

ECE-656: Fall 2011

Lecture 33:

Heterostructures

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L32 outline

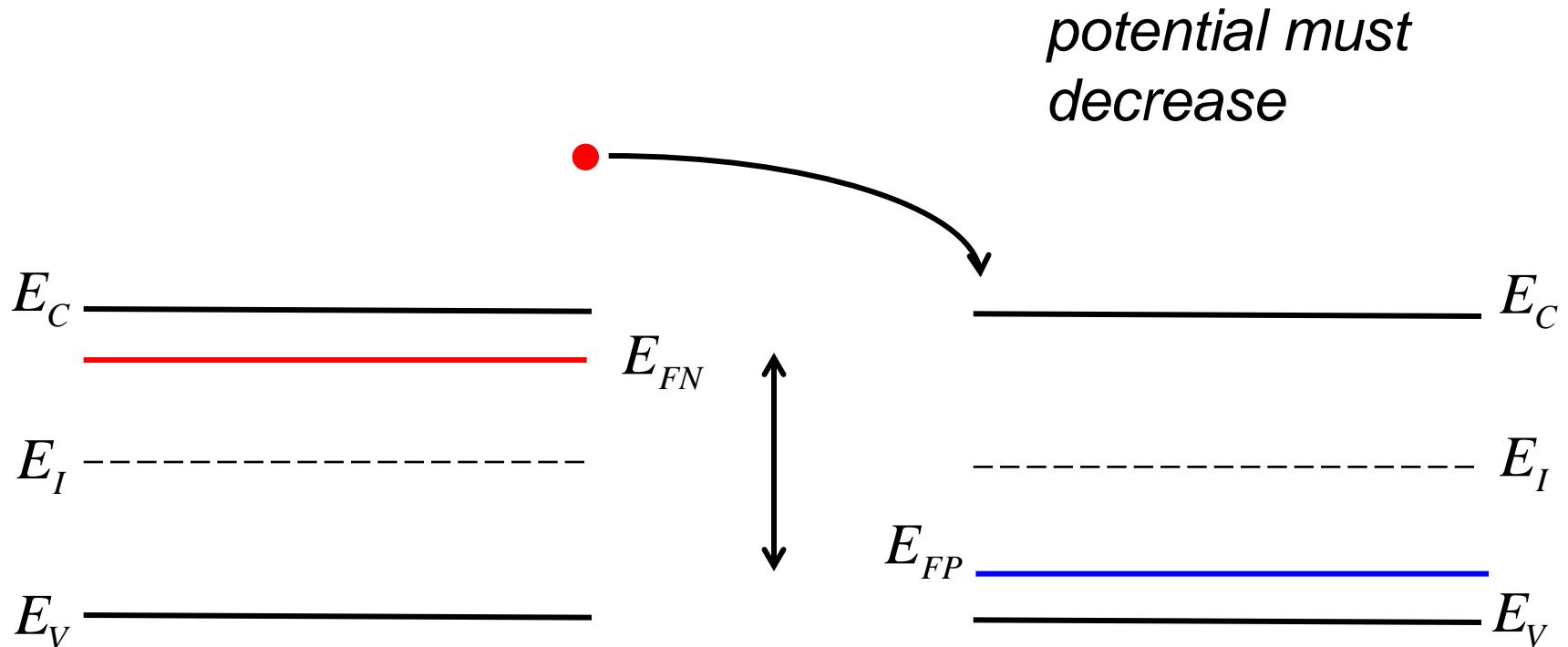
- 1) Review of L31
- 2) Carrier temperature and heat flux
- 3) Heterostructures**
- 4) Summary

(Reference: Chapter 5, Lundstrom, FCT)



