

ECE-656: Fall 2009

**Lecture 30:
Balance Equation Approach: III**

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general approach

$$\frac{\partial f}{\partial t} + v_x \frac{\partial f}{\partial x} - \frac{q\mathcal{E}_x}{\hbar} \frac{\partial f}{\partial k_x} = \hat{C} f \quad \Rightarrow \quad \frac{\partial n_\phi}{\partial t} = -\nabla \cdot \vec{F}_\phi + G_\phi - R_\phi$$

$$n_\phi(x, t) = \frac{1}{L} \sum_p \phi(p) f(x, p, t) \quad G_\phi = -q\mathcal{E}_x \left\{ \frac{1}{L} \sum_{p_x} \frac{\partial \phi}{\partial p_x} f \right\}$$

$$F_\phi \equiv \frac{1}{L} \sum_{p_x} \phi(p_x) v_x f(x, p_x, t) \quad R_\phi \equiv \frac{n_\phi - n_\phi^0}{\langle \tau_\phi \rangle}$$

the four balance equations

$$\frac{\partial n_L(x, t)}{\partial t} = - \frac{d[I_{nx}/(-q)]}{dx}$$

0th moment of BTE

$$\langle \tau_m \rangle \frac{\partial I_x(x, t)}{\partial t} + I_x = n_L q \mu_n \mathcal{E}_x + 2 \mu_n \frac{dW}{dx}$$

1st moment of BTE

$$\frac{\partial W(x, t)}{\partial t} = - \frac{dF_W}{dx} + I_x \mathcal{E}_x - \frac{(W - W_0)}{\langle \tau_E \rangle}$$

2nd moment of BTE

$$\frac{\partial F_W(x, t)}{\partial t} = - \frac{dX}{dx} - \frac{3q}{m^*} n_L u \mathcal{E}_x - \frac{F_W}{\langle \tau_{F_W} \rangle}$$

3rd moment of BTE

outline

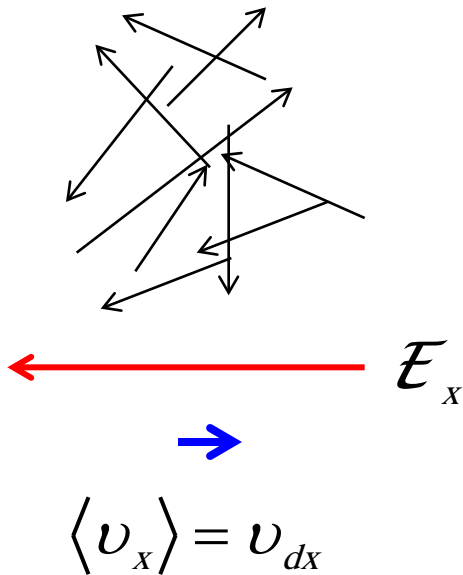
- 1) Carrier Temperature and Heat Flux**
- 2) Balance equations in 3D
- 3) Heterostructures
- 4) Summary

electron temperature

- Electrons gain energy from the electric field.
- Electrons lose energy to the lattice by inelastic scattering.
- If the electric field is high, electrons will gain energy faster than they lose it to the lattice.
- Under such conditions, the electron energy $>$ lattice energy.
- If we measure the electron energy by a “temperature,” then the electrons are *hot*.

electron temperature

random thermal motion of electrons



$$W = \frac{1}{L} \sum_{k_x} E(k) f = \frac{1}{L} \sum_{k_x} \frac{1}{2} m^* v_x^2 f$$

$$W = n_L \frac{1}{2} m^* \langle v_x^2 \rangle$$

$$v_x = v_{dx} + c \quad \langle c \rangle \equiv 0$$

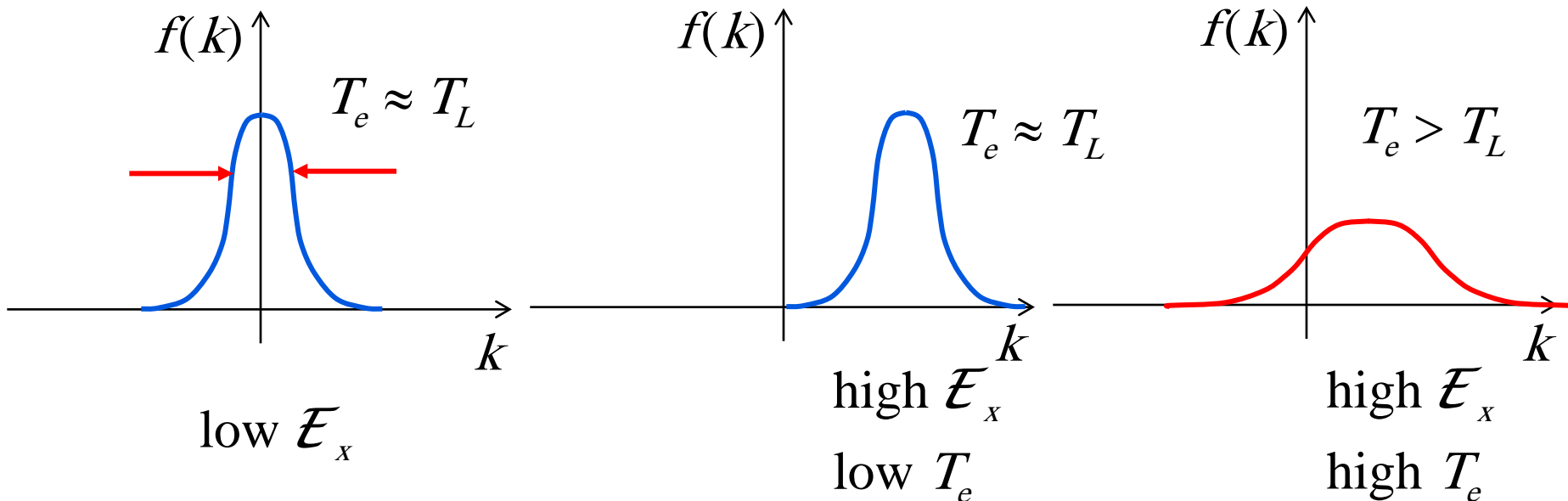
$$v_x^2 = v_{dx}^2 + 2cv_{dx} + c^2$$

$$W = n_L \left(\frac{1}{2} m^* v_{dx}^2 + \frac{1}{2} m^* \langle c^2 \rangle \right)$$

