

# Sums and Integrals

Mark Lundstrom

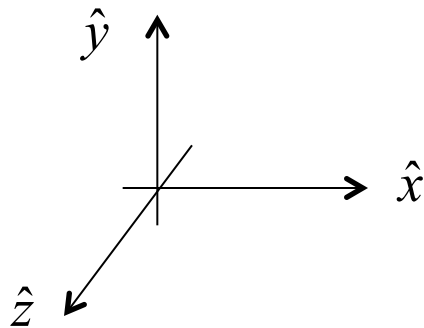
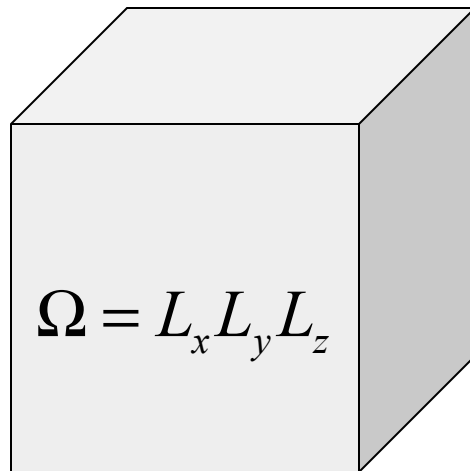
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# Outline

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- 1) **Carrier densities in 3D, 2D, and 1D**
- 2) Working in k-space
- 3) Working in energy space
- 4) Other important moments
- 5) Alternative definition of DOS
- 5) Summary

# 3D bulk semiconductor



$$N = \sum_{\vec{k}} f_0 [E(\vec{k})]$$

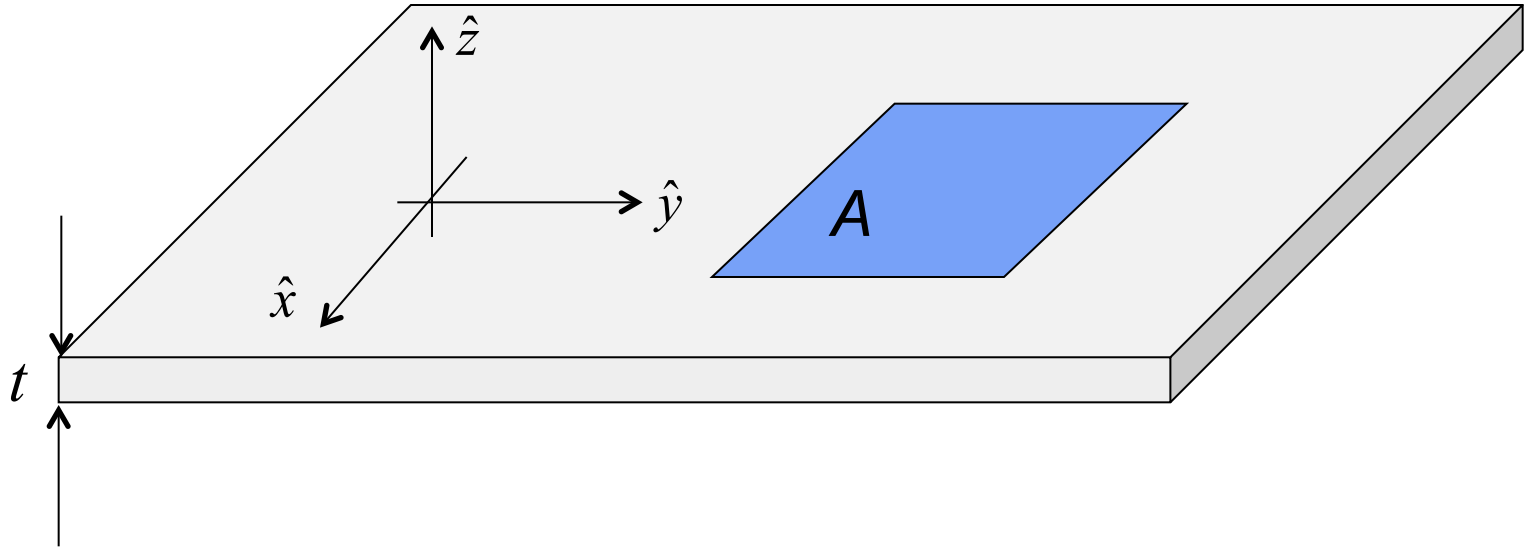
$$n = \frac{1}{\Omega} \sum_{\vec{k}} f_0(E_k)$$

$$n = \frac{1}{\Omega} \int f_0(E_k) N_k d^3k$$

$$N_k = 2 \times \left( \frac{\Omega}{8\pi^3} \right) = \frac{\Omega}{4\pi^3}$$

$$\sum_{\vec{k}} \bullet \rightarrow \frac{\Omega}{4\pi^3} \int_{BZ} \bullet d^3k$$

