



# **ECE695: Reliability Physics of Nano-Transistors**

## **Lecture 19: Spin-Dependent Recombination and Electrically Detected Magnetic Resonance**

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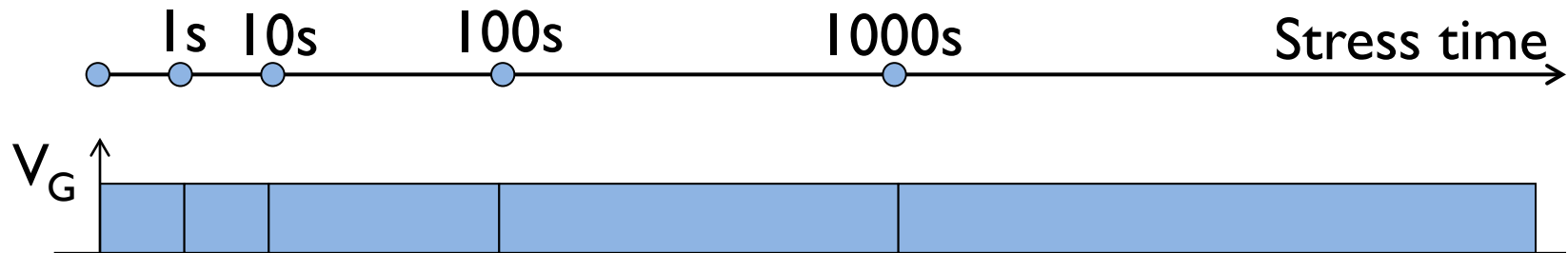
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# Outline

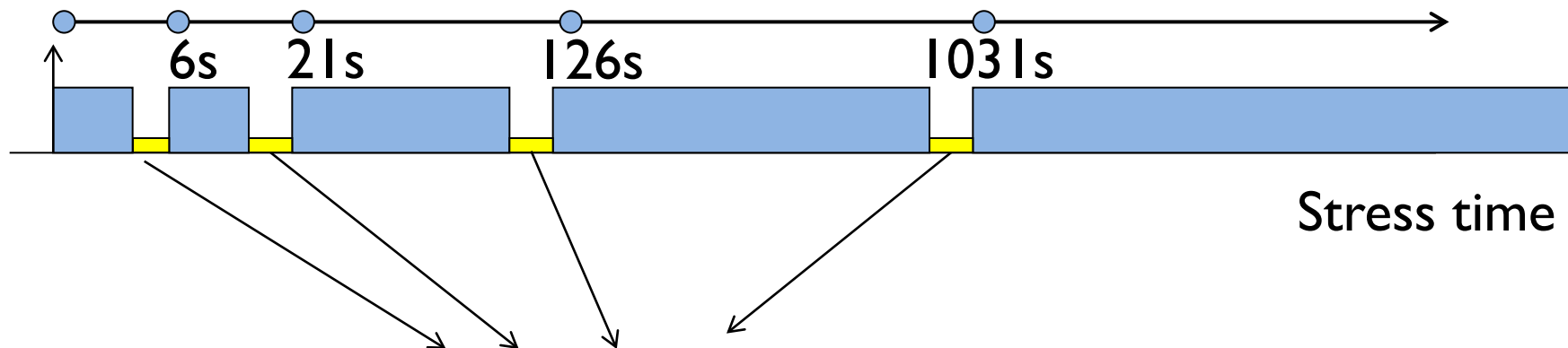
1. Importance of measuring interface damage
2. Electronic Spin Resonance (A quick review)
3. Spin Dependent Recombination
4. Electrically detected spin-resonance and noise-spectroscopy
5. Comparing the approaches
6. Conclusions

# Measurement is a complex process

We periodically stop the stress and measure defects ...

