

# Modern Physics

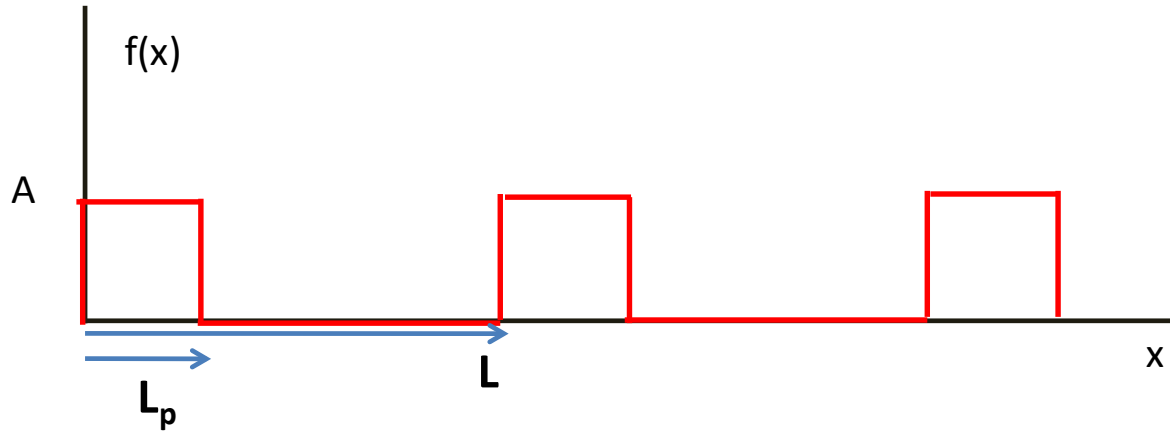
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Unit 4: Heisenberg's Uncertainty Principle

Lecture 4.4: The Fourier Integral

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# Fourier Series



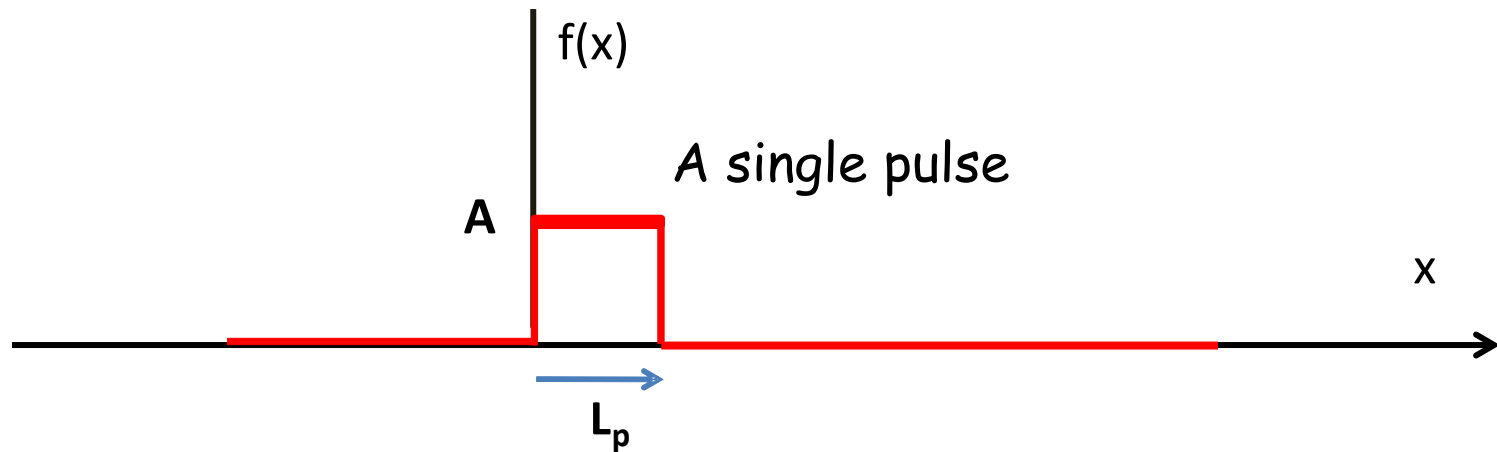
What happens when  $L \gg L_p$ ?

$$f(x) = a_o + \sum_{n=1}^N \left[ a_n \cos\left(\frac{2\pi nx}{L}\right) + b_n \sin\left(\frac{2\pi nx}{L}\right) \right]$$

Progressively more terms are required in the Fourier Series before the waveform is adequately approximated.

