

Optical Phonon Scattering (ODP)

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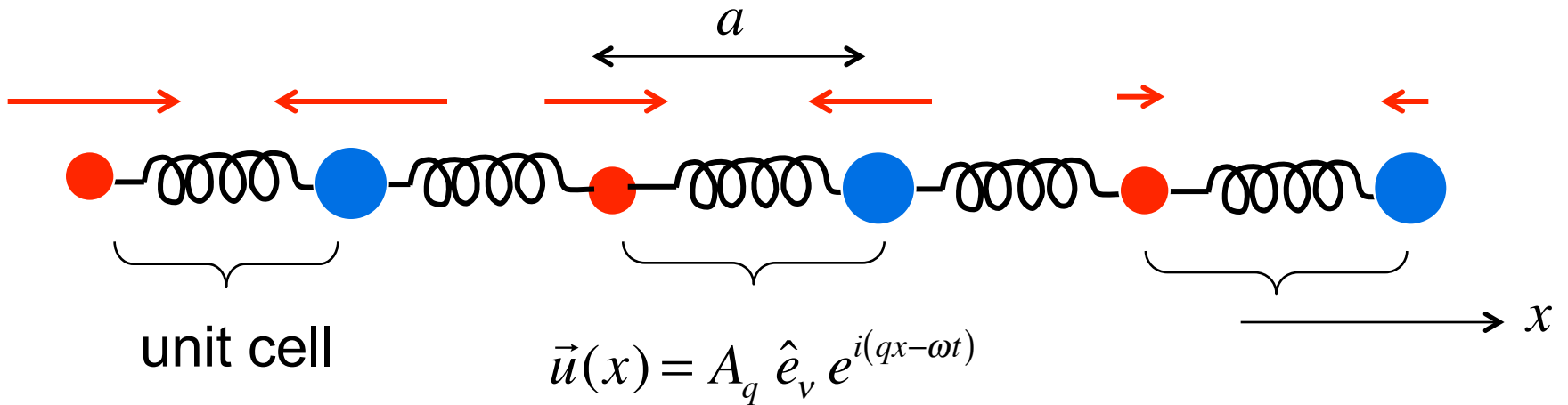
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Outline

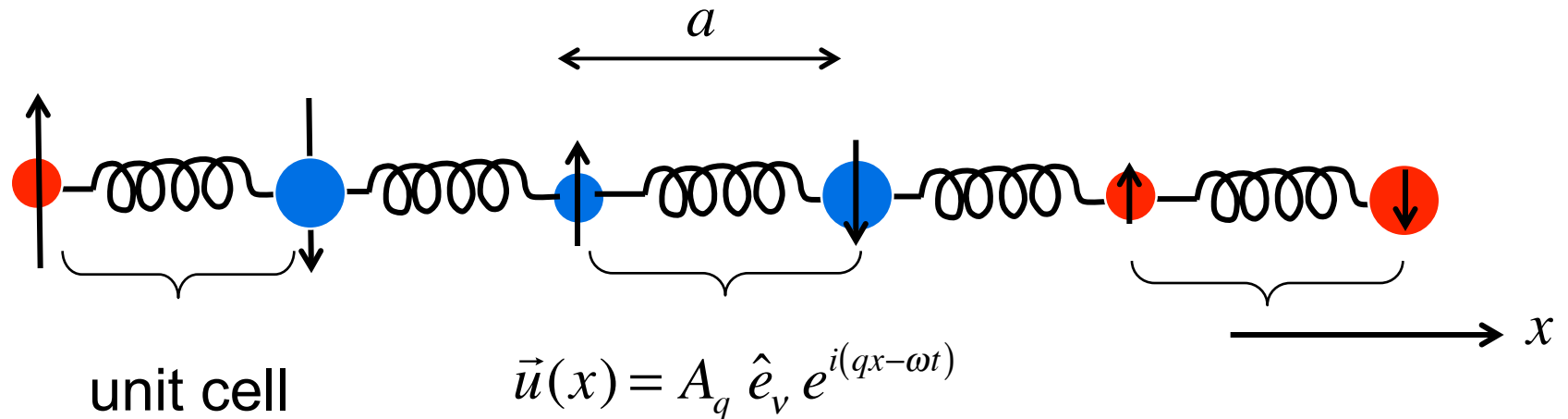
- 1) Quick review of optical phonons
- 2) Electron-phonon coupling (ODP)
- 3) Phonon amplitude
- 4) Matrix element and transition rate
- 5) Momentum and energy conservation
- 6) Scattering time
- 7) Summary

