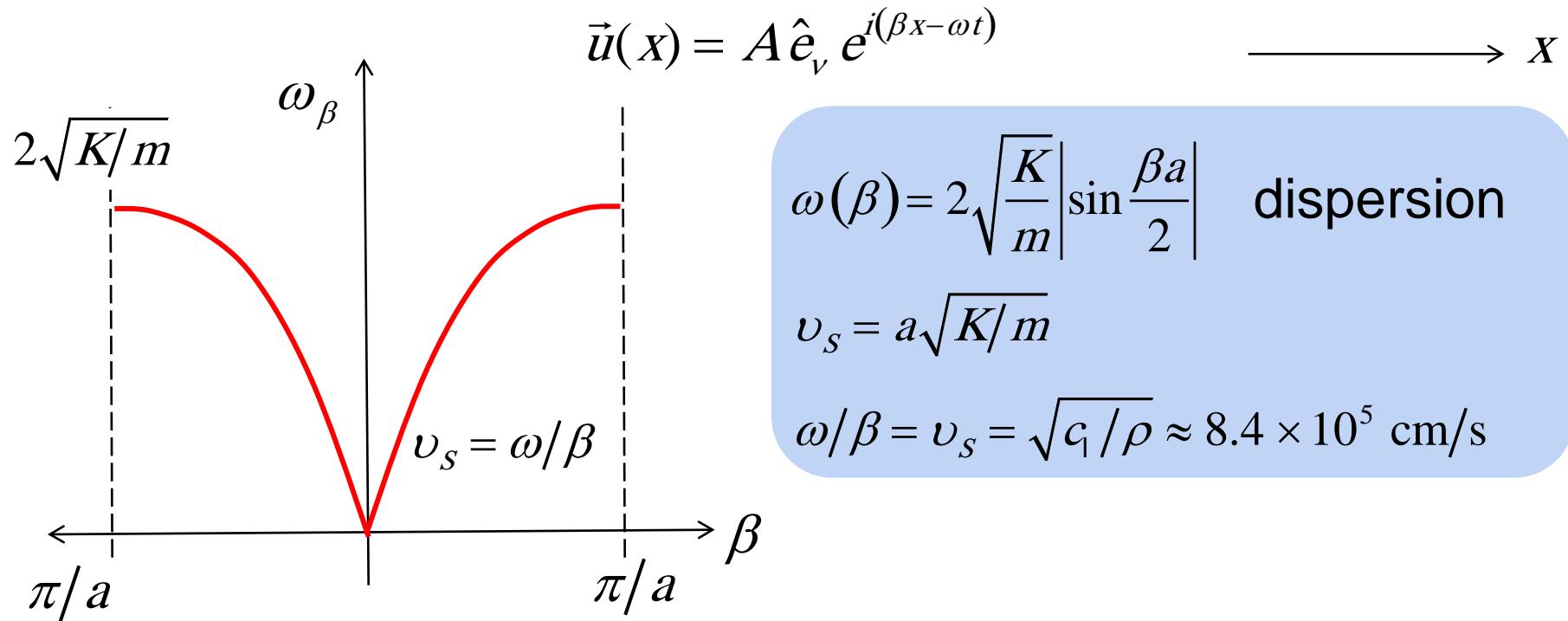
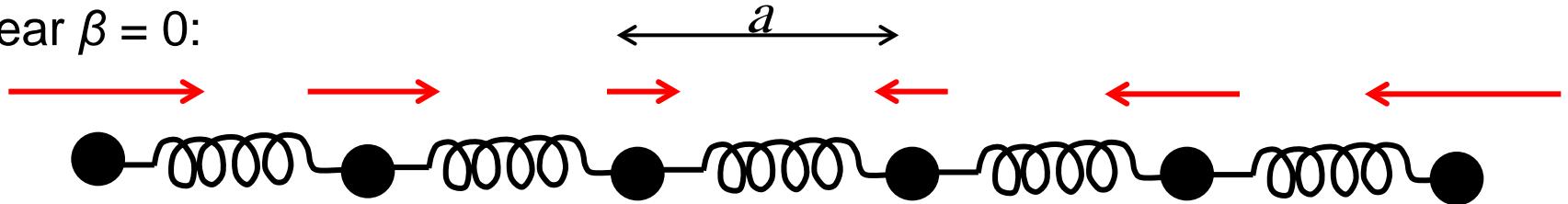
$\hat{e}_\nu = \hat{x}$ longitudinal