# 2D “bulk” semiconductor



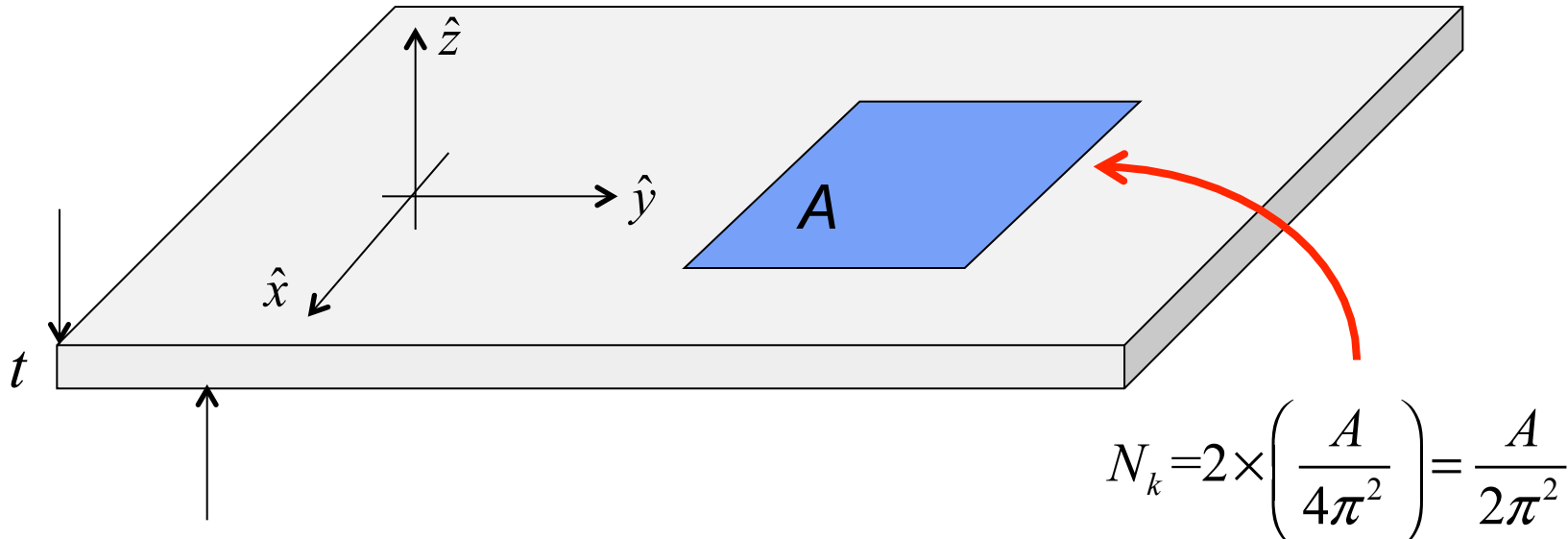
$$N = \sum_{\vec{k}} f_0(E_k)$$

$$n_S = \frac{1}{A} \sum_{\vec{k}} f_0(E_k) \text{ cm}^{-2}$$

$$n_S = \sum_j \frac{1}{A} \sum_{\vec{k}} f_0(E_k) \text{ cm}^{-2}$$

sum over subbands

# 2D bulk semiconductor: single subband

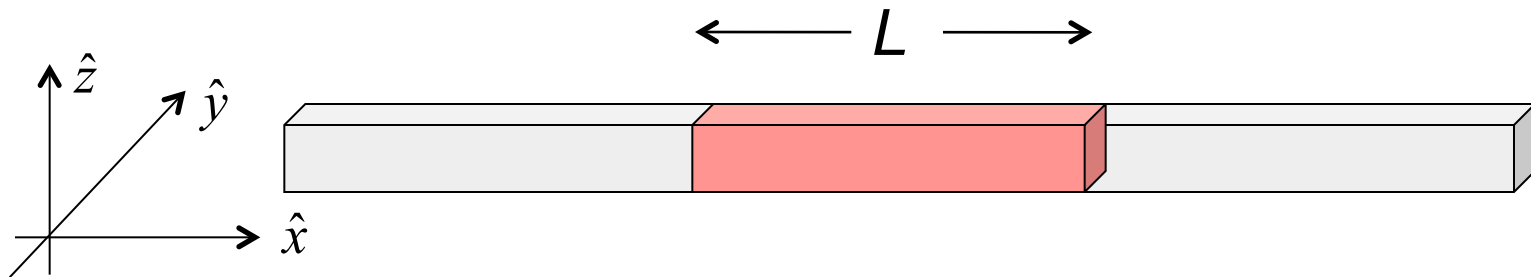


$$n_S = \frac{1}{A} \sum_{\vec{k}} f_0(E_k) \text{ cm}^{-2}$$

$$n_S = \frac{1}{A} \int_k f_0(E_k) N_k d^2k \text{ cm}^{-2}$$

$$\sum_{\vec{k}} \bullet \rightarrow \frac{A}{2\pi^2} \int_{BZ} \bullet d^2k$$

# 1D “bulk” semiconductor



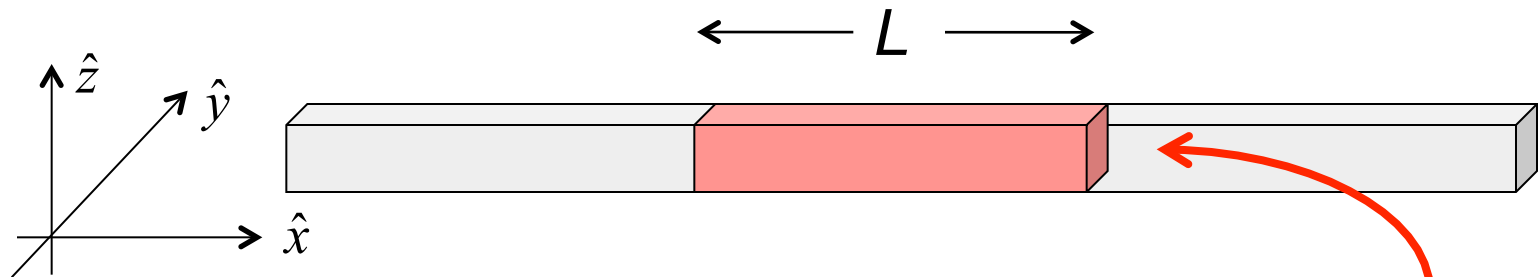
$$N = \sum_{k_x} f_0(E_k)$$

$$n_L = \frac{1}{L} \sum_{k_x} f_0(E_k) \text{ cm}^{-1}$$

$$n_L = \sum_j \frac{1}{L} \sum_{k_x} f_0(E_k) \text{ cm}^{-1}$$

sum over subbands

