

Thermoelectric Effects in Bipolar Devices: Internally Cooled Semiconductor Lasers

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Heating in semiconductor lasers is detrimental to device performance, affecting such characteristic parameters as threshold current, efficiency, and lifetime. Previous work has identified several sources of heating (Joule heating, radiation absorption, contact heating, and non-radiative recombination) and has modeled heat transport through thermal conduction [1]. These models have neglected macroscopic heat transport mechanisms such as convection and have likewise neglected several microscopic heat exchange processes such as thermoelectric effects that occur when carriers move across heterojunctions and diffuse against built-in fields. These microscopic effects have been demonstrated to produce net cooling of more than ten degrees in bulk thin films [2], and their magnitude in (bipolar) laser diodes has been shown to be of the same order as conventional heating processes, with predicted cooling under certain bias conditions and device geometries [3].

In Fig. 1a we show measured data taken with a $25 \times 25 \mu\text{m}^2$ NIST-calibrated microthermocouple (accurate to 10mK) on a $20 \times 500 \mu\text{m}^2$ $\lambda=980\text{nm}$ InGaP/InGaAs/GaAs laser mounted on a heatsunk copper block. A thermistor was stabilized at 17°C and was located 2mm from the laser. Temperature measurements were taken both on the laser surface and on the heat sink approximately $10 \mu\text{m}$ from the bottom contact to the $100 \mu\text{m}$ -thick GaAs substrate. A temperature difference ΔT was measured at zero bias, indicating the effect of convection from ambient air (21.4°C) warming the laser. This convective power is of greater magnitude than the laser's optical power, and must be included in a thermal model [4]. As the laser reaches threshold, a fraction of the electrical bias power is dissipated through radiation rather than through heat, appearing as a kink in ΔT . By carefully applying a thermal impedance model for heat conduction and by considering convection, this can be used to derive the laser's optical power without recourse to an optical detector [5].

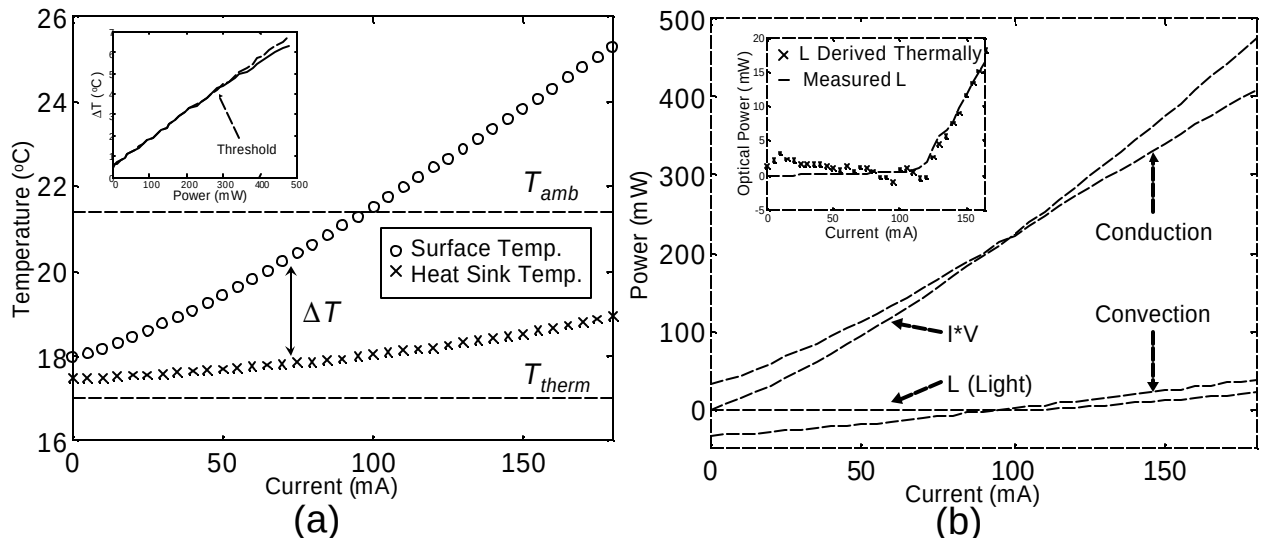


Fig. 1a: Measured data; INSET: slope discontinuity in ΔT at threshold. **Fig. 1b:** Calculated rates of conduction and convection in comparison to measured bias power and optical power; INSET: calculated optical power.