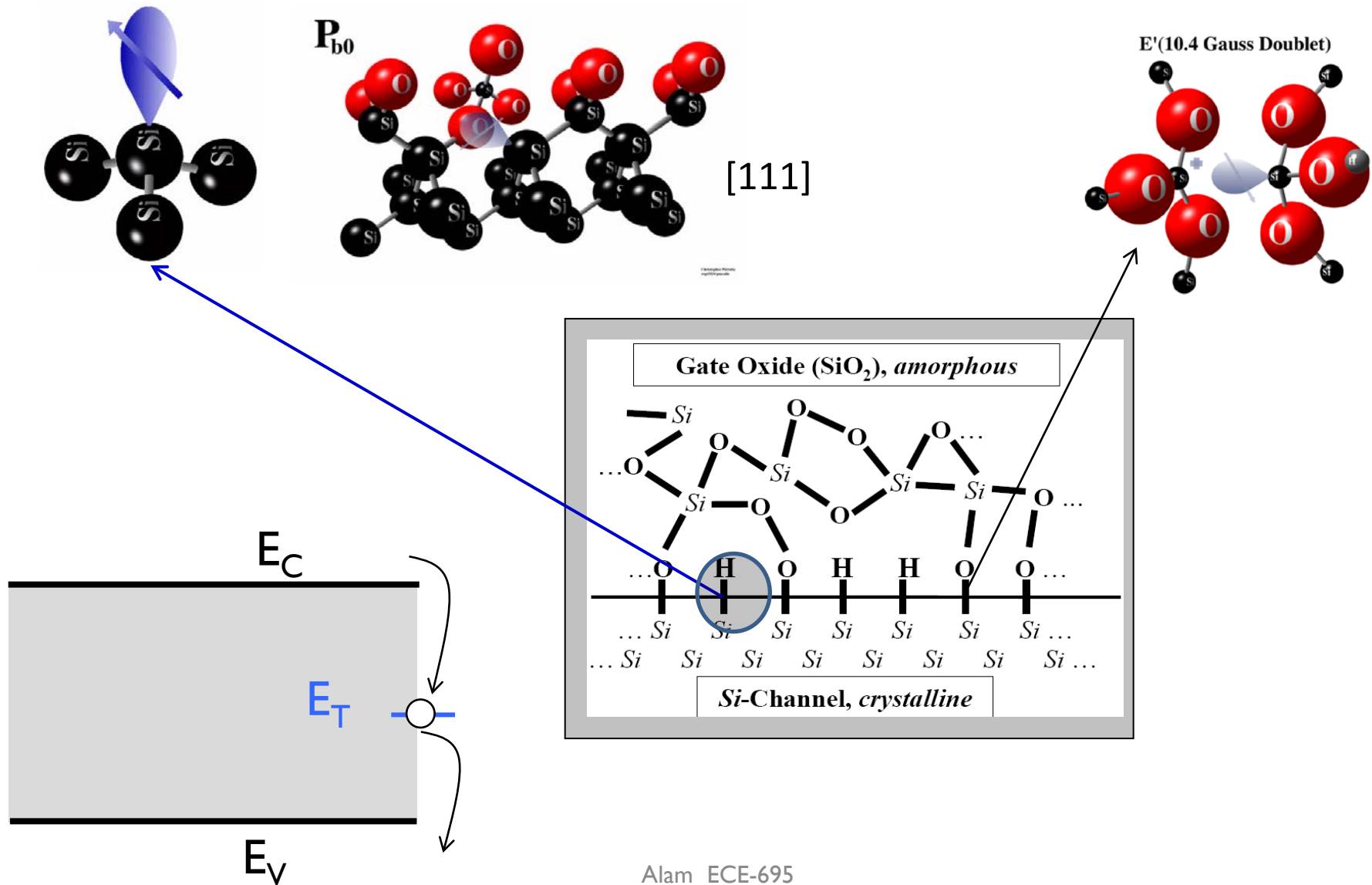


The measurements are often complex and interpretation of data depends on our interpretation of measurement

