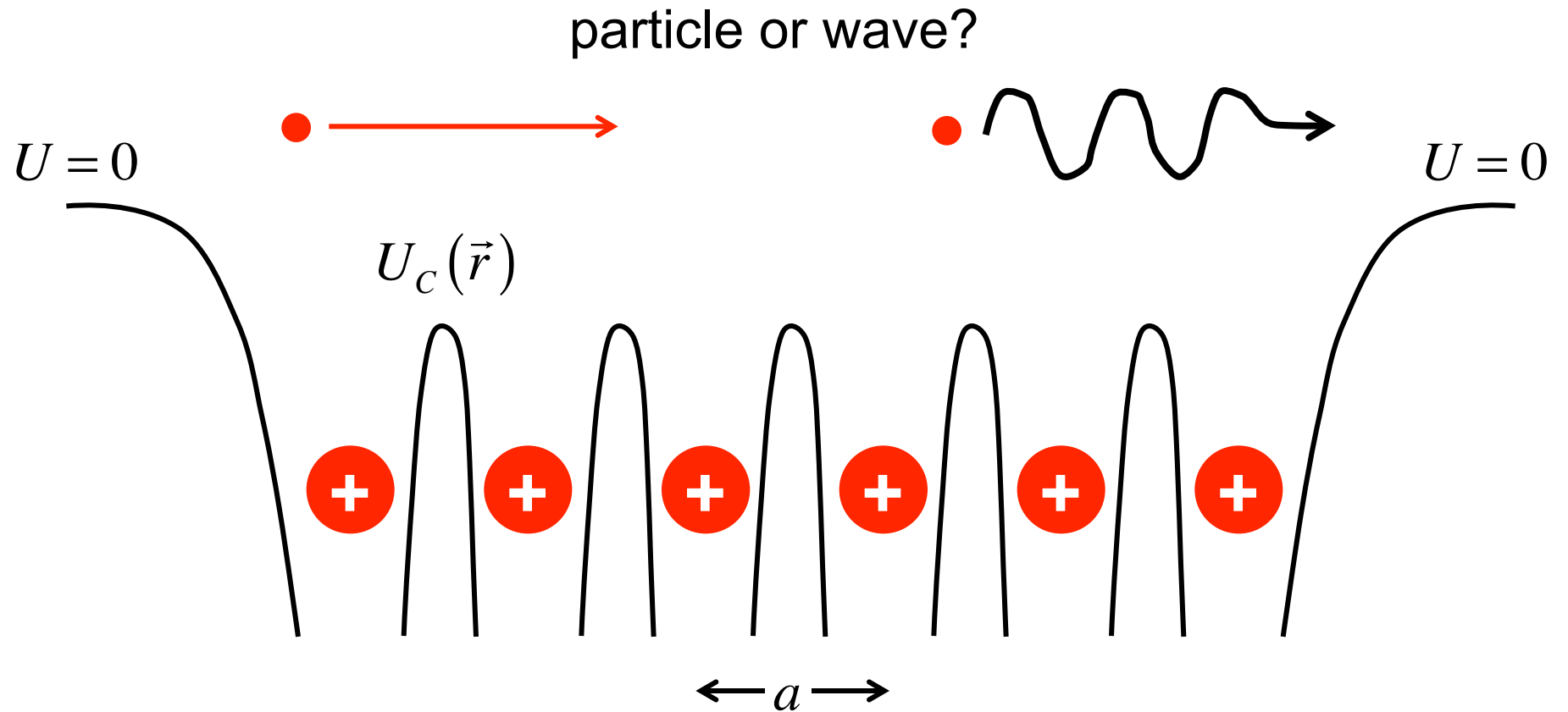


Semiclassical Transport

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Mobile electrons in crystals