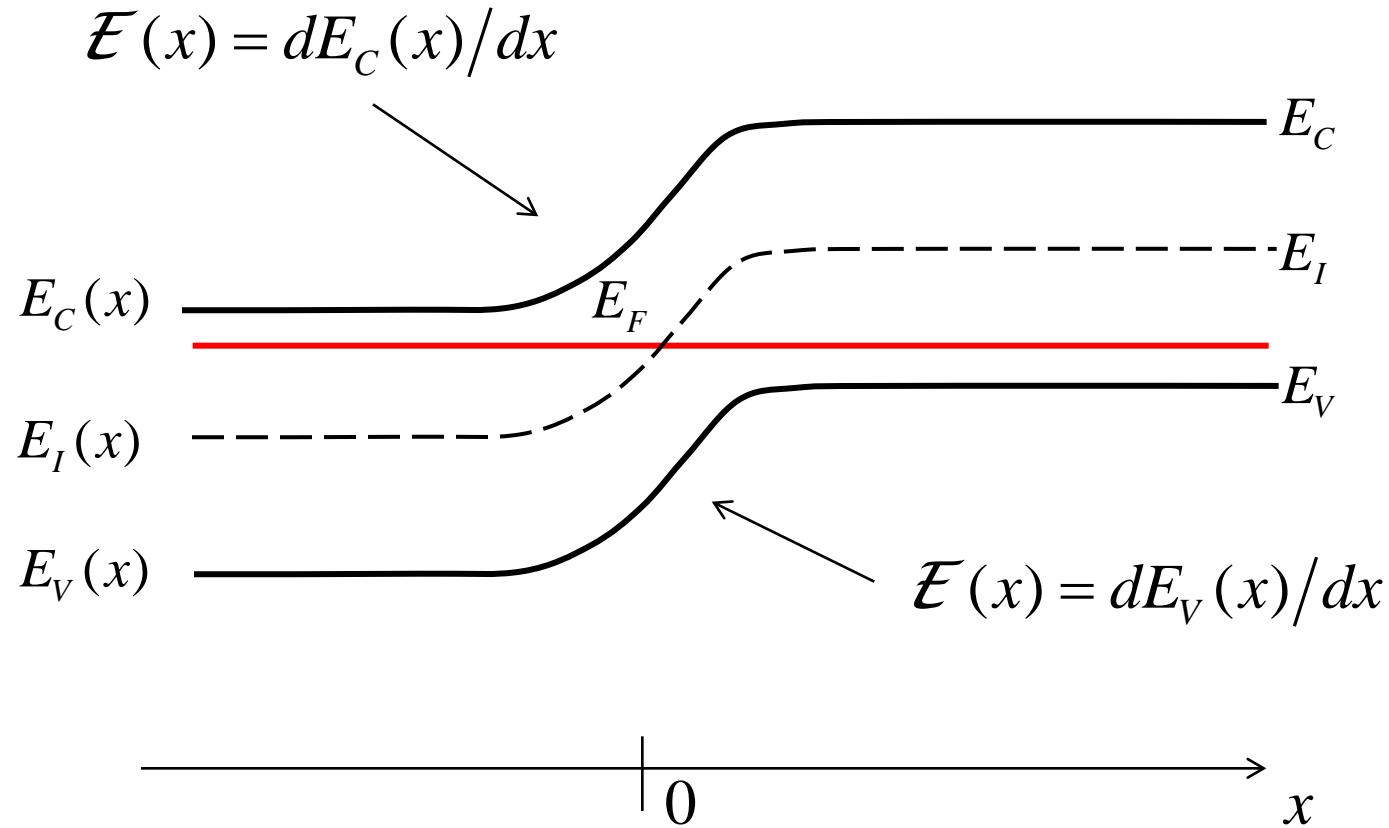
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review: pn homojunctions

