

From Atoms to Materials: Predictive Theory and Simulations

Week 3: Dynamics of Atoms – Classical Mechanics and MD Simulations

Lecture 3.2: Interatomic Potentials for Molecular Materials: Covalent Interactions

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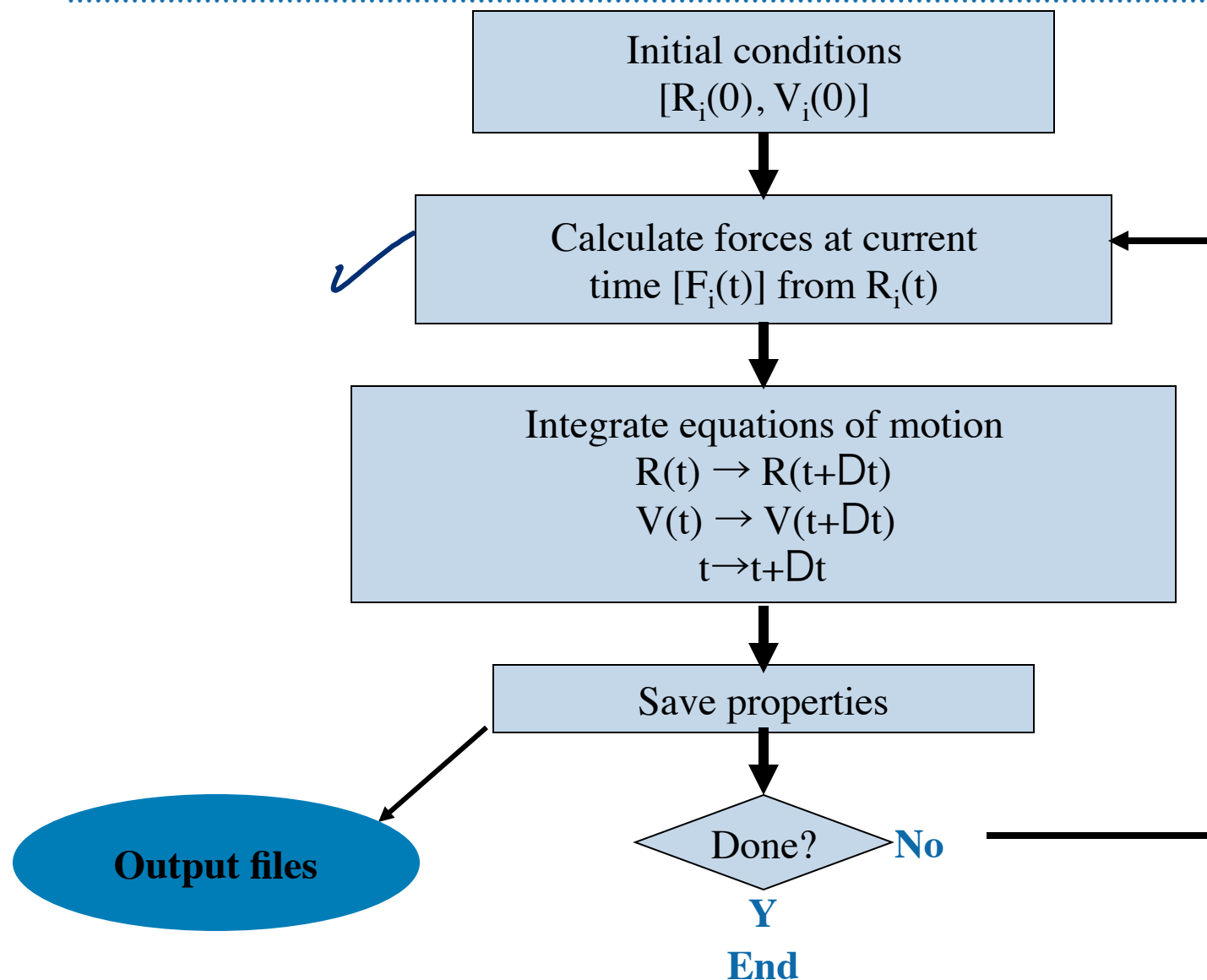
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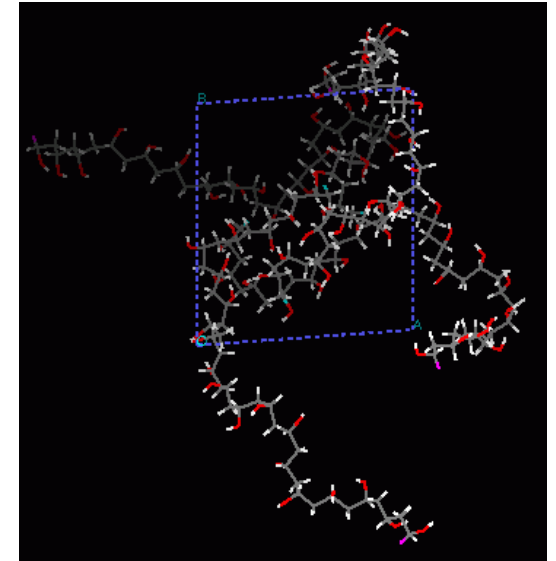
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Structure of a minimalist MD code