De Broglie Wavelength

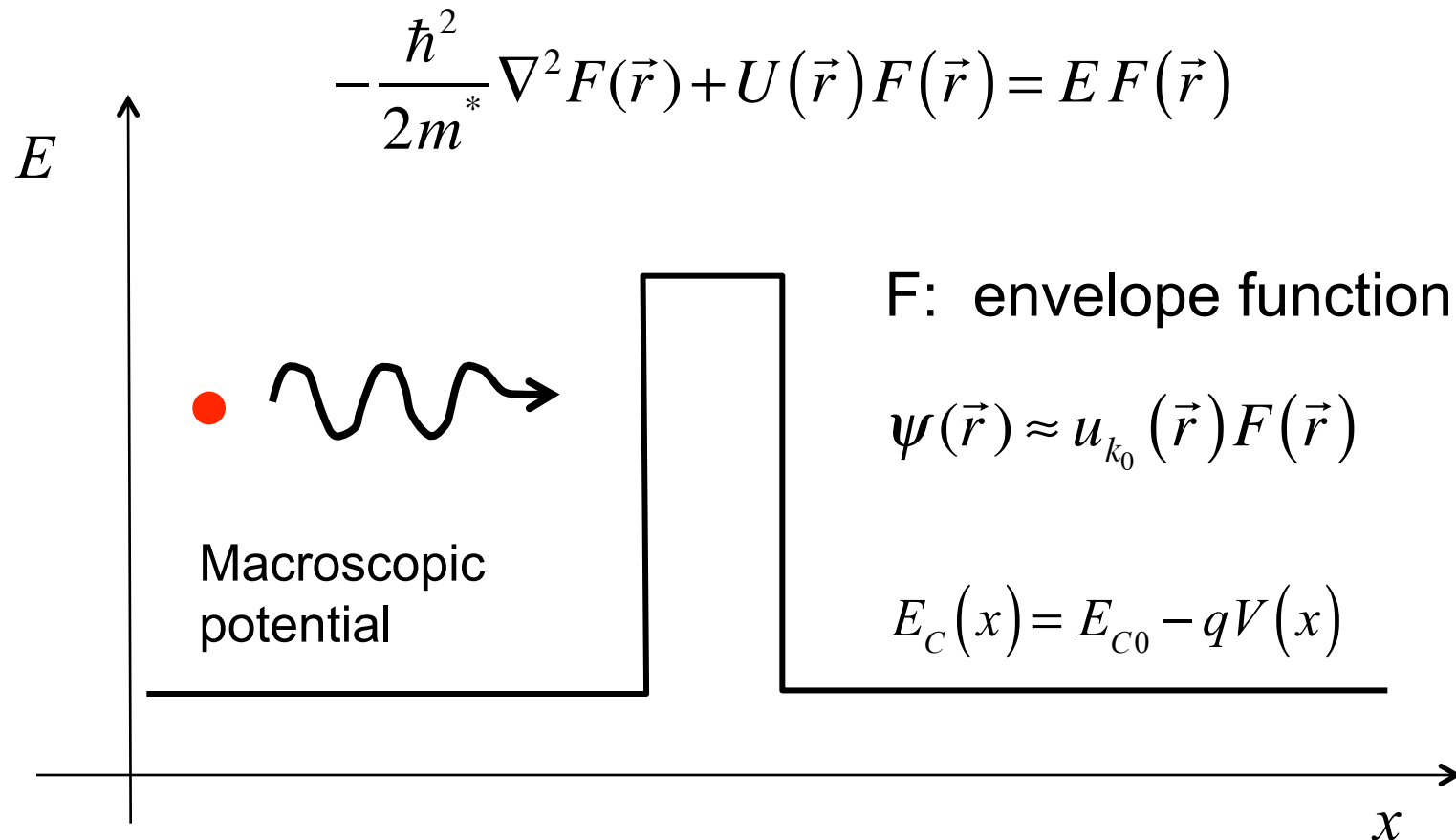
The crystal potential varies on an atomic scale. It gives us the band structure when we solve the Schrodinger equation. Quantum mechanics is necessary.

What is the wavelength of a free (mobile) electron?