In the limit, when the waveform consists of only ONE pulse (what's the periodicity?), a continuous distribution of sines and cosines are required.

# The Fourier Transform



In this case,  $f(x)$  is no longer periodic. To reconstruct  $f(x)$ , we now require a continuous distribution of sines and cosines ( $k$  is no longer discrete!). The sums become integrals. Instead of the  $a_n$  and  $b_n$ ,  $f(x)$  is now specified by some function  $g(k)$  according to

$$f(x) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} g(k) e^{ikx} dk$$

Once  $f(x)$  is known, the function  $g(k)$  can be calculated from

$$g(k') = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} f(x) e^{-ik'x} dx$$

**Note: This is now called a Fourier Transform, not a Fourier Series**

# Points to Ponder

$$g(k') = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} f(x) e^{-ik'x} dx$$

1. Remember, you know  $f(x)$  AND you must specify a value for  $k'$
2. In general,  $g(k')$  will be a complex number
3. The Fourier Transform is really an “ideal” model because
  - it requires a function  $f(x)$  that goes from  $-\infty$  to  $+\infty$
  - it requires an infinite number of  $k'$
4. For every  $k'$  you specify, you need to perform the integral to get the value of  $g(k')$  - so you need to perform an infinite number of integrals
5. Sometimes, the integration works out so you have an **analytical form** for  $g(k')$  - this is somewhat of a special case
6. In a Fourier Series, you can understand one term, but in a Fourier transform, it's very difficult to comprehend a single term because there are an infinite number of them, they span both positive and negative values of  $k'$  AND  $g(k')$  for a specific  $k'$  can be a complex number

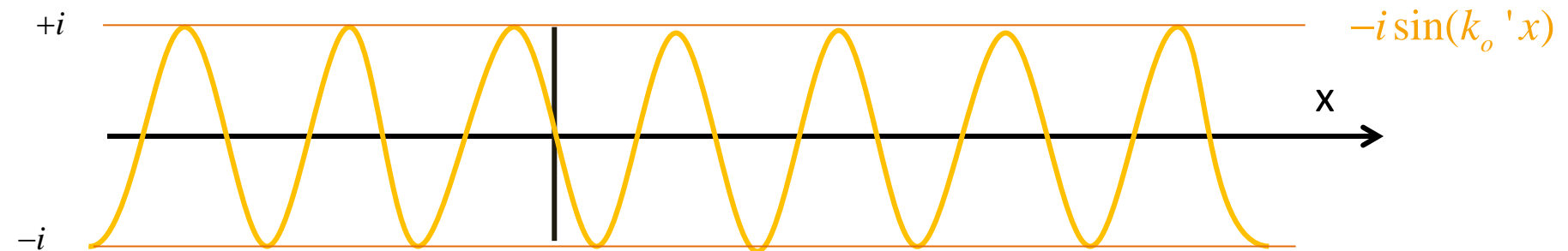
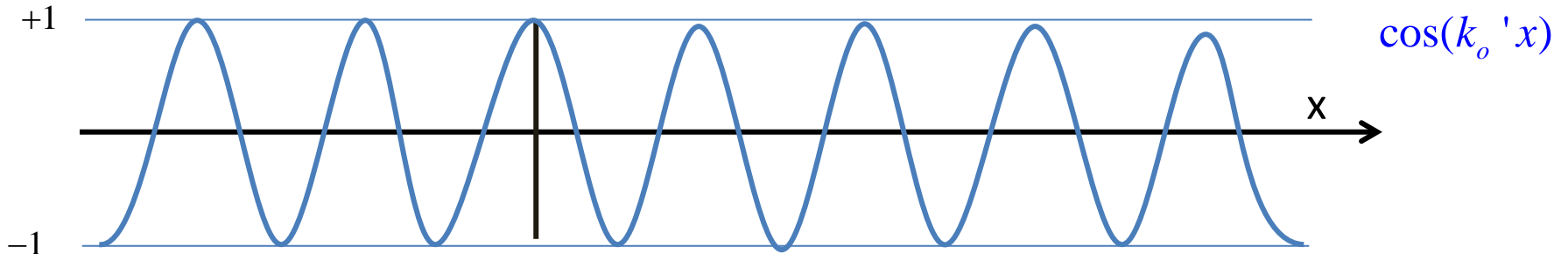
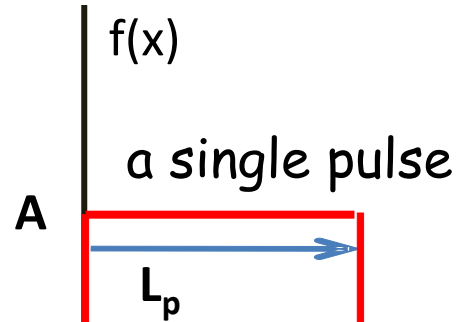
} like a  
double  
infinity