How can we compute interatomic potential?

- *Ab initio* electronic structure methods
 - Density Functional Theory, Hartree-Fock
 - Quantum Monte Carlo
 - Very accurate but computationally intensive (applicable to small systems – 100s-1000s atoms)
- Interatomic potentials:



$$V\left(\left\{R_i\right\}\right)=V_{\text{cov/met}}\left(\left\{R_i\right\}\right)+V_{\text{electr}}\left(\left\{R_i\right\}\right)+V_{\text{vdW}}\left(\left\{R_i\right\}\right)$$

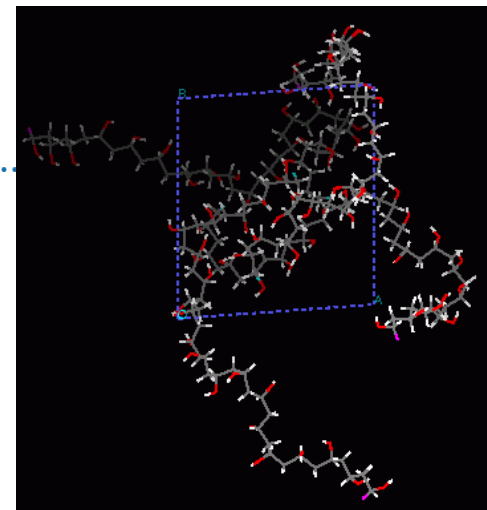
Parameterized to describe specific materials

- *Ab initio* data
- Experiments

Covalent interactions

Represent the effect of chemical bonds

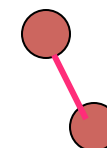
- 2-body, 3-body and 4-body interactions



2-body – bond stretch

- Pairs of atoms separated by a chemical bond)

$$\phi_{bond}(R) = \frac{1}{2} k_{\alpha\beta} \left(R - R_{\alpha\beta}^0 \right)^2 \quad \sim 700 \text{ kcal/mol / \AA}^2$$



Advantages of harmonic bonds

- Stable (bonds can't be broken) and simple

Disadvantages

- Symmetric and compression and tension (no thermal expansion)

Covalent interactions

Other functional forms for bond stretch:

Lennard-Jones (6-12)