This macroscopic model of heat transfer assumes a uniform temperature throughout the laser that is often approximately valid for typical devices. However, if a laser has both small thermal conductivity and large band offsets (with respect to the photon energy), large internal temperature gradients can develop due to thermoelectric effects.

Transport in semiconductor lasers is distinguished by its bipolar nature. Current injection into the active region of a diode structure requires the charged carriers to move against the built-in electric field at the junction. Heat exchange with the crystal lattice is required for this injection to take place in steady-state. This heat exchange can be described by a bias dependent Seebeck coefficient for both majority and minority carriers. This heat exchange can be used to dramatically reduce the junction temperature with respect to the electrical contacts [6] (Figure 2).

During the 150-year history of thermoelectric devices only *passive* electronic devices have been used for thermoelectric cooling. A typical thermoelectric cooler consists of a polycrystalline semiconductor sandwiched between metallized ceramic plates. Since the electrical characteristics of these devices are those of simple resistors these devices cannot be used for switching, amplification, or light generation. Thermoelectric effects in bipolar devices represent a new class of *active thermoelectric* devices. This bipolar thermoelectric device can generate large temperature gradients that can enhance the ability of these diode lasers to generate photons at longer wavelengths and higher temperatures.

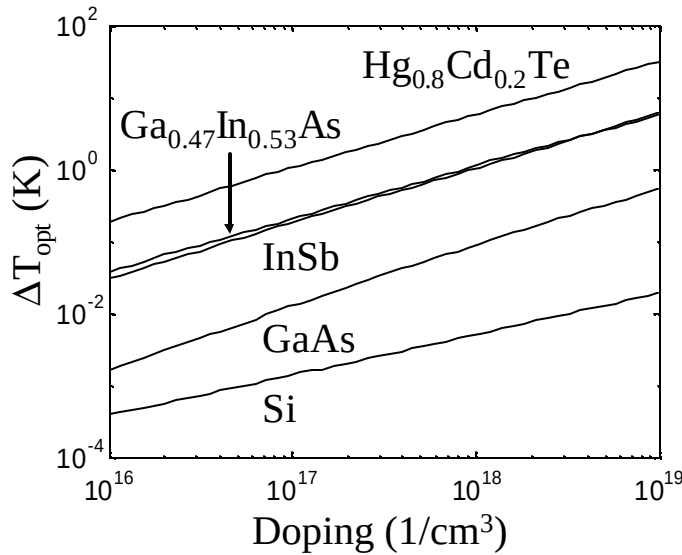


Figure 2 The maximum internal temperature difference that can be generated within a diode structure. Calculations are shown for symmetrically doped diodes with various materials parameters.

Fig. 3 shows a $\lambda=2.64\mu\text{m}$ GaSb/Ga_{0.7}In_{0.3}As_{0.26}Sb_{0.74}/Ga_{0.8}In_{0.2}As_{0.17}Sb_{0.83} device in which the cladding doping has been modified to create either heating or cooling as injected carriers enter the core region [6]. These heating and cooling effects can be large with respect to other heating processes; Figure 4 shows the results of a finite-element simulation which takes into account thermoelectric effects as well as other thermal sources such as Joule heating and non-radiative recombination (contact resistance is typically very low in GaSb due to Fermi-level pinning). While leakage current (~8% in these devices) produces thermoelectric cooling in both devices, injection current cools rather than heats the optimized device, producing an overall lower

temperature in the quantum well. Increased removal of heat from the contacts is expected to magnify this effect.

By examining for the first time these macroscopic and microscopic heat exchange terms, we uncover new design strategies for the internal and external cooling of lasers, as well as a new characterization tool.