$$W = W_{drift} + W_{thermal}$$

electron temperature

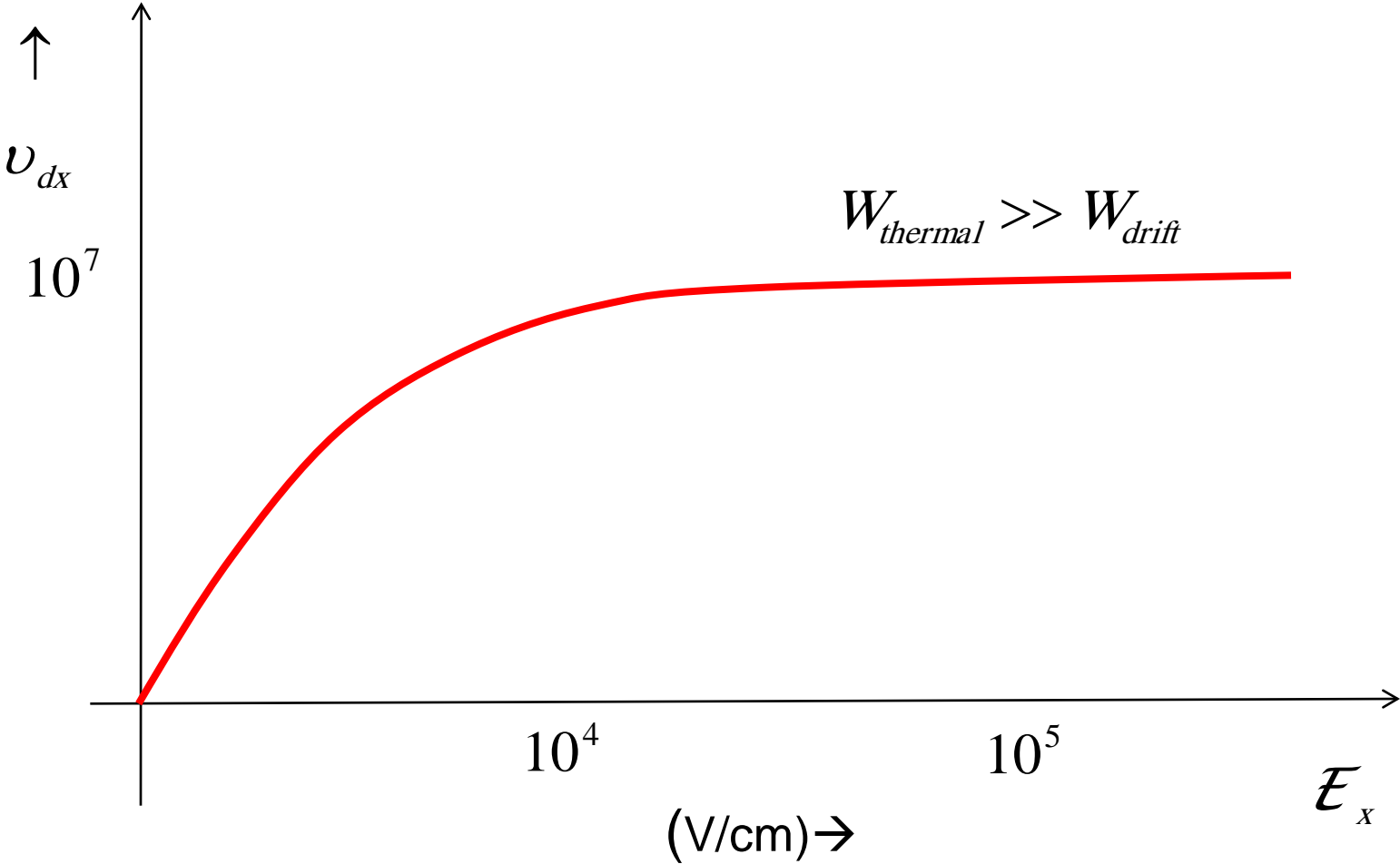
$$W = n_L \left(\frac{1}{2} m^* v_{dx}^2 + \frac{1}{2} m^* \langle c^2 \rangle \right) \quad \frac{1}{2} k_B T_e \equiv \frac{1}{2} m^* \langle c^2 \rangle$$



$$f(k) \sim e^{-\hbar^2 k^2 / 2 m^* k_B T_L}$$

$$f(k) \sim e^{-\hbar^2 k^2 / 2 m^* k_B T_e}$$

example



heat flux

$$F_W = \frac{1}{L} \sum_{p_x} E(p_x) v_x f(x, p_x, t)$$

$$F_W = \frac{1}{L} \sum_{p_x} \left(\frac{1}{2} m^* v_x^2 \right) v_x f(x, p_x, t)$$

$$F_W = \frac{1}{2} m^* n_L \langle v_x^2 v_x \rangle$$

$$F_W = \frac{1}{2} m^* n_L \langle v_x^2 (v_{dx} + c) \rangle$$

$$F_W = \frac{1}{2} m^* n_L \langle v_x^2 \rangle v_{dx} + \frac{1}{2} m^* n_L \langle v_x^2 c \rangle$$

$$v_x = v_{dx} + c$$

heat flux

$$F_W = \frac{1}{2} m^* n_L \langle v_x^2 \rangle v_{dx} + \frac{1}{2} m^* n_L \langle v_x^2 c \rangle \quad v_x = v_{dx} + c$$

$$F_W = W v_{dx} + \frac{1}{2} m^* n_L \langle (v_{dx}^2 + 2v_{dx}c + c^2)c \rangle$$

$$F_W = W v_{dx} + \frac{1}{2} m^* n_L \left[v_{dx}^2 \langle c \rangle + 2 \langle c^2 \rangle v_{dx} + \langle c^2 c \rangle \right]$$

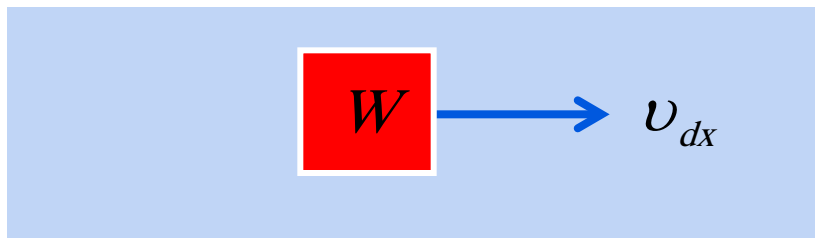
$$F_W = W v_{dx} + n_L m^* \langle c^2 \rangle v_{dx} + n_L \frac{1}{2} m^* \langle c^2 c \rangle$$

heat flux

$$F_W = Wv_{dx} + n_L m^* \langle c^2 \rangle v_{dx} + n_L \frac{1}{2} m^* \langle c^2 c \rangle \quad v_x = v_{dx} + c$$

$$k_B T_e \equiv m^* \langle c^2 \rangle \quad Q_x \equiv n_L \frac{1}{2} m^* \langle c^2 c \rangle \quad \text{“heat flux”}$$