$\hat{e}_\nu = \hat{y}$ transverse

$\hat{e}_\nu = \hat{z}$ transverse

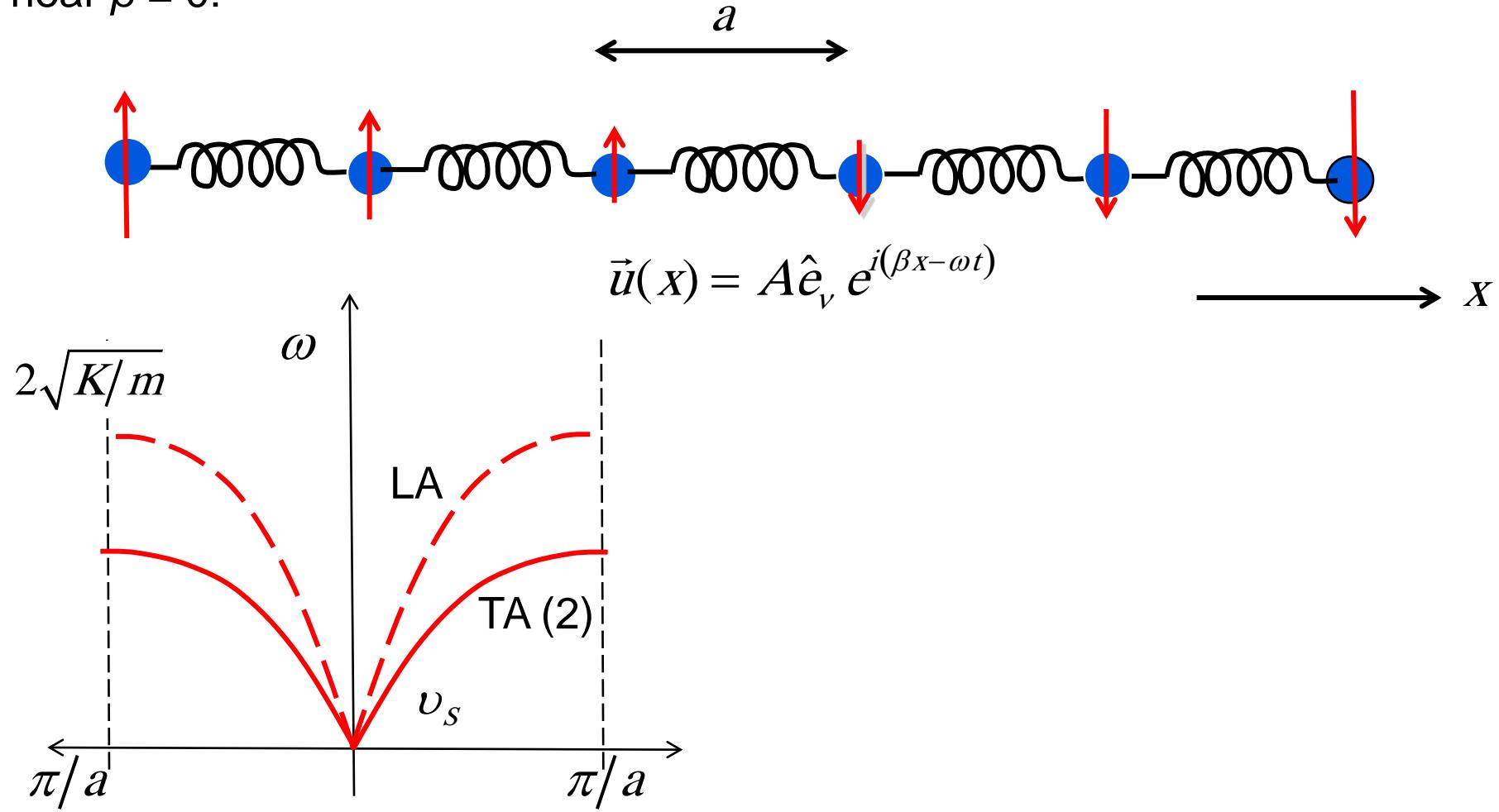
LA phonons (1)

near $\beta = 0$:

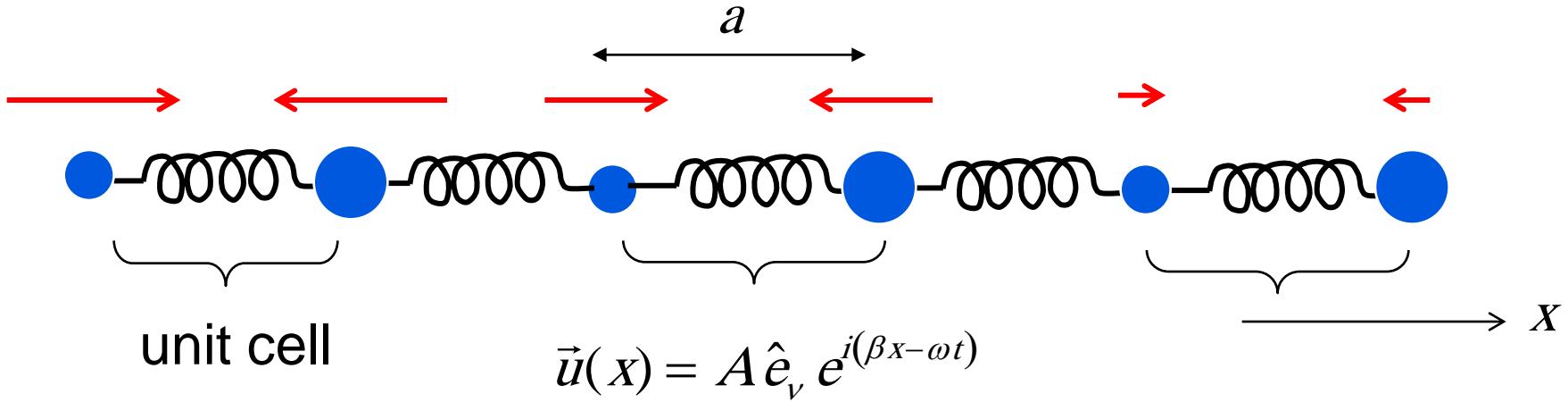


TA phonons (2)

near $\beta = 0$:

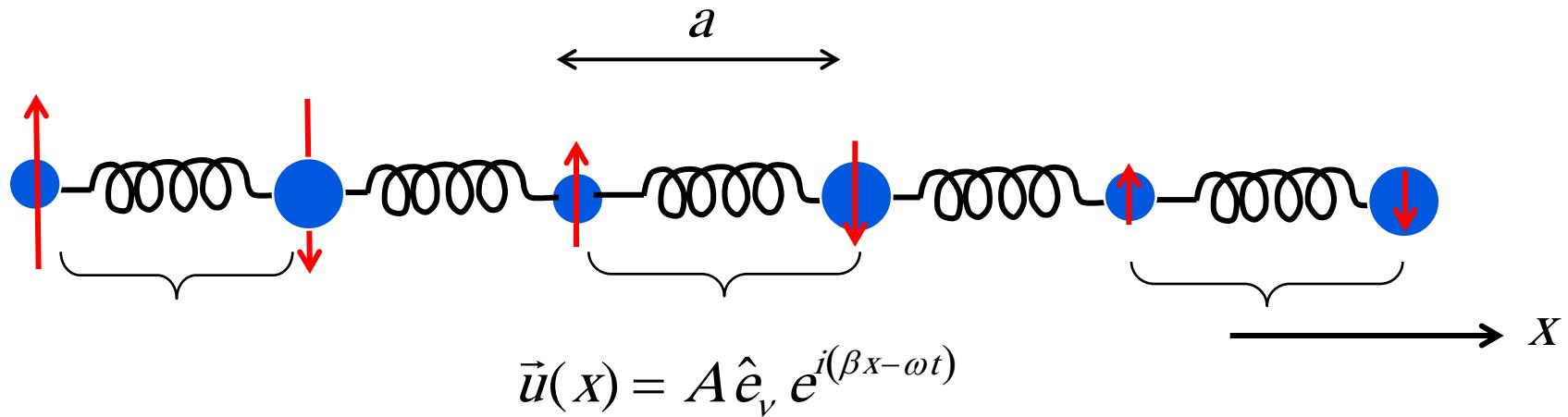


LO phonons (1)



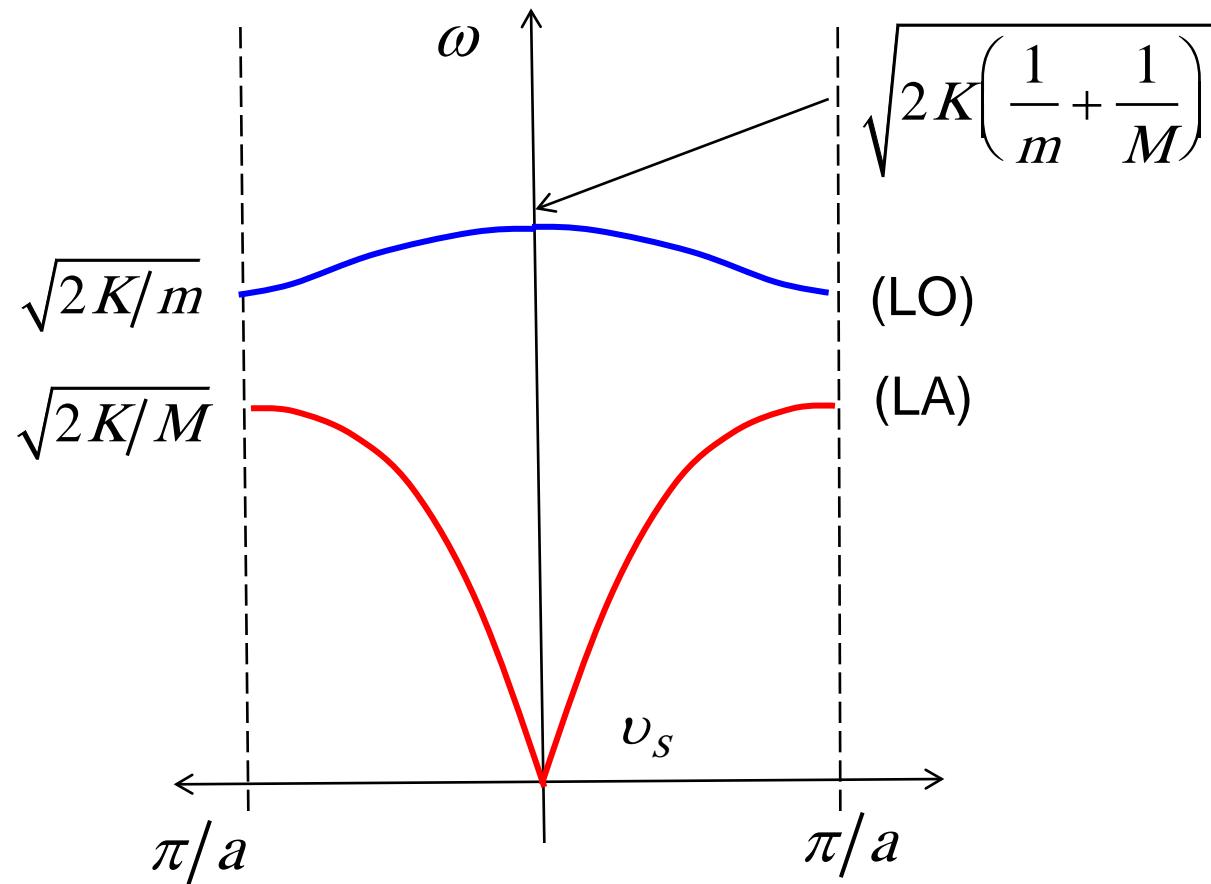
the two atoms in a unit cell oscillate out of phase