# What the Fourier Transform does

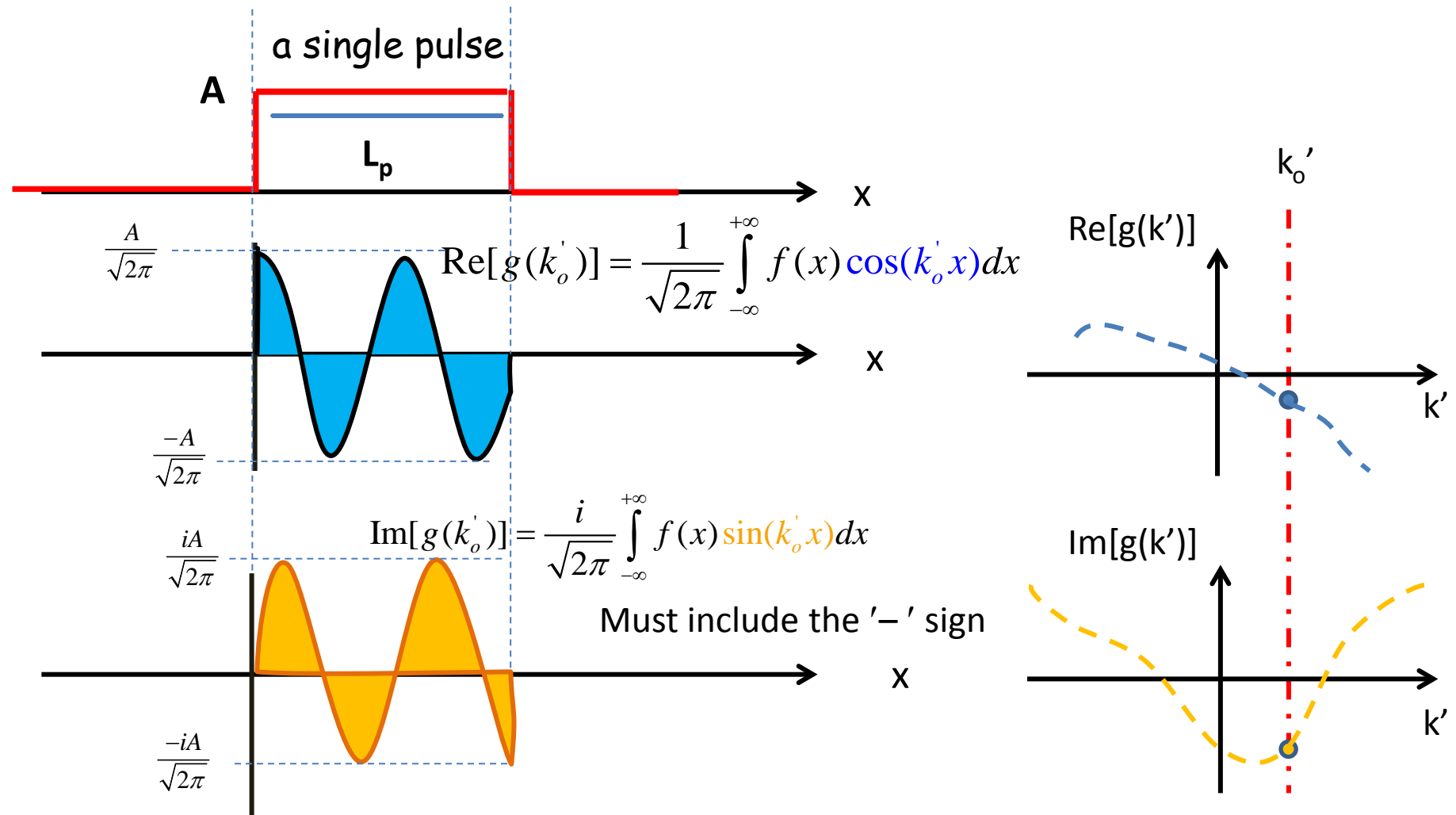
Pick a  $k' = k'_o$



$$e^{-ik'x} = \cos(k_o'x) - i \sin(k_o'x)$$



# Integrating Graphically



# Working out the math

$$g(k') = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} f(x) e^{-ik'x} dx = \frac{1}{\sqrt{2\pi}} \int_0^{L_p} A e^{-ik'x} dx$$

$$= \frac{A}{\sqrt{2\pi}} \frac{e^{-ik'x}}{-ik'} \Big|_0^{L_p} = \frac{A}{\sqrt{2\pi}} \left[ \frac{e^{-ik'L_p} - 1}{-ik'} \right]$$

$$= \frac{A}{\sqrt{2\pi}} e^{-ik'L_p/2} \left[ \frac{e^{-ik'L_p/2} - e^{ik'L_p/2}}{-ik'L_p/2} \right] \frac{L_p}{2} = \frac{A}{\sqrt{2\pi}} e^{-ik'L_p/2} \left[ \frac{e^{ik'L_p/2} - e^{-ik'L_p/2}}{ik'L_p/2} \right] \frac{L_p}{2}$$

$$= \frac{AL_p}{\sqrt{2\pi}} e^{-ik'L_p/2} \left[ \frac{e^{ik'L_p/2} - e^{-ik'L_p/2}}{2i} \right] \frac{1}{k'L_p/2}$$

$$= \frac{AL_p}{\sqrt{2\pi}} e^{-ik'L_p/2} \left[ \frac{\sin\left(\frac{k'L_p}{2}\right)}{\left(\frac{k'L_p}{2}\right)} \right]$$

Note that now,  
there is NO  
restriction on  $k'$

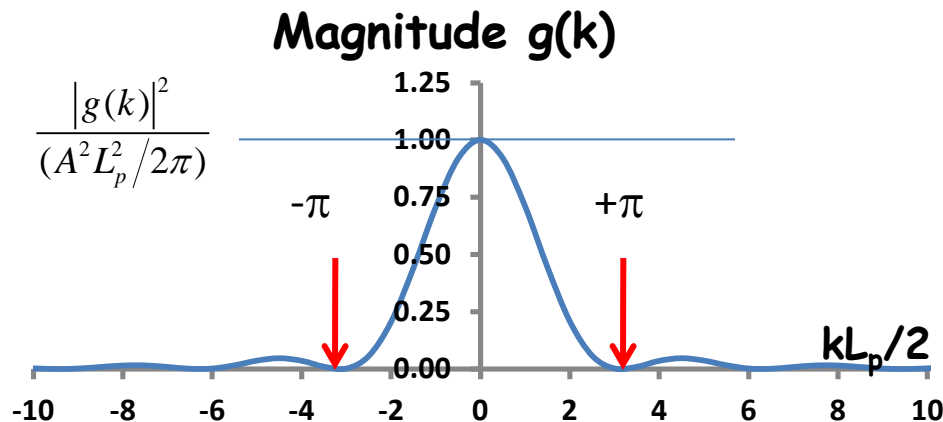
$$g(k) = \frac{AL_p}{\sqrt{2\pi}} e^{-ikL_p/2} \left[ \frac{\sin\left(\frac{kL_p}{2}\right)}{\left(\frac{kL_p}{2}\right)} \right]$$