LO phonons



the two atoms in a unit cell oscillate out of phase

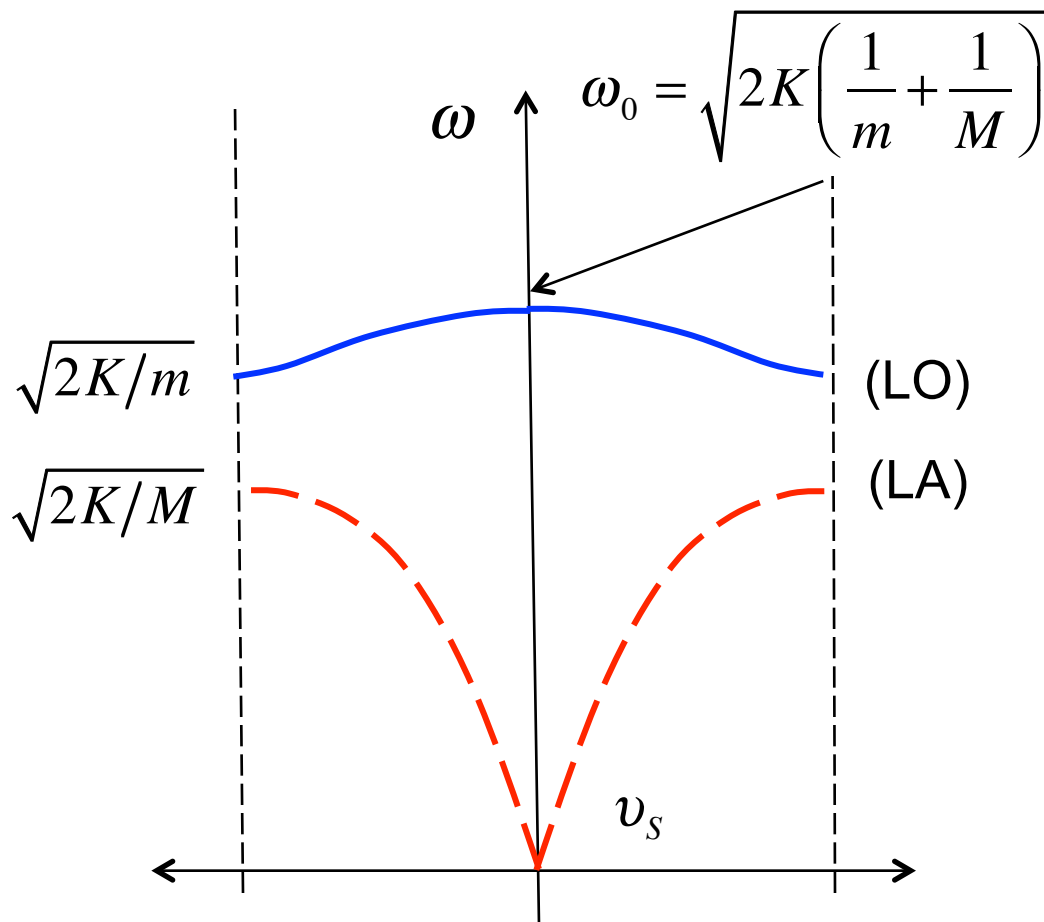
TO phonons



The two atoms in a unit cell oscillate out of phase