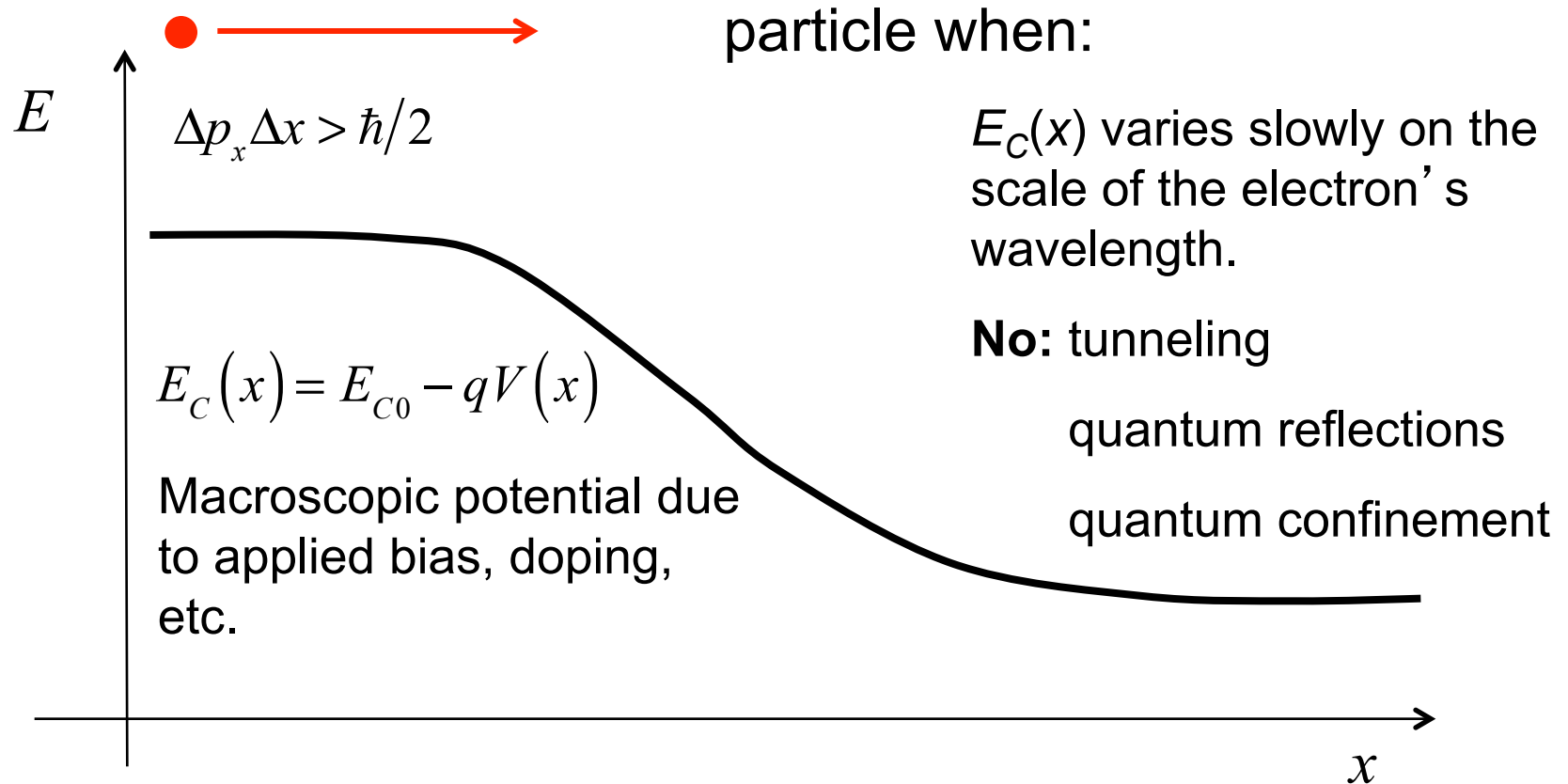
$$p = \hbar k = \hbar \frac{2\pi}{\lambda_B} \qquad E = \frac{p^2}{2m^*} \approx \frac{3}{2} k_B T$$

$$\lambda_B = \sqrt{\frac{4\pi^2 \hbar^2}{3m^* k_B T}} \simeq 10\text{nm (electrons in Si at 300K)}$$