TO phonons (2)



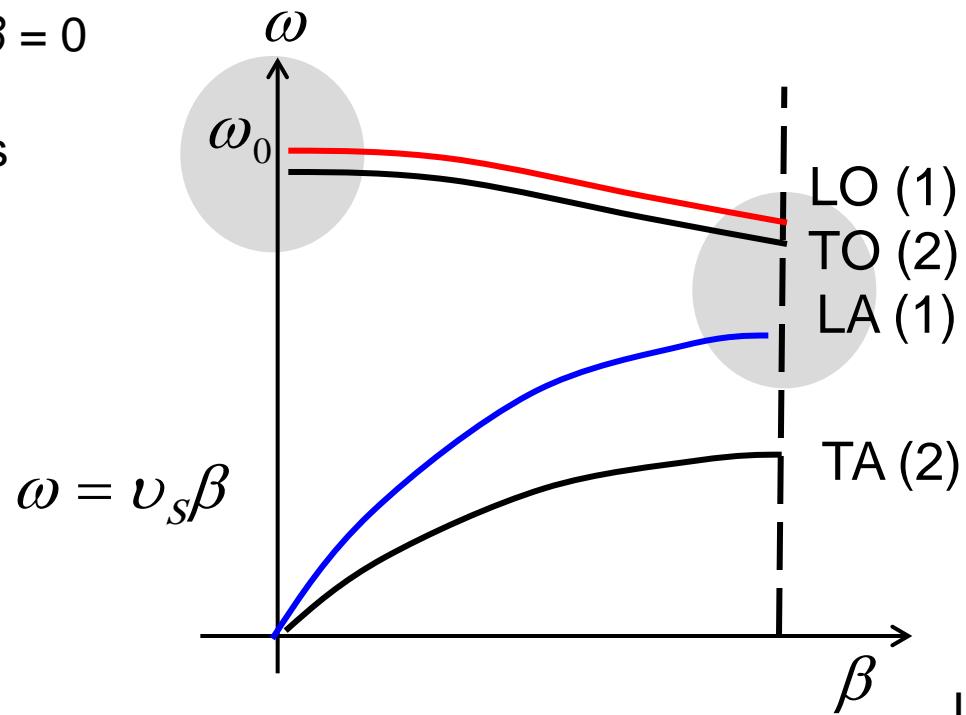
the two atoms in a unit cell oscillate out of phase

1D spring model (longitudinal modes)