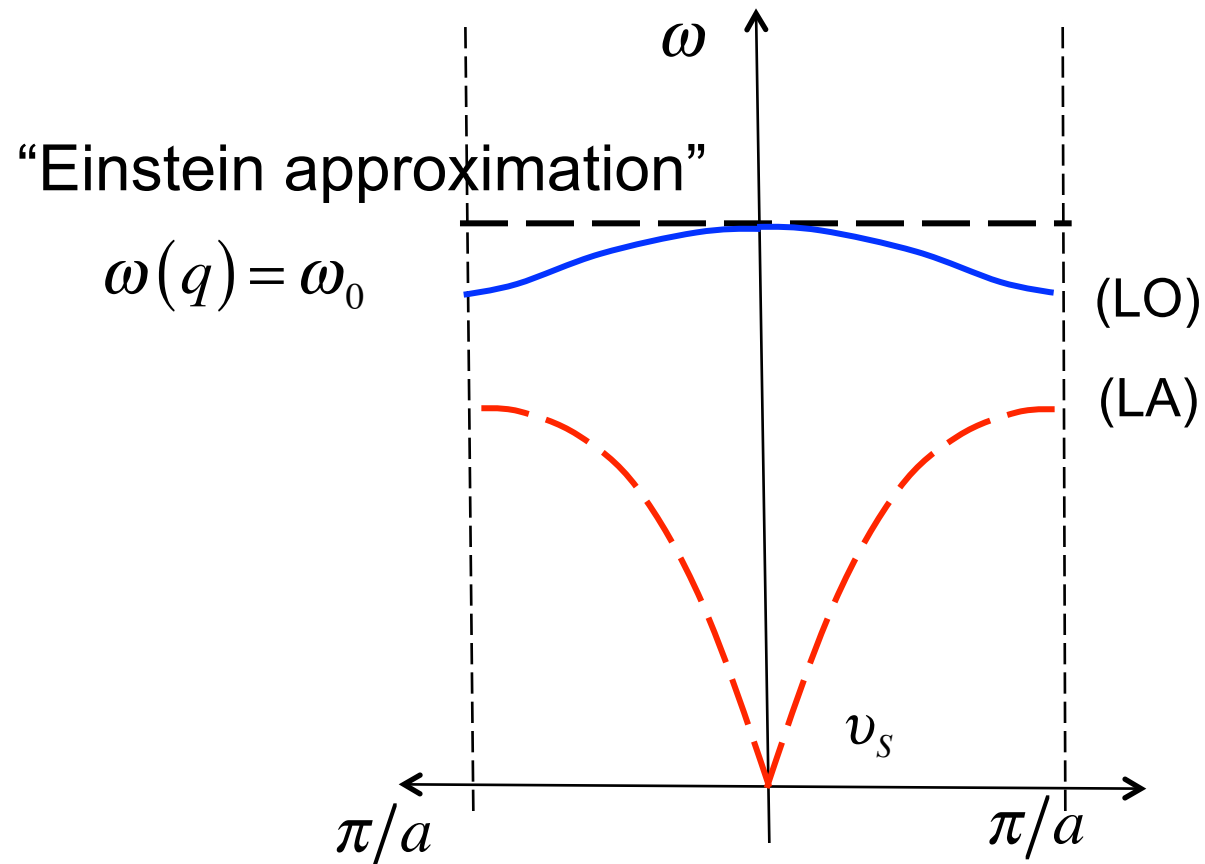
Coupling to electrons is not as strong as for LO phonons.

LO phonons



Optical phonons have **low group velocity**.