Effective mass wave equation

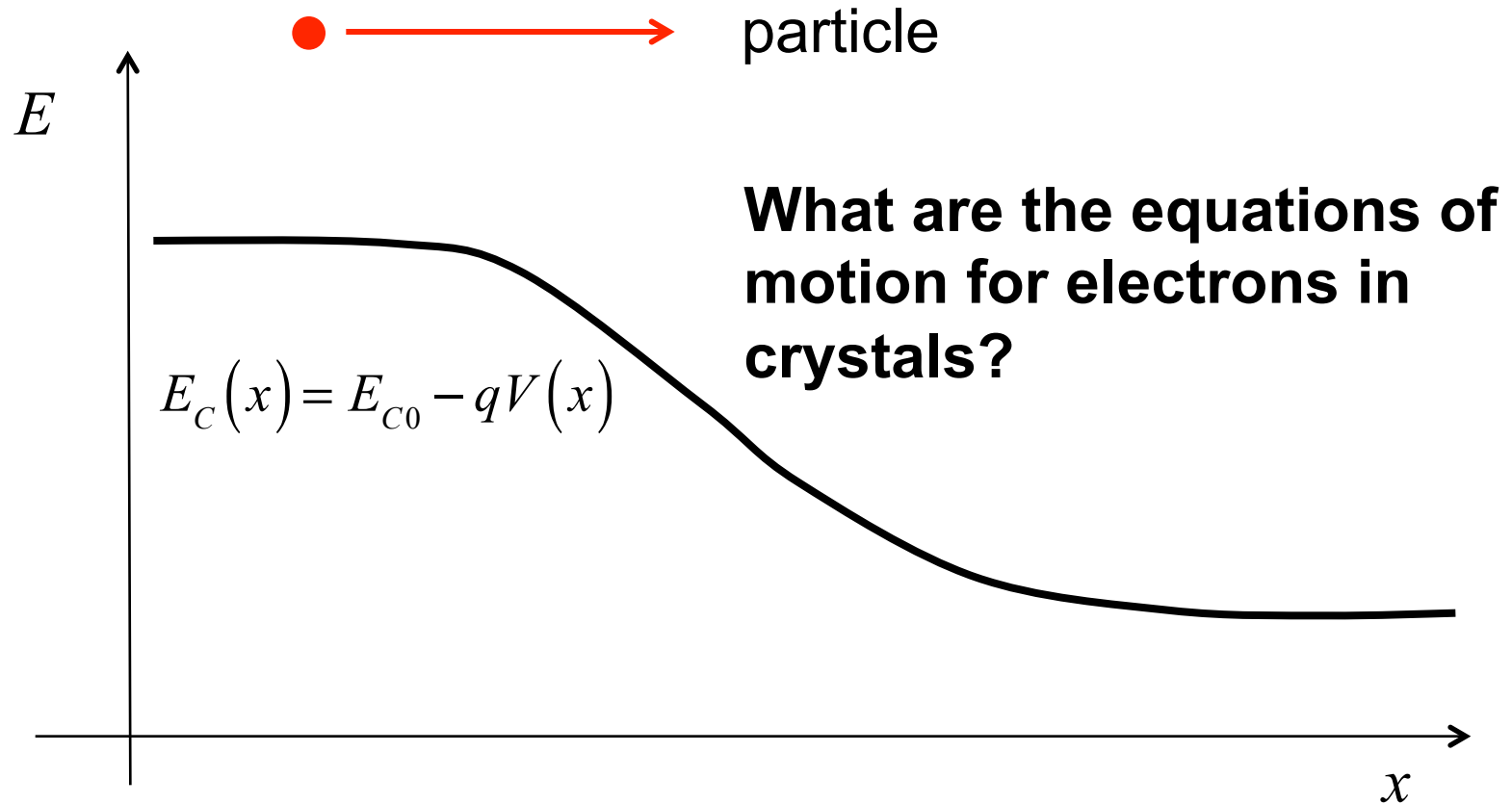


Semi-classical transport

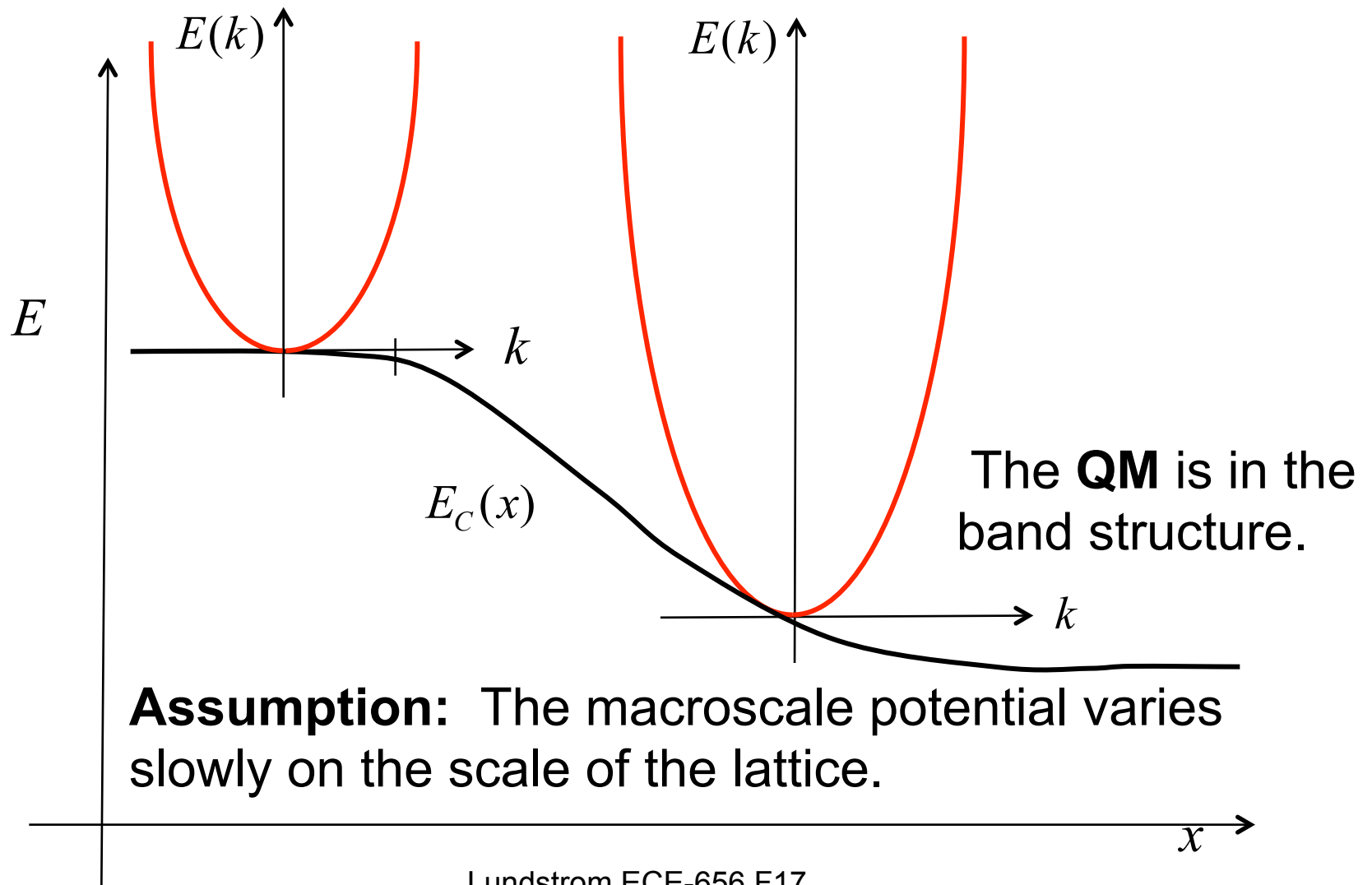


But: Quantum mechanics determines the band structure (e.g. effective mass)