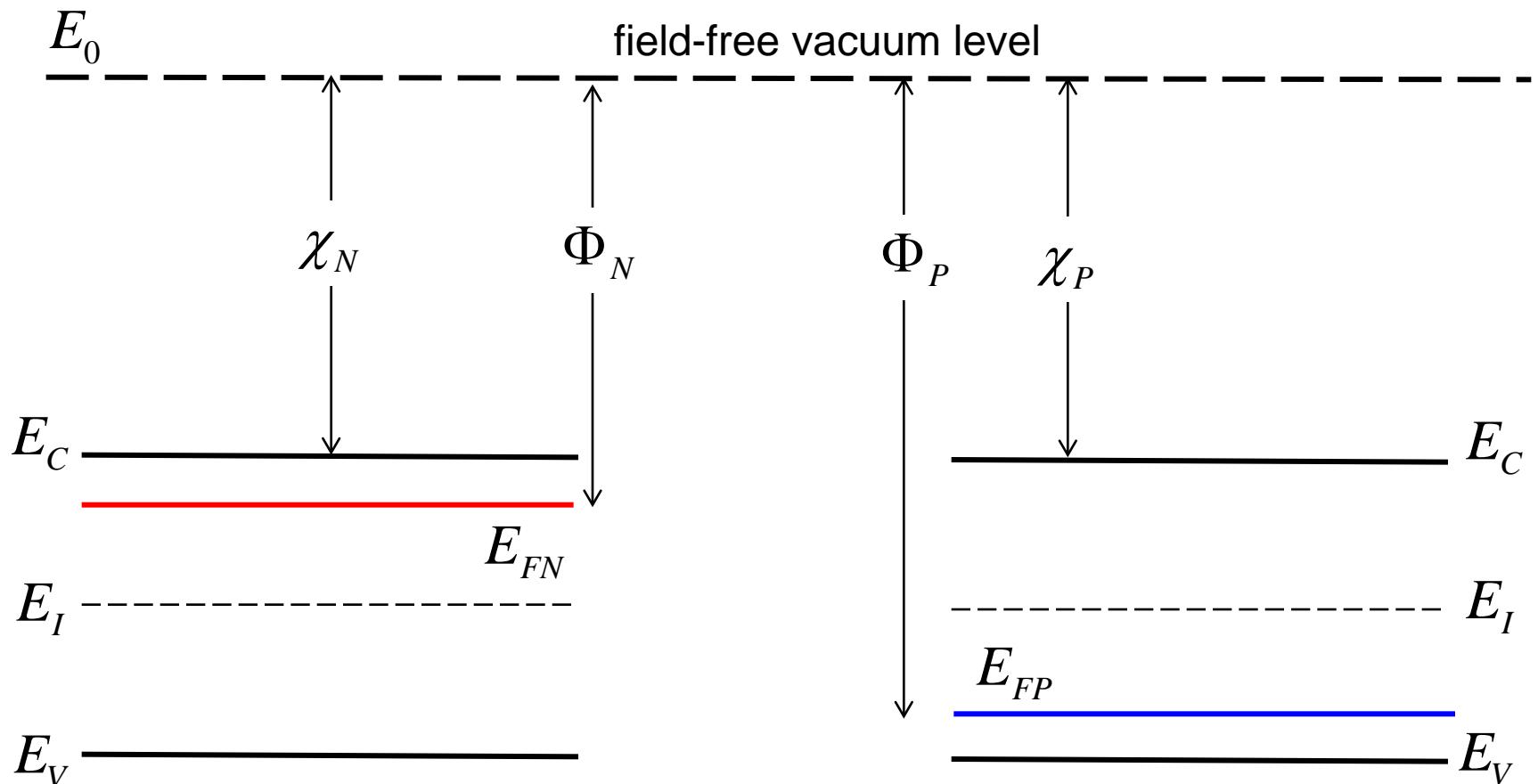


$$qV_{BI} = (E_{FN} - E_{FP})/q$$

review: pn homojunctions



reference for the energy bands