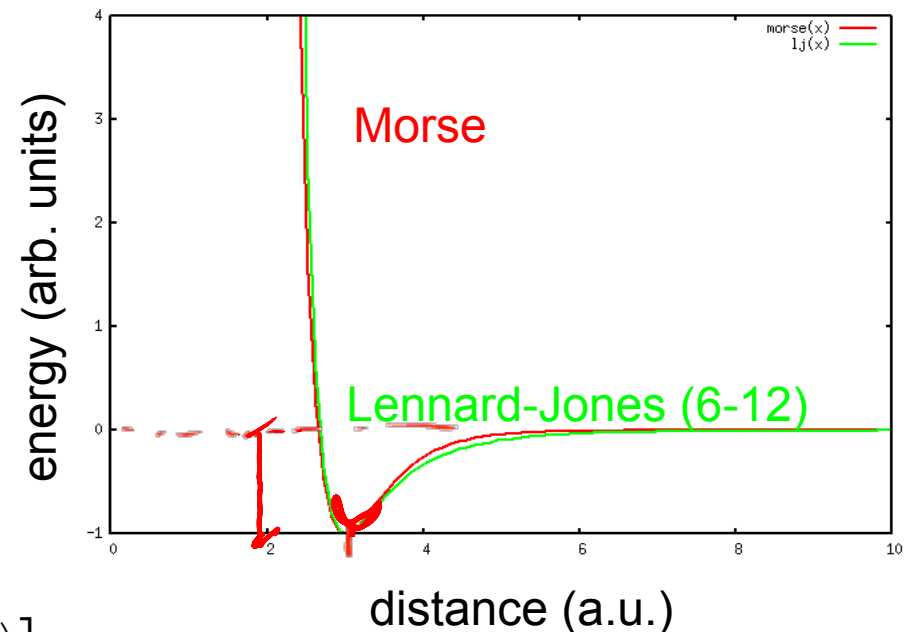
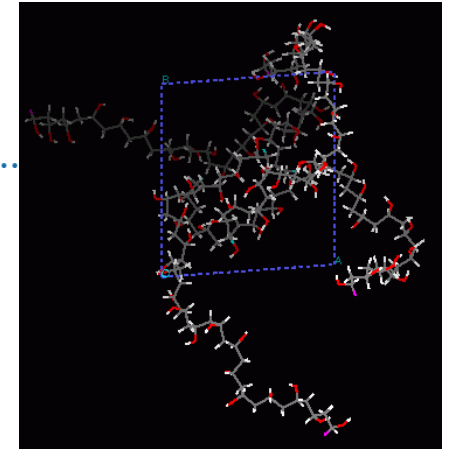
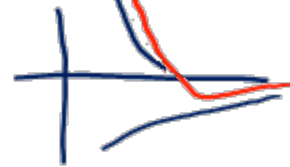
$$\phi_{LJ}(R) = 4\epsilon \left[\left(\frac{\sigma}{R} \right)^{12} - \left(\frac{\sigma}{R} \right)^6 \right]$$

Exponential-6

$$\phi_{\text{exp-6}}(R) = \epsilon \left[e^{-\gamma R} - AR^{-6} \right]$$

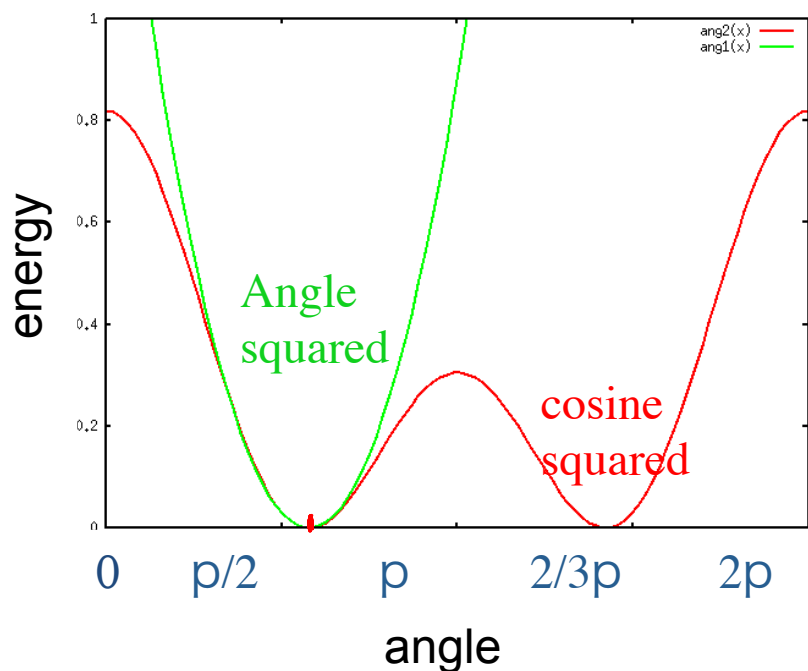
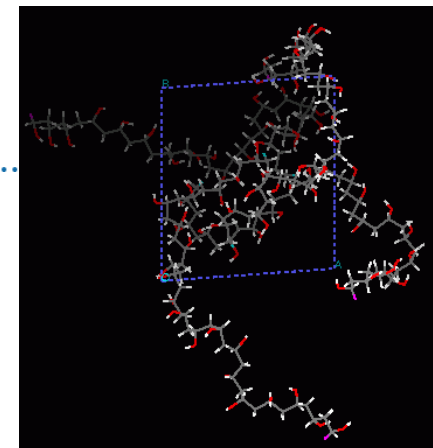
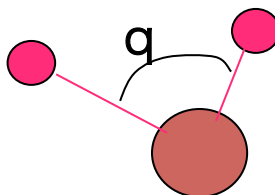
Morse

$$\phi_{\text{Morse}}(r) = \epsilon \left[\exp \left(\gamma \left(1 - \frac{R}{R_0} \right) \right) - 2 \exp \left(\frac{\gamma}{2} \left(1 - \frac{R}{R_0} \right) \right) \right]$$



