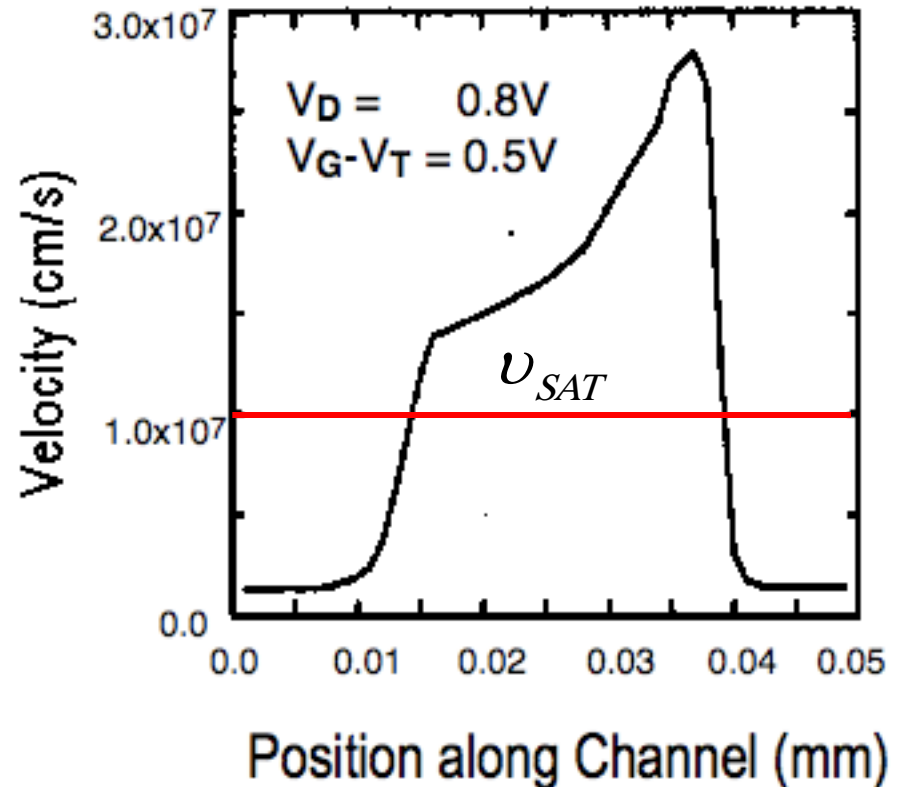
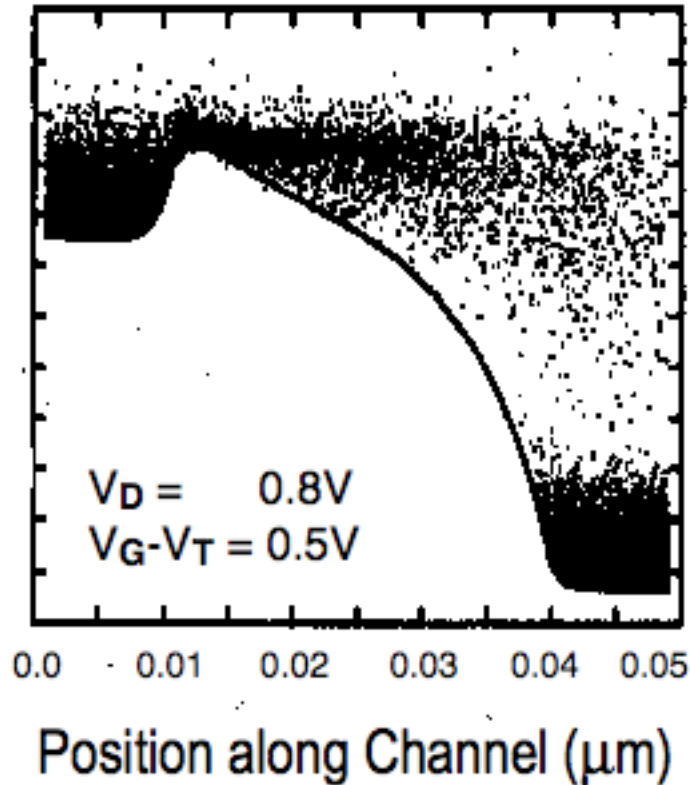


ECE-656: Fall 2009

**Lecture 31:
Monte Carlo Simulation**

Professor Mark Lundstrom
Electrical and Computer Engineering
Purdue University, West Lafayette, IN USA

Damocles

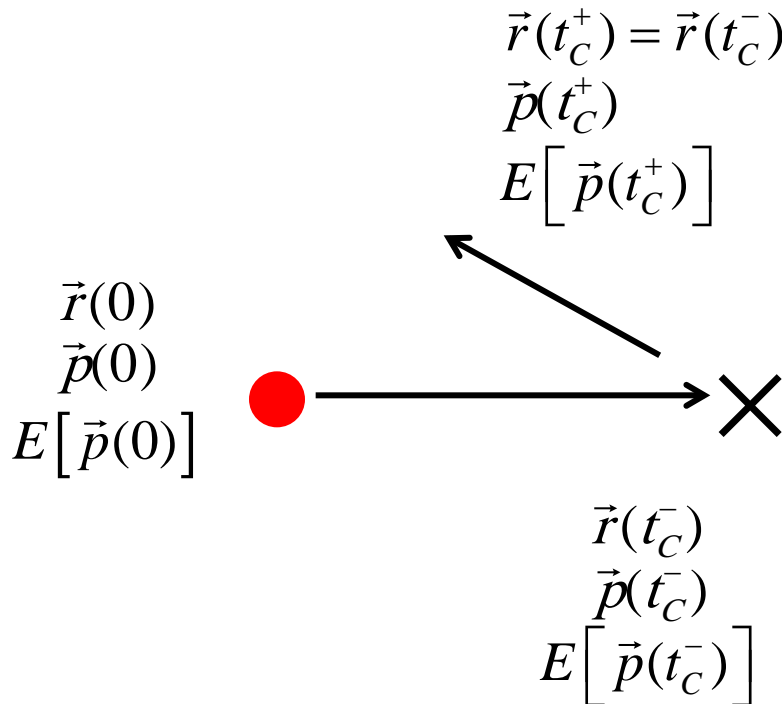


Frank, Laux, and Fischetti, IEDM Tech. Dig., p. 553, 1992

outline

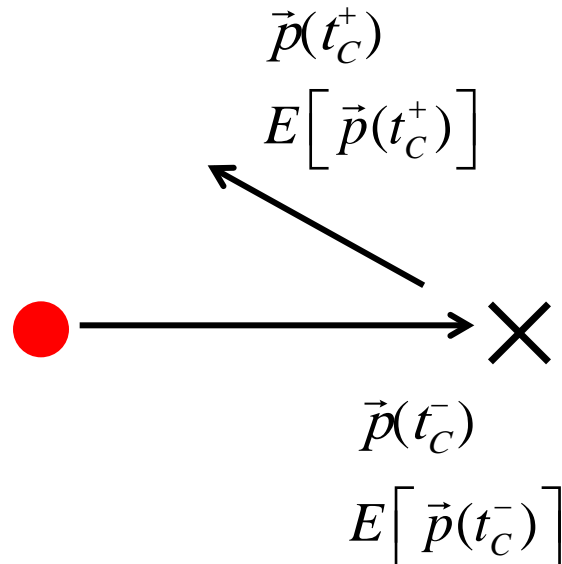
- 1) Introduction
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carrier scattering



- 1) Scattering events are assumed to be instantaneous - scattering changes the carrier's momentum (and energy) but not position.
- 2) Scattering events are treated quantum mechanically. Carrier motion between scattering events is treated semi-classically

carrier scattering



Elastic scattering:

$$E[\vec{p}(t_c^+)] = E[\vec{p}(t_c^-)]$$

Inelastic scattering:

$$E[\vec{p}(t_c^+)] \neq E[\vec{p}(t_c^-)]$$

Isotropic scattering:

-no preferred direction

Scattering rate:

$$\frac{1}{\tau(E_i)} : D(E_f)$$

total scattering rate

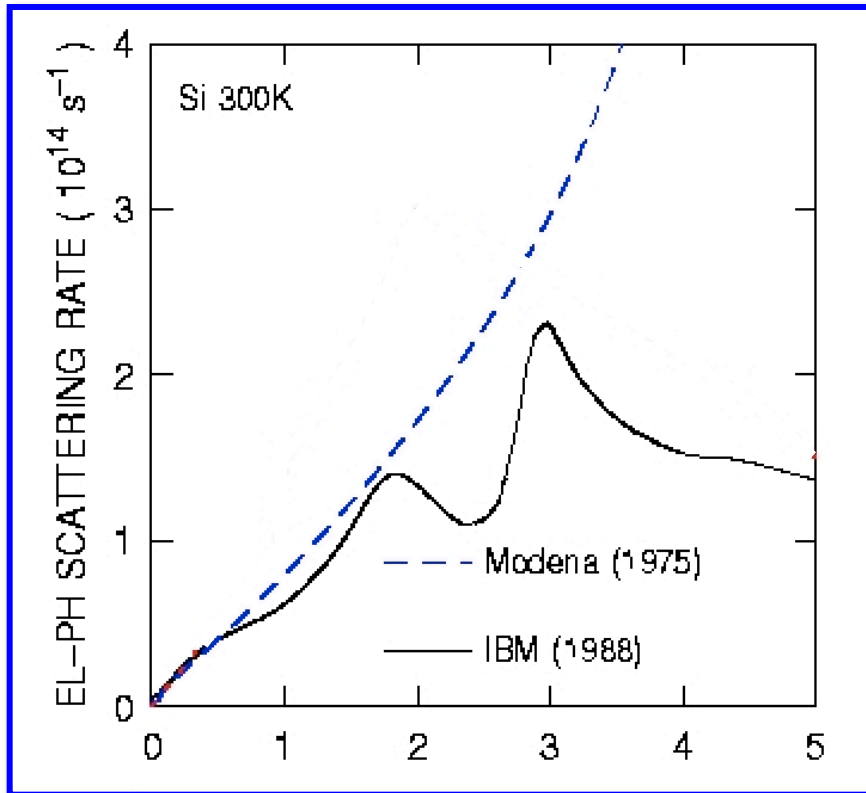
$$\frac{1}{\tau(E_i)} = \sum_{k=1}^N \frac{1}{\tau_k(E_i)} = \Gamma(E_i)$$

scattering mechanisms:

- ionized impurity
- acoustic phonon emission
- acoustic phonon absorption
- optical phonon ems
- optical phonon abs
- intervalley
- electron-electron
- electron-hole
- “polar optical phonon”
- surface roughness
- etc.

total scattering rate

typical scattering rates vs. energy



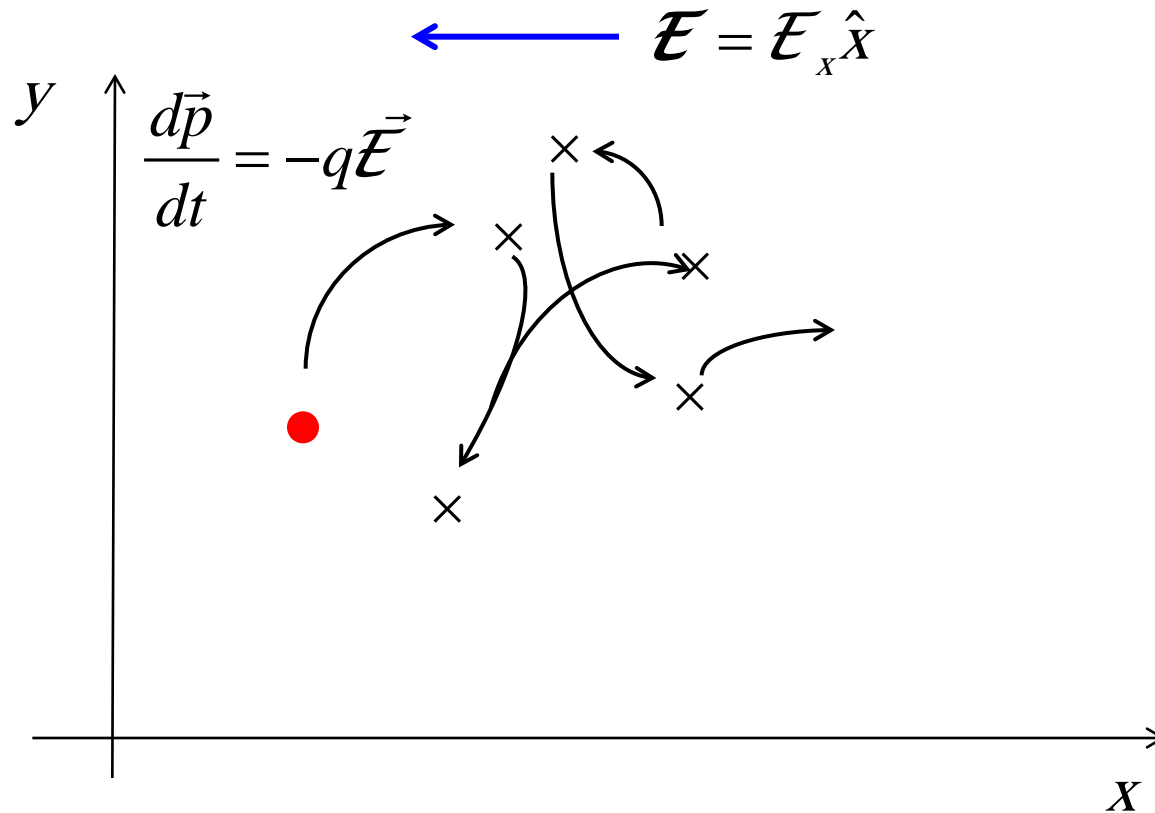
http://www.research.ibm.com/DAMOCLES/html_files/sirates.html#history

Lundstrom ECE-656 F09

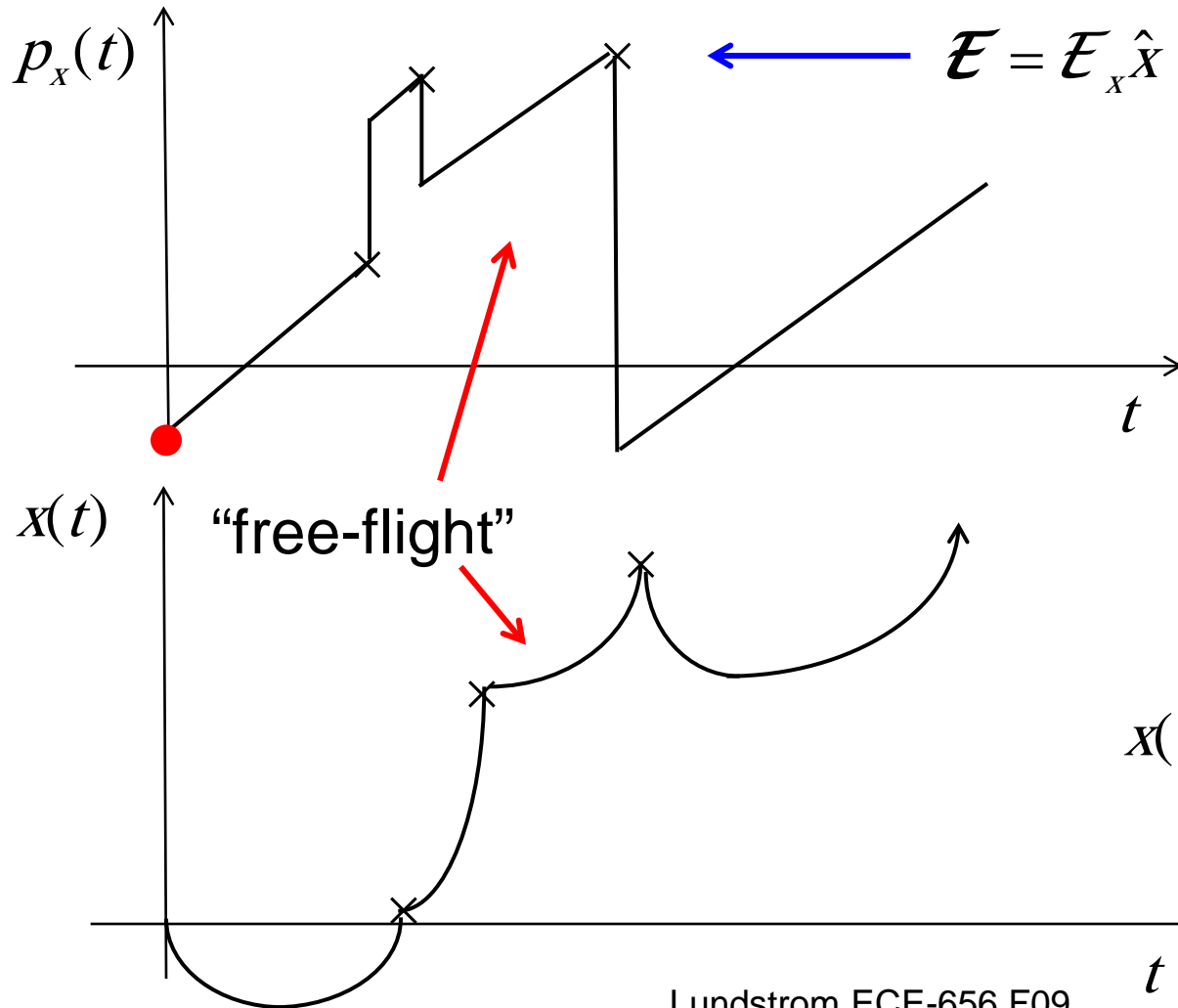
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carrier trajectories in 2D



carrier trajectories in phase space



$$\frac{dp_x}{dt} = -q\mathcal{E}_x$$

$$\frac{dp_y}{dt} = 0$$

$$x(t) = x(0) + \int_0^{t_c} v_x(t) dt$$

MC algorithm

- 1) “free flight” for t_C seconds.
- 2) update $E(t_C^-)$ and $r(t_C^-)$
- 3) identify collision
- 4) update $E(t_C^+)$ and $p(t_C^+)$
- 5) Set $t = 0$ and repeat

