

# The Coupled Current Equations (Again)

Mark Lundstrom

Electrical and Computer Engineering  
Purdue University  
West Lafayette, IN USA

# The balance equation prescription

$$\frac{\partial n_\phi}{\partial t} = -\nabla \cdot \vec{F}_\phi + G_\phi - R_\phi$$

$$n_\phi(\vec{r}, t) = \frac{1}{\Omega} \sum_{\vec{p}} \phi(\vec{p}) f(\vec{r}, \vec{p}, t)$$

$$G_\phi = -q\vec{E} \cdot \left\{ \frac{1}{\Omega} \sum_{\vec{p}} \vec{\nabla}_p \phi f \right\}$$

$$\vec{F}_\phi \equiv \frac{1}{\Omega} \sum_{\vec{p}} \phi(\vec{p}) \vec{v} f(\vec{r}, \vec{p}, t)$$

$$R_\phi \equiv \frac{n_\phi - n_\phi^0}{\langle \tau_\phi \rangle}$$

# The balance equations

$$\frac{\partial n}{\partial t} = -\nabla \cdot \frac{\vec{J}_n}{(-q)}$$

$$\frac{\partial \vec{J}_n}{\partial t} = \frac{2q}{m^*} \nabla W_{xx} + \frac{nq^2}{m^*} \vec{\mathcal{E}} - \frac{\vec{J}_n}{\langle \tau_m \rangle}$$

$$\frac{\partial W}{\partial t} = -\nabla \vec{F}_W - \vec{J}_n \vec{\mathcal{E}} - \frac{W - W_0}{\langle \tau_E \rangle}$$

$$\frac{\partial \vec{F}_W}{\partial t} = \dots$$

← Terminate here  
(find a way to estimate  
 $W$  in terms of known  
quantities).

# The DD equation

$$\frac{\partial n}{\partial t} = -\nabla \cdot \frac{\vec{J}_n}{(-q)}$$

$$\frac{\partial \vec{J}_n}{\partial t} = \frac{2q}{m^*} \nabla W_{xx} + \frac{nq^2}{m^*} \vec{\mathcal{E}} - \frac{\vec{J}_n}{\langle \tau_m \rangle} \quad \leftarrow \text{Terminate here.}$$

$$J_{nx} = nq\mu_n \mathcal{E}_x + 2\mu_n \frac{\partial W_{xx}}{\partial x} \quad J_{nx} = nq\mu_n \mathcal{E}_x + \frac{2}{3}\mu_n \frac{\partial W}{\partial x} \quad W = \frac{3}{2}nk_B T_e$$

$$T_e = T_L = T \quad J_{nx} = nq\mu_n \mathcal{E}_x + qD_n \frac{\partial n}{\partial x} \quad \frac{D_n}{\mu_n} = \frac{k_B T}{q}$$

# Coupled current equations

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Recall that we derived the coupled current equations for the flow of charge and heat by solving the BTE in the Relaxation Time Approximation.

## Coupled current equations

$$J = \sigma \mathcal{E} - \sigma S dT/dx$$

$$J_Q = T \sigma S \mathcal{E} - \kappa_0 dT/dx$$

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$$\mathcal{E} = \rho J + S \frac{dT}{dx}$$

$$J_Q = \pi J - \kappa_e \frac{dT}{dx}$$

(diffusive transport)

$$\sigma = \int \sigma'(E) dE$$

$$\sigma'(E) = \frac{2q^2}{h} \lambda(E) \frac{M(E)}{A} \left( -\frac{\partial f_0}{\partial E} \right)$$

$$S = -\frac{k_B}{q} \int \left( \frac{E - E_F}{k_B T} \right) \sigma'(E) dE \bigg/ \int \sigma'(E) dE$$

$$\pi = TS$$

$$\kappa_0 = T \left( \frac{k_B}{q} \right)^2 \int \left( \frac{E - E_F}{k_B T} \right)^2 \sigma'(E) dE$$

$$\kappa_e = \kappa_0 - \pi S \sigma$$

## Balance equation prescription

$$\frac{1}{\Omega} \sum_{\vec{p}} \phi(\vec{p}) \left\{ \frac{\partial f}{\partial t} + \vec{v} \cdot \nabla_r f - q \vec{E} \cdot \nabla_p f = \hat{C} f \right\}$$

$$\frac{\partial n_\phi}{\partial t} = -\nabla \cdot \vec{F}_\phi + G_\phi - R_\phi$$

$$n_\phi(\vec{r}, t) = \frac{1}{\Omega} \sum_{\vec{p}} \phi(\vec{p}) f(\vec{r}, \vec{p}, t)$$

$$G_\phi = -q \vec{E} \cdot \left\{ \frac{1}{\Omega} \sum_{\vec{p}} \vec{\nabla}_p \phi f \right\}$$

$$\vec{F}_\phi \equiv \frac{1}{\Omega} \sum_{\vec{p}} \phi(\vec{p}) \vec{v} f(\vec{r}, \vec{p}, t)$$

$$R_\phi \equiv \frac{n_\phi - n_\phi^0}{\langle \tau_\phi \rangle}$$

# Current equations

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$$\frac{\partial n_\phi}{\partial t} = -\nabla \cdot \vec{F}_\phi + G_\phi - R_\phi$$