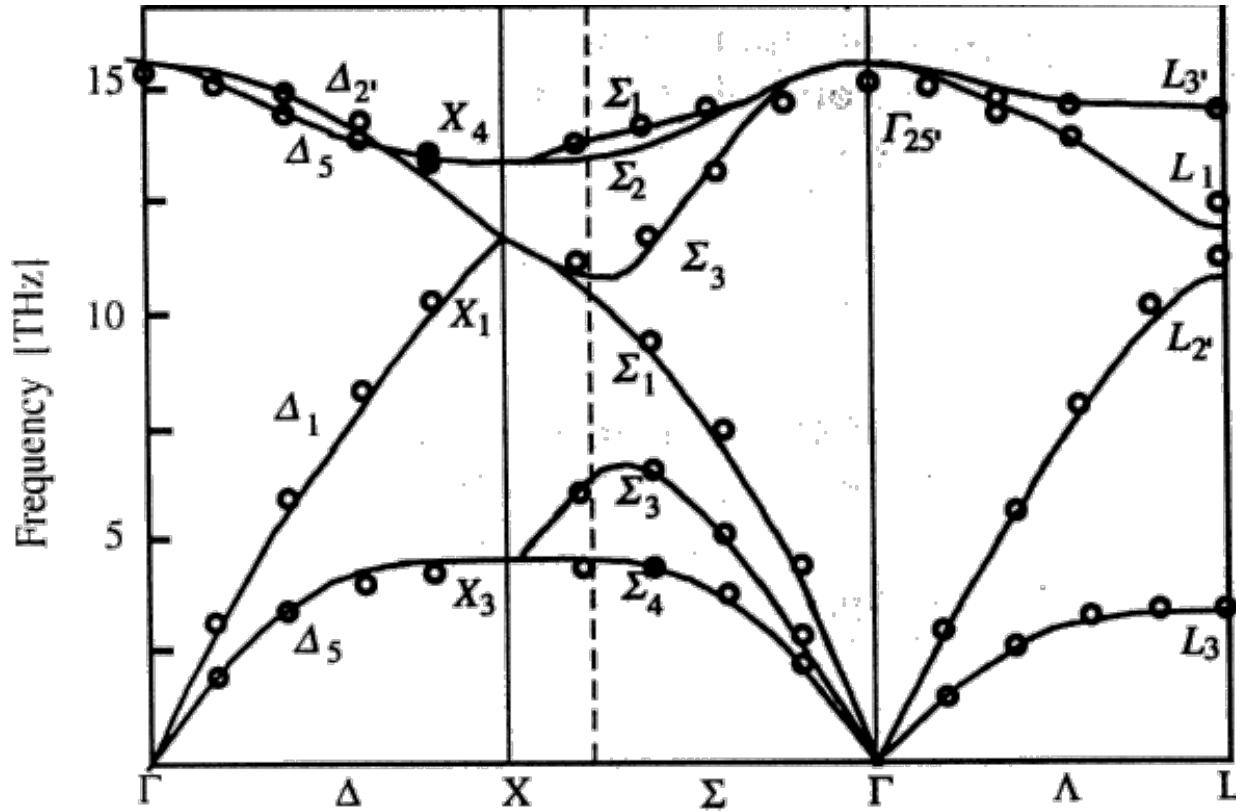
phonon dispersion characteristics

LO and TO
degenerate at $\beta = 0$
for non polar
semiconductors



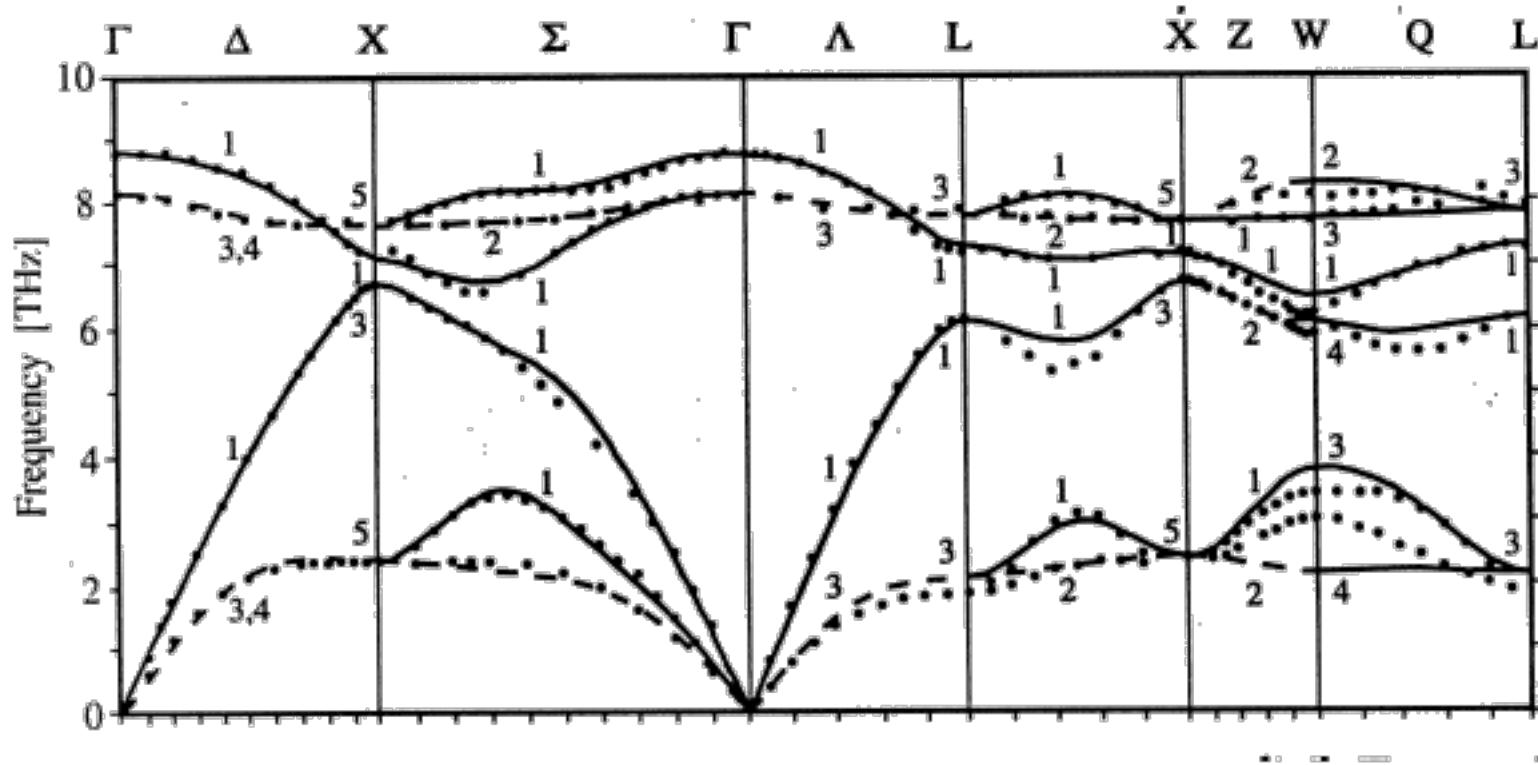
LO and LA
degenerate at zone
boundary for non
polar semiconductors