Simplified dispersion for optical phonons



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Electron-phonon coupling (LA)

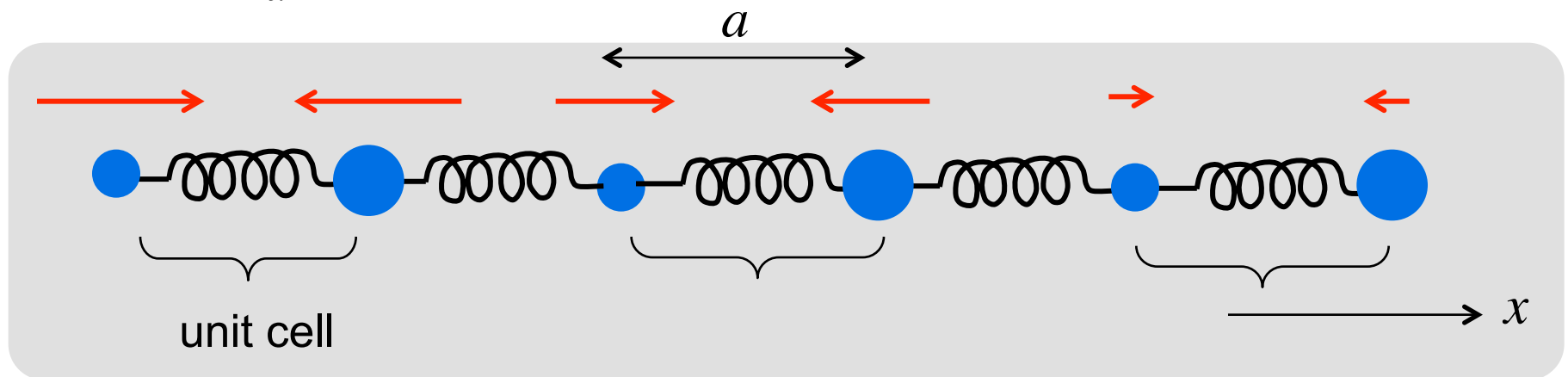
The bandgap depends on lattice constant: $\delta E_G = D \frac{\delta a}{a}$

$$\delta E_C = D_C \frac{\delta a}{a} \quad \text{“deformation potential”}$$

Optical deformation potential (ODP) scattering

Electron-phonon coupling (LO)

$$\delta E_c = D_c \frac{\delta a}{a} \quad \text{deformation potential}$$



$$\delta a(x) = u_q(x)$$

$$u_q(x, t) = A_q e^{\pm i(qx - \omega t)}$$

$$U_S \approx D_O u_q = K_q u_q$$

$$|K_q|^2 = D_O^2$$

“optical deformation potential scattering (ODP)”

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Phonon amplitude

$$S(\vec{p}, \vec{p}') = \frac{2\pi}{\hbar} |K_q|^2 |A_q^{a,e}|^2 \delta(p' - \vec{p} \mp \hbar\vec{q}) \delta(E' - E \mp \Delta E)$$

$$u_q(\vec{r}, t) = A_q e^{\pm i(\vec{q}\cdot\vec{r} - \omega t)}$$

$$|A_q^{a,e}|^2 \rightarrow \frac{\hbar}{2\rho\Omega\omega} \left(N_\omega + \frac{1}{2} \mp \frac{1}{2} \right)$$