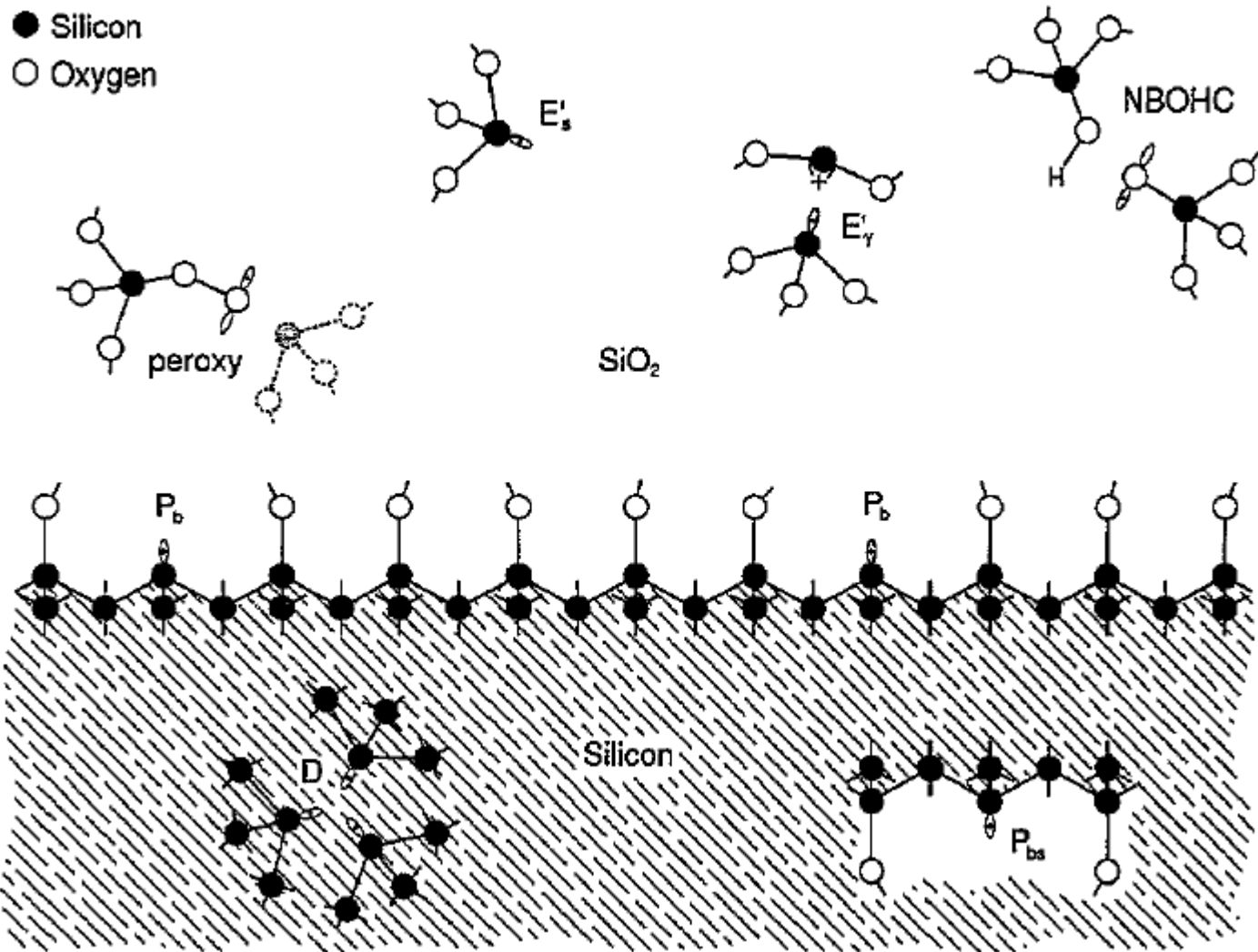


C-V, SILC, DCIV, CP, **SDR**, Idlin methods, ....