# 1D bulk semiconductor: single subband



$$N_k = 2 \times \left( \frac{L}{2\pi} \right) = \frac{L}{\pi}$$

$$n_L = \frac{1}{L} \sum_{k_x} f_0(E_k) \text{ cm}^{-1}$$

$$n_L = \frac{1}{L} \int_{BZ} f_0(E_k) N_k dk_x$$

$$\sum_{k_x} \bullet \rightarrow \frac{L}{\pi} \int_{BZ} \bullet dk$$

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# Density-of-states in k-space

1D:

$$N_k = 2 \times \left( \frac{L}{2\pi} \right) = \frac{L}{\pi} \quad dk$$

2D:

$$N_k = 2 \times \left( \frac{A}{4\pi^2} \right) = \frac{A}{2\pi^2} \quad dk_x dk_y$$

3D:

$$N_k = 2 \times \left( \frac{\Omega}{8\pi^3} \right) = \frac{\Omega}{4\pi^3} \quad dk_x dk_y dk_z$$

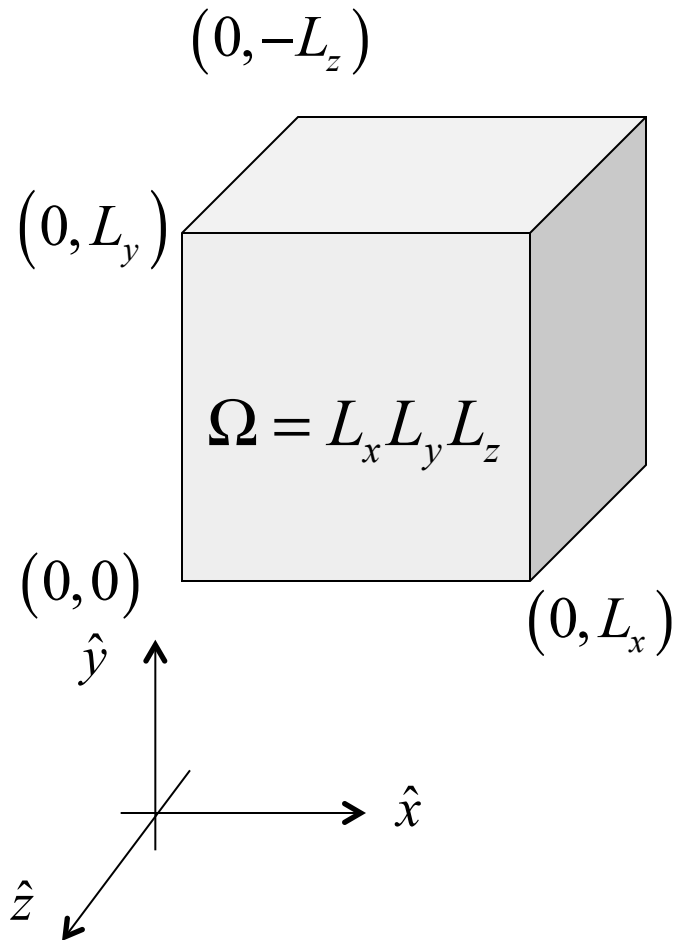
$N_k$   
independent of  $E(k)$

$$\sum_{\vec{k}} \bullet \rightarrow N_k \int_{BZ} \bullet dk$$

(Should include valley degeneracy factor,  $g_v$ .)