$$S^{abs}(\vec{p}, \vec{p}') \propto N^\omega$$

ABS

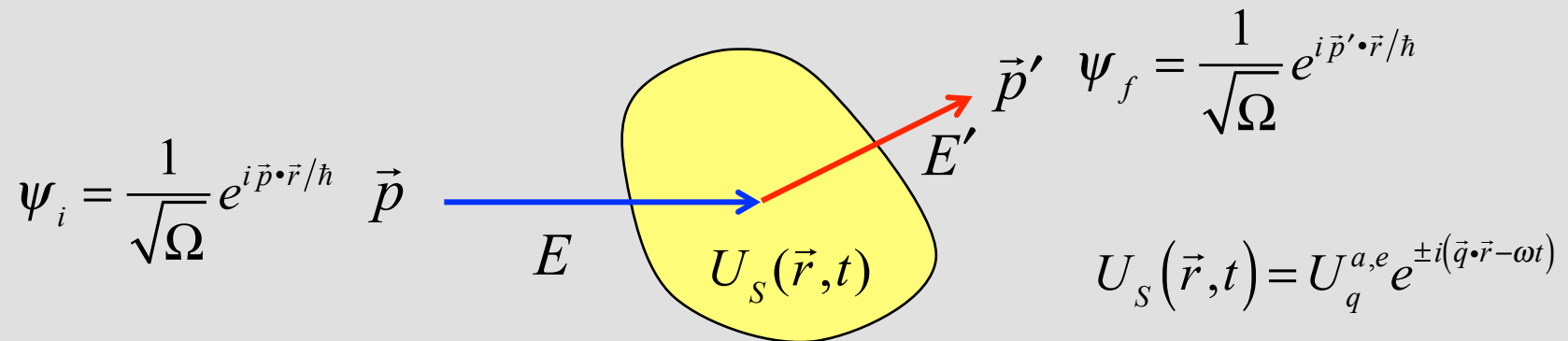
$$S^{ems}(\vec{p}, \vec{p}') \propto N_\omega + 1$$

EMS

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ODP Phonon scattering



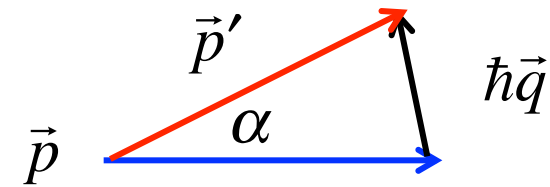
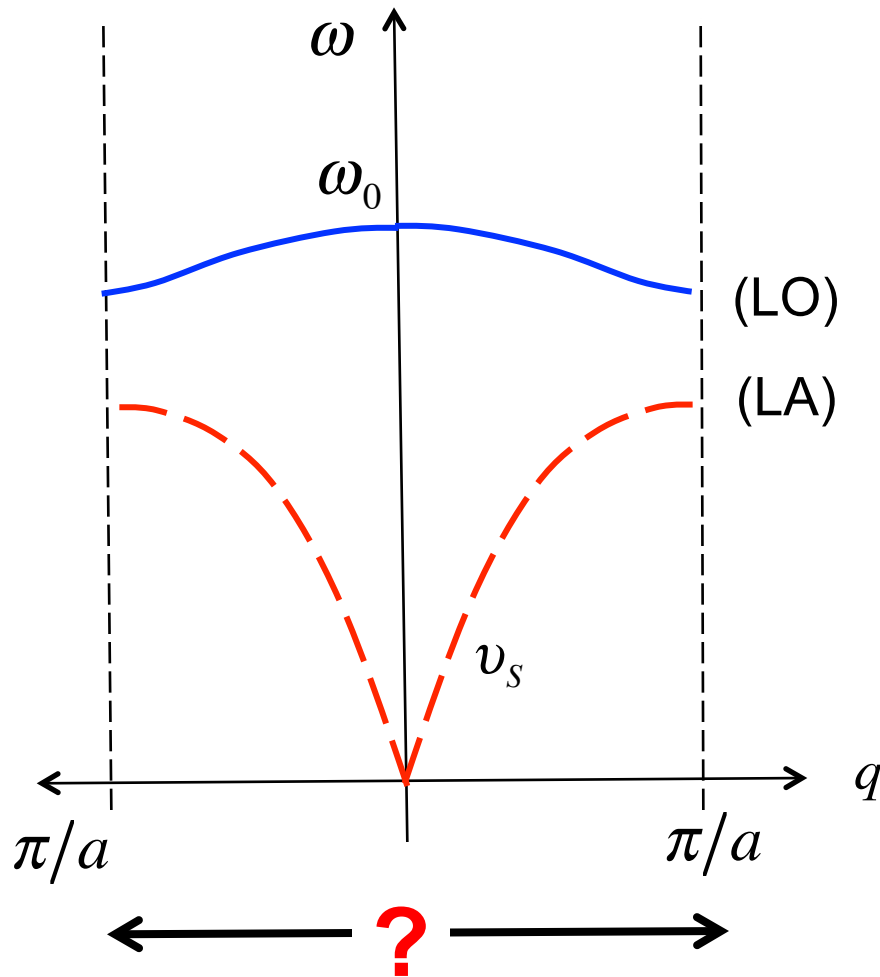
$$S(\vec{p}, \vec{p}') = \frac{2\pi}{\hbar} |U_q^{a,e}|^2 \delta(p' - \vec{p} \mp \hbar\vec{q}) \delta(E' - E \mp \hbar\omega_0)$$