$$g(k) = \frac{AL_p}{\sqrt{2\pi}} \left[ \frac{\sin\left(\frac{kL_p}{2}\right)}{\left(\frac{kL_p}{2}\right)} \right] \cos\left(\frac{kL_p}{2}\right) - i \frac{AL_p}{\sqrt{2\pi}} \left[ \frac{\sin\left(\frac{kL_p}{2}\right)}{\left(\frac{kL_p}{2}\right)} \right] \sin\left(\frac{kL_p}{2}\right)$$

$g(k)$  has real and imaginary parts

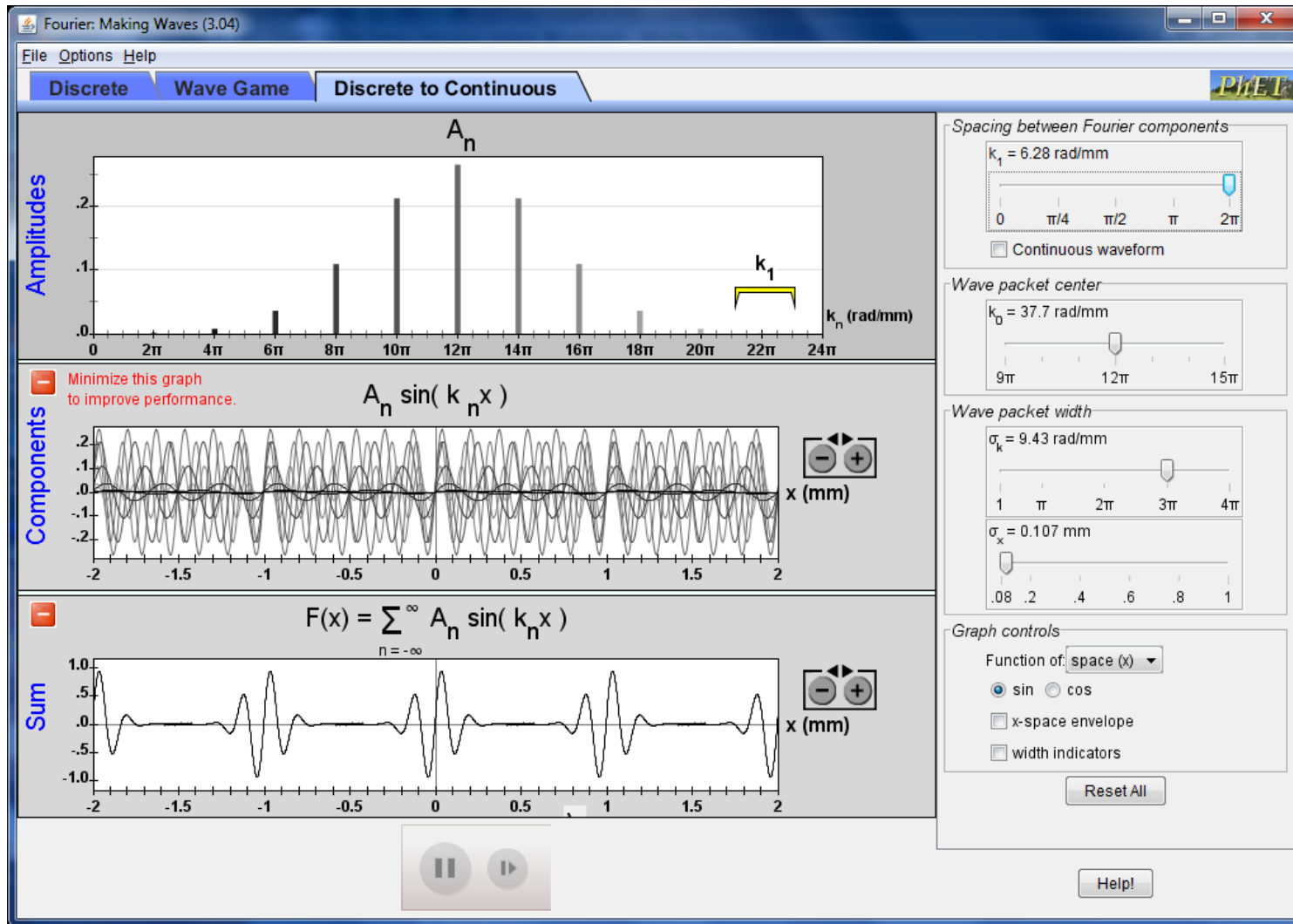
$$|g(k)|^2 = g^*(k)g(k) = \frac{A^2 L_p^2}{2\pi} \left[ \frac{\sin\left(\frac{kL_p}{2}\right)}{\left(\frac{kL_p}{2}\right)} \right]^2$$