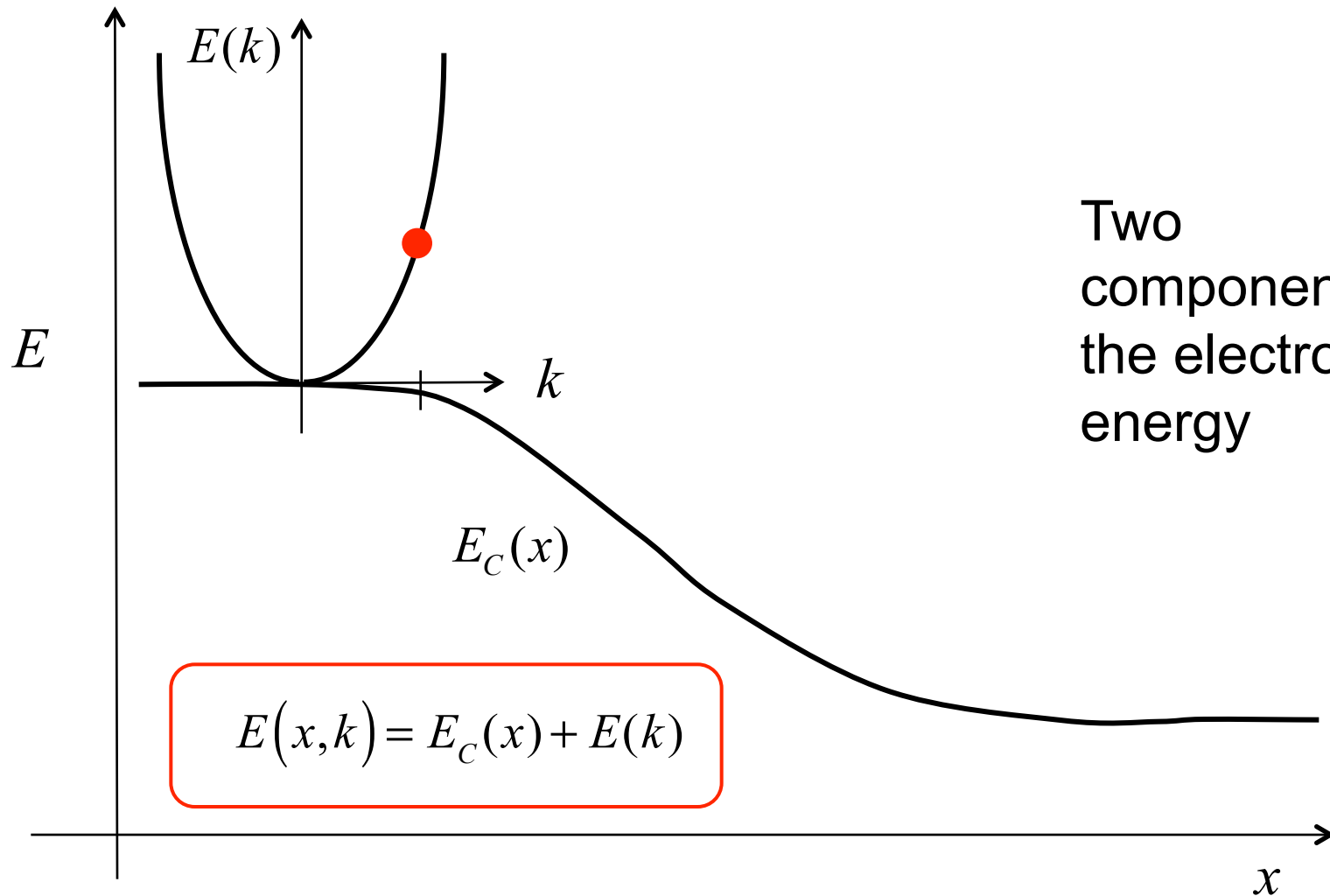
Semi-classical transport



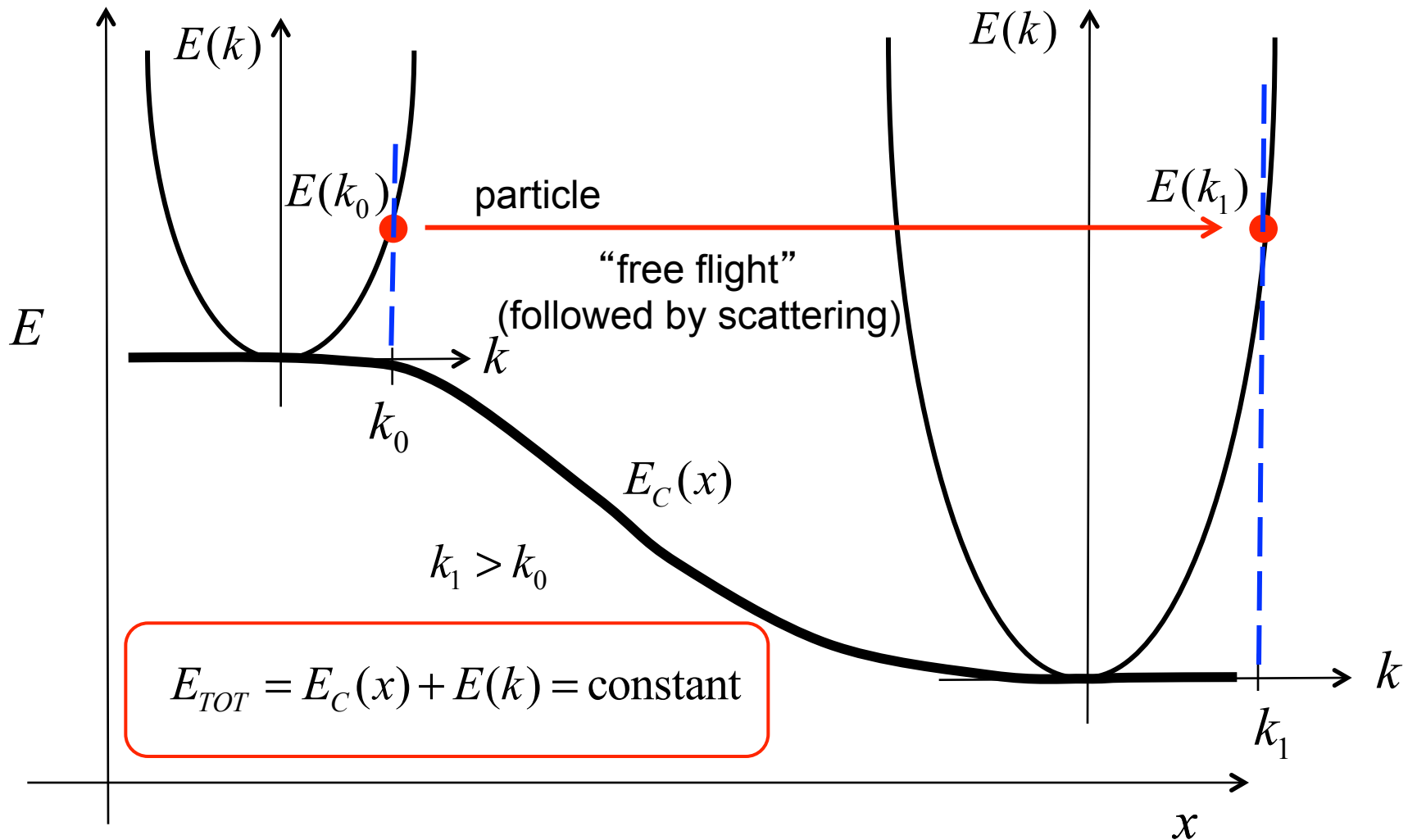
Semi-classical transport



Semi-classical transport



Semi-classical transport



Semi-classical transport

$$E_{TOT} = E_C(x) + E(k)$$

$$\frac{dE_{TOT}(x,k)}{dt} = 0 = \frac{dE_C(x)}{dx} \frac{dx}{dt} + \frac{dE(k)}{dk_x} \frac{dk_x}{dt}$$

$$0 = \frac{dE_C(x)}{dx} v_x + \frac{1}{\hbar} \frac{dE}{dk_x} \frac{d(\hbar k_x)}{dt}$$

$$0 = \frac{dE_C(x)}{dx} v_x + v_x \frac{d(\hbar k_x)}{dt}$$

$$\frac{d(\hbar k_x)}{dt} = F_e = -\frac{dE_C(x)}{dx}$$

$$p_x = \hbar k_x$$

$$\frac{dp_x}{dt} = F_e$$

“Newton’s Law”