$$|U_q^{a,e}|^2 = |K_q^{a,e}|^2 |A_q^{a,e}|^2 = D_0^2 \left(N_\omega + \frac{1}{2} \mp \frac{1}{2} \right) \frac{\hbar}{2\Omega\rho\omega_0}$$

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Which phonons scatter?



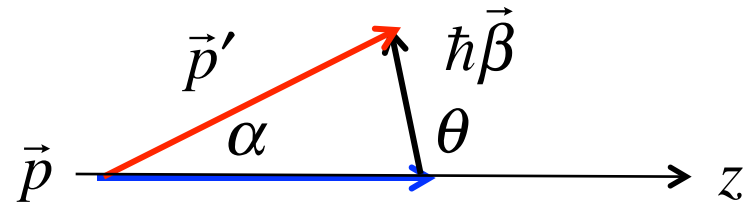
$$E' = E \pm \hbar\omega_0$$

$$\vec{p}' = \vec{p} \pm \hbar\vec{q}$$

Momentum conservation

$$\frac{p'^2}{2m^*} = \frac{p^2}{2m^*} \pm \hbar\omega_\beta$$

$$\vec{p}' = \vec{p} \pm \hbar\vec{\beta}$$



$$\hbar q = 2p \left[\mp \cos\theta \pm \frac{\omega_0}{qv} \right]$$

top sign: ABS

$$\hbar q = 2p \left[\mp \cos\theta \pm \frac{\omega_0}{qv} \right]$$

$$\hbar q = p \left[-\cos\theta + \sqrt{\cos^2\theta \pm \frac{\hbar\omega_0}{E}} \right]$$

$$\hbar q_{\max} = p \left[1 + \sqrt{1 \pm \frac{\hbar\omega_0}{E}} \right]$$

(inelastic)

Any electron can absorb an optical phonon, but to emit an optical phonon, its energy must be > optical phonon energy.

Summary

- 1) Energy-momentum conservation restrict the range of phonon wavevectors that can scatter electrons.

$$q_{\min} < q < q_{\max}$$

- 2) Intra-valley optical phonon scattering is inelastic at room temperature and involves phonons near the zone center.

$$\hbar q_{\min} = p \left[\mp 1 \pm \left(1 \pm \hbar \omega_0 / E_k \right)^{1/2} \right] \quad (\text{ABS/EMS})$$

$$\hbar q_{\max} = p \left[1 + \left(1 \pm \hbar \omega_0 / E_k \right)^{1/2} \right] \quad (\text{For EMS}) \quad E_k > \hbar \omega_0$$

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Scattering times

$$S(\vec{p}, \vec{p}') = \frac{2\pi}{\hbar} |K_q^{a,e}|^2 |A_q^{a,e}|^2 \delta(p' - \vec{p} \mp \hbar \vec{q}) \delta(E' - E \mp \hbar \omega_0)$$

$$|K_q^{a,e}| = D_o^2 \quad |A_q^{a,e}|^2 = \left(N_\omega + \frac{1}{2} \mp \frac{1}{2} \right) \frac{\hbar}{2\Omega \rho \omega_0}$$

$$\frac{1}{\tau(\vec{p})} = \sum_{\vec{p}', \uparrow} S(\vec{p}, \vec{p}') \quad \frac{1}{\tau_m(\vec{p})} = \sum_{\vec{p}', \uparrow} S(\vec{p}, \vec{p}') \frac{\Delta p_z}{p_{z0}}$$