$$F_W = Wv_{dx} + n_L k_B T_e v_{dx} + Q_x$$



$$P\Omega = Nk_B T$$

energy balance equations

$$\frac{\partial n_L(x, t)}{\partial t} = - \frac{d[I_{nx}/(-q)]}{dx}$$

electron continuity equation

$$I_x = n_L q \mu_n \mathcal{E}_x + 2 \mu_n \frac{dW}{dx}$$

drift-diffusion equation

$$\frac{\partial W(x, t)}{\partial t} = - \frac{dF_W}{dx} + I_x \mathcal{E}_x - \frac{(W - W_0)}{\tau_E}$$

energy-balance equation

$$F_W = -3 \mu_E W \mathcal{E}_x - 3 \mu_E \frac{d(W k_B T / q)}{dx}$$

energy-flux equation

energy balance equations

$$\frac{\partial n_L(x, t)}{\partial t} = - \frac{d[I_{nx}/(-q)]}{dx}$$

electron continuity equation

$$I_x = n_L q \mu_n \mathcal{E}_x + 2 \mu_n \frac{dW}{dx}$$

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energy-balance equation

$$F_W = W v_{dx} + n_L k_B T_e v_{dx} + Q_x$$

energy-flux equation

heat flux

$$F_W = Wv_{dx} + n_L k_B T_e v_{dx} + Q_x$$

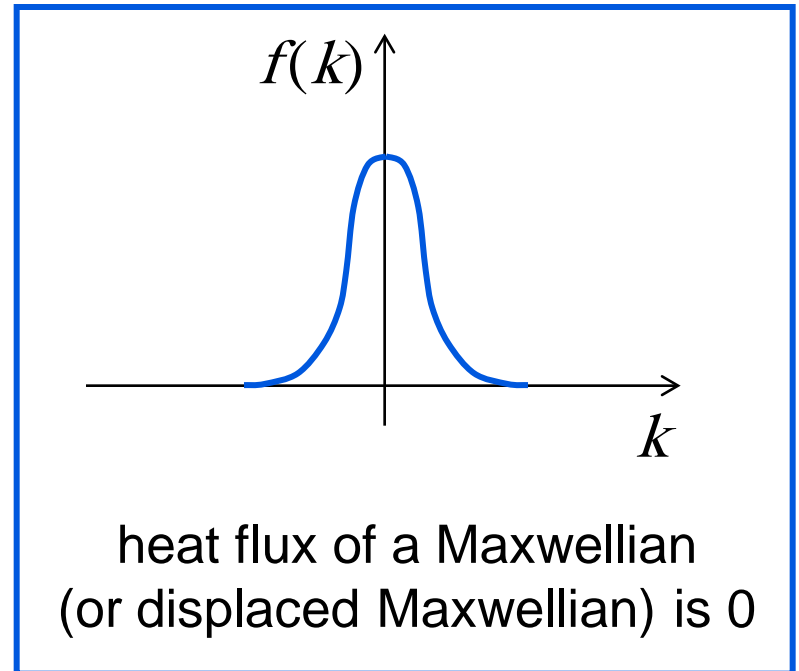
$$Q_x \equiv n_L \frac{1}{2} m^* \langle c^2 c \rangle$$

$$W = n_L \left(\frac{1}{2} m^* v_{dx}^2 + \frac{1}{2} k_B T_e \right)$$

$$Q_x \approx -\kappa_e \frac{dT}{dx}$$

but more generally

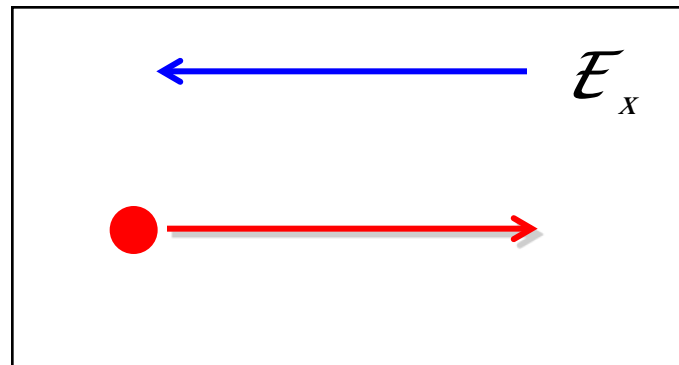
$$I_{Q_x} \approx \pi I_x - \kappa_e \frac{dT_e}{dx}$$



Q and I_q

does $Q_x = I_q$?