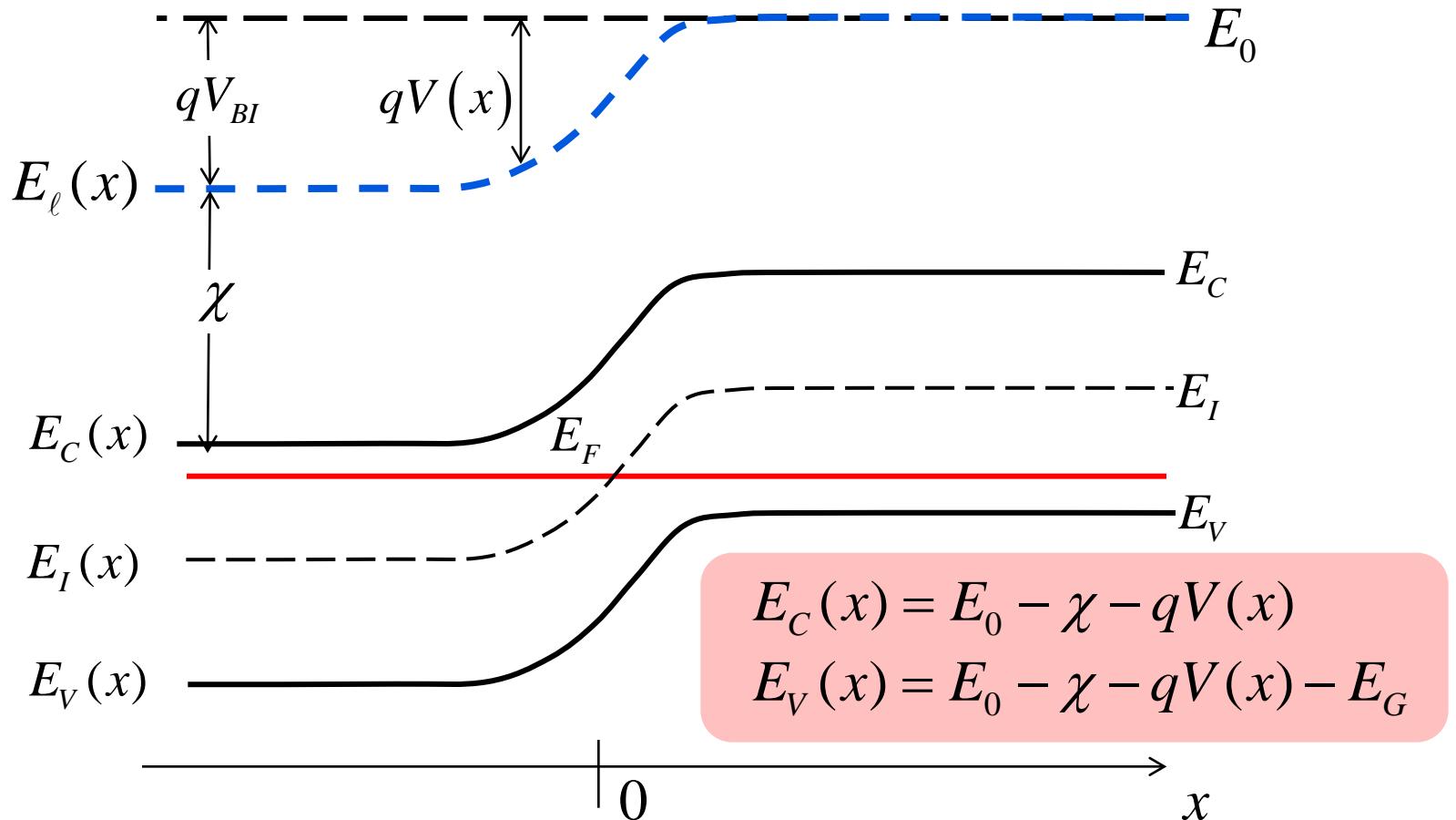
$$E_C = E_0 - \chi$$

$$E_V = E_0 - \chi - E_G$$