Simple approach

$$S(\vec{p}, \vec{p}') = \frac{2\pi}{\hbar} |K_q^{a,e}|^2 |A_q^{a,e}|^2 \delta(p' - \vec{p} \mp \hbar\vec{q}) \delta(E' - E \mp \hbar\omega_0)$$

If there is no q dependence here, then the scattering is **isotropic**, and we don't need to explicitly account for momentum conservation.

$$\frac{1}{\tau} = \frac{1}{\tau_m}$$

Simple approach

$$S(\vec{p}, \vec{p}') = \frac{2\pi}{\hbar} |K_q^{a,e}|^2 |A_q^{a,e}|^2 \delta(E' - E \mp \hbar\omega_0)$$

$$|K_q^{a,e}| = D_0^2 \quad |A_q^{a,e}|^2 = \left(N_\omega + \frac{1}{2} \mp \frac{1}{2} \right) \frac{\hbar}{2\Omega\rho\omega_0} \quad N_\omega = \frac{1}{e^{\hbar\omega_0/k_B T} - 1}$$

$$S(\vec{p}, \vec{p}') = \frac{2\pi}{\hbar} D_0^2 \frac{\hbar}{2\Omega\rho\omega_0} \left(N_\omega + \frac{1}{2} \mp \frac{1}{2} \right) \delta(E' - E \mp \hbar\omega_0)$$

ODP scattering

$$S(\vec{p}, \vec{p}') = \frac{2\pi}{\hbar} U_{op}^{a,e} \frac{1}{\Omega} \delta(E' - E \mp \hbar\omega_0) \quad U_{op}^{a,e} = D_0^2 \frac{\hbar}{2\Omega\rho\omega_0} \left(N_\omega + \frac{1}{2} \mp \frac{1}{2} \right)$$

$$\frac{1}{\tau} = \frac{1}{\tau_m} = \sum_{\vec{p}' \uparrow} S(\vec{p}, \vec{p}')$$

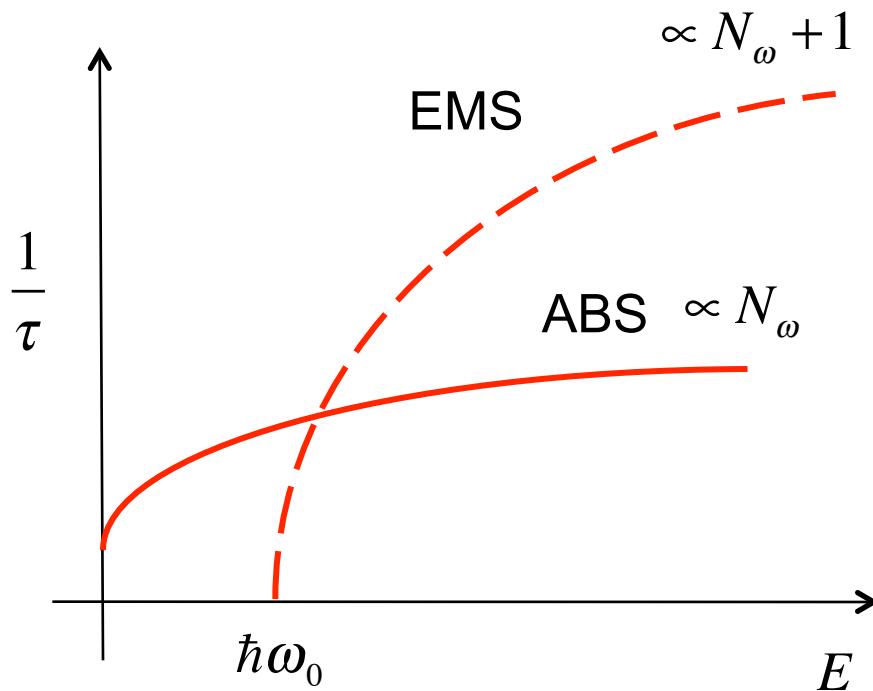
$$\frac{1}{\tau_m(E)} = \frac{2\pi}{\hbar} U_{op}^{a,e} \frac{D_{3D}(E \pm \hbar\omega_0)}{2}$$