N-type semiconductor



—————→ x

$$I_q = \pi I > 0$$

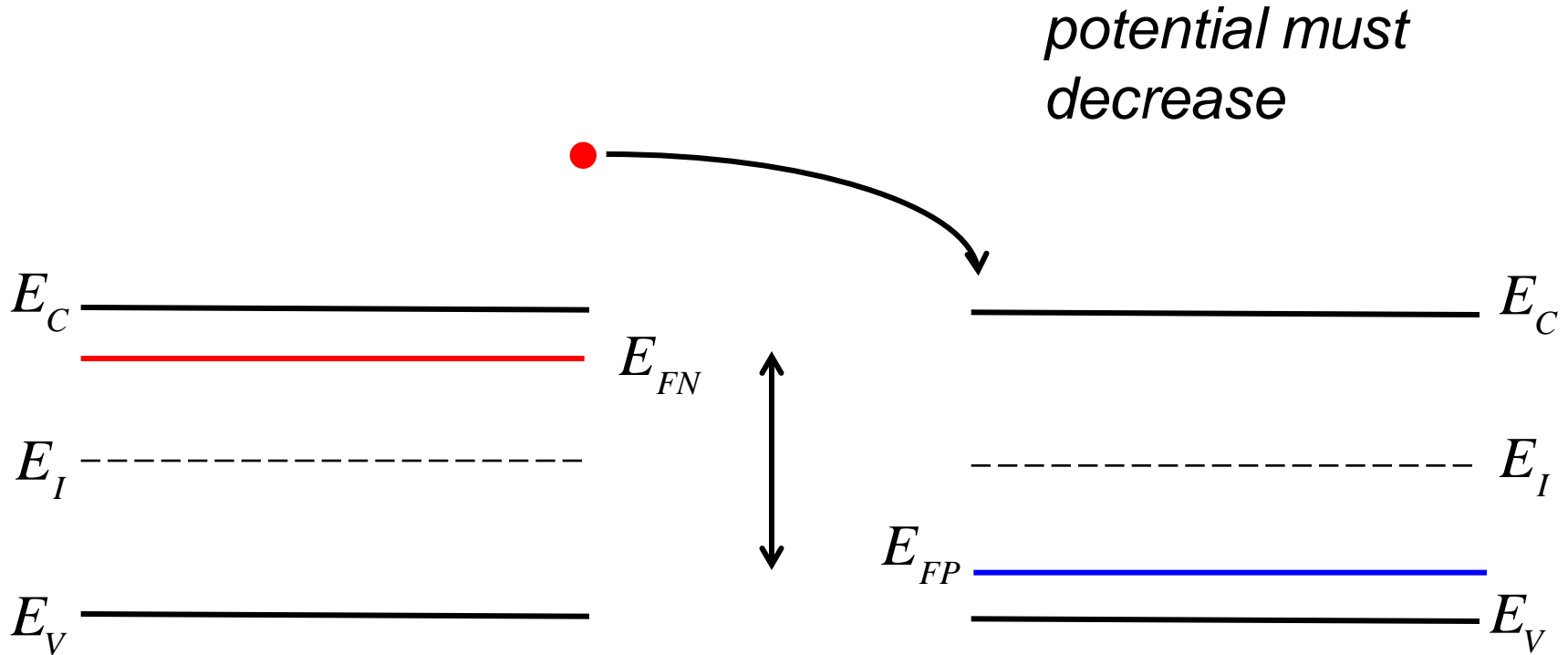
$$Q_x < 0$$

Mark A. Stettler, Muhammad A. Alam, and Mark S. Lundstrom, "A Critical Examination of the Assumptions Underlying Macroscopic Transport Equations for Silicon Devices," *IEEE Trans. on Electron Devices*, **40**, 733, 1993.

outline

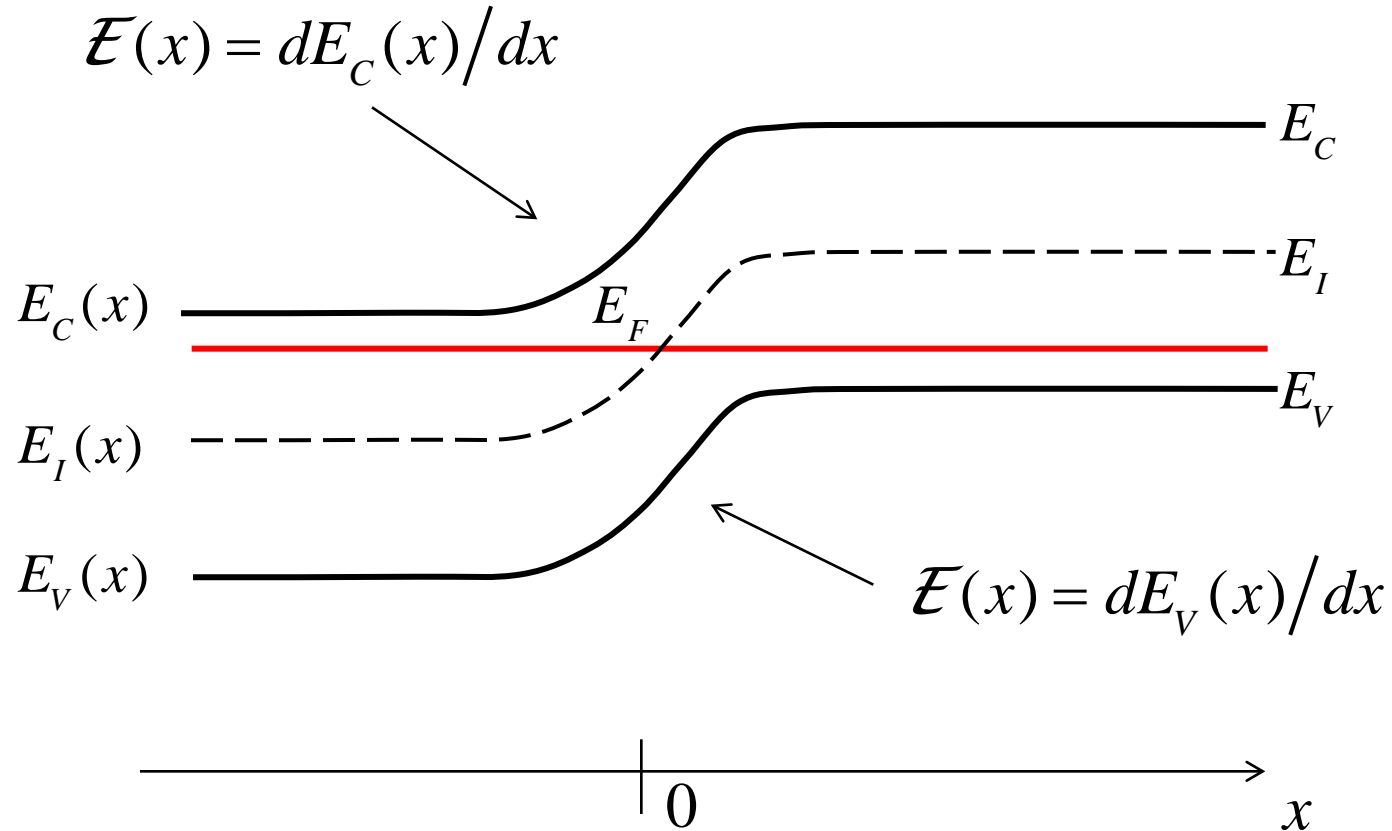
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- 4) Summary