# Example 1: Electron density in 3D

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$$n = \frac{1}{\Omega} \sum_{\vec{k}} f_0(E_k) \text{ cm}^{-3}$$

$$n = \frac{1}{\Omega} \frac{\Omega}{4\pi^3} \int_{BZ} f_0(E_k) d^3k \text{ cm}^{-3}$$

$$f_0 = \frac{1}{1 + e^{(E - E_F)/k_B T}}$$

## Example: cont.

$$n = \frac{1}{4\pi^3} \int_{BZ} f_0(E_k) d^3k$$

Note: We extend the integral to infinity because we can usually assume that the higher energy states at large  $k$  have  $E \gg E_F$ , so they are not occupied.

$$n = \frac{1}{4\pi^3} \int_0^\infty \frac{4\pi k^2 dk}{1 + e^{(E-E_F)/k_B T}}$$

$$n = \frac{1}{\pi^2} \int_0^\infty \frac{k^2 dk}{1 + e^{(E_C + \hbar^2 k^2 / 2m^* - E_F)/k_B T}}$$

$$n = \frac{1}{\pi^2} \int_0^\infty \frac{k^2 dk}{1 + e^{-\eta_F} e^{\hbar^2 k^2 / 2m^* k_B T}}$$

$$f_0 = \frac{1}{1 + e^{(E-E_F)/k_B T}}$$

$$E = E_C + E(k) = E_C + \frac{\hbar^2 k^2}{2m^*}$$

$$\eta_F \equiv (E_F - E_C)/k_B T$$

## Example: cont.

$$n = \frac{1}{\pi^2} \int_0^{\infty} \frac{k^2 dk}{1 + e^{-\eta_F} e^{\hbar^2 k^2 / 2m^* k_B T}}$$

$$\eta_F = (E_F - E_C) / k_B T$$

$$n = \frac{(2m^* k_B T)^{3/2}}{2\pi^2 \hbar^3} \int_0^{\infty} \frac{\eta^{1/2} d\eta}{1 + e^{\eta - \eta_F}}$$

$$\eta \equiv (E - E_C) / k_B T$$

$$\eta = \hbar^2 k^2 / 2m^* k_B T$$

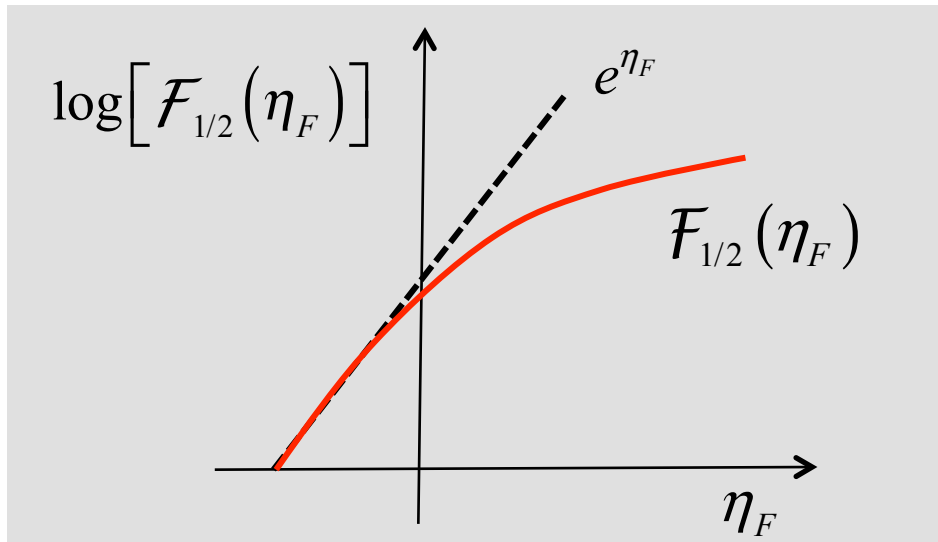
$$n = \frac{(2m^* k_B T)^{3/2}}{4\pi^{3/2} \hbar^3} \frac{2}{\sqrt{\pi}} \int_0^{\infty} \frac{\eta^{1/2} d\eta}{1 + e^{\eta - \eta_F}}$$

$$k^2 dk = \frac{(2m^* k_B T)^{3/2}}{2\hbar^3} \eta^{1/2} d\eta$$

$$n = N_C \mathcal{F}_{1/2}(\eta_F)$$

## Example: cont.

$$n = \frac{1}{\Omega} \sum_{\vec{k}} f_0(E_k) \rightarrow N_C \mathcal{F}_{1/2}(\eta_F) \text{ cm}^{-3}$$



$$\eta_F = (E_F - E_C)/k_B T$$

$$N_C = \frac{1}{4} \left( \frac{2m^* k_B T}{\pi \hbar^2} \right)^{3/2}$$

$$\mathcal{F}_{1/2}(\eta_F) \equiv \frac{2}{\sqrt{\pi}} \int_0^{\infty} \frac{\eta^{1/2} d\eta}{1 + e^{\eta - \eta_F}}$$

$$\eta_F \ll 0 \quad E_F \ll E_C \quad \mathcal{F}_{1/2}(\eta_F) \rightarrow e^{\eta_F} \quad n = N_C e^{\eta_F} \text{ cm}^{-3}$$