magnitude of  $g(k)$





# Check it out!



[http://phet.colorado.edu/simulations/sims.php?sim=Fourier\\_Making\\_Waves](http://phet.colorado.edu/simulations/sims.php?sim=Fourier_Making_Waves)

Appendix: Using the Fourier Transform, you can now find the momentum eigenfunctions for the infinite square well

**Application:** Find the Momentum Eigenfunction for the  $n=3$  quantum state in the Infinite Square Well

$$\Phi_n(k) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} \Psi_n(x) e^{-ikx} dx = \frac{1}{\sqrt{2\pi}} \int_{-L/2}^{+L/2} \Psi_n(x) e^{-ikx} dx$$

The Infinite Square Well - as an example

| n | $\Psi_n(x)$  | $E_n$            | Levels      |
|---|--|------------------|-------------|
| 1 | $\sqrt{\frac{2}{L}} \cos\left(\frac{\pi x}{L}\right)$  | $h^2/8m_e L^2$   | $E_1=E_0$   |
| 2 | $\sqrt{\frac{2}{L}} \sin\left(\frac{2\pi x}{L}\right)$ | $4h^2/8m_e L^2$  | $E_2=4E_0$  |
| ③ | $\sqrt{\frac{2}{L}} \cos\left(\frac{3\pi x}{L}\right)$ | $9h^2/8m_e L^2$  | $E_3=9E_0$  |
| 4 | $\sqrt{\frac{2}{L}} \sin\left(\frac{4\pi x}{L}\right)$ | $16h^2/8m_e L^2$ | $E_4=16E_0$ |

Choose  $n=3$  for the Infinite Square Well - calculate the momentum eigenfunction

$$\Phi_3(k) = \frac{1}{\sqrt{2\pi}} \int_{-L/2}^{+L/2} \Psi_3(x) e^{-ikx} dx = \frac{1}{\sqrt{2\pi}} \int_{-L/2}^{+L/2} \Psi_3(x) e^{-ikx} dx$$