# Review and background



# Different types of ESR-visible defects



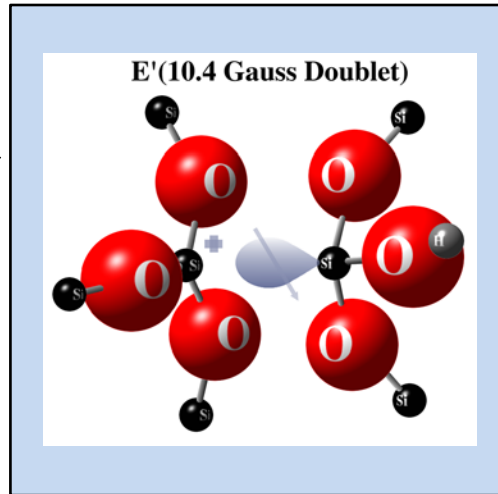
Nonbridging oxygen hole center

*Helms and Poindexter, 1994.*

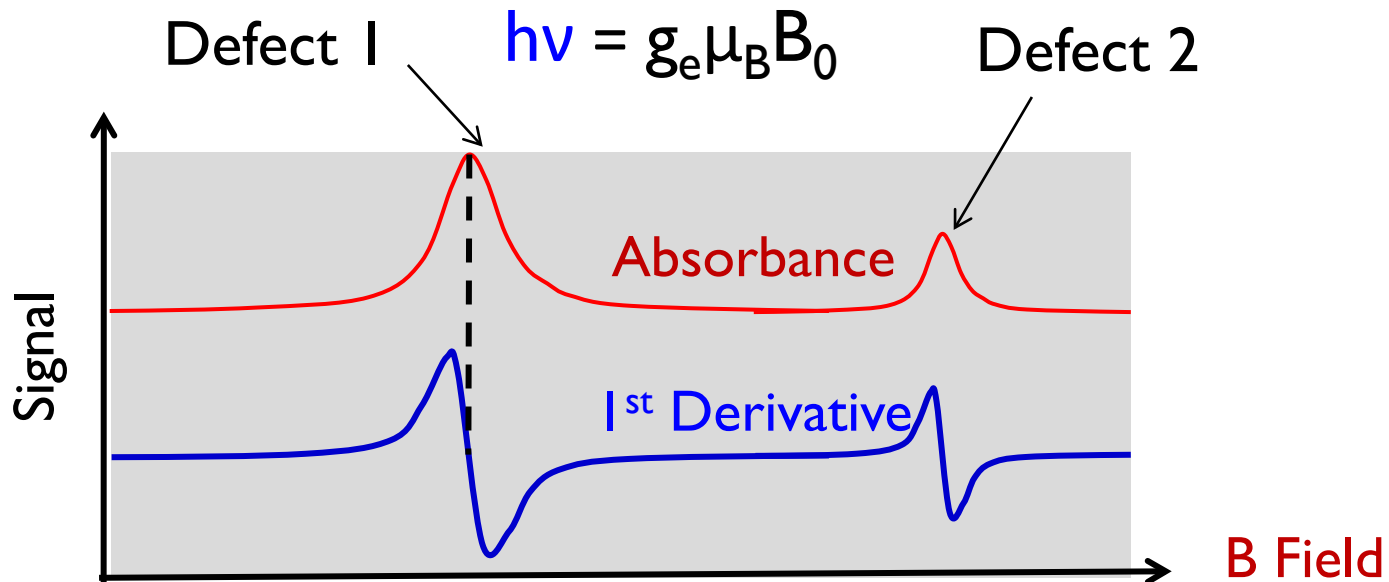
# ESR: a 'microscope' for defects

10 GHz Microwave

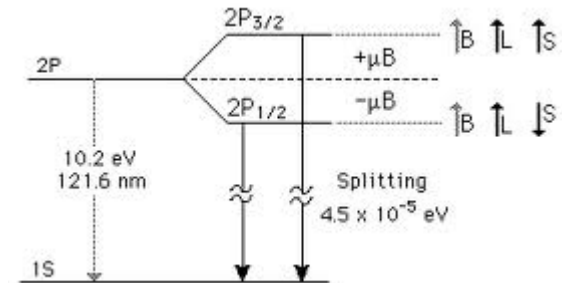
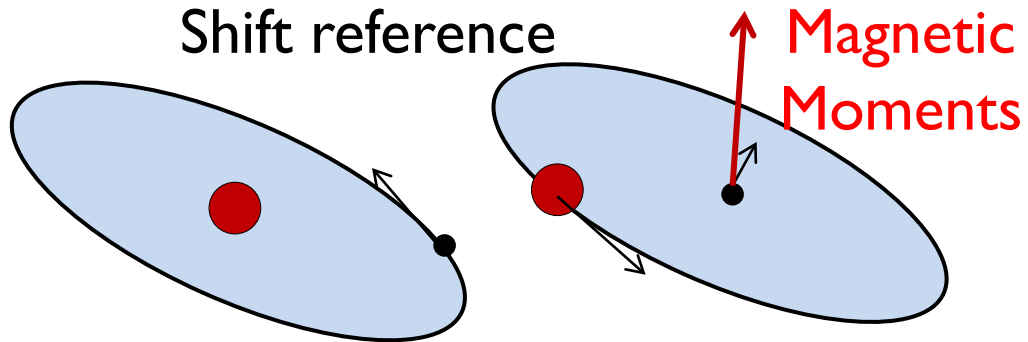
Variable B-field



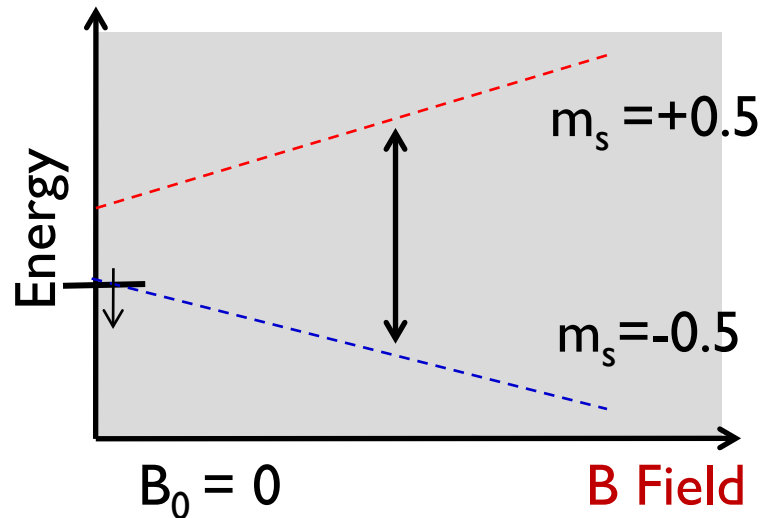
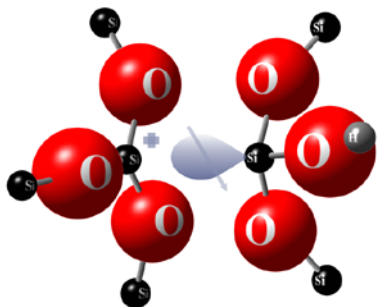
Absorption spectra