(non-degenerate semiconductor)

# Fermi-Dirac integrals

$$\mathcal{F}_j(\eta_F) = \frac{1}{\Gamma(j+1)} \int_0^{\infty} \frac{\eta^j d\eta}{1 + e^{\eta - \eta_F}}$$

$$\Gamma(n) = (n-1)! \quad (n \text{ integer})$$

$$\Gamma(1/2) = \sqrt{\pi}$$

$$\Gamma(p+1) = p\Gamma(p)$$

$$\mathcal{F}_j(\eta_F) \rightarrow e^{\eta_F} \quad \eta \ll 1$$

$$(E_F - E_C)/k_B T \ll 1$$

$$\frac{d\mathcal{F}_j}{d\eta_F} = \mathcal{F}_{j-1}$$

don't confuse with....  $F_j(\eta) = \int_0^{+\infty} \frac{x^j dx}{1 + e^{x-\eta}}$

For an introduction to Fermi-Dirac integrals, see: "Notes on Fermi-Dirac Integrals," 3rd Ed., by R. Kim and M. Lundstrom) <https://www.nanohub.org/resources/5475>

# outline

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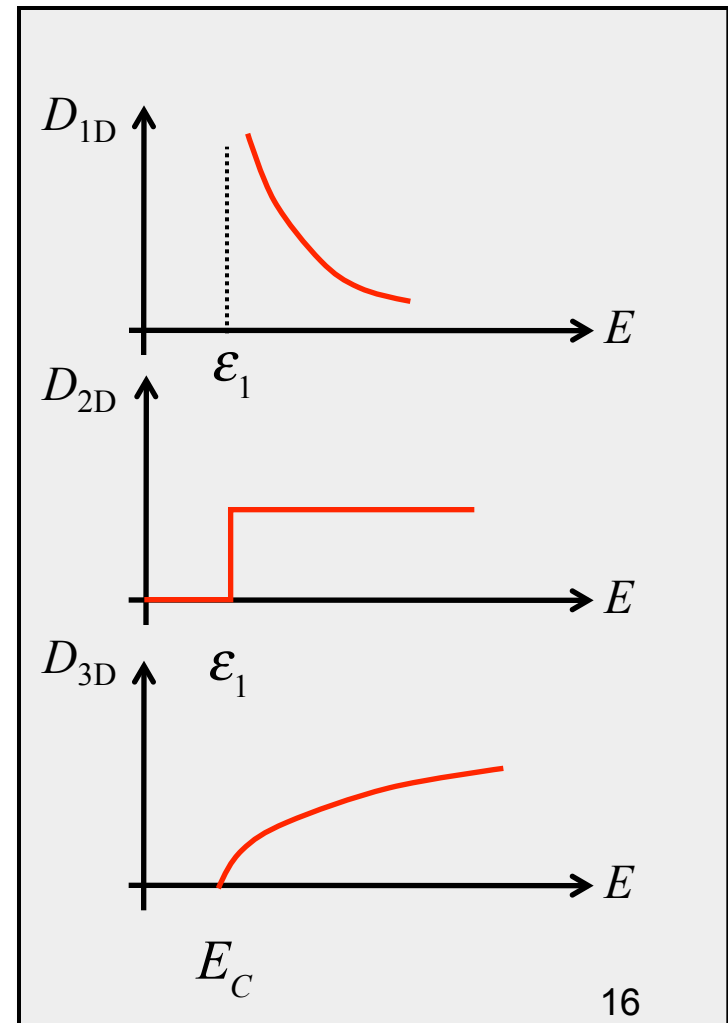
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# Parabolic bands: 1D, 2D, and 3D

$$D_{1D}(E) = g_V \frac{1}{\pi \hbar} \sqrt{\frac{2m^*}{(E - \varepsilon_1)}} \Theta(E - \varepsilon_1)$$

$$D_{2D}(E) = g_V \frac{m^*}{\pi \hbar^2} \Theta(E - \varepsilon_1)$$

$$D_{3D}(E) = g_V \frac{m^* \sqrt{2m^*(E - E_C)}}{\pi^2 \hbar^3} \Theta(E - E_C)$$



$$(E(k) = E_C + \hbar^2 k^2 / 2m^*)$$



# Energy space (3D)

$$n = \int_{E_C}^{\infty} f_0(E) D(E) dE$$

$$n = \frac{(2m^*)^{3/2}}{2\pi^2 \hbar^3} \int_{E_C}^{\infty} \frac{(E - E_C)^{1/2}}{1 + e^{(E - E_F)/k_B T}} dE$$