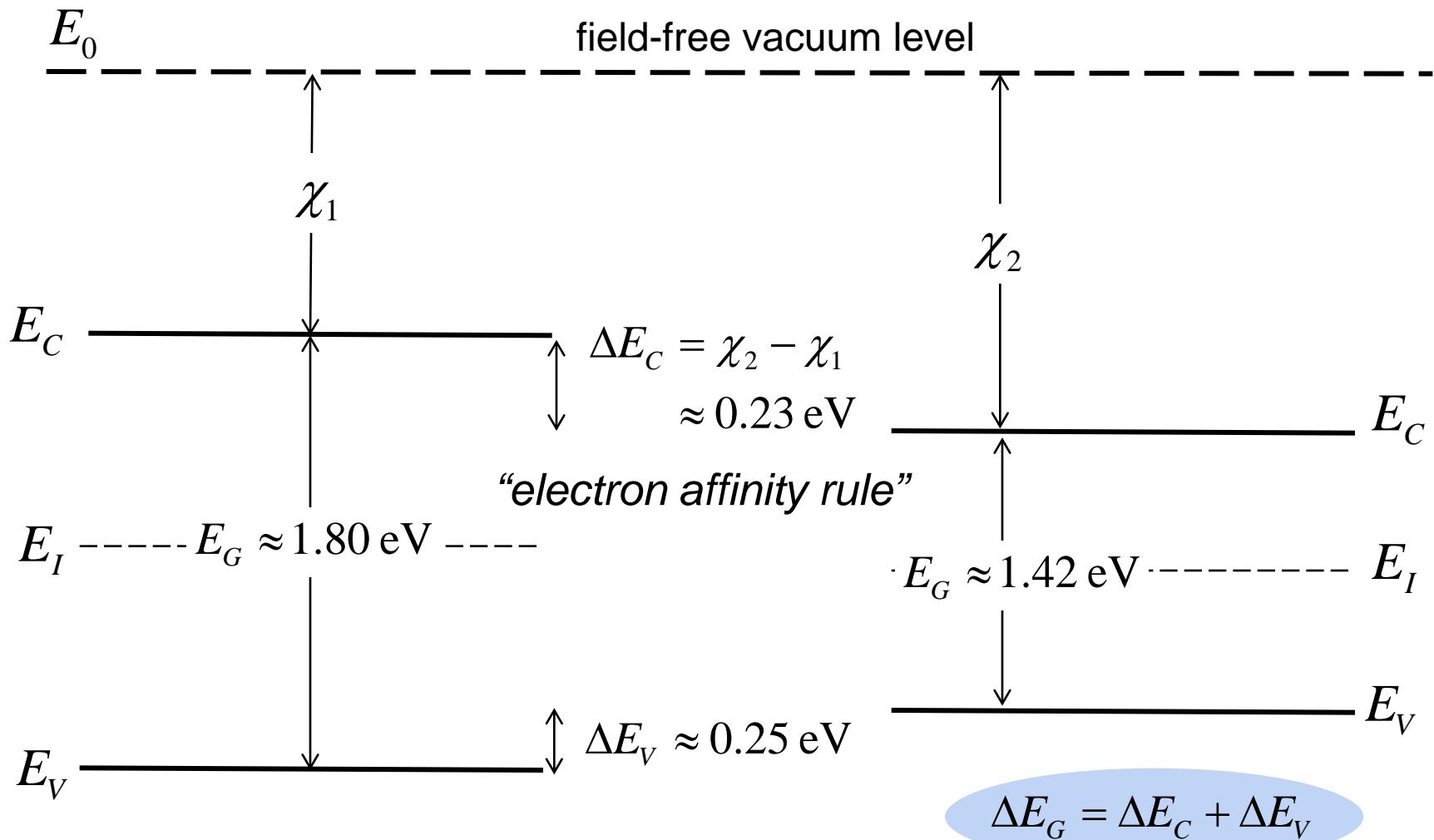
$$qV_{BI} = (\Phi_P - \Phi_N)$$

Lundstrom ECE-656 F11

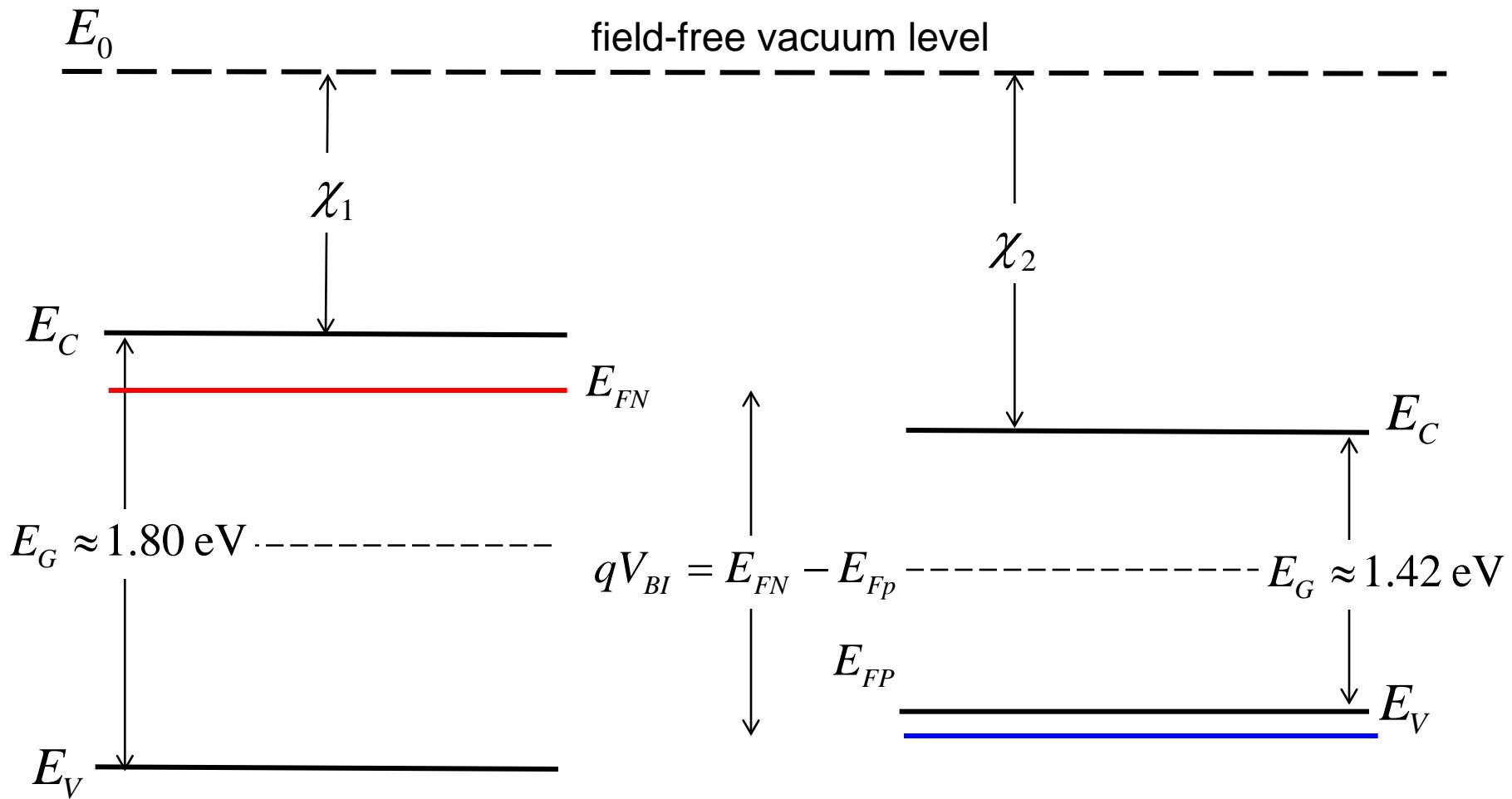