review: pn homojunctions

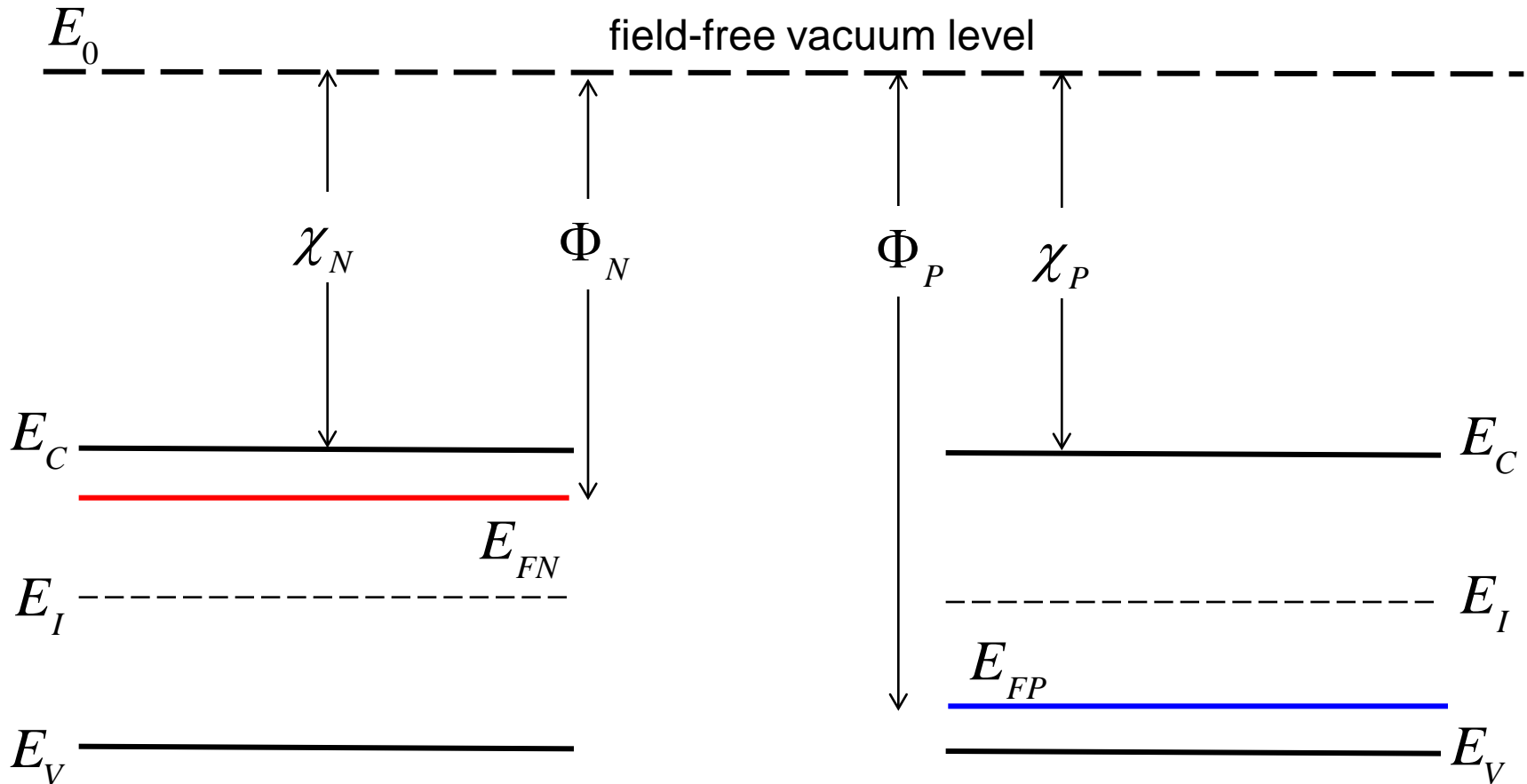


$$qV_{BI} = (E_{FN} - E_{FP})/q$$

review: pn homojunctions



reference for the energy bands

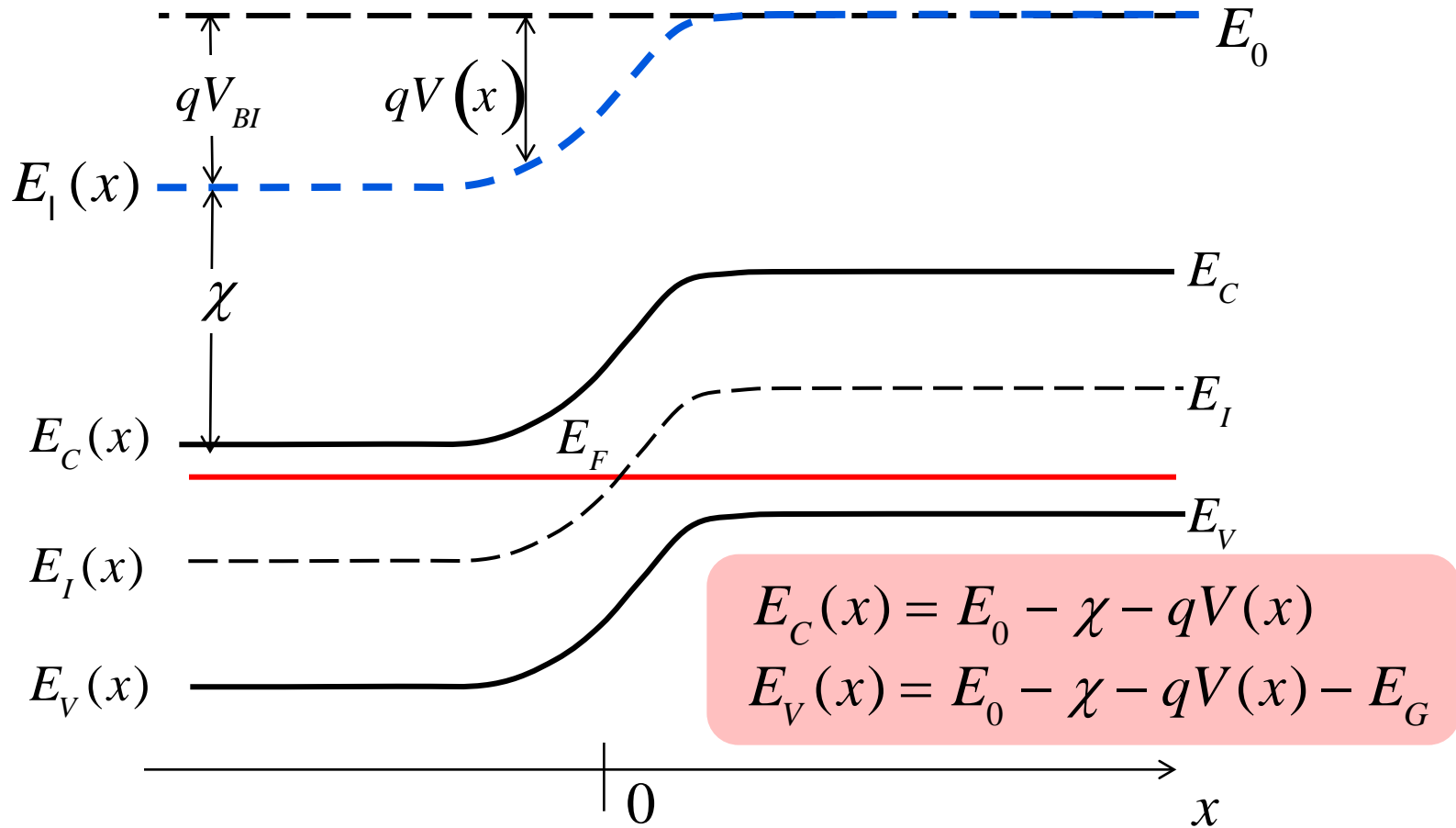