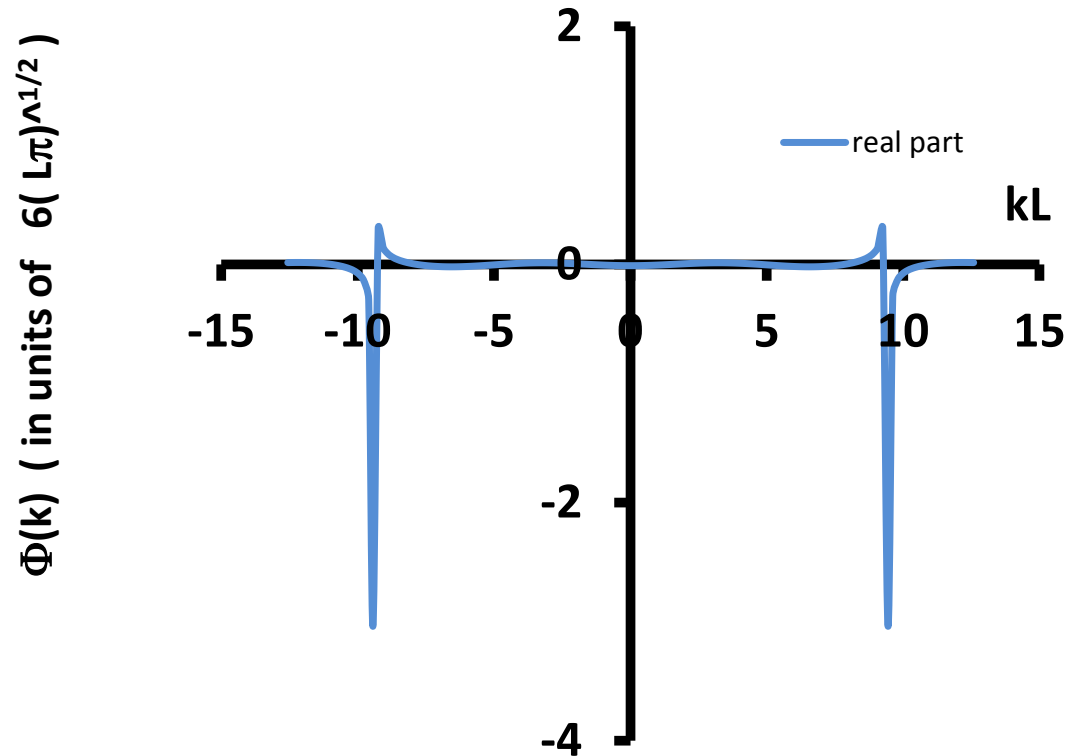
$$\Psi_3(x) = \sqrt{\frac{2}{L}} \cos\left(\frac{3\pi x}{L}\right)$$

$$\begin{aligned} \Phi_3(k) &= \frac{1}{\sqrt{2\pi}} \sqrt{\frac{2}{L}} \int_{-L/2}^{+L/2} \cos\left(\frac{3\pi x}{L}\right) e^{-ikx} dx \\ &= 6\sqrt{L\pi} \left[ \left( \frac{1}{[kL - 3\pi][kL + 3\pi]} \right) \cos\left(\frac{kL}{2}\right) \right] \end{aligned}$$

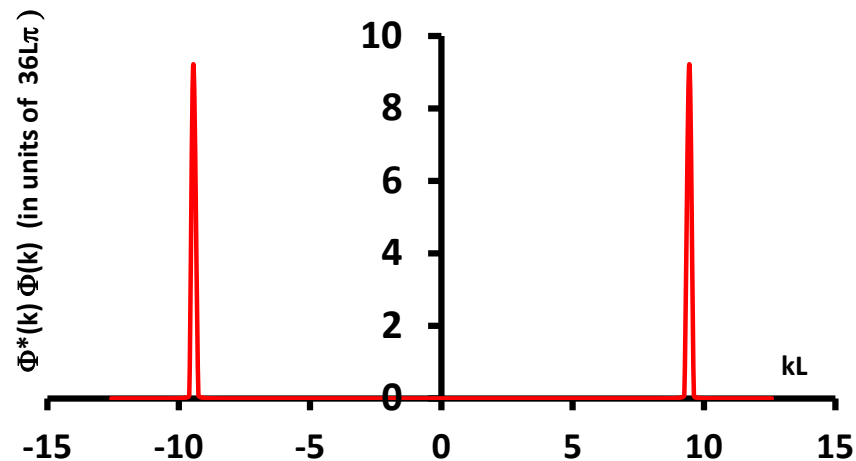
Note: math details included at end of Appendix

# PLOTTING the RESULTS (n=3 for infinite square well)

Eigenfunction of momentum  $\Phi(k)$



Probability  $\Phi^*(k) \Phi(k)$   
(sharp peaks at momentum eigenvalues)