local vacuum level



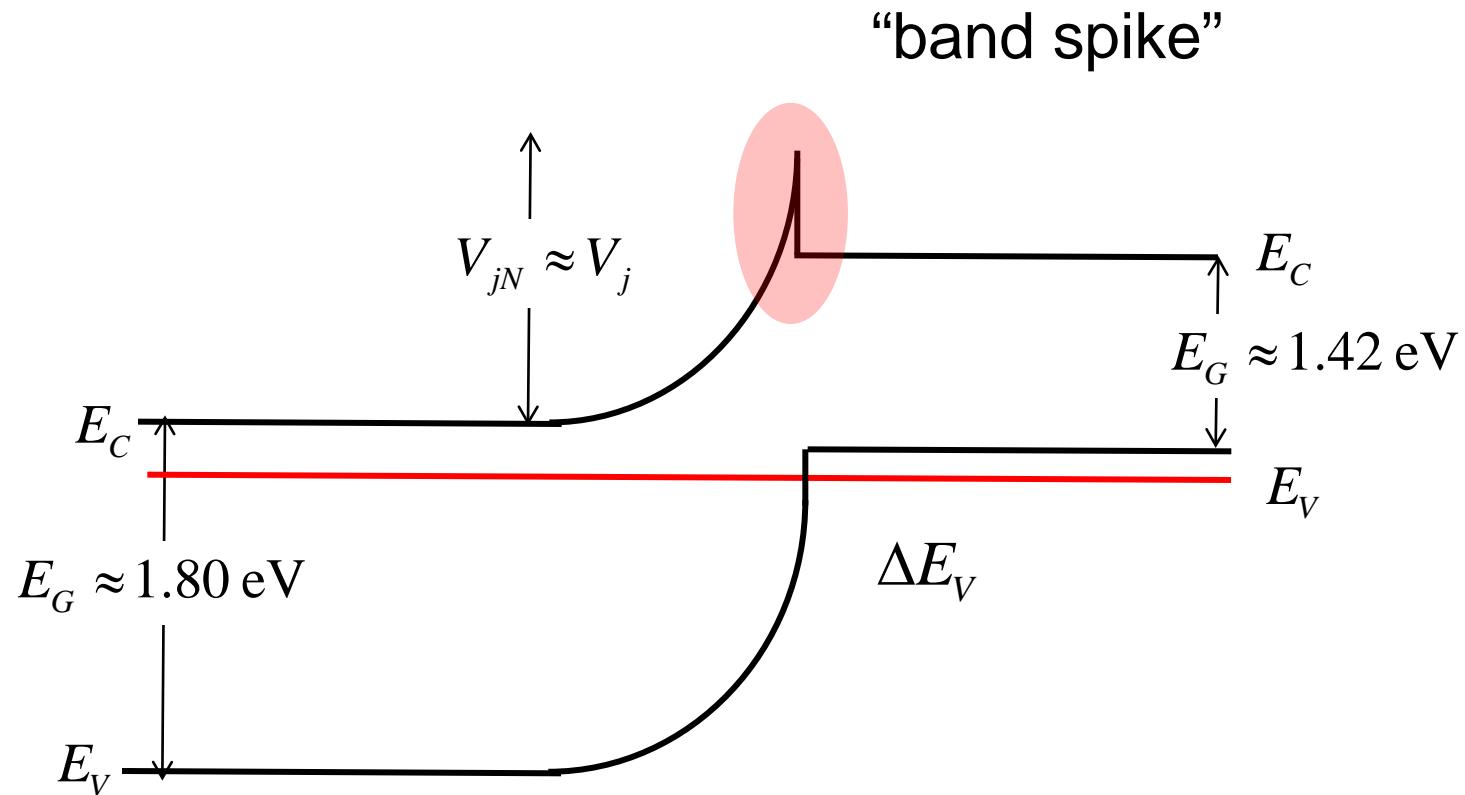
$\text{Al}_{0.3}\text{Ga}_{0.7}\text{As : GaAs}$ (Type I HJ)