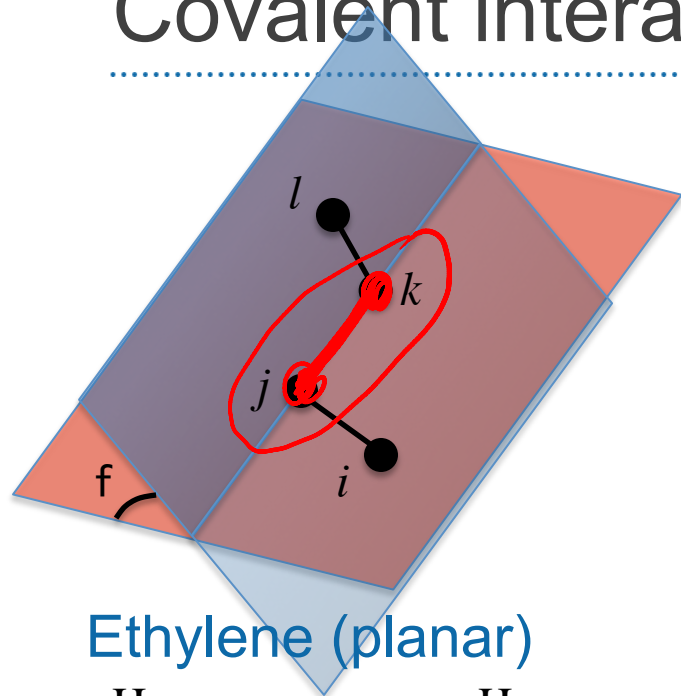
Covalent interactions: 3-body

Angle bending:

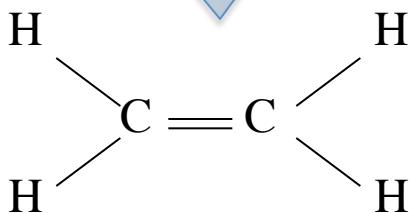


$$\phi_{angle}(\theta) = \begin{cases} \frac{1}{2} k_{ijk} (\theta - \theta_{ijk}^0)^2 & \checkmark \\ \frac{1}{2} \frac{k_{ijk}}{\sin^2 \theta_{ijk}^0} [\cos \theta - \cos \theta_{ijk}^0]^2 & \checkmark \end{cases}$$

Covalent interactions: 4 body



Ethylene (planar)



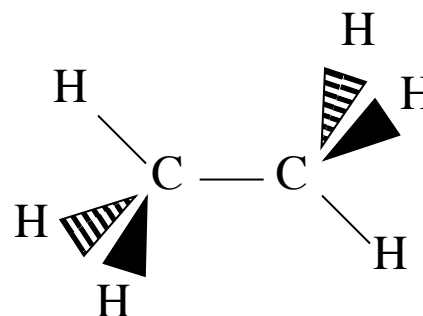
sp² hybridization

- Equilibrium at 0° & 180°
- Maxima at: 90° & 270°
- High torsional barrier
 - Rotating around C-C bond breaks the PI bond

Torsion bending (dihedral angles)

$$\varphi_{\text{torsions}}(\phi) = k_{ijkl} \left[1 - \cos \left(n_{jk} \left(\phi - \phi_{jk}^0 \right) \right) \right]$$