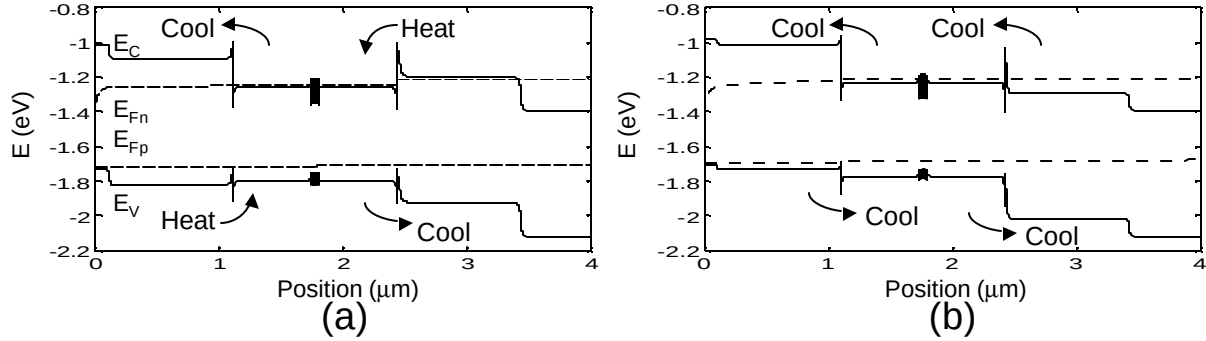


Fig. 3: Simulated band structure for GaInAsSb (a) conventional and (b) optimized laser structures biased at 475 A/cm^2 (threshold), with associated thermoelectric heat exchange illustrated. $J_{th} < 100 \text{ A/cm}^2$ has been shown for this material.

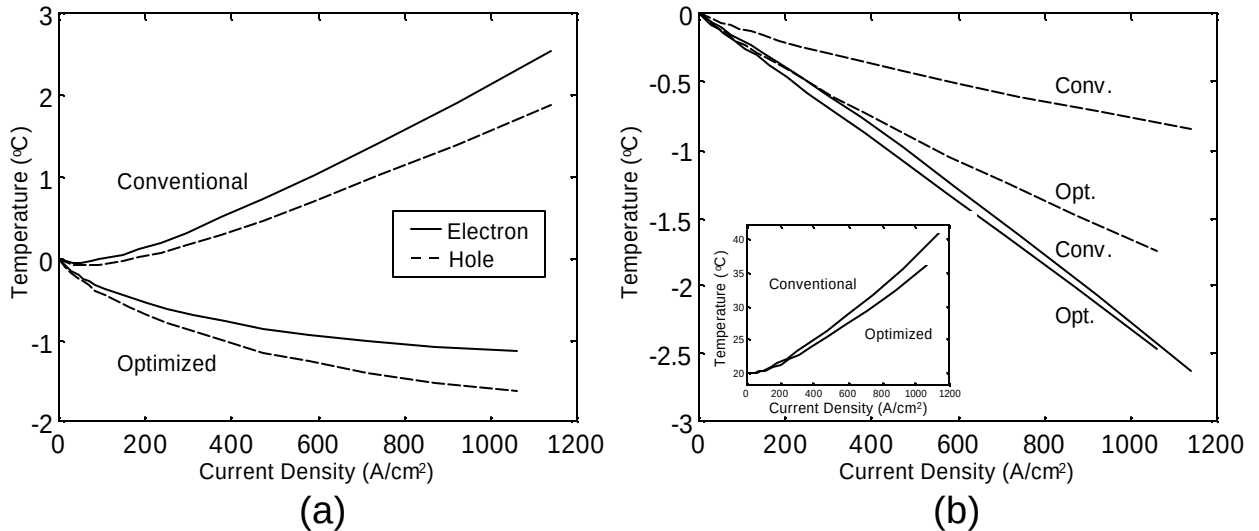


Fig. 4: Contribution by (a) injection and (b) leakage thermoelectric terms to the quantum well temperature. INSET: Overall quantum well temperature, including all thermoelectric terms as well as Joule and recombination heating.

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