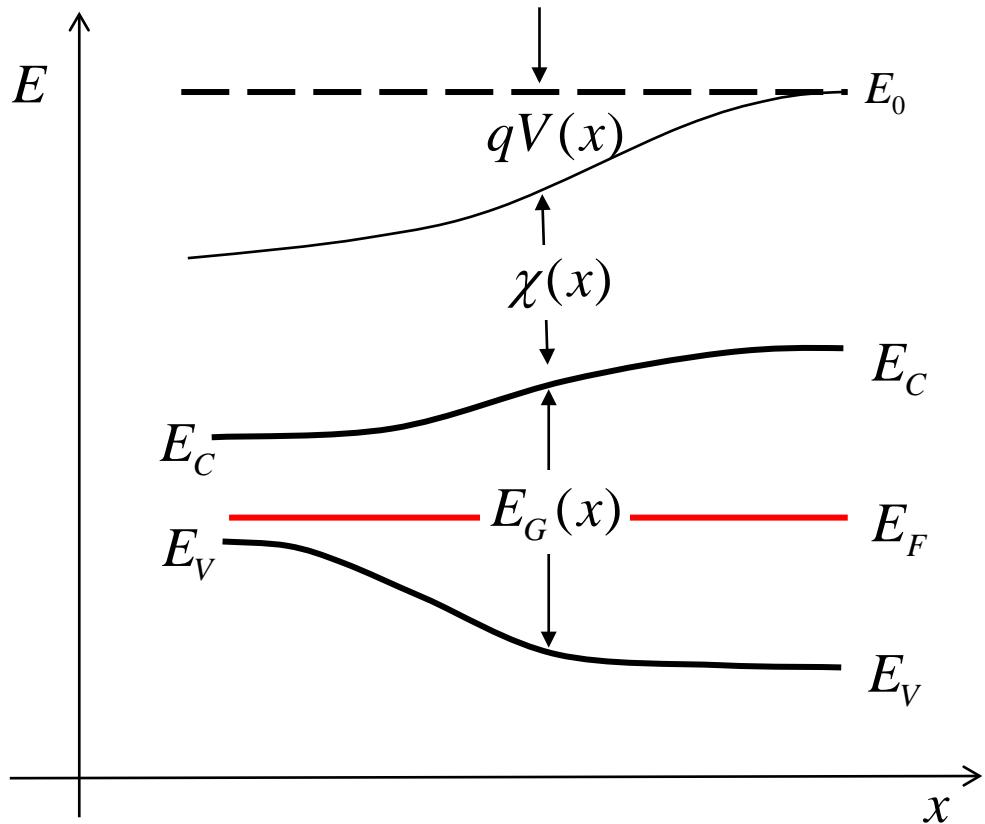
N-Al_{0.3}Ga_{0.7}As : p⁺-GaAs (Type I HJ)



N-Al_{0.3}Ga_{0.7}As : p⁺-GaAs (Type I HJ)



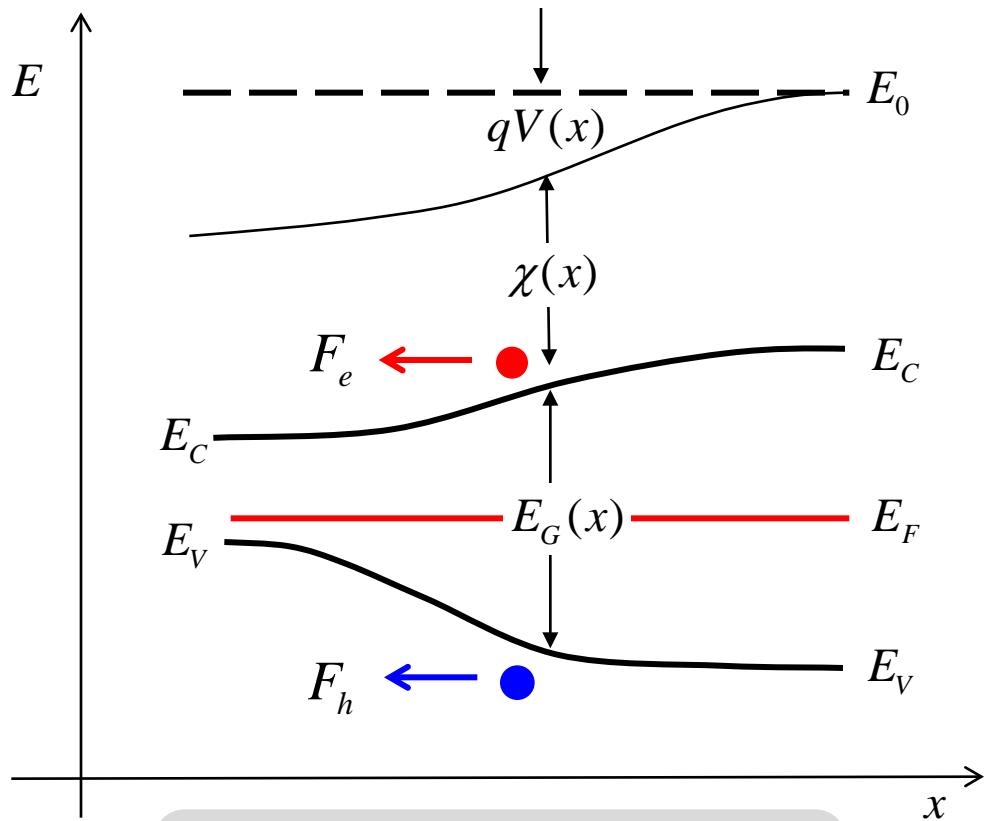
general, graded heterostructure



$$E_C(x) = E_0 - \chi(x) - qV(x)$$

$$E_V(x) = E_C(x) - E_G(x)$$

“quasi-electric fields”