# The Math Details

$$\Phi_3(k) = \frac{1}{\sqrt{2\pi}} \int_{-L/2}^{+L/2} \Psi_3(x) e^{-ikx} dx = \frac{1}{\sqrt{2\pi}} \int_{-L/2}^{+L/2} \Psi_3(x) e^{-ikx} dx$$

$$\Psi_3(x) = \sqrt{\frac{2}{L}} \cos\left(\frac{3\pi x}{L}\right)$$

$$\Phi_3(k) = \frac{1}{\sqrt{2\pi}} \sqrt{\frac{2}{L}} \int_{-L/2}^{+L/2} \cos\left(\frac{3\pi x}{L}\right) e^{-ikx} dx$$

$$= \frac{1}{\sqrt{2\pi}} \sqrt{\frac{2}{L}} \left[ e^{-ikx} \left( \frac{-\left(\frac{3\pi}{L}\right) \sin\left(\frac{3\pi}{L} x\right) + ik \cos\left(\frac{3\pi}{L} x\right)}{k^2 - \left(\frac{3\pi}{L}\right)^2} \right) \right]_{-L/2}^{+L/2}$$

$$= \frac{1}{\sqrt{2\pi}} \sqrt{\frac{2}{L}} \left[ e^{-ikL/2} \left( \frac{-\left(\frac{3\pi}{L}\right) \sin\left(\frac{3\pi}{L} \frac{L}{2}\right) + ik \cos\left(\frac{3\pi}{L} \frac{L}{2}\right)}{k^2 - \left(\frac{3\pi}{L}\right)^2} \right) - e^{-ik(-L/2)} \left( \frac{-\left(\frac{3\pi}{L}\right) \sin\left(\frac{3\pi}{L} \left(-\frac{L}{2}\right)\right) + ik \cos\left(\frac{3\pi}{L} \left(-\frac{L}{2}\right)\right)}{k^2 - \left(\frac{3\pi}{L}\right)^2} \right) \right]$$

$$\begin{aligned}
&= \frac{1}{\sqrt{2\pi}} \sqrt{\frac{2}{L}} \left[ e^{-ikL/2} \left( \frac{-\left(\frac{3\pi}{L}\right)(-1) + ik \times 0}{k^2 - \left(\frac{3\pi}{L}\right)^2} \right) - e^{+ik(L/2)} \left( \frac{-\left(\frac{3\pi}{L}\right)(1) + ik \times 0}{k^2 - \left(\frac{3\pi}{L}\right)^2} \right) \right] \\
&= \frac{1}{\sqrt{2\pi}} \sqrt{\frac{2}{L}} \left[ e^{-ikL/2} \left( \frac{\left(\frac{3\pi}{L}\right)}{k^2 - \left(\frac{3\pi}{L}\right)^2} \right) + e^{+ik(L/2)} \left( \frac{\left(\frac{3\pi}{L}\right)}{k^2 - \left(\frac{3\pi}{L}\right)^2} \right) \right] \\
&= \frac{2}{\sqrt{L\pi}} \left[ \left( \frac{\left(\frac{3\pi}{L}\right)}{k^2 - \left(\frac{3\pi}{L}\right)^2} \right) \cos\left(\frac{kL}{2}\right) \right] \\
&= \frac{2}{\sqrt{L\pi}} \left( \frac{3\pi}{L} \right) \left[ \left( \frac{1}{\left[ k - \left(\frac{3\pi}{L}\right) \right] \left[ k + \left(\frac{3\pi}{L}\right) \right]} \right) \cos\left(\frac{kL}{2}\right) \right] \\
&= \frac{2}{\sqrt{L\pi}} \left( \frac{3\pi}{L} \right) \left[ \left( \frac{L^2}{[kL - 3\pi][kL + 3\pi]} \right) \cos\left(\frac{kL}{2}\right) \right] \\
&= \frac{6\pi L}{\sqrt{L\pi}} \left[ \left( \frac{1}{[kL - 3\pi][kL + 3\pi]} \right) \cos\left(\frac{kL}{2}\right) \right] \\
&= 6\sqrt{L\pi} \left[ \left( \frac{1}{[kL - 3\pi][kL + 3\pi]} \right) \cos\left(\frac{kL}{2}\right) \right]
\end{aligned}$$