r_1

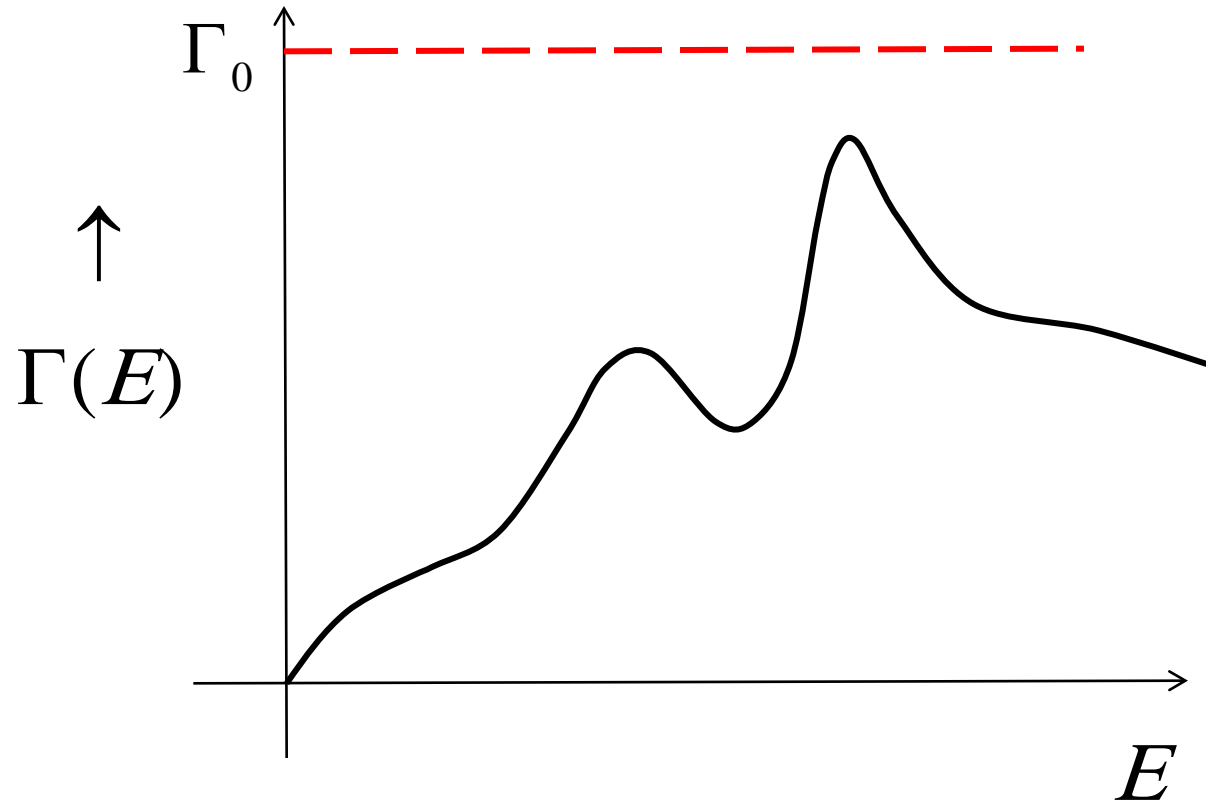
r_2

r_3, r_4

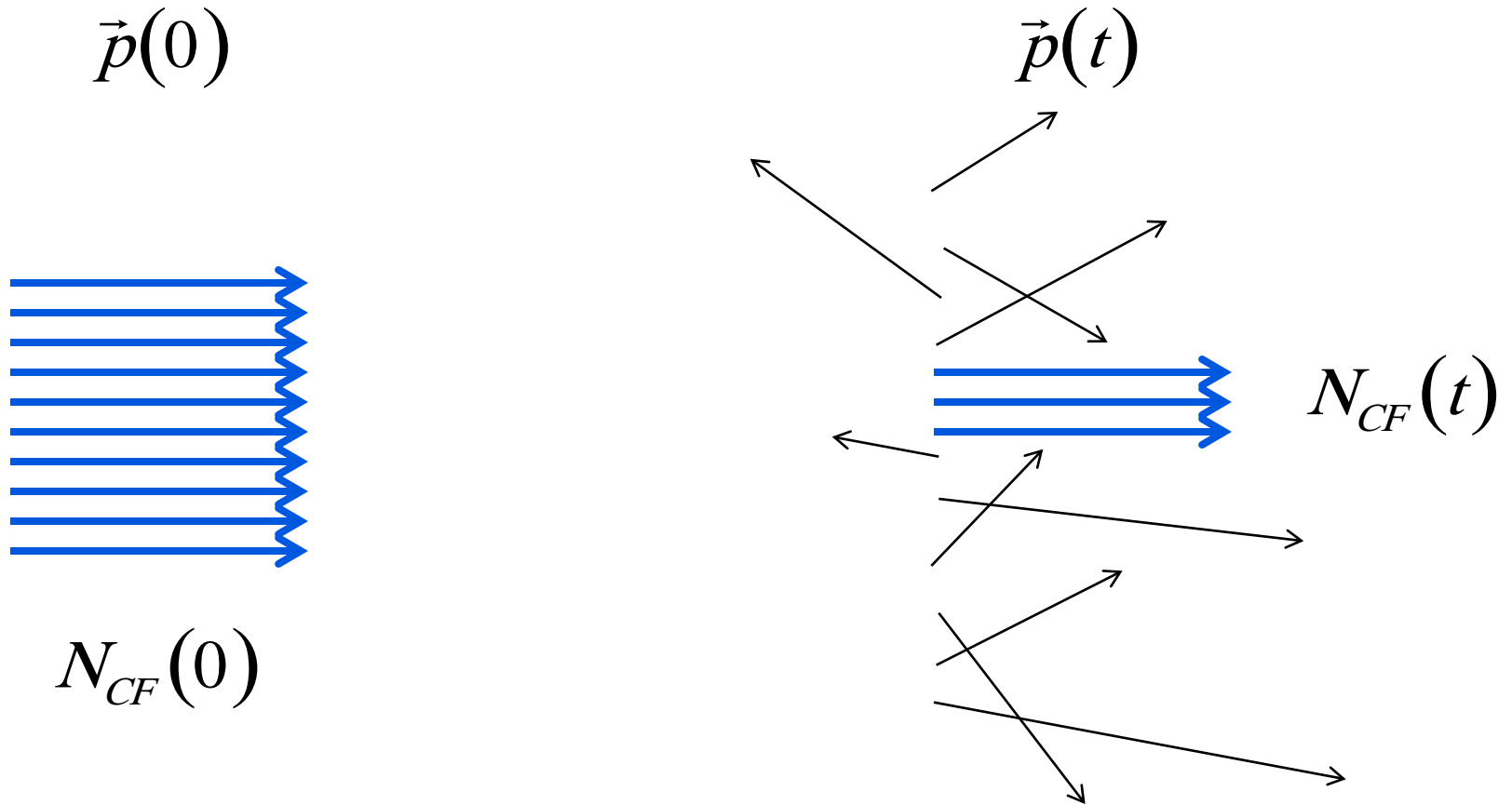
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free flight



free flight



free flight time

How many carriers survive until time $t + dt$ without scattering?

$$N_{CF}(t + dt) = N_{CF}(t) - N_{CF}(t)\Gamma_0 dt$$

$$\frac{dN_{CF}(t)}{dt} = -N_{CF}(t)\Gamma_0$$

$$N_{CF}(t) = N_{CF}(0)e^{-\Gamma_0 t}$$

$$\frac{N_{CF}(t)}{N_{CF}(0)} = e^{-\Gamma_0 t} \quad \langle \tau \rangle = 1/\Gamma_0$$

free flight time

What is the probability that a carrier survives until time, t , and then scatters between t and $t + dt$?

$$\mathcal{P}(t)dt = \left(\frac{N_{CF}(t)}{N_{CF}(0)} \right) \Gamma_0 dt = e^{-\Gamma_0 t} \Gamma_0 dt$$

free-flight time
probability density
function

If we have a random number generator that produces random numbers with a probability distribution, $\mathcal{P}(r)$, how do we use it to choose free flight times?

$$\mathcal{P}(t)dt = e^{-\Gamma_0 t} \Gamma_0 dt = \mathcal{P}(r)dr$$

free flight times

Assume a random number generator that produces random numbers uniformly distributed between 0 and 1.

$$\mathcal{P}(r)dr = e^{-\Gamma_0 t} \Gamma_0 dt$$

$$dr = e^{-\Gamma_0 t} \Gamma_0 dt$$

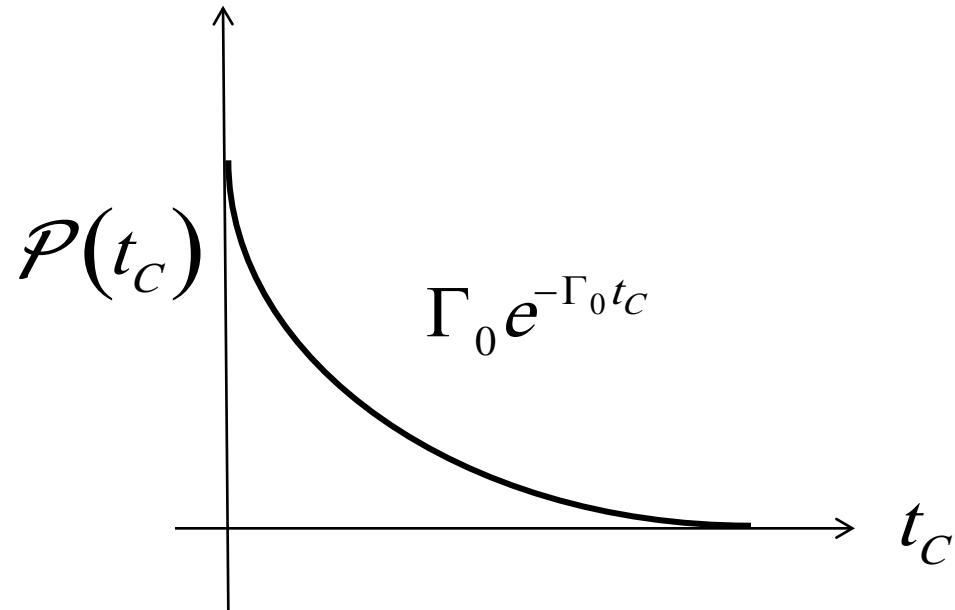
$$\int_0^{r_c} dr = r_c = \int_0^{t_c} e^{-\Gamma_0 t} \Gamma_0 dt = -e^{-\Gamma_0 t} \Big|_0^{t_c} = (1 - e^{-\Gamma_0 t_c})$$