measured phonon dispersion: Si



<http://www.personal.psu.edu/pce3/Research%20Topics.htm>

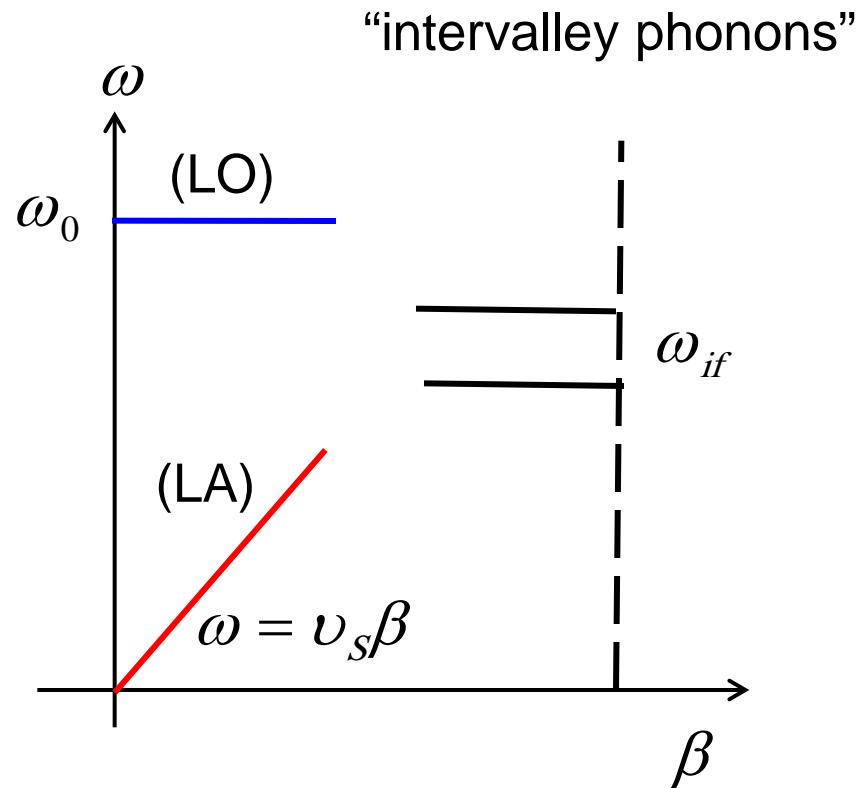
measured phonon dispersion: GaAs



<http://www.personal.psu.edu/pce3/Research%20Topics.htm>

simplified phonon dispersion

- 1) Longitudinal modes couple most strongly with electrons.
- 2) Intravalley scattering requires small β .
- 3) Intervalley requires β near the Brillouin zone boundary.



amplitude of vibration

$$\vec{u}(x) = A \hat{e}_\nu \left[e^{i(\beta x - \omega t)} + e^{-i(\beta x - \omega t)} \right]$$

$$u(x_j) = 2|A| \cos(\beta aj - \omega_\beta t) \quad x_j = aj$$

$$\frac{du_j}{dt} = 2|A|\omega \sin(\beta aj - \omega_\beta t)$$

$$\langle KE \rangle = \frac{1}{2} m \left\langle \left| \frac{du_j}{dt} \right|^2 \right\rangle = 2m|A|^2 \omega^2 \left\langle \sin^2(\beta aj - \omega_\beta t) \right\rangle = m|A|^2 \omega^2$$

$$\langle E \rangle = \langle KE \rangle + \langle PE \rangle = 2m|A|^2 \omega^2$$

$$\langle E_{TOT} \rangle = N_a 2m|A|^2 \omega^2 = \Omega \rho 2|A|^2 \omega^2$$

amplitude of vibration (ii)

$$\langle E_{TOT} \rangle = \Omega \rho 2|A|^2 \omega^2$$

quantum mechanics says: $\langle E_{TOT} \rangle = \hbar\omega$

$$|A|^2 = \frac{\hbar}{2\Omega\rho\omega}$$

$$u_j = (A e^{i(\beta x - \omega t)} + cc) = 2|A| \cos(\beta aj - \omega_\beta t)$$

$$u_j = \sqrt{\frac{\hbar}{2\Omega\rho\omega}} \cos(\beta aj - \omega_\beta t) \quad x = aj \quad j = 1, 2, 3, \dots$$

outline

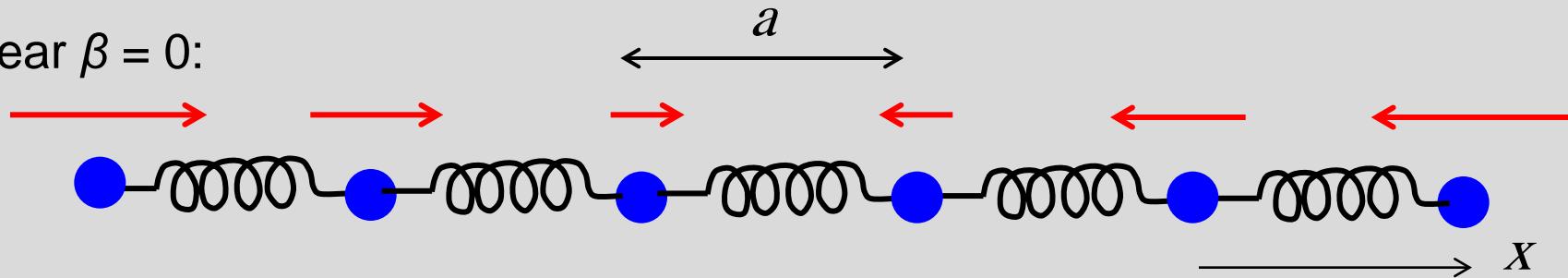
- 1) About phonons
- 2) Electron-phonon coupling**
- 3) Summary

electron-phonon coupling (LA)

the bandgap depends on lattice constant: $\delta E_G = D \frac{\delta a}{a}$
 $\delta E_C = D_C \frac{\delta a}{a}$ ‘deformation potential’