Semi-classical transport: k-space

$$\frac{d(\hbar\vec{k})}{dt} = -\nabla_r E_C(\vec{r}) = -q\vec{E}(\vec{r}) \quad \left\{ \frac{d\vec{p}}{dt} = \vec{F}_e \right\}$$

$$\hbar\vec{k}(t) = \hbar\vec{k}(0) + \int_0^t -q\vec{E}(t') dt'$$

equation of motion for
“semi-classical transport”
in k-space.

No band structure (or effective mass) is involved!

Semi-classical transport: real-space

$$\frac{d(\hbar\vec{k})}{dt} = -\nabla_r E_C(\vec{r}) = -q\vec{E}(\vec{r})$$

$$\hbar\vec{k}(t) = \hbar\vec{k}(0) + \int_0^t -q\vec{E}(t') dt'$$

$$\vec{v}_g(t) = \frac{1}{\hbar} \nabla_k E[\vec{k}(t)]$$

$$\vec{r}(t) = \vec{r}(0) + \int_0^t \vec{v}_g(t') dt'$$

} equations of motion for semi-classical transport in real-space.

Motion in real space brings in the band structure!

Semi-classical transport: parabolic bands

Equations of motion for semi-classical transport.

Parabolic energy bands:

$$\frac{d(\hbar\vec{k})}{dt} = -\nabla_r E_C(\vec{r}) = -q\vec{\mathcal{E}}(\vec{r})$$

$$\hbar\vec{k}(t) = \hbar\vec{k}(0) + \int_0^t -q\vec{\mathcal{E}}(t') dt'$$

$$\vec{v}_g(t) = \frac{1}{\hbar} \nabla_k E[\vec{k}(t)]$$

$$\vec{r}(t) = \vec{r}(0) + \int_0^t \vec{v}_g(t') dt'$$

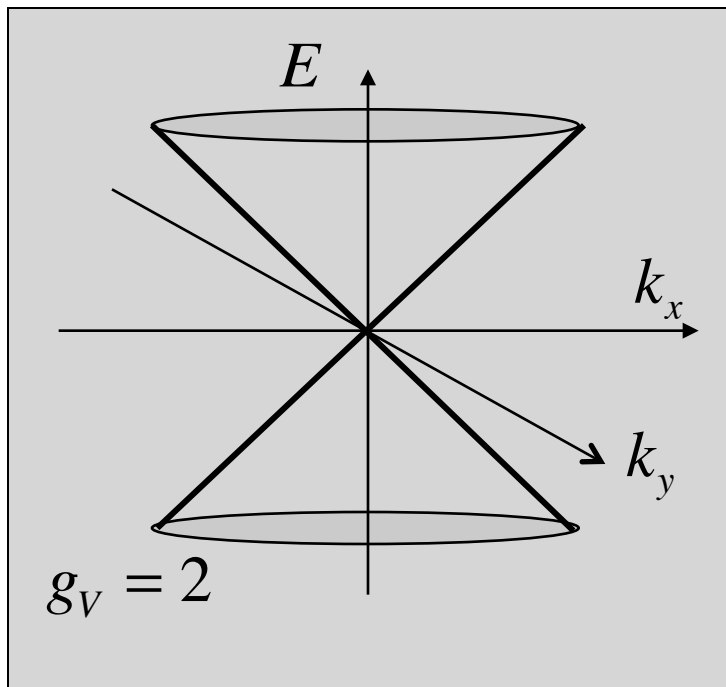
$$E(k) = \frac{\hbar^2 k^2}{2m^*}$$

$$\vec{v}_g(t) = \frac{\hbar\vec{k}}{m^*}$$



Exercise 1: graphene

What are the equations of motion for electrons in graphene?