$$\frac{1}{\tau_m^{abs}} \propto N_{\omega_0}$$

$$\frac{1}{\tau_m^{ems}} \propto N_{\omega_0} + 1$$

ODP scattering: Summary

$$\frac{1}{\tau_m(E)} = \frac{\pi D_o^2}{\rho \omega_0} \left(N_\omega + \frac{1}{2} \mp \frac{1}{2} \right) \frac{D_{3D}(E \pm \hbar \omega_0)}{2}$$

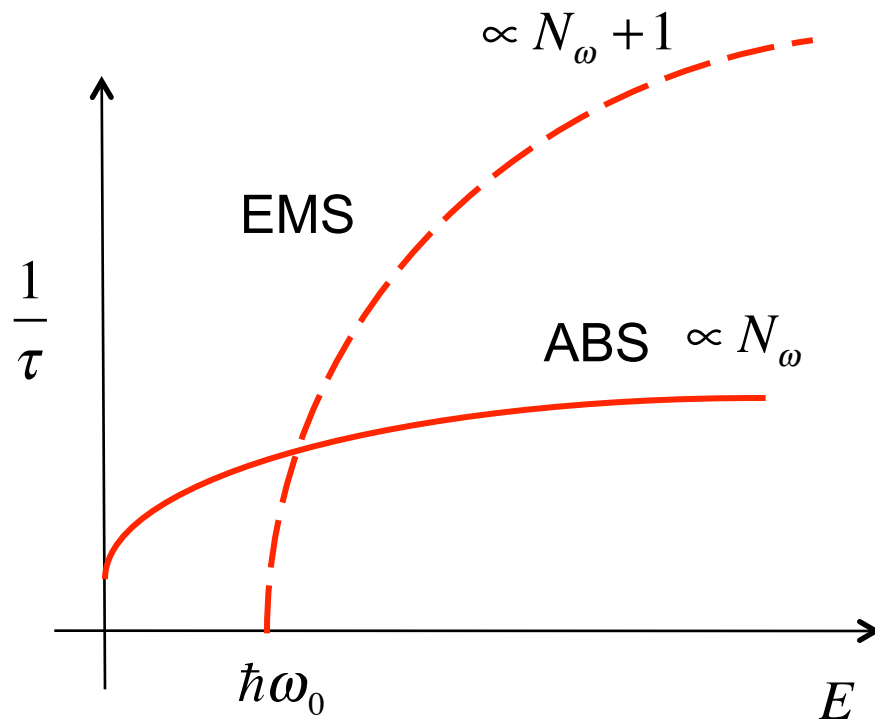


Power law scattering?

$$\tau \neq \tau_0 (E/k_B T)^s$$

Absorption vs. emission

$$\frac{1}{\tau_m} = \frac{\pi D_0^2}{\rho \omega_0} \left(N_\omega + \frac{1}{2} \mp \frac{1}{2} \right) \frac{D_{3D}(E)}{2}$$



Compare N and N+1
for GaAs at 300K:

$$\hbar\omega_0 = 0.033 \text{ eV}$$

$$N_{\omega_0} = \frac{1}{e^{0.033/0.026} - 1} = 0.4$$

$$N_{\omega_0} + 1 = 1.4$$

Questions?

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