$$t_c = \frac{1}{\Gamma_0} \ln(1 - r_c)$$

free flight distribution

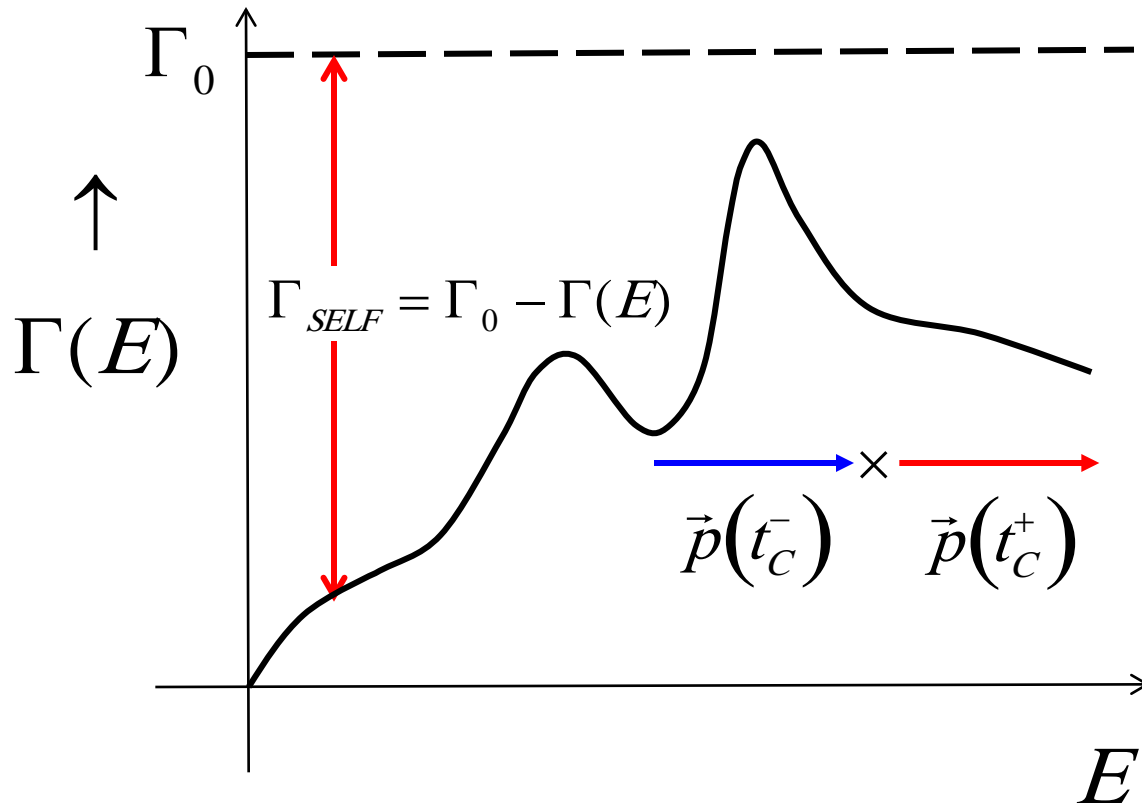
$$t_C = \frac{1}{\Gamma_0} \ln(r_1)$$

$$\langle t_C \rangle = \frac{1}{\Gamma_0}$$



but, the scattering rate is not constant with energy!

free flight



$$\Gamma(E) = \sum_{k=1}^N \frac{1}{\tau_k(E)}$$

$$\Gamma_0 = \sum_{k=1}^{N+1} \frac{1}{\tau_k(E)}$$

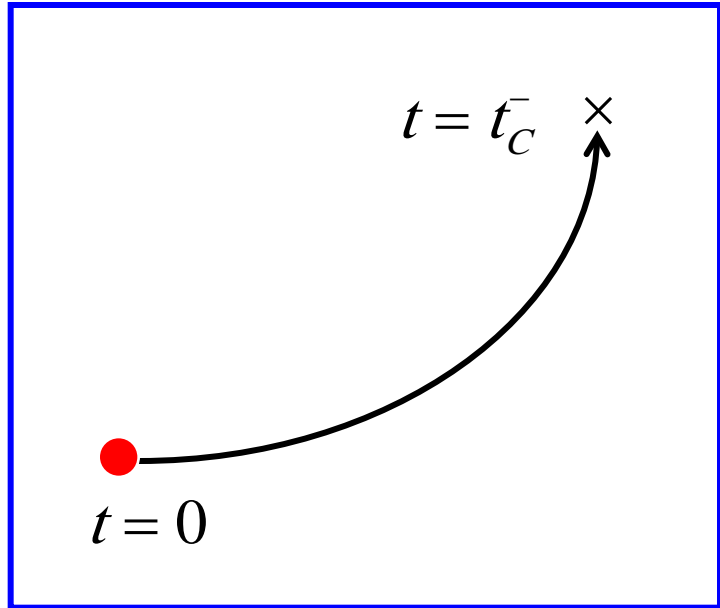
$$t_c = \frac{1}{\Gamma_0} \ln(r_1)$$

update at end of free-flight

$$\vec{p}(t_c^-) = \vec{p}(0) - q\vec{E}t_c$$

$$E(t_c^-) = E[\vec{p}(t_c^-)]$$

$$\vec{r}(t_c^-) = r(0) + \vec{v}(0)t_c$$



- 1) what collision terminated the free flight?
- 2) how does the collision affect the carrier's momentum and energy?

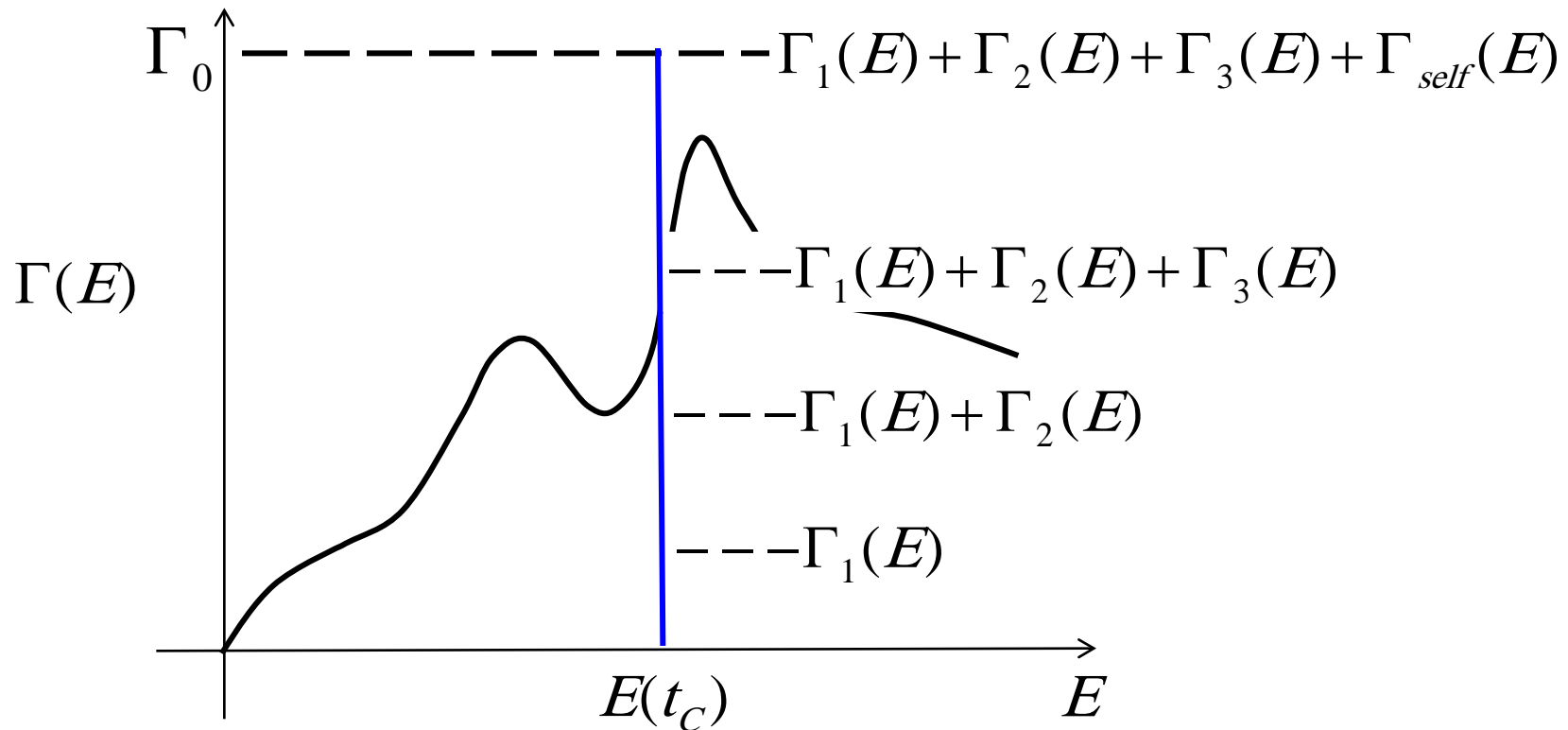
outline

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collision

As an example, assume that there are 3 scattering mechanisms, plus self-scattering.