$$n = \frac{(2m^* k_B T)^{3/2}}{2\pi^2 \hbar^3} \int_0^{\infty} \frac{\eta^{1/2} d\eta}{1 + e^{\eta - \eta_F}}$$

$$n = N_C \mathcal{F}_{1/2}(\eta_F) \text{ cm}^{-3}$$

$$f_0 = \frac{1}{1 + e^{(E - E_F)/k_B T}}$$

$$D(E) = \frac{(2m^*)^{3/2}}{2\pi^2 \hbar^3} (E - E_C)^{1/2}$$

$$\eta_F = (E_F - E_C)/k_B T$$

$$\eta = (E - E_C)/k_B T$$

$$N_C = \frac{1}{4} \left( \frac{2m^* k_B T}{\pi \hbar^2} \right)^{3/2}$$

# k-space vs. energy-space

k-space:

$$n = \frac{1}{\Omega} \sum_k f_0(E_k) = \frac{1}{\Omega} \int_{BZ} f_0(E_k) N_k d^3k \text{ cm}^{-3}$$

$$N_k d^3k = \frac{\Omega}{4\pi^3}$$

$$n = N_C \mathcal{F}_{1/2}(\eta_F) \text{ cm}^{-3}$$

energy-space:

$$n = \int_{E_C}^{\infty} f_0(E) D(E) dE$$

$$D(E) dE = \frac{(2m^*)^{3/2}}{2\pi^2 \hbar^3} (E - E_C)^{1/2} dE$$

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# Moments of the Fermi function

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In equilibrium, physical quantities are various moments of the Fermi function.

$$n = \frac{1}{\Omega} \sum_{\vec{k}} \phi(\vec{k}) f_0(E_k) \quad \text{cm}^{-3}$$

$$\phi(\vec{k}) = v_x(\vec{k})$$

$$\phi(\vec{k}) = 1$$

$$\langle v_x \rangle = \frac{1}{\Omega} \sum_{\vec{k}} v_x(\vec{k}) f_0(E_k) = 0$$

$$n = \frac{1}{\Omega} \sum_{\vec{k}} f_0(E_k) \quad \text{cm}^{-3}$$

$$\frac{\langle v_x \rangle}{n} = \frac{\frac{1}{\Omega} \sum_{\vec{k}} v_x(\vec{k}) f_0(E_k)}{\frac{1}{\Omega} \sum_{\vec{k}} f_0(E_k)} = 0$$

# Other moments

$$n = \frac{1}{\Omega} \sum_{\vec{k}} \phi(\vec{k}) f_0(E_k) \quad \text{cm}^{-3}$$

$$\phi(\vec{k}) = 1 \quad n \quad \text{cm}^{-3}$$

$$\phi(\vec{k}) = v_x(\vec{k}) \quad \langle v_x \rangle = 0$$

$$\phi(\vec{k}) = E(\vec{k}) \quad \langle E_k \rangle$$

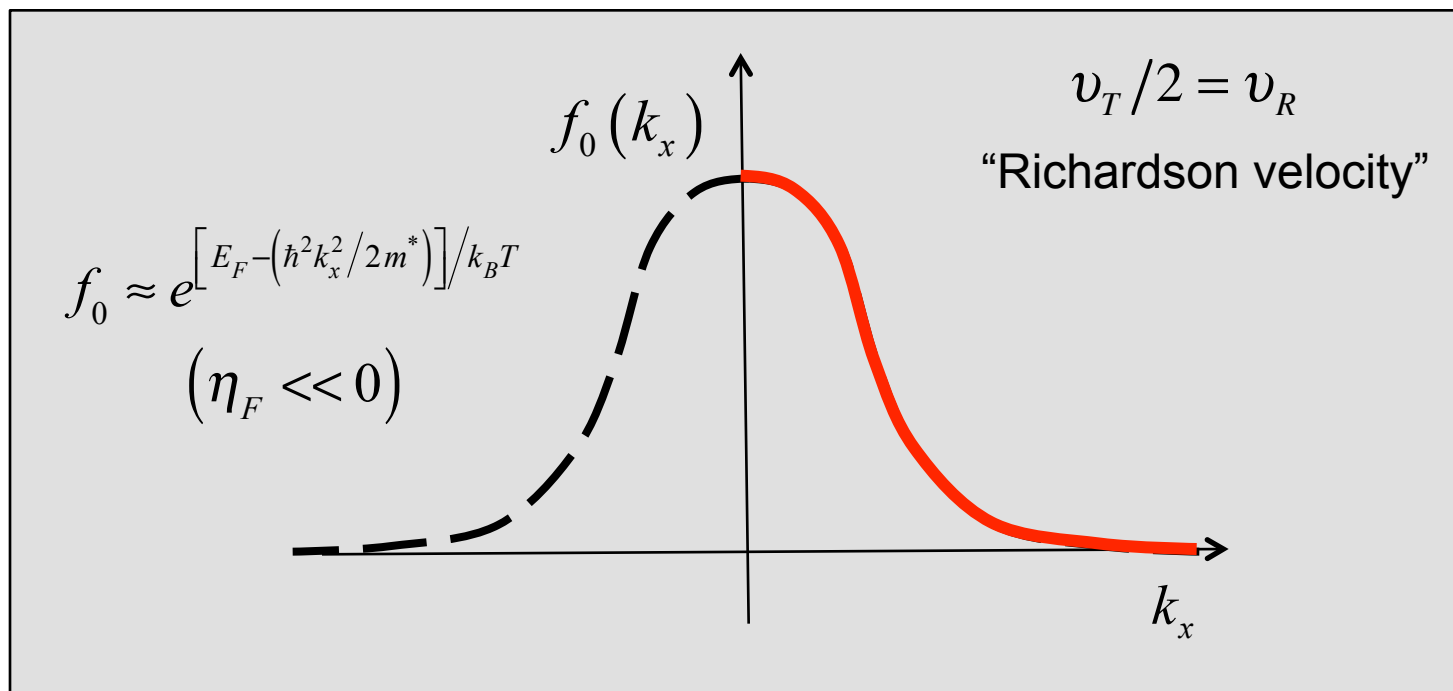
$$\frac{\langle E_k \rangle}{n} = u = \frac{\frac{1}{\Omega} \sum_{\vec{k}} E_k f_0(E_k)}{\frac{1}{\Omega} \sum_{\vec{k}} f_0(E_k)} \quad \text{J}$$