LA phonons:

near $\beta = 0$:



$$\delta a = u(x) - u(x-a) = u(x) - \left\{ u(x) - \frac{\partial u}{\partial x} a \right\} = \frac{\partial u}{\partial x} a$$

$$\frac{\delta a}{a} = \frac{\partial u}{\partial x} \quad \text{“strain”}$$

$$U_s = \delta E_C = D_C \frac{\delta a}{a}$$

deformation potential scattering

$$\frac{\delta a}{a} \approx \frac{\partial u}{\partial x} \quad \text{“strain”} \quad U_s = \delta E_c = D_c \frac{\delta a}{a} = D_c \frac{\partial u_\beta}{\partial x}$$

$$u_\beta(x, t) = A_\beta e^{\pm i(\beta z - \omega t)}$$

$$U_s = D_A \frac{\partial u_\beta}{\partial x} = \pm i\beta D_A u_\beta = K_\beta u_\beta$$

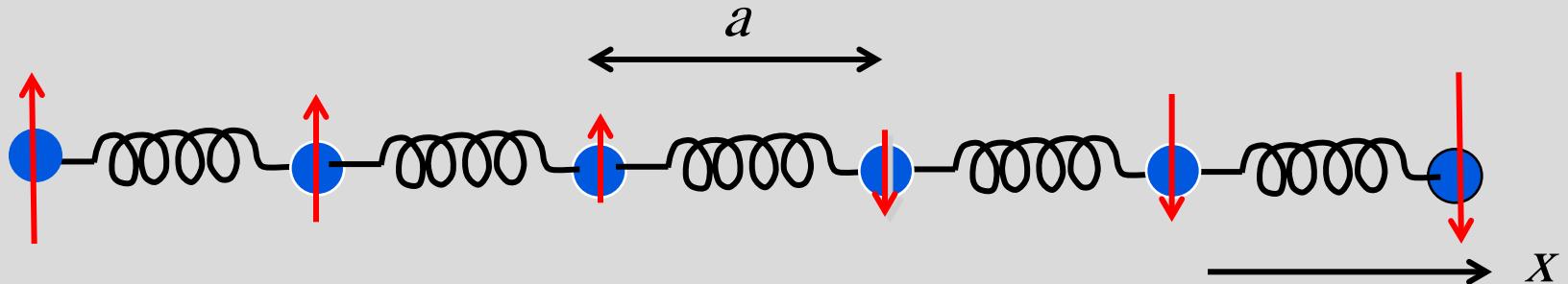
$$|K_\beta|^2 = \beta^2 D_A^2$$

“acoustic deformation potential scattering (ADP)”

electron-phonon coupling (TA)

bandgap depends on lattice constant: $\delta E_G = D \frac{\delta a}{a}$

near $\beta = 0$:

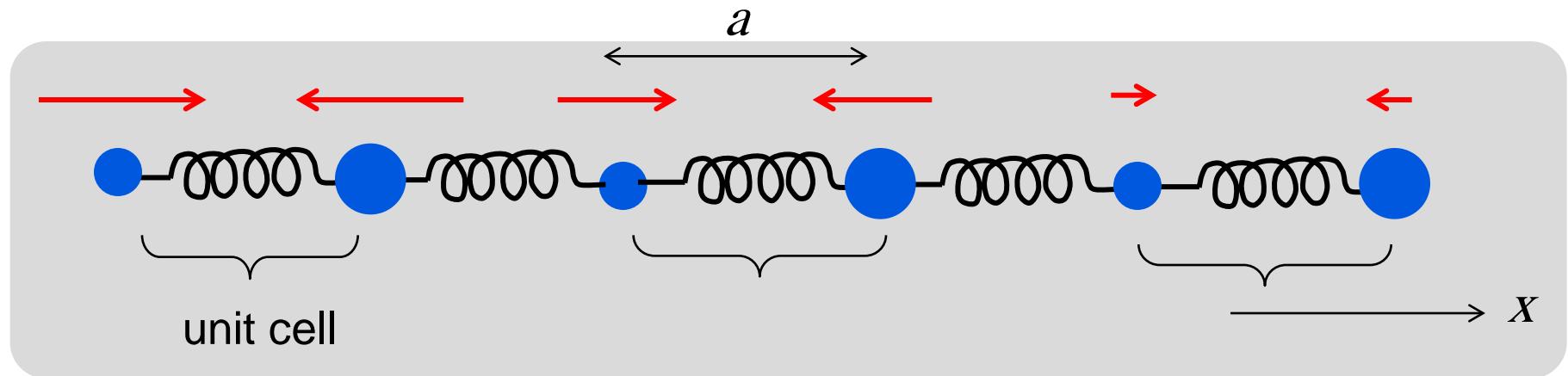


$$\frac{\delta a}{a} \propto \left(\frac{\partial u}{\partial X} \right)^2$$

To first order, *only* LA phonons scatter electrons

electron-phonon coupling (LO)