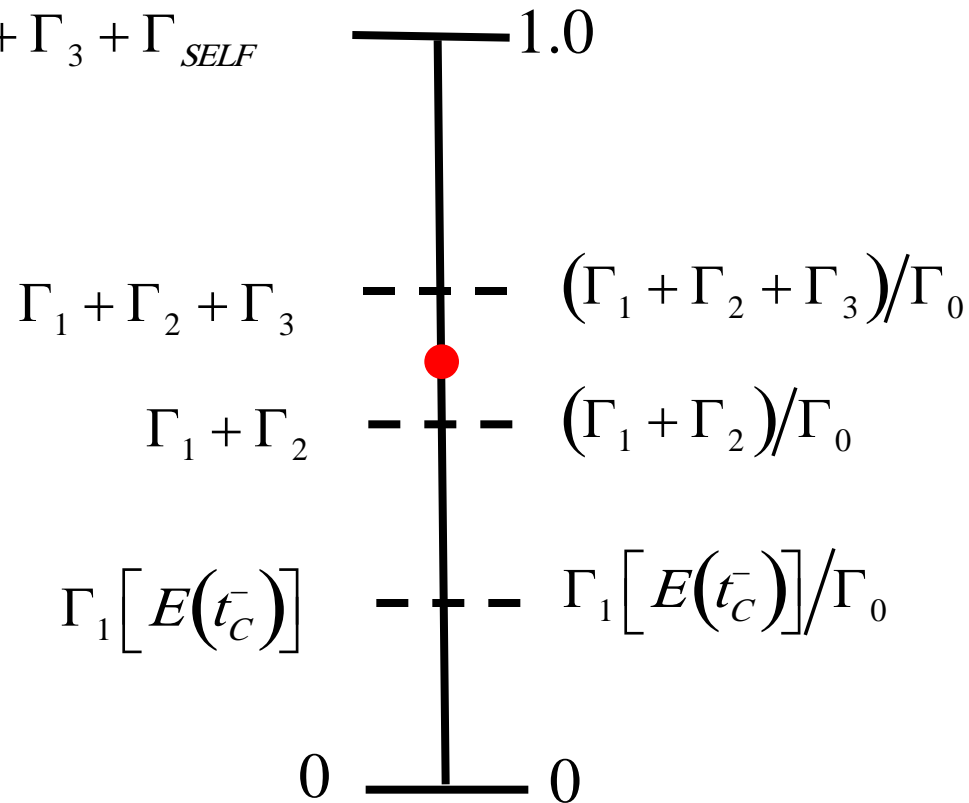
$$\Gamma(E) = \Gamma_1(E) + \Gamma_2(E) + \Gamma_3(E) + \Gamma_{SELF}(E) = \Gamma_0$$



collision

$$\Gamma(E) = \Gamma_1(E) + \Gamma_2(E) + \Gamma_3(E) + \Gamma_{SELF}(E) = \Gamma_0$$

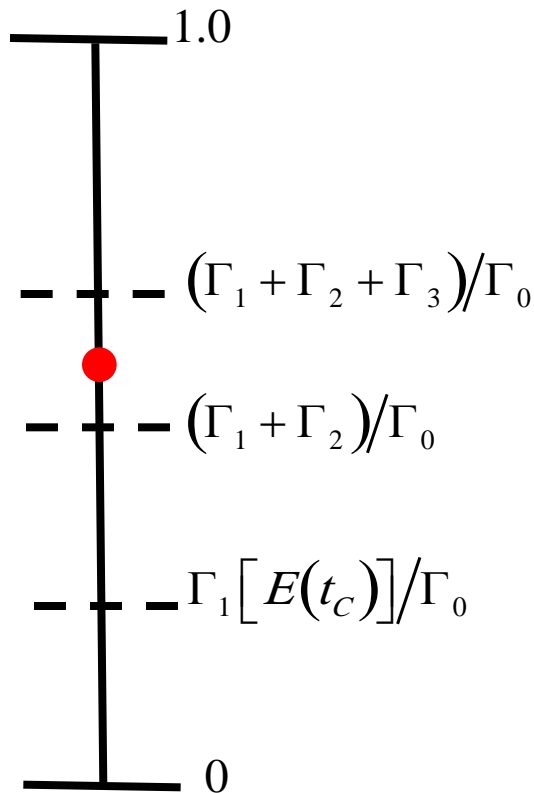
$$\Gamma_0 = \Gamma_1 + \Gamma_2 + \Gamma_3 + \Gamma_{SELF}$$



choose random number between 0 and 1

collision

$$\Gamma(E) = \Gamma_1(E) + \Gamma_2(E) + \Gamma_3(E) + \Gamma_{SELF}(E) = \Gamma_0$$



$$0 < r_2 < \Gamma_1/\Gamma_0$$

process 1

$$\Gamma_1/\Gamma_0 < r_2 < (\Gamma_1 + \Gamma_2)/\Gamma_0$$

process 2

$$(\Gamma_1 + \Gamma_2)/\Gamma_0 < r_2 < (\Gamma_1 + \Gamma_2 + \Gamma_3)/\Gamma_0$$

process 3

$$(\Gamma_1 + \Gamma_2 + \Gamma_3)/\Gamma_0 < r_2 < 1.0$$

self-scattering

choose random number uniformly distributed between 0 and 1

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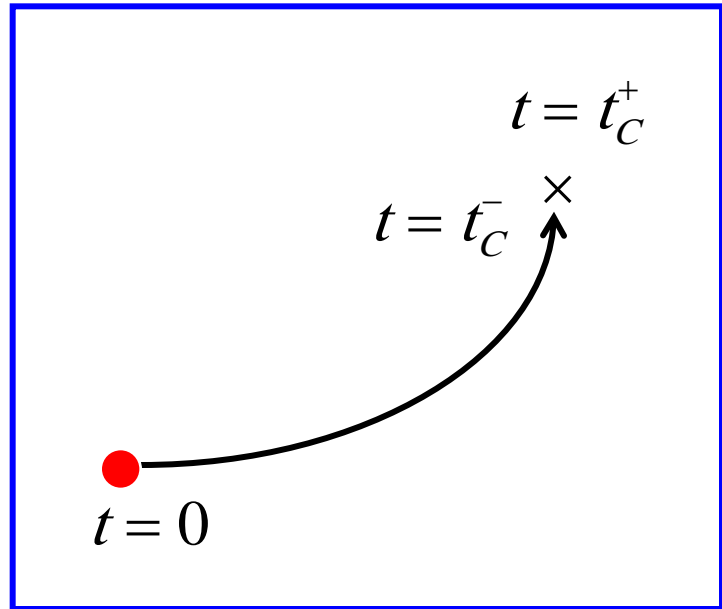
update after scattering

$$\vec{r}(t_C^+) = \vec{r}(t_C^-)$$

$$E(t_C^+) = E(t_C^-) + \Delta E_i$$

$$\Delta E_i = 0, \pm \hbar \omega_0, \text{ etc}$$

$$\vec{p}(t_C^+) = ?$$



We know the magnitude of p from the known energy after the collision, but how do we determine the direction?

update after scattering: 1D

$$\vec{p}(t_c) \longrightarrow$$

i) forward scattering by phonon absorption

$$\vec{p}(t_c^+) \longrightarrow$$

ii) backward scattering by phonon absorption

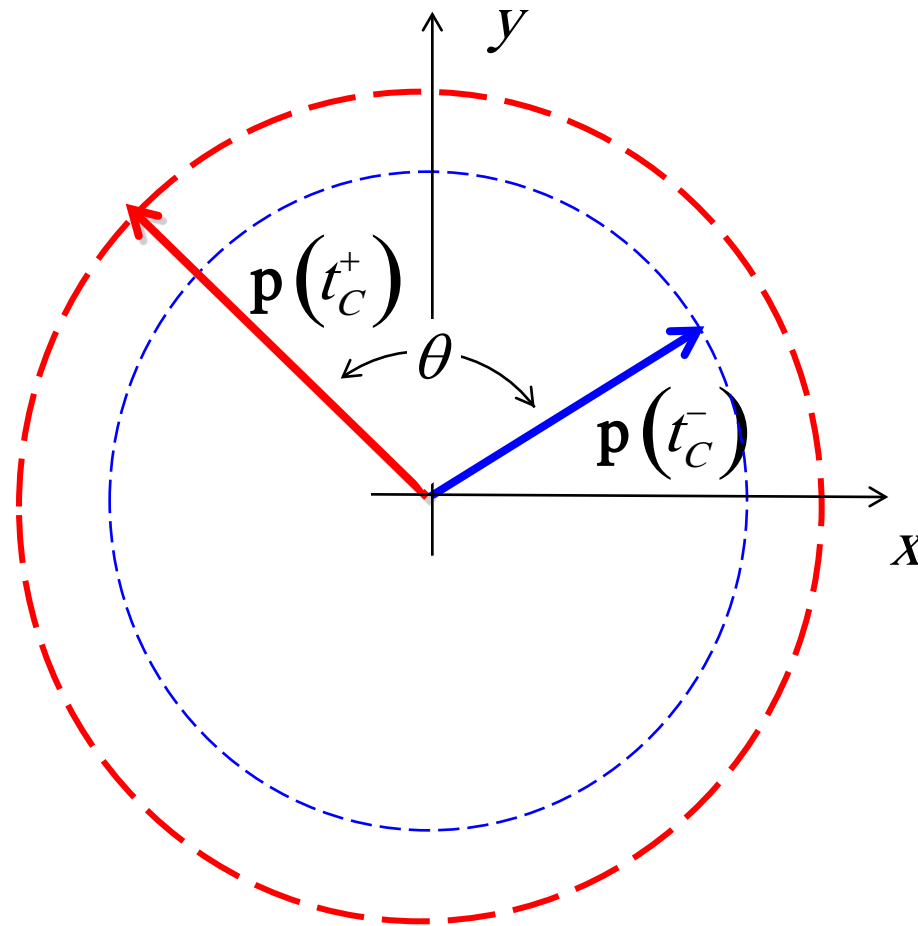
$$\vec{p}(t_c^+) \longleftarrow$$

If scattering is isotropic, then forward and backward scattering are equally probable.