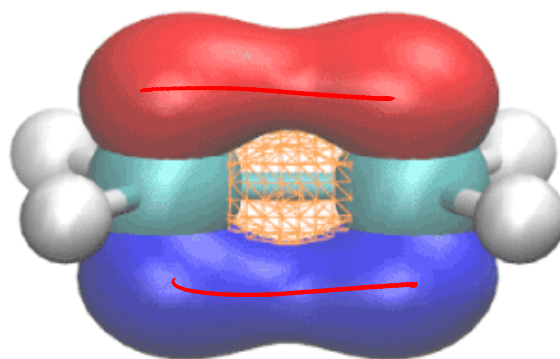
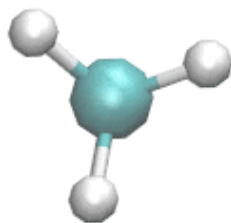
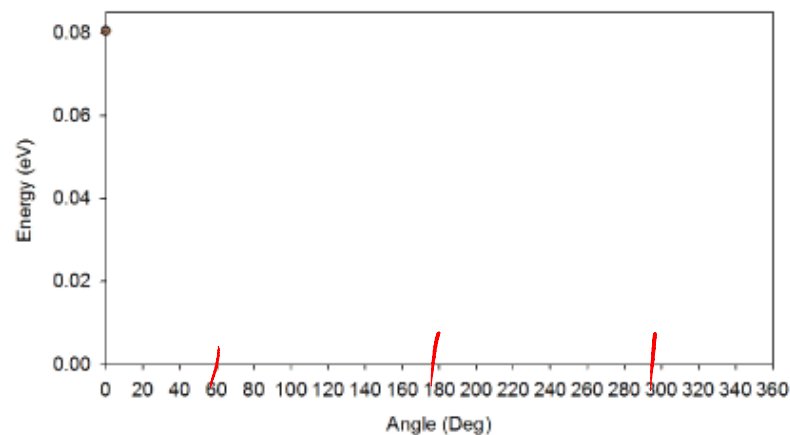
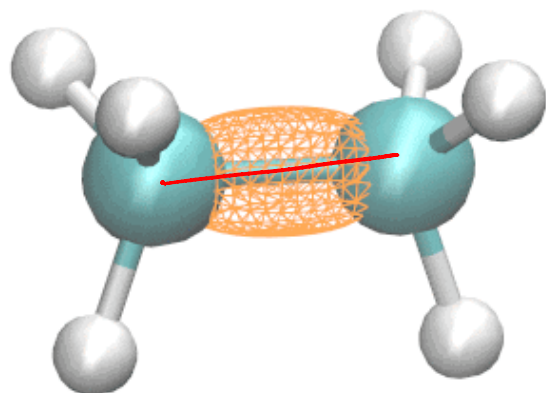
Ethane (staggered)



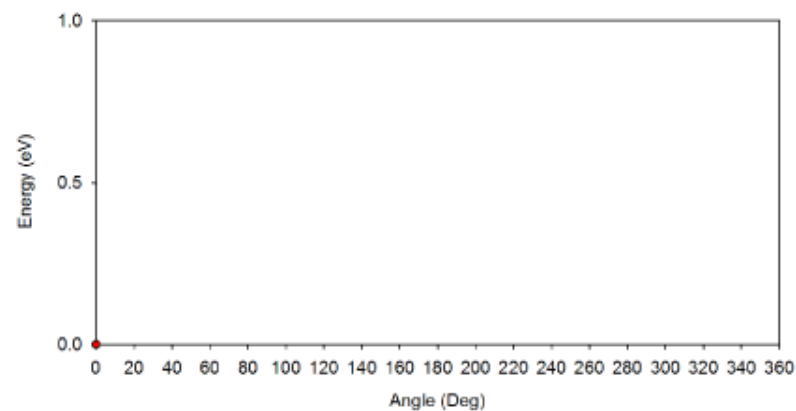
sp³ hybridization

- Equilibrium at 60°, 180°, & 300°
- Maxima at: 0°, 120°, & 240°
- Low torsional barrier
 - Dominated by vdW between H

Hybrid orbitals and torsional potentials



Bond Length=1.33 Ang



Atom types in molecular force fields

- Element alone is not enough to determine the interactions
 - E.g.: sp³ carbon and sp² carbon behave differently
- In traditional molecular force fields the user specifies *force field types* to establish the character of the atom
 - Force field types & connectivities need to be determined prior to MD simulation
 - FF types and connectivities are kept fixed during run
- Reactive force fields (like ReaxFF and REBO) describe interactions purely in terms of atomic positions and allow for atoms to change coordination and environment

Force fields for molecular simulations

CHARMM: Brooks, Bruccoleri, Olafson, States, Swaminathan, and Karplus. J. Comput. Chem., **4**, 187 (1983).

DREIDING: Mayo, Olafson, and Goddard, III, J. Chem. Phys. 94, 8897 (1990).

AMBER: Cornell, Cieplak, Bayly, Gould, Merz, Ferguson, Spellmeyer, Fox, Caldwell, and Kollman. J. Am. Chem. Soc., **117**, 5179 (1995).

COMPASS: H. Sun, J. Phys. Chem. B, **102**, 7338–7364 (1998)

OPLS: G. A. Kaminski, R. A. Friesner, J. Tirado-Rives, and W. L. Jorgensen J. Phys. Chem. B **105** 6474-6487 (2001).