Charge current:

$$\phi(\vec{p}) = (-q)v_x$$

$$n_\phi(\vec{r}, t) = J_{nx}(\vec{r}, t) = \sum_{\vec{p}} \phi(\vec{p}) f(\vec{r}, \vec{p}, t)$$

Heat current:

$$\phi(\vec{p}) = (E(\vec{p}) + E_C - F_n)v_x$$

$$n_\phi(\vec{r}, t) = J_{Qx}(\vec{r}, t) = \sum_{\vec{p}} \phi(\vec{p}) f(\vec{r}, \vec{p}, t)$$



# Charge current

$$J_{nx} = nq\mu_n \mathcal{E}_x + \frac{2}{3} \mu_n \frac{\partial W}{\partial x}$$

(assumes  $W_{xx} = W/3$ )

Now assume:  $W = \frac{3}{2} nk_B T$

$$J_{nx} = nq\mu_n \mathcal{E}_x + \frac{2}{3} \mu_n \frac{\partial \left( \frac{3}{2} nk_B T \right)}{\partial x}$$

$$J_{nx} = nq\mu_n \mathcal{E}_x + \mu_n \frac{\partial (nk_B T)}{\partial x}$$

## Seebeck coefficient

$$J_{nx} = nq\mu_n \mathcal{E}_x + \mu_n \frac{d}{dx} \{ nk_B T \}$$

$$\mu_n \frac{d}{dx} (nk_B T) = \mu_n \left[ k_B T \frac{\partial n}{\partial x} + nk_B \frac{dT}{dx} + k_B T \frac{\partial n}{\partial T} \frac{dT}{dx} \right]$$

$$\mu_n \frac{d}{dx} (nk_B T) = k_B T \mu_n \frac{\partial n}{\partial x} + nq\mu_n \left( \frac{k_B}{q} \right) \left\{ 1 + T \frac{1}{n} \frac{\partial n}{\partial T} \right\} \frac{dT}{dx}$$

## Seebeck coefficient

$$\mu_n \frac{d}{dx}(nk_B T) = k_B T \mu_n \frac{\partial n}{\partial x} + nq\mu_n \left( \frac{k_B}{q} \right) \left\{ 1 + T \frac{1}{n} \frac{\partial n}{\partial T} \right\} \frac{dT}{dx}$$

$$n = N_C e^{(F_n - E_C)/k_B T} \quad N_C = \frac{1}{4} \left( \frac{2m^* k_B T}{\pi \hbar^2} \right)^{3/2}$$

$$\mu_n \frac{d}{dx}(nk_B T) = k_B T \mu_n \frac{\partial n}{\partial x} + nq\mu_n \left( \frac{k_B}{q} \right) \left\{ \frac{3}{2} - \frac{F_n - E_C}{k_B T} \right\} \frac{dT}{dx}$$

# Seebeck coefficient

$$J_{nx} = nq\mu_n \mathcal{E}_x + \mu_n \frac{d}{dx} \{nk_B T\}$$

$$J_{nx} = nq\mu_n \mathcal{E}_x + qD_n \frac{dn}{dx} + nq\mu_n \left( \frac{k_B}{q} \right) \left\{ \frac{E_C - F_n}{k_B T} + \frac{3}{2} \right\} \frac{dT}{dx}$$

# Seebeck coefficient

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$$J_{nx} = nq\mu_n \mathcal{E}_x - \sigma_n S_n \frac{dT}{dx}$$

$$S_n = \left( -\frac{k_B}{q} \right) \left\{ \frac{E_C - F_n}{k_B T} + \frac{3}{2} \right\}$$

# Comparison

## Landauer or BTE

$$J_{nx} = \sigma_n \mathcal{E}_x - \sigma_n S_n \frac{dT_L}{dx}$$

$$\sigma_n = nq\mu_n$$

$$\mu_n = \frac{v_T \lambda_0 \Gamma(r+2)}{2k_B T / q}$$

$$S_n = -\left(\frac{k_B}{q}\right) \left( \frac{E_C - F_n}{k_B T} + (r+2) \right)$$

## Moment of BTE

$$J_{nx} = \sigma_n \mathcal{E}_x - \sigma_n S_n \frac{dT}{dx}$$

$$\sigma_n = nq\mu_n$$

$$\mu_n = \frac{q \langle \tau_m \rangle}{m^*}$$

$$S_n = -\left(\frac{k_B}{q}\right) \left\{ \frac{E_C - F_n}{k_B T} + \frac{3}{2} \right\}$$

## Landauer / BTE vs. moments

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See *Fundamentals of Carrier Transport*, Sec. 5.5.2 for a discussion of how to make the moment equations consistent with the Landauer/BTE approach.

# Charge and heat flow balance equations

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$$\mathcal{E}_x = \rho J_x + S \frac{dT}{dx}$$

$$J_{Qx} = \pi J_x - \kappa_e \frac{dT}{dx}$$

$$J_x = \sigma \mathcal{E}_x - S\sigma \frac{dT}{dx}$$

$$\phi(\vec{p}) = (-q)\vec{v}(\vec{p})$$

$$J_{Qx} = \pi J_x - \kappa_e \frac{dT}{dx}$$

$$\phi(\vec{p}) = (E - F_n)\vec{v}(\vec{p})$$



# Heat current

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Repeat the process for the heat current.

Is the Kelvin Relation satisfied?

What is the Wiedemann-Franz Law in this approach?

# Questions

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