$$0 < r_3 \leq 0.5: \quad \vec{p}(t_c^+) = \vec{p}(t_c^-)$$

$$0.5 < r_2 \leq 1.0: \quad \vec{p}(t_c^+) = -\vec{p}(t_c^-)$$

update after scattering: 2D

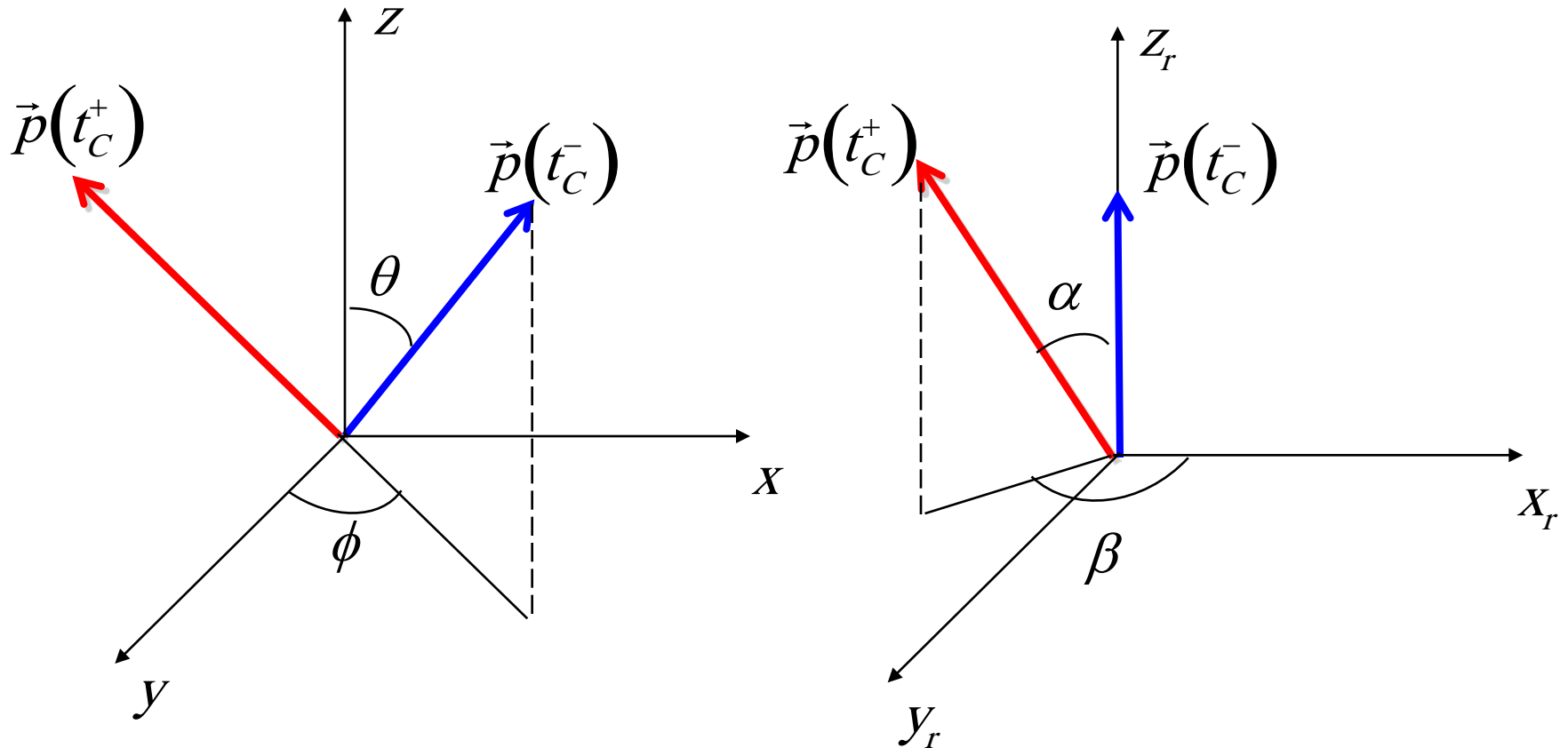


assume isotropic scattering
(e.g. optical phonon abs)

$$\theta = 2\pi r_3$$

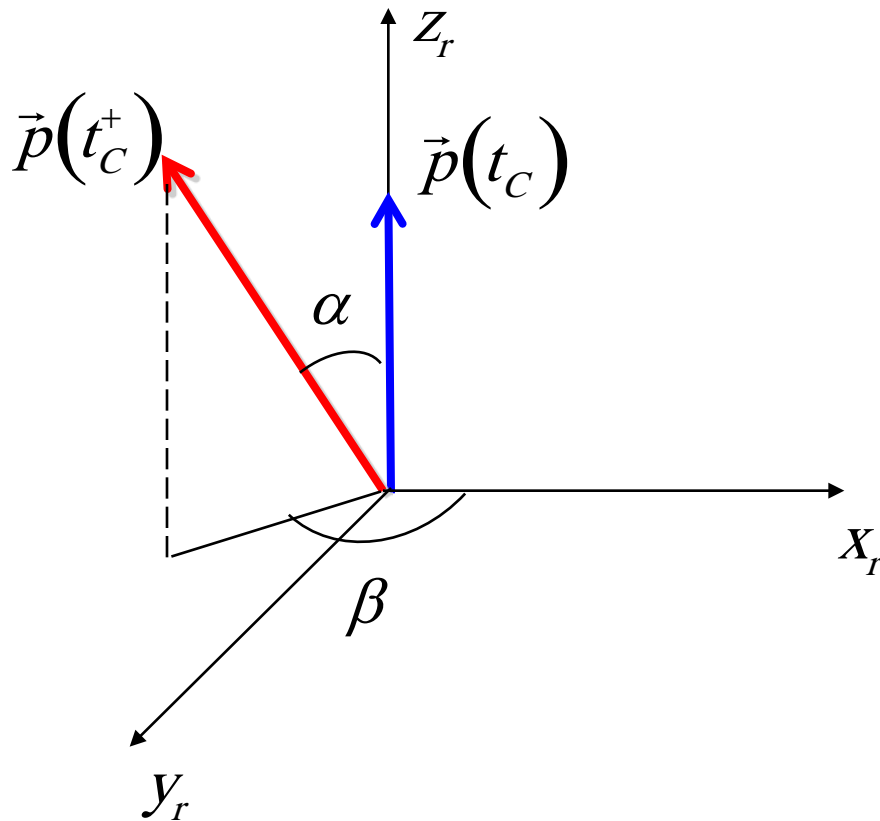
r_3 is a random number uniformly
distributed between 0 and 1

update after scattering: 3D



need two random numbers to select $(\alpha, \beta) \rightarrow (\theta, \phi)$

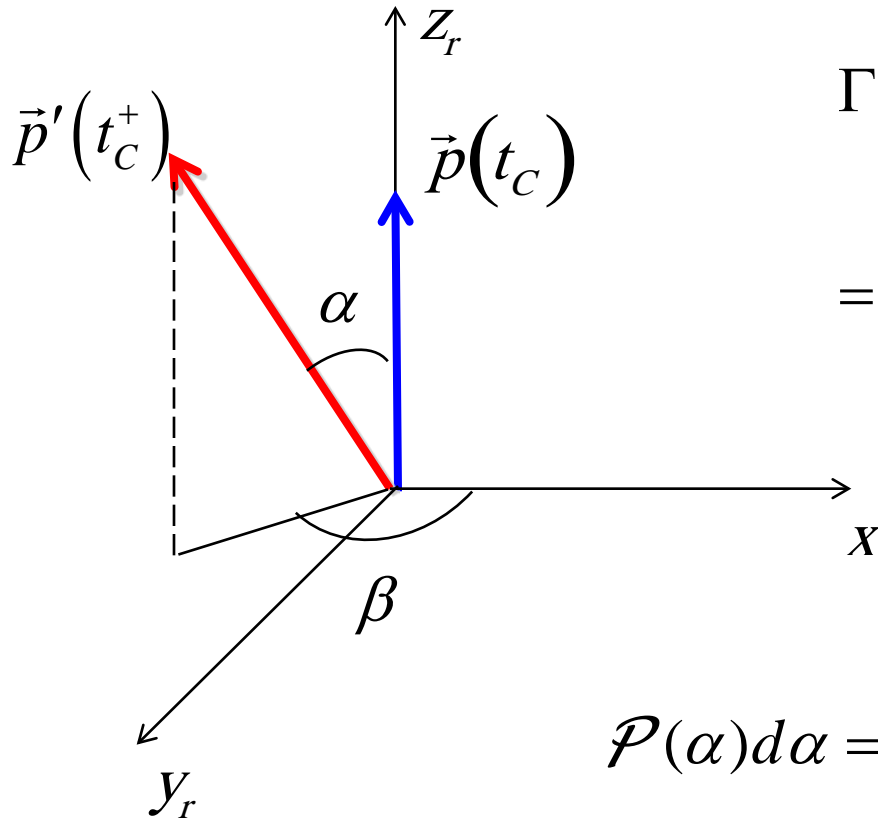
update after scattering: 3D



$$\beta = 2\pi r_3$$

r_3 is a random number uniformly distributed between 0 and 1

update after scattering: 3D

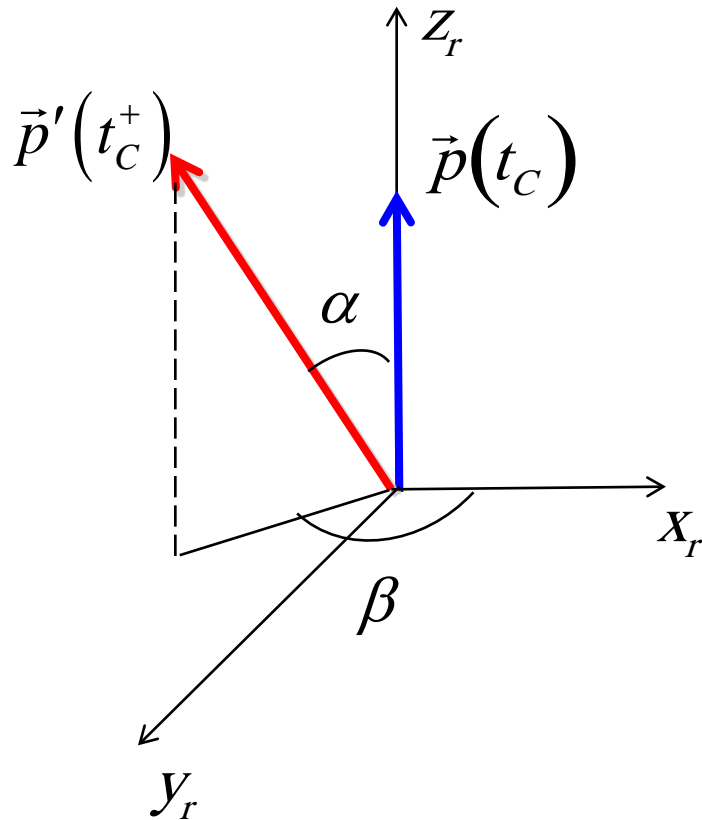


$$\Gamma_i = \frac{1}{\tau_i} = \sum_{\vec{p}'} S(\vec{p}, \vec{p}')$$

$$= C \int_0^{2\pi} d\beta \int_0^{\pi} \int_0^{\infty} S(\vec{p}, \vec{p}') \sin \alpha d\alpha p'^2 dp'$$