$$\delta E_C = D_C \frac{\delta a}{a} \quad \text{deformation potential}$$



$$\delta a(x) = u_\beta(x)$$

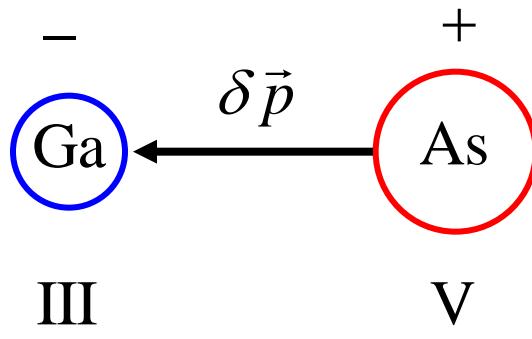
$$u_\beta(x, t) = A_\beta e^{\pm i(\beta z - \omega t)}$$

$$U_S \approx D_O u_\beta = K_\beta u_\beta$$

$$|K_\beta|^2 = D_O^2$$

“optical deformation potential scattering (ODP)”

polar semiconductors



$$u_\beta(x, t) = A_\beta e^{\pm i(\beta z - \omega t)}$$

$$U_s = q \int \frac{\delta p_x}{\epsilon_0 V_u} dx$$

$$\delta \vec{P} = \frac{\delta \vec{p}}{V_u}$$

$$D_x = \epsilon_0 \mathcal{E}_x + P_x$$

$$\delta D = \epsilon_0 \delta \mathcal{E} + \delta P$$

$$\nabla \bullet \delta D = 0$$

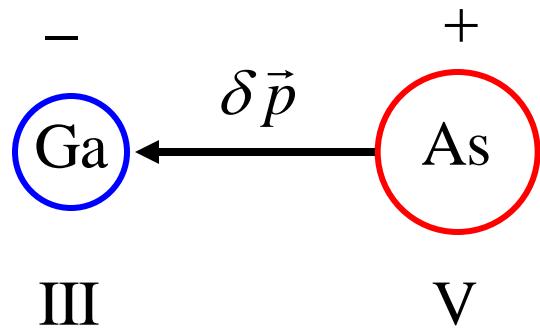
$$\delta D = \delta D_x e^{i \beta_x x} \rightarrow \delta D = 0$$

$$\delta \mathcal{E}_x = - \frac{\delta P_x}{\epsilon_0} = - \frac{\delta p_x}{\epsilon_0 V_u}$$

$$U_s = -q \int \mathcal{E}_x dx$$

optical phonons in polar semiconductors

$$\delta p = q^* u_\beta$$



$$u_\beta(x, t) = A_\beta e^{\pm i(\beta x - \omega t)}$$

$$U_s = \frac{q q^* u_\beta}{i \beta \epsilon_0 V_u}$$

$$\left(\frac{q^*}{V_u} \right)^2 = \frac{\epsilon_0 \rho \omega_0^2}{\kappa_0} \left(\frac{\kappa_0}{\kappa_\infty} - 1 \right)$$

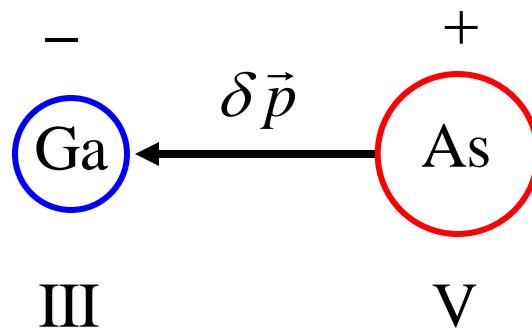
$$U_s = K_\beta u_\beta$$

$$U_s = q \int \frac{\delta p_x}{\epsilon_0 V_u} dx$$

$$|K_\beta|^2 = \frac{\rho q^2 \omega_0^2}{\beta^2 \kappa_0 \epsilon_0} \left(\frac{\kappa_0}{\kappa_\infty} - 1 \right)$$

small angle scattering dominates

acoustic phonons in polar semiconductors



$$u_\beta(x, t) = A_\beta e^{\pm i(\beta x - \omega t)}$$

$$\delta p = q^* \frac{\partial u_\beta}{\partial x} = q^* i \beta u_\beta$$

$$U_s = \frac{q q^*}{\epsilon_0 V_u} u_\beta = \frac{q e_{PZ}}{\kappa_0 \epsilon_0} u_\beta$$

$$U_s = K_\beta u_\beta$$

$$U_s = q \int \frac{\delta p_x}{\epsilon_0 V_u} dx$$

$$|K_\beta|^2 = \frac{q^2 e_{PZ}^2}{\kappa_0^2 \epsilon_0^2}$$

scattering potentials

$$u_\beta(\vec{r}, t) = A_\beta e^{\pm i(\vec{\beta} \cdot \vec{r} - \omega_\beta t)} \quad U_s = K_\beta u_\beta$$

ADP $|K_\beta|^2 = \beta^2 D_A^2$

ODP $|K_\beta|^2 = D_o^2$

PZ $|K_\beta|^2 = (qe_{PZ}/\kappa_s \epsilon_0)^2$

POP $|K_\beta|^2 = \frac{\rho q^2 \omega_0^2}{\beta^2 \kappa_0 \epsilon_0} \left(\frac{\kappa_0}{\kappa_\infty} - 1 \right)$

other scattering mechanisms

- 1) Neutral impurity
- 2) Alloy scattering
- 3) Surface / edge roughness scattering
- 4) Plasmon scattering
- 5) Electron-electron scattering
- 6) Electron-hole

questions

- 1) About phonons
- 2) Electron-phonon coupling
- 3) Summary