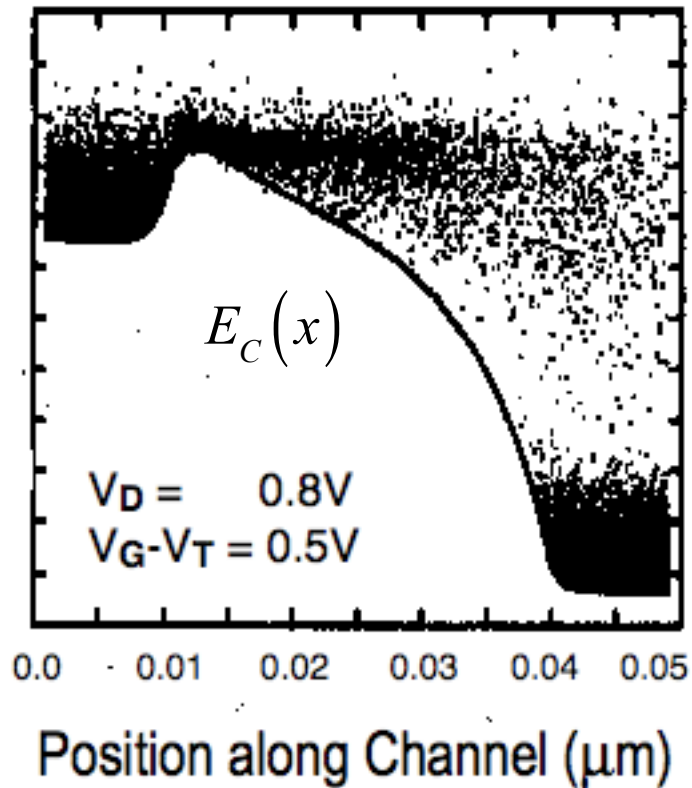
$$E(k) = \pm \hbar v_F \sqrt{k_x^2 + k_y^2} = \pm \hbar v_F k$$

Exercise 2: equations of motion for $m^*(x)$

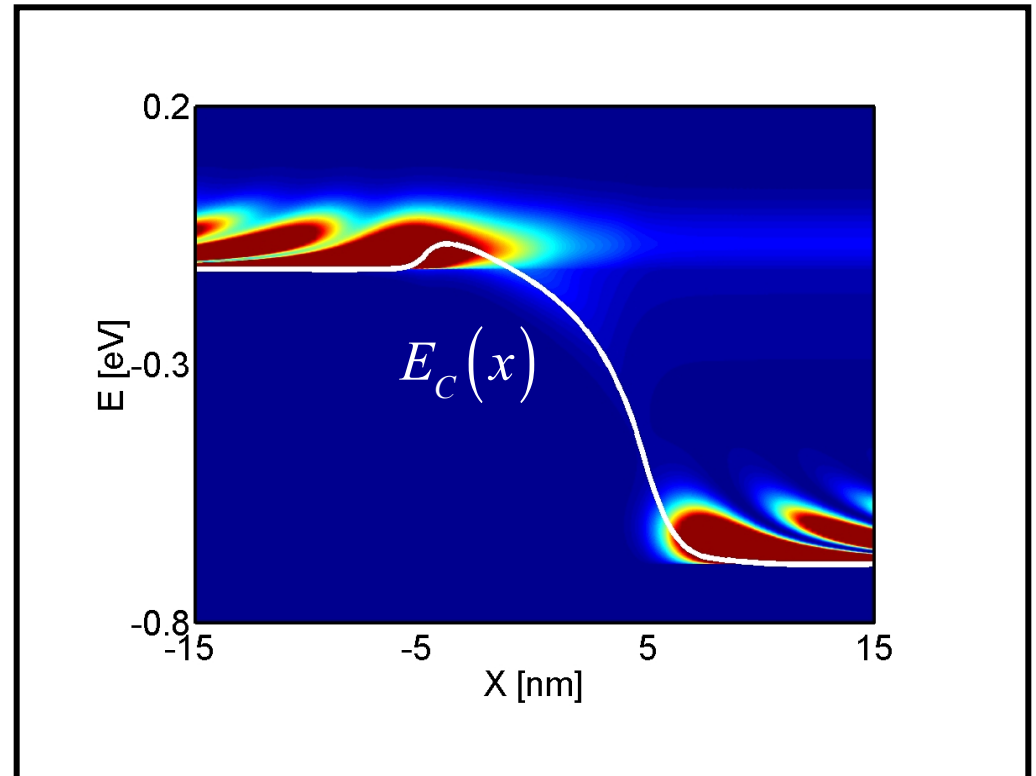
What are the equations of motion for electrons in a parabolic band semiconductor with a position dependent effective mass?

$$E(k, \vec{r}) \approx \frac{\hbar^2 k^2}{2m^*(\vec{r})}$$

Semiclassical vs. Quantum Transport



Frank, Laux, and Fischetti,
IEDM Tech. Dig., p. 553,
1992.



nanoMOS (nanoHUB.org)

Summary

- 1) Semiclassical transport means that the potential changes slowly – and that the uncertainty in position and energy is small enough.
- 2) Motion in k-space is simple.
- 3) Motion in real space involves the bandstructure.