$$\mathcal{P}(\alpha) d\alpha = \frac{C \int_0^{2\pi} d\beta \int_0^{\infty} S(\vec{p}, \vec{p}') \sin \alpha d\alpha p'^2 dp'}{C \int_0^{2\pi} d\beta \int_0^{\pi} \int_0^{\infty} S(\vec{p}, \vec{p}') \sin \alpha d\alpha p'^2 dp'}$$

update after scattering: 3D

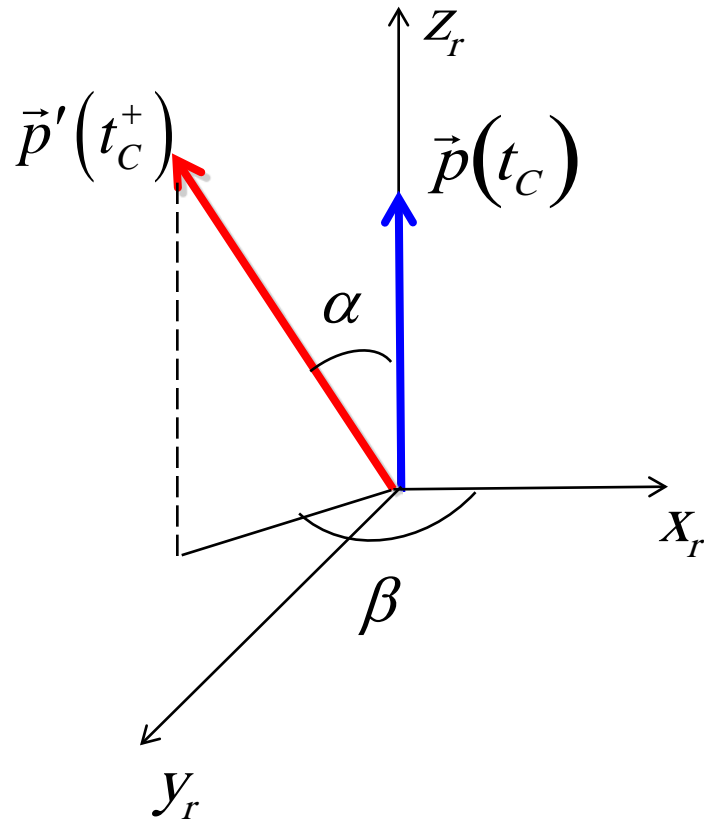


$$\mathcal{P}(\alpha)d\alpha = \frac{\int_0^{\infty} S(\vec{p}, \vec{p}') \sin \alpha d\alpha p'^2 dp'}{\int_0^{\pi} \int_0^{\infty} S(\vec{p}, \vec{p}') \sin \alpha d\alpha p'^2 dp'}$$

example: isotropic scattering

$$\mathcal{P}(\alpha)d\alpha = \frac{\sin \alpha d\alpha}{\int_0^{\pi} \sin \alpha d\alpha} = \frac{\sin \alpha d\alpha}{-\cos \alpha \Big|_0^{\pi}}$$

update after scattering: 3D



example: isotropic scattering

$$\mathcal{P}(\alpha) d\alpha = \frac{\sin \alpha d\alpha}{2} = \mathcal{P}(r) dr$$

$$\int_0^{r_4} dr = r_4 = \int_0^{\alpha} \frac{\sin \alpha d\alpha}{2} = \frac{-\cos \alpha}{2} \Big|_0^{\alpha}$$

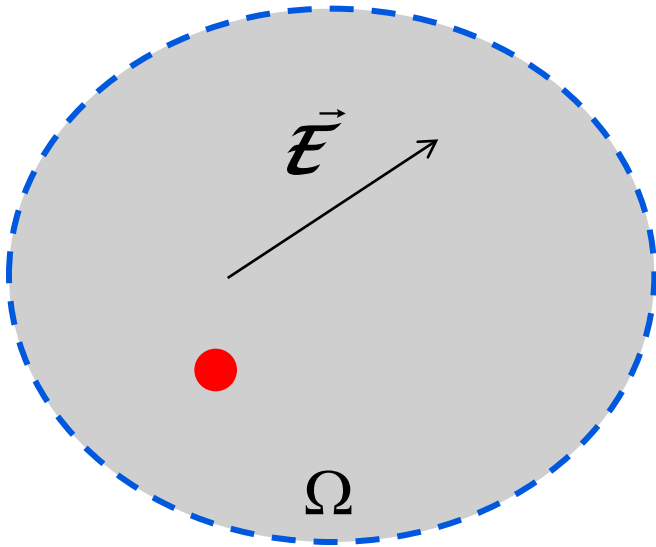
$$2r_4 = 1 - \cos \alpha$$

$$\cos \alpha = 1 - 2r_4$$

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MC simulation of a bulk semiconductor



initialize $\vec{p}(t=0)$

- 1) select: $t_C = -\frac{1}{\Gamma_0} \ln r_1$
- 2) update: $\vec{r}(t_C^-), \vec{p}(t_C^-), E(t_C^-)$
- 3) identify scattering event: r_2
- 4) update: $E(t_C^+), \vec{p}(t_C^+), \theta, \phi \quad r_3, r_4$
- 5) repeat

MC simulation of a device

$$Q_{TOT} = qALN_D^+$$

$$Q_{super} = \frac{N_{electrons}}{N_{sim}} \times (-q)$$

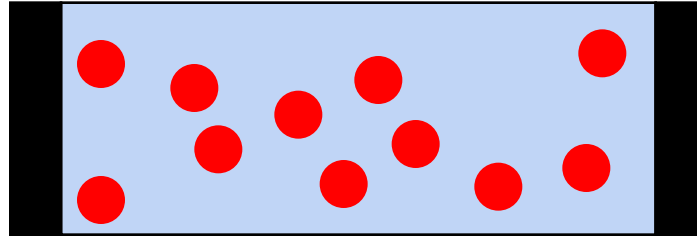
- 1) Load up device with N “super electrons”

example:

$$A = 1\mu\text{m} \times 1\mu\text{m} \quad L = 10\mu\text{m} \quad N_D = 10^{18} \text{ cm}^{-3}$$

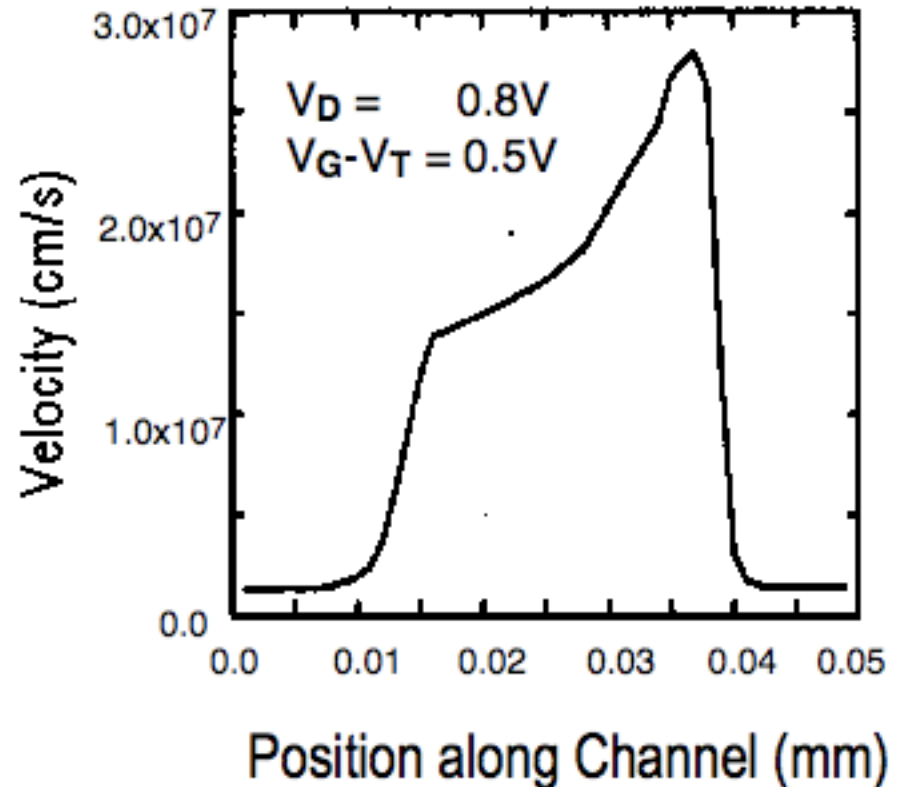
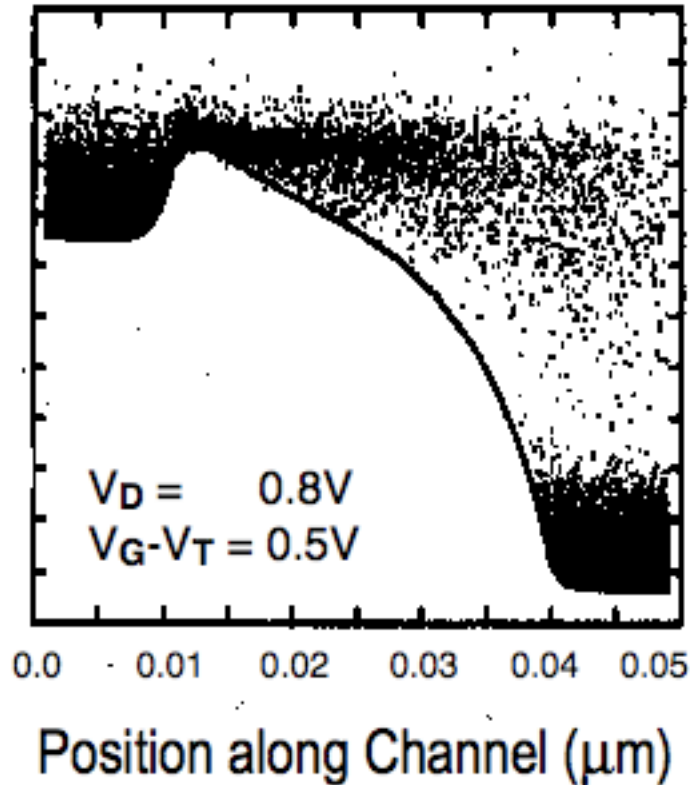
$$Q_{TOT} = 10^7 \times (-q)$$

