

# Quantum Transport:

Atom to Transistor

**Prof. Supriyo Datta**  
ECE 659  
Purdue University

02.14.2003

## Lecture 14: Basis Functions, Density Matrix I

Ref. Chapter 4.3



Network for Computational Nanotechnology



## Retouch on Concepts

00:00

- Summary of Equations:

$$E\Phi = H_{op}\Phi$$

$$\Phi(\vec{r}) = \sum_m f_m u_m(\vec{r})$$

$$\Phi(\vec{r}) = \sum_i f_i u'_i(\vec{r})$$

and

$$u'_i = \sum_m C_{mi} u_m(\vec{r})$$

$$\{\Phi\} = [C]\{\Phi'\}$$

$$[A'] = C^+ A C$$

- Recall the concept of a Hilbert Space, a function space conceptually similar to a vector space. Basis functions in Hilbert space are like the unit vectors  $\hat{x}$ ,  $\hat{y}$ ,  $\hat{z}$  in vector space.
- Proper use of basis functions facilitate matrix reduction and so are very useful computationally.
- Computationally we often use a non-orthogonal basis set (for example recall  $H_2$ ). Conceptually this can cause many problems. The usual procedure is to transform the oblique sets to orthogonal sets.
- From now on it will be assumed that we are working with orthogonal basis sets:

$$\int d\vec{r} u_n^* u_m = d_{mn}$$

## Basis Transformations

04:08

- Basis transformations are defined by the relation:

$$u'_i = \sum_m C_{mi} u_m(\vec{r})$$

Where:

$$\Phi(\vec{r}) = \sum_m f_m u_m(\vec{r}) = \sum_i f'_i u'_i(\vec{r})$$

So to transform a vector to a new basis set we'll have:

$$\{\Phi\} = [C] \{\Phi'\}$$

- In matrix notation, basis transformations are defined by:

$$[A'] = [C]^\dagger A [C]$$

- To show this, we have:

$$A'_{ij} = \int d\vec{r} u_i^*(\vec{r}) A_{op} u'_j(\vec{r})$$

$$\begin{aligned} &= \int d\vec{r} \sum_m C_{mi}^* u_m^*(\vec{r}) A_{op} \sum_n C_{nj} u_n(\vec{r}) \\ &= \sum_m \sum_n C_{mi}^* A_{mn} C_{nj} \end{aligned}$$

$$\text{But } C_{mi}^* = (C^+)_{im}$$

$$\therefore A'_{ij} = \sum_m \sum_n (C^+)_{im} A_{mn} C_{nj} = [C^+ A C]_{ij}$$

Thus, we see that C can be used to transform matrix operators as well as vectors.

## Hamiltonian Matrix Transformation

11:10

- First, recall the definition of a unitary transformation:
  - Vector length is preserved
  - Proviso on unitary matrix:  $C^\dagger C = C C^\dagger = I$
- Consider the process of finding the eigenvalues and eigenvectors of a Hamiltonian matrix.
- In matlab, we invoke:  
 **$[V,D] = \text{eig}(H)$**   
where D is a **diagonal** matrix with the eigenvalues on the diagonal and V is a square matrix with the eigenvectors as its columns.

- One way to visualize this process is to consider  $[H] \rightarrow [D]$  as a basis transformation from the real space basis to the eigenvector basis. Formally this is expressed as:

$$[D] = V^\dagger [H] V$$

$$V = \begin{bmatrix} | & | & & | \\ 1 & 2 & & n \\ | & | & \cdots & | \\ | & | & & | \end{bmatrix}$$

## Ordering of the Hamiltonian Matrix

18:16

- How do we know it is  $D = V^+HV$  and not  $D = VHV^+$ ?
- Look at the organization of old and new basis sets in  $V$ ,  $H$  and  $D$ ...

$$V = \begin{bmatrix} \xrightarrow{\text{new basis}} \\ \downarrow \text{old basis} \end{bmatrix} \quad H = \begin{bmatrix} \xrightarrow{\text{old basis}} \\ \downarrow \text{old basis} \end{bmatrix} \quad D = \begin{bmatrix} \xrightarrow{\text{new basis}} \\ \downarrow \text{new basis} \end{bmatrix}$$

- Rules of matrix multiplication require that the columns of one matrix match the rows of next matrix.
- So by observation it must be  $D = V^+HV$

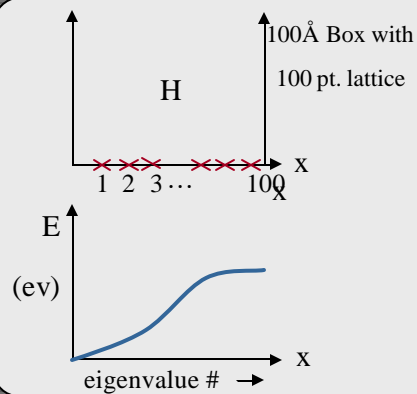
$$\begin{bmatrix} \xrightarrow{\text{new basis}} \\ \downarrow \text{new basis} \end{bmatrix} \underset{D}{=} \begin{bmatrix} \xrightarrow{\text{old basis}} \\ \downarrow \text{new basis} \end{bmatrix} \times \begin{bmatrix} \xrightarrow{\text{old basis}} \\ \downarrow \text{old basis} \end{bmatrix} \times \begin{bmatrix} \xrightarrow{\text{new basis}} \\ \downarrow \text{old basis} \end{bmatrix} \underset{H}{}$$