$$u = \frac{3}{2} k_B T$$

$$(\eta_F \ll 0)$$

# Uni-directional thermal velocity

$$\langle v^+ \rangle = \frac{\frac{1}{\Omega} \sum_{k_x > 0, k_y, k_z} v_k f_0(E_k)}{\frac{1}{\Omega} \sum_{k_x > 0, k_y, k_z} f_0(E_k)} \quad \text{cm/s} = \tilde{v}_T \quad v_T = \sqrt{\frac{2k_B T}{\pi m^*}} \quad (\eta_F \ll 0)$$



# About thermal velocities

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$$u = \frac{3}{2} k_B T \qquad u = \frac{1}{2} m^* \langle v^2 (E) \rangle$$
$$\frac{1}{2} m^* \langle v^2 \rangle = \frac{3}{2} k_B T$$

$$\langle v^2 \rangle^{1/2} = v_{rms} = \sqrt{\frac{3k_B T}{m^*}} \neq v_T = \sqrt{\frac{2k_B T}{\pi m^*}}$$

- i) unidirectional thermal velocity
- ii) Richardson thermal velocity
- iii) rms thermal velocity

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# An alternative way to define the DOS

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$$D_{1D}(E) = \frac{1}{L} \sum_k \delta(E - E_k) \quad \frac{\#}{\text{J-m}}$$

$$D_{2D}(E) = \frac{1}{A} \sum_{\vec{k}} \delta(E - E_{\vec{k}}) \quad \frac{\#}{\text{J-m}^2}$$

$$D_{3D}(E) = \frac{1}{\Omega} \sum_{\vec{k}} \delta(E - E_{\vec{k}}) \quad \frac{\#}{\text{J-m}^3}$$

# Proof (1D)

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in k-space, we know:

$$n_L = \frac{1}{L} \sum_k f_0(E_k)$$

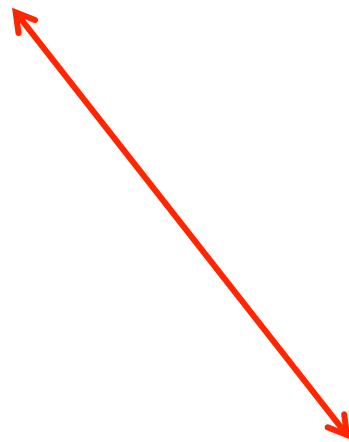
can also work in energy-space:

$$n_L = \int f_0(E) D_{1D}(E) dE$$

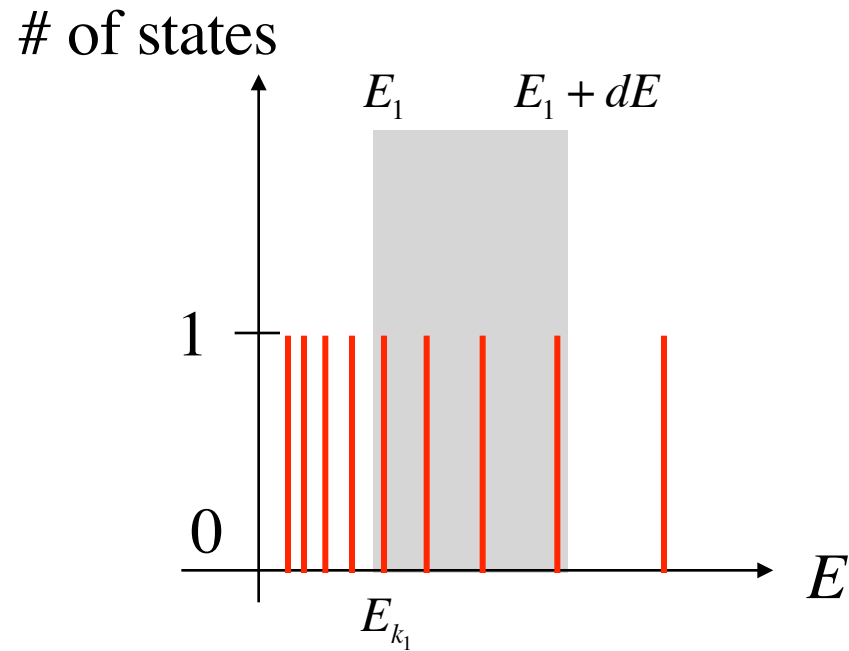
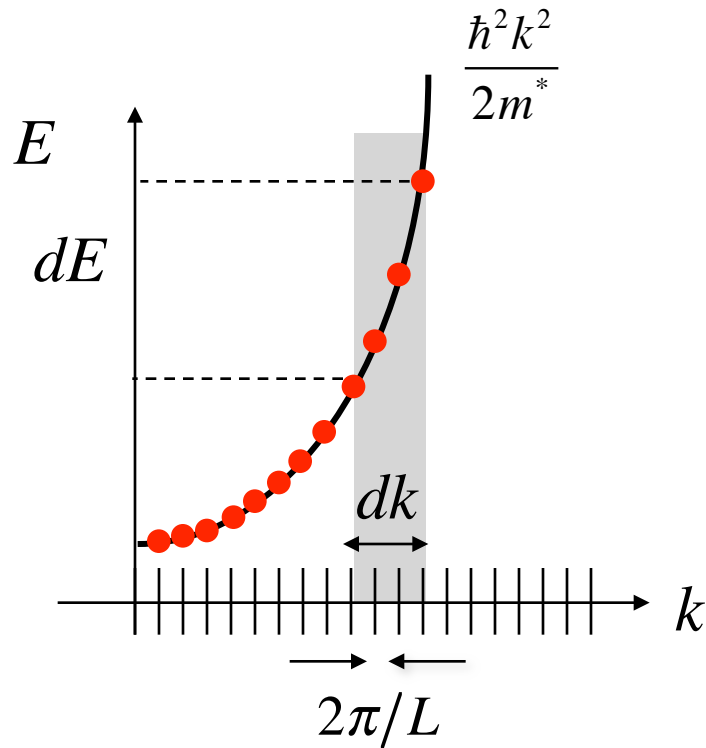
$$n_L = \int f_0(E) \frac{1}{L} \sum_k \delta(E - E_k) dE$$

$$n_L = \frac{1}{L} \sum_k \int f_0(E) \delta(E - E_k) dE$$

$$n_L = \frac{1}{L} \sum_k f_0(E_k)$$



# Interpretation



$$\int_{E_1}^{E_1+dE} D_{1D}(E) dE = \int_{E_1}^{E_1+dE} \frac{1}{L} \sum_k \delta(E - E_k) dE = \frac{1}{L} \sum_k \int_{E_1}^{E_1+dE} \delta(E - E_k) dE$$

**counts the states between  $E$  and  $E + dE$**

# Example: 2D DOS for parabolic energy bands

$$D_{2D}(E_1) = \frac{1}{A} \sum_{\vec{k}} \delta(E_1 - E_{\vec{k}}) \frac{\#}{\text{J-m}^2}$$

$$D_{2D}(E_1) = \frac{1}{A} \sum_{\vec{k}} \delta(E_1 - E_{\vec{k}}) \rightarrow \frac{1}{A} g_V \frac{A}{(2\pi)^2} \times 2 \int_0^{\infty} \delta(E_1 - E_{\vec{k}}) 2\pi k dk$$

$$D_{2D}(E_1) = g_V \frac{1}{\pi} \times \int_0^{\infty} \delta(E_1 - E(k)) k dk$$

$$E(k) = \frac{\hbar^2 k^2}{2m^*}$$

$$dE = \frac{\hbar^2 2k dk}{2m^*}$$

$$k dk = \frac{m^*}{\hbar^2} dE$$

$$D_{2D}(E_1) = g_V \frac{m^*}{\pi \hbar^2} \times \int_0^{\infty} \delta(E_1 - E(k)) dE$$

$$D_{2D}(E_1) = g_V \frac{m^*}{\pi \hbar^2} \quad \checkmark$$

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# Summary

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- 1) Be familiar with how to work out sums in  $k$ -space and integrals in energy space.
- 2) Become familiar with Fermi-Dirac integrals.
- 3) Understand the uni-directional thermal velocity.
- 4) Understand the delta-function definition of the DOS.