MC simulation of a device



- 1) load up device with N “super electrons”
- 2) track all electron trajectories for $\otimes T$ (replace electrons that leave the device)
- 3) collect statistics
- 4) solve Poisson equation for new self-consistent potential
- 5) repeat until steady-state

Damocles



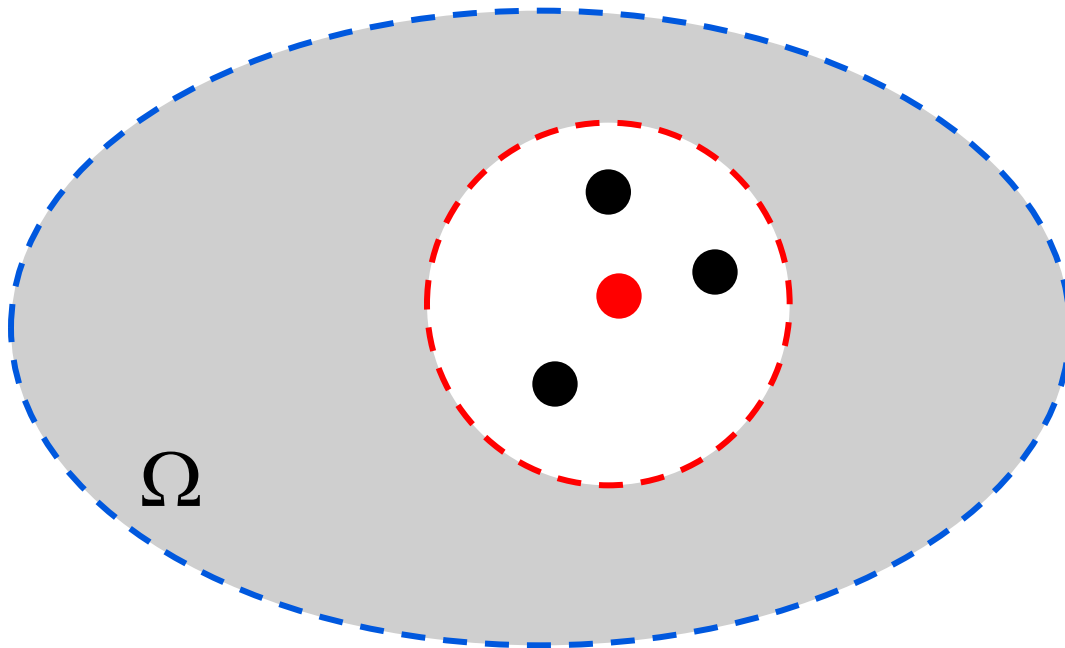
Frank, Laux, and Fischetti, IEDM Tech. Dig., p. 553, 1992

comments

$$\frac{\partial f}{\partial t} + \vec{v} \cdot \nabla_r f - q \vec{\mathcal{E}} \cdot \nabla_p f = \hat{C} f$$

a six-dimensional, time-dependent problem in 3D

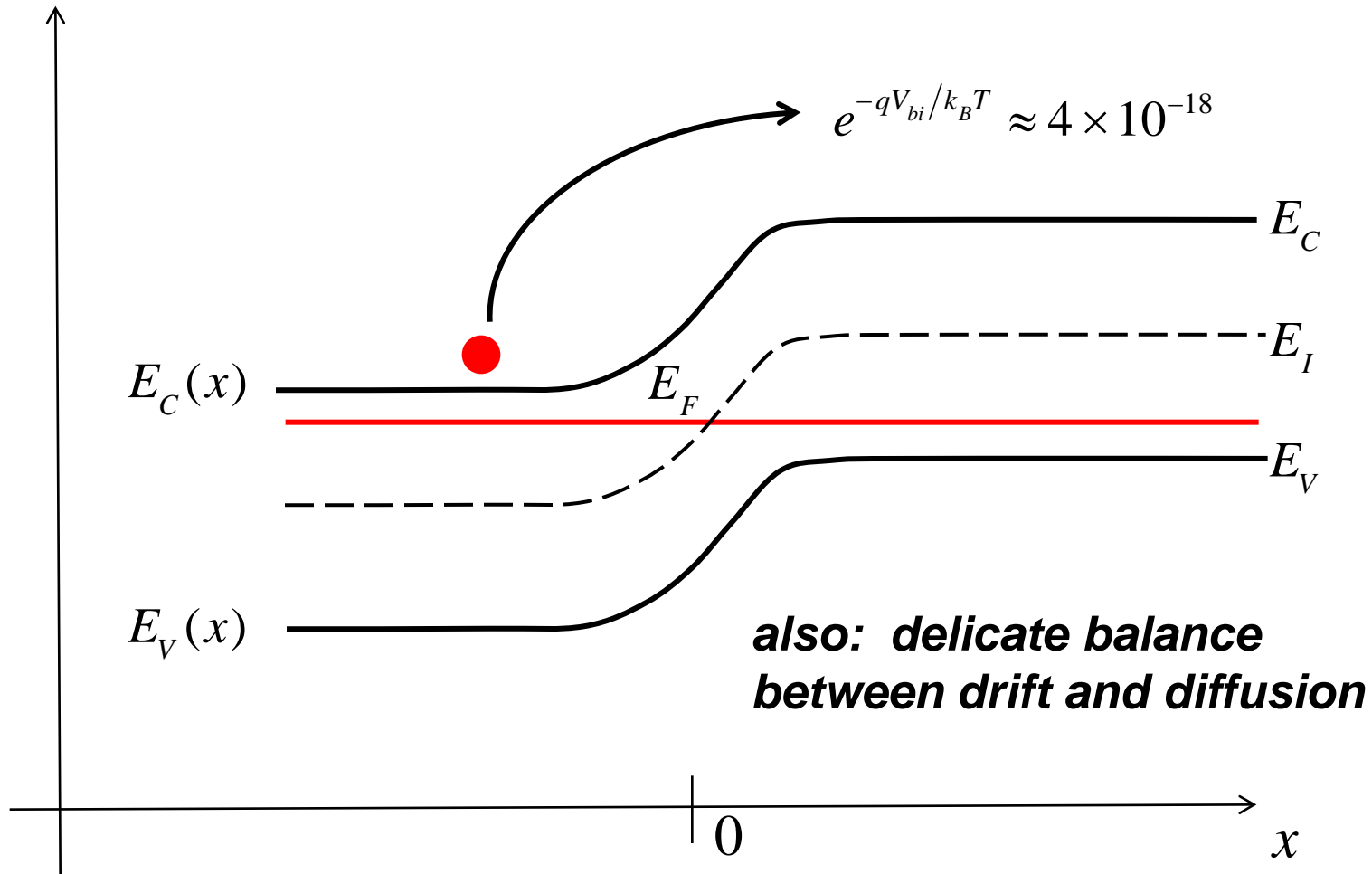
“curse of dimensionality”



- 1) long range forces due to the average charge in the device (from Poisson's equation)
- 2) short-range (Coulomb) forces

$$F = \frac{q_1 q_2}{4 \pi \epsilon r_{12}^2}$$

near-equilibrium



Monte Carlo simulation and the BTE

One can show (see Lundstrom, Sec. 6.8) that Monte Carlo simulation provides a solution to the BTE -

except when short range e-e scattering is included in which case it goes beyond the BTE by treating particle-particle correlations.

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summary

- 1) Monte Carlo simulation is a technique to solve the BTE (can go beyond by treating correlations)
- 2) Details of bandstructure and scattering are readily included
- 3) But rare events can be hard to treat
- 4) When applicable, MC simulation typically provides the most accurate solutions of the BTE
- 5) For a description of a state-of-the-art MC simulator:

<http://www.research.ibm.com/DAMOCLES/home.html>

questions

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