$$E_C(x) = E_0 - \chi(x) - qV(x)$$

$$E_V(x) = E_C(x) - E_G(x)$$

$$F_e = -\frac{dE_C}{dx} = q \frac{dV}{dx} + \frac{d\chi}{dx}$$

$$F_e = -q\mathcal{E}(x) - q\mathcal{E}_{QN}(x)$$

$$\mathcal{E}_{QN} \equiv -\frac{1}{q} \frac{d\chi}{dx}$$

$$F_h = +\frac{dE_V}{dx} = -q \frac{dV}{dx} - \frac{d(\chi + E_G)}{dx}$$

$$F_h = +q\mathcal{E}(x) + q\mathcal{E}_{QP}(x)$$

$$\mathcal{E}_{QP} \equiv -\frac{1}{q} \frac{d(\chi + E_G)}{dx}$$

BTE

$$\frac{\partial f}{\partial t} + \vec{v} \bullet \nabla_r f + \frac{d \vec{p}}{dt} \bullet \nabla_p f = 0$$

$$\frac{d \vec{p}}{dt} = \frac{d(\hbar \vec{k})}{dt} = -\nabla_r E_C(\vec{r}) = -q \vec{E}(\vec{r})$$

(constant effective mass)

These equations do not hold when the effective mass is position dependent. Lundstrom, FCT, Sec. 5.8.

alternative approach: hole current

$$J_p = p \mu_p \frac{dF_p}{dx}$$

$$p = N_V(x) e^{(E_V - F_p)/k_B T_L} \quad F_p = E_V(x) - k_B T_L \ln(p/N_V)$$

$$\frac{dF_p}{dx} = \frac{dE_V(x)}{dx} - k_B T_L \left[\frac{1}{p} \frac{dp}{dx} - \frac{1}{N_V} \frac{dN_V}{dx} \right]$$

$$J_p = p \mu_p \left[\frac{dE_V(x)}{dx} + \frac{k_B T_L}{N_V} \frac{dN_V}{dx} \right] - k_B T_L \mu_p \frac{dp}{dx}$$

$$\frac{dE_V(x)}{dx} = \frac{d}{dx} [E_0 - \chi(x) - qV(x) - E_G(x)] = q(\mathcal{E}(x) + \mathcal{E}_{QP})$$

hole and electron currents

$$J_p = pq\mu_p \left[\mathcal{E} + \mathcal{E}_{QP} + \frac{k_B T_L}{q} \frac{1}{N_V} \frac{dN_V}{dx} \right] - qD_p \frac{dp}{dx}$$

“DOS effect”

$$J_n = nq\mu_n \left[\mathcal{E} + \mathcal{E}_{QN} - \frac{k_B T_L}{q} \frac{1}{N_C} \frac{dN_C}{dx} \right] + qD_n \frac{dn}{dx}$$

quasi-electric fields

$$\mathcal{E}_{QP} \equiv -\frac{1}{q} \frac{d(\chi + E_G)}{dx} \quad \mathcal{E}_{QN} \equiv -\frac{1}{q} \frac{d\chi}{dx}$$

outline

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balance equation summary

- 1) The four balance equations can be reduced to two continuity equations and two constitutive relations.
- 2) We can write them as two equations in two unknowns.
- 3) The unknowns are n and W or n and T_e .

the simplified equations

$$\frac{\partial n}{\partial t} = \frac{1}{q} \nabla \cdot \vec{J}_n$$

$$\vec{J}_n = n q \mu_n \vec{\mathcal{E}} + \frac{2}{3} \mu_n \nabla W$$

$$\frac{\partial W}{\partial t} = -\nabla \cdot \vec{F}_W + \vec{J}_n \cdot \vec{\mathcal{E}} - \frac{(W - W_0)}{\langle \tau_E \rangle}$$

$$\vec{F}_W = ?$$

cont. eqn. for electrons

current equation

continuity eqn. for energy

current eqn. for energy

energy current equation

First approach:

$$\vec{F}_W = -\frac{2}{3} \mu_E W \vec{\mathcal{E}} - \langle \tau_{F_W} \rangle \nabla \bullet \vec{X} \quad X_{ij} \approx \frac{5}{3} \frac{k_B T_e}{q} \mu_E W \delta_{ij}$$

Second approach:

$$\vec{F}_W = W \vec{v}_d + n k_B T_e \vec{v}_d + \vec{Q} \quad \vec{Q} \approx -\kappa_e \nabla T_e$$

questions?

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