# Electrons in a Box

22:38

- As an example, we will calculate the electron density of the 'electrons in a box'

## Electrons in a box



- Remember, by the method of finite differences (using box boundary conditions).

$$H = \begin{pmatrix} +2t_0 & -t_0 & & \\ -t_0 & +2t_0 & \ddots & \\ & \ddots & \ddots & \ddots \\ & & & -t_0 & +2t_0 \end{pmatrix}$$

- We want to find the electron density

$$n(x) = \sum_{occ. a} |f_a(x)|^2$$

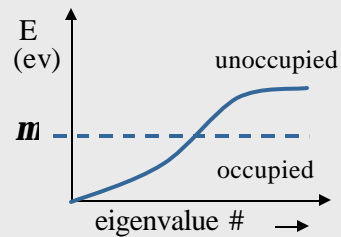
Note: 'occ. a' refers to the sum over all **occupied** states. It is very important to not that here the states are either full or empty.

## Electron Density

25:48

- We can redefine  $n(x)$  by applying the Fermi function to all states. Where the Fermi function provides the “degree of occupation” between 0 and 1 of a given state at a known electrochemical potential  $\mu$ .

Occupied and unoccupied levels at  $\mu$



- So the electron density is:

$$n(x) = \sum_a |\mathbf{f}_a(x)|^2 f_a$$

$$= \sum_a \mathbf{f}_a(x) f_a \mathbf{f}_a^*(x)$$

- We can also re-write this as:

$$n(x) = \sum_a \sum_b \mathbf{f}_a(x) \mathbf{r}_{ab} \mathbf{f}_b^*(x)$$

where

$$\mathbf{r}_{ab} = \begin{cases} f_a & \text{if } a = b \\ 0 & \text{if } a \neq b \end{cases}$$

- Note:  $\mathbf{r}_{ab}$  forms a diagonal matrix called the density matrix.

## Density Matrix

31:00

- Generalizing we write:

$$\tilde{r}(x, x') = \sum_a \sum_b f_a(x) r_{ab} f_b^*(x')$$

Where  $n(x)$  is the diagonal of  $\tilde{r}(x, x')$

- This relation can be seen to represent a unitary transformation from the eigenvector basis to **real** space. Note:  $\Phi_a(x)$  are given by columns of  $V$ :

$$f_a(x) = x \begin{pmatrix} \vdots \\ \vdots \\ \vdots \end{pmatrix} = V_{x,a}$$

- In other words:

$$\Phi_a(x) = (V)_{a,x}$$

$$\Phi_b^*(x') = (V^+)_{b,x'} = V_{x',b}^*$$

- Summarizing,

$$\tilde{r} = V r V^+$$

$$r = V^+ \tilde{r} V$$

- $\tilde{r}$  is in **real** space
- $r$  is in the eigenstate space
- The diagonal elements of  $\tilde{r}$  are equal to the electron density  $n(x)$ .



- So, in the eigenstate basis or “space”  $\rho$  is a diagonal matrix with elements

$$\rho_{aa} = \frac{1}{1 + e^{(E_a - \mu)/k_B T}}$$

$$= f_0(E_a - \mu)$$

- Now in general let us denote the density matrix for any space as  $\rho$ , where  $\rho$  is given by:

$$\rho = f_0([H] - \mu[I])$$

- What is meant by  $f_0([H] - \mu[I])$  ?

More Generally, how is the ‘function’ of a matrix calculated? For a diagonal matrix it is simply the ‘function’ operated on all elements. How about matrices with off diagonal elements?

- Example: Given  $[H]$  with off diagonal elements calculate  $\sin[H]$ . To do this we must first diagonalize  $[H]$ , then operate  $\sin()$  upon the diagonalized form of  $[H]$ , and then finally transform  $[H]$  back into its original space.

## Equilibrium Density Matrix

42:00

Example: continued.

(1) Diagonalize  $[H]$

$$\begin{pmatrix} H_1 & & \\ & H_2 & \\ & & \ddots \end{pmatrix}$$

(2) Operate  $\sin()$

$$\begin{pmatrix} \sin(H_1) & & \\ & \sin(H_2) & \\ & & \ddots \end{pmatrix}$$

(3) Transform back to original space

$$\mathbf{V} \begin{pmatrix} \sin(H_1) & & \\ & \sin(H_2) & \\ & & \ddots \end{pmatrix} \mathbf{V}^+$$

- Note: In matlab matrix functions and element by element functions are differentiated by addition of an 'm' for matrix functions. e.g. :  $\sin()$  : represents element by element operation  $\sinm()$ : represents matrix operations.

- Finally we see the expression for  $\rho$  in real space is:

$$\tilde{\rho} = V \begin{pmatrix} f_0(E_1 - \mu) & & \\ & f_0(E_2 - \mu) & \\ & & \ddots \\ & & & f_0(E_n - \mu) \end{pmatrix} V^+$$

- Interestingly,  $\rho$  is only diagonal in the eigenvector basis. Off diagonal elements of  $\rho$  in alternate basis sets are used in some calculations, but more often than not only the diagonal elements of  $\rho$  (which in any space provide electron density) are of interest.