$$E_C = E_0 - \chi$$

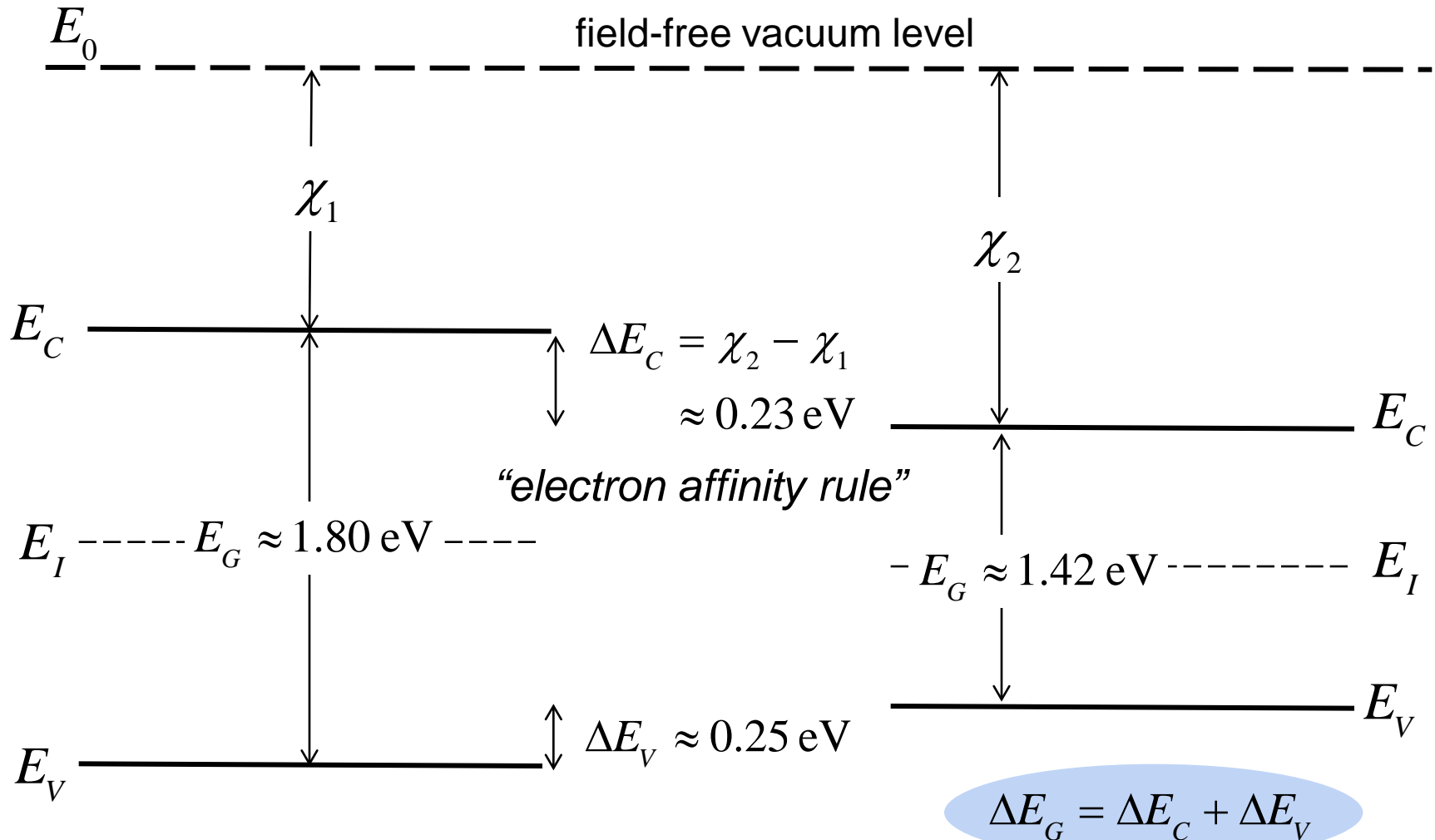
$$E_V = E_0 - \chi - E_G$$

$$qV_{BI} = (\Phi_P - \Phi_N)$$

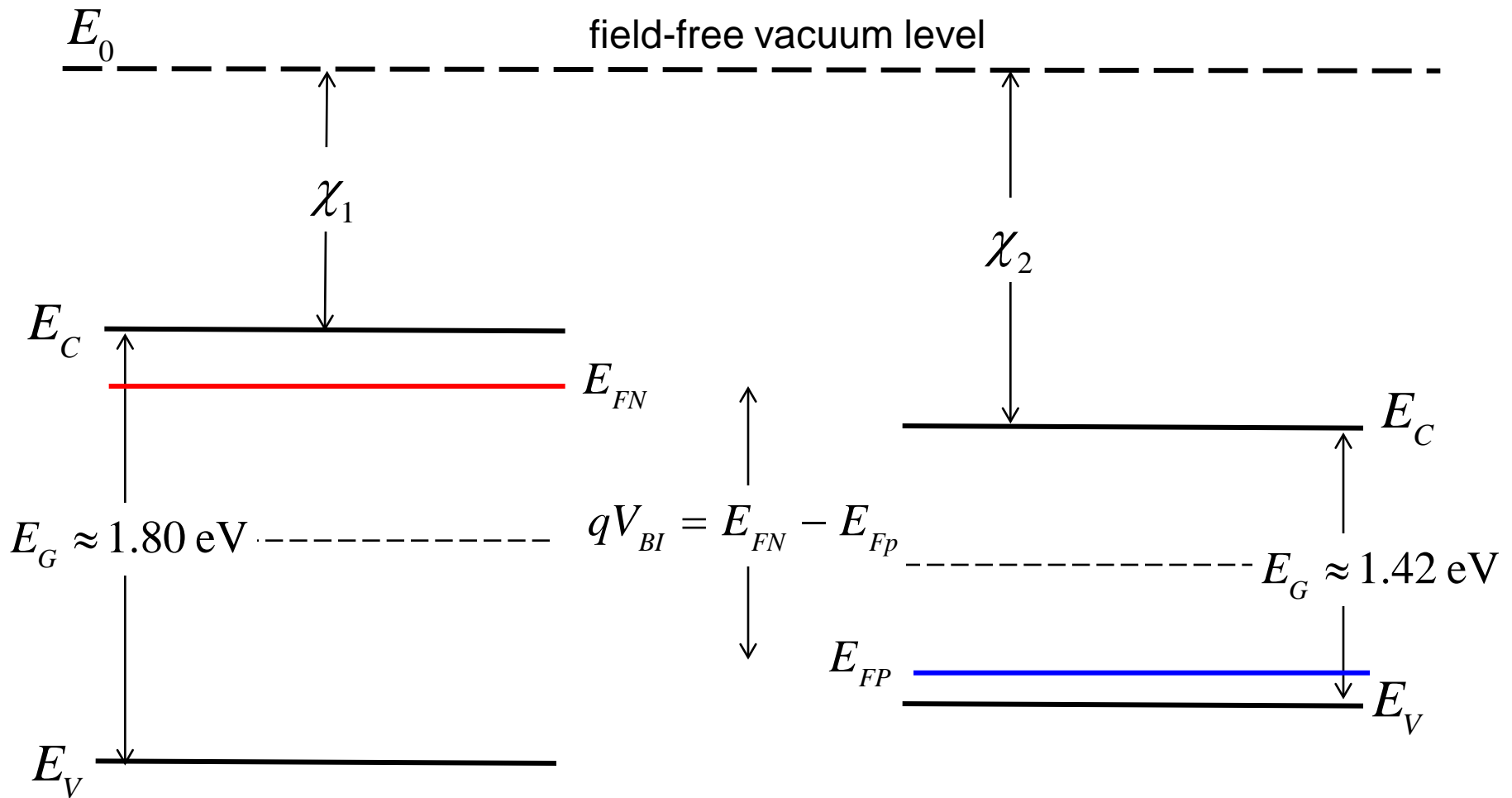
local vacuum level



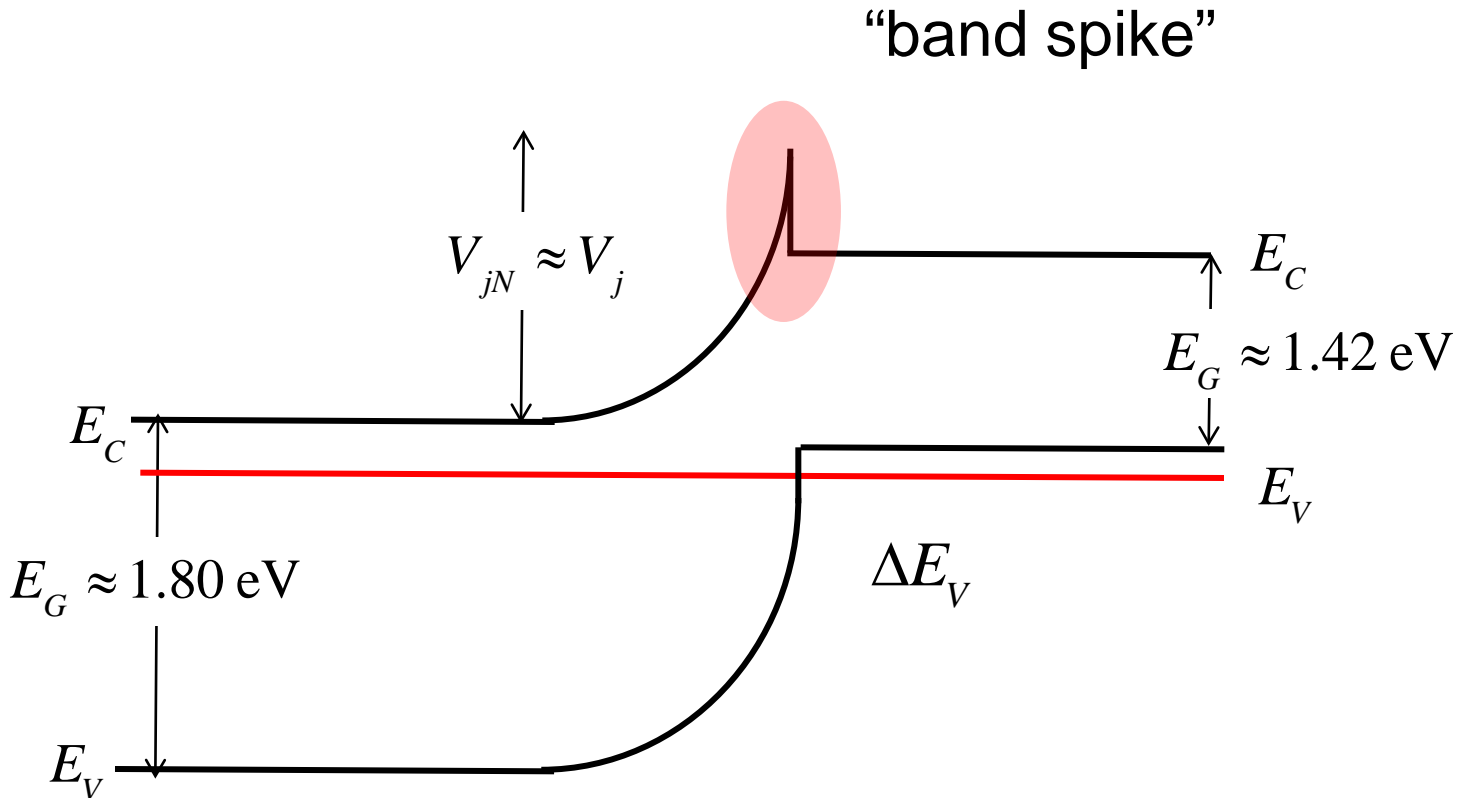
Al_{0.3}Ga_{0.7}As : GaAs (Type I HJ)



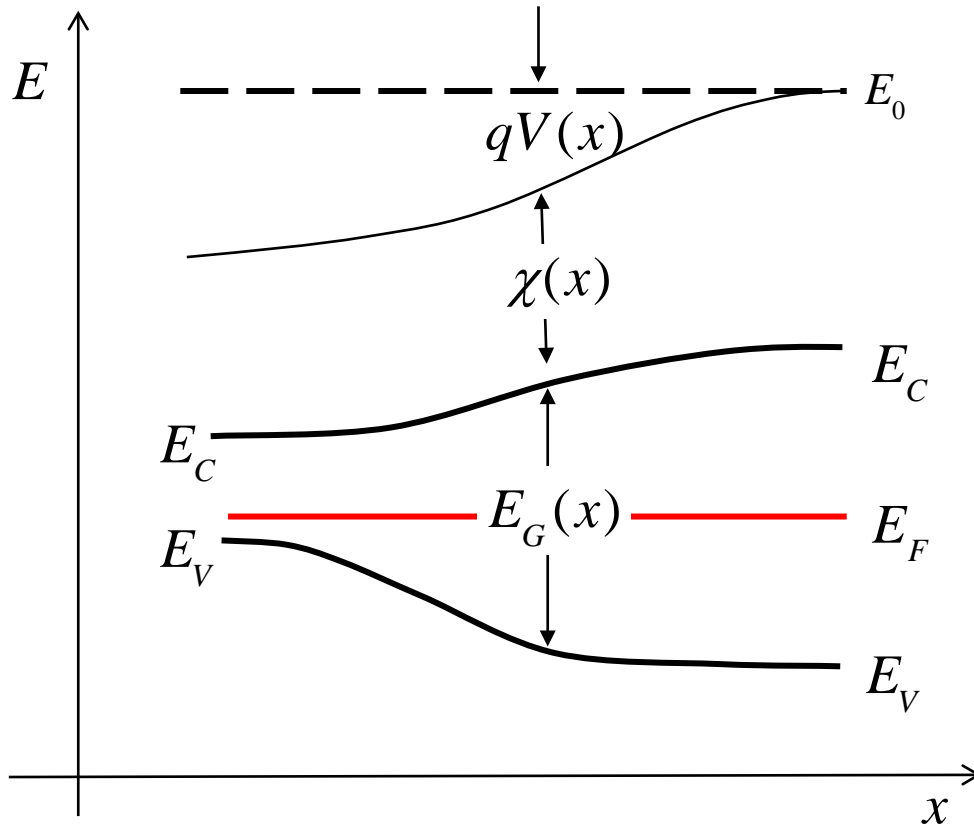
N-Al_{0.3}Ga_{0.7}As : p⁺-GaAs (Type I HJ)



N-Al_{0.3}Ga_{0.7}As : p⁺-GaAs (Type I HJ)



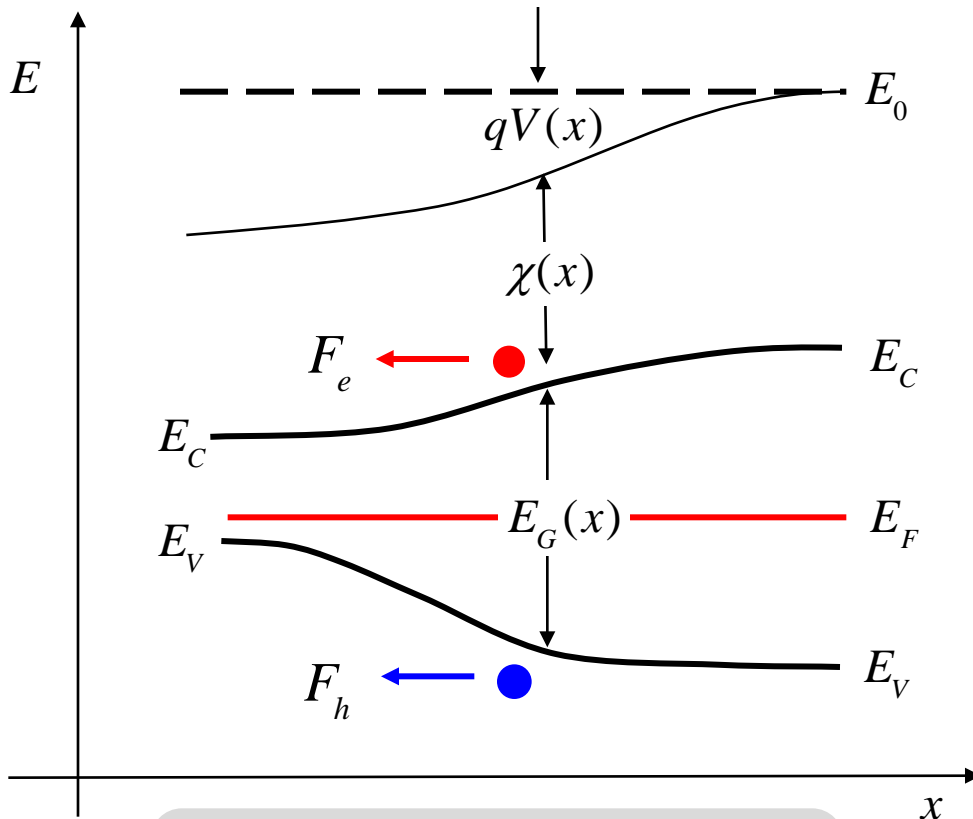
general, graded heterostructure



$$E_C(x) = E_0 - \chi(x) - qV(x)$$

$$E_V(x) = E_C(x) - E_G(x)$$

quasi-electric fields



$$F_e = -\frac{dE_C}{dx} = q \frac{dV}{dx} + \frac{d\chi}{dx}$$

$$F_e = -q\mathcal{E}(x) - q\mathcal{E}_{QN}(x)$$

$$\mathcal{E}_{QN} \equiv -\frac{1}{q} \frac{d\chi}{dx}$$

$$F_h = +\frac{dE_V}{dx} = -q \frac{dV}{dx} - \frac{d(\chi + E_G)}{dx}$$

$$F_h = +q\mathcal{E}(x) + q\mathcal{E}_{QP}(x)$$

$$\mathcal{E}_{QP} \equiv -\frac{1}{q} \frac{d(\chi + E_G)}{dx}$$

$$E_C(x) = E_0 - \chi(x) - qV(x)$$

$$E_V(x) = E_C(x) - E_G(x)$$

BTE

$$\frac{\partial f}{\partial t} + \vec{v} \cdot \nabla_r f + \frac{d\vec{p}}{dt} \cdot \nabla_p f = 0$$

$$\frac{d\vec{p}}{dt} = \frac{d(\hbar\vec{k})}{dt} = -\nabla_r E_C(\vec{r}) = -q\vec{\mathcal{E}}(\vec{r})$$

(constant effective mass)

This equation does not hold when the effective mass is position dependent. See HW 12.

alternative approach: hole current

$$J_p = p\mu_p \frac{dF_p}{dx} \quad p = N_V(x) e^{(E_V - F_p)/k_B T} \quad F_p = E_V(x) - k_B T \ln(p/N_V)$$

$$\frac{dF_p}{dx} = \frac{dE_V(x)}{dx} - k_B T \left[\frac{1}{p} \frac{dp}{dx} - \frac{1}{N_V} \frac{dN_V}{dx} \right]$$

$$J_p = p\mu_p \left[\frac{dE_V(x)}{dx} + \frac{k_B T}{N_V} \frac{dN_V}{dx} \right] - k_B T \mu_p \frac{dp}{dx}$$

$$\frac{dE_V(x)}{dx} = \frac{d}{dx} \left[E_0 - \chi(x) - E_G(x) - qV(x) \right] = q \left(\mathcal{E}(x) + \mathcal{E}_{QP} \right)$$

hole and electron currents

$$J_p = pq\mu_p \left[\mathcal{E} + \mathcal{E}_{QP} + \frac{k_B T}{q} \frac{1}{N_V} \frac{dN_V}{dx} \right] - qD_p \frac{dp}{dx}$$

“DOS effect”

$$J_p = nq\mu_n \left[\mathcal{E} + \mathcal{E}_{QN} - \frac{k_B T}{q} \frac{1}{N_C} \frac{dN_C}{dx} \right] + qD_n \frac{dn}{dx}$$

quasi-electric fields

$$\mathcal{E}_{QP} \equiv -\frac{1}{q} \frac{d(\chi + E_G)}{dx} \quad \mathcal{E}_{QN} \equiv -\frac{1}{q} \frac{d\chi}{dx}$$

questions

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- 2) Balance equations in 3D
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