# Hyperfine interaction & paramagnetic resonance



$$\Delta E = g_e \mu_B B_0$$

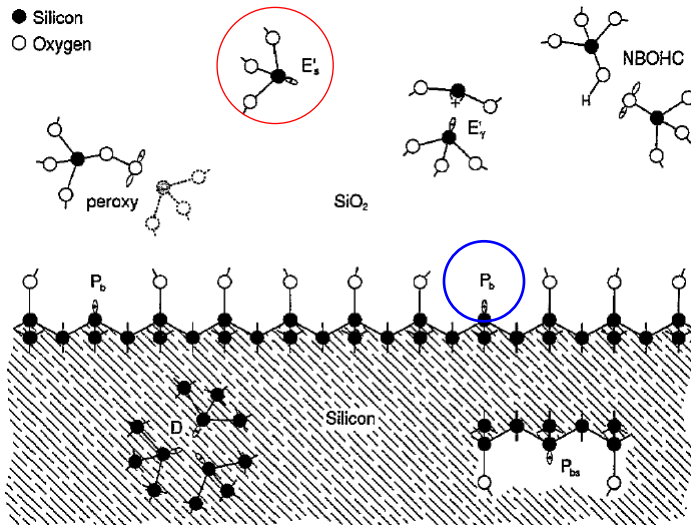
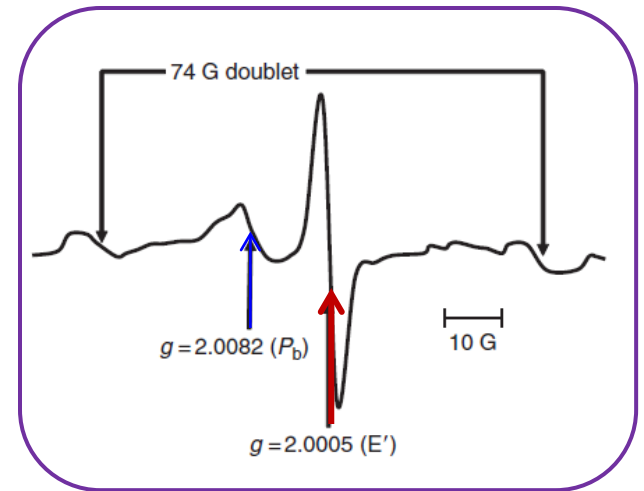
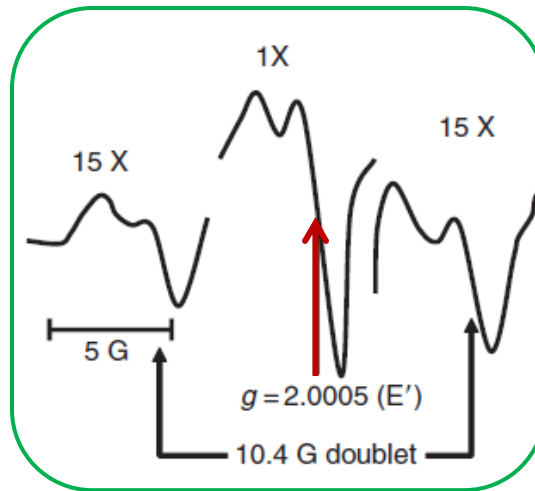
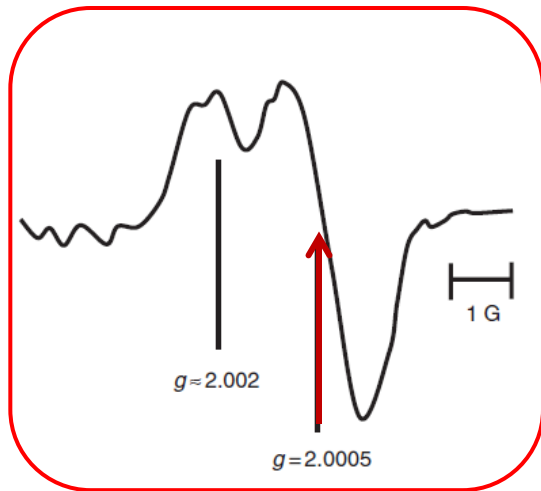


B-value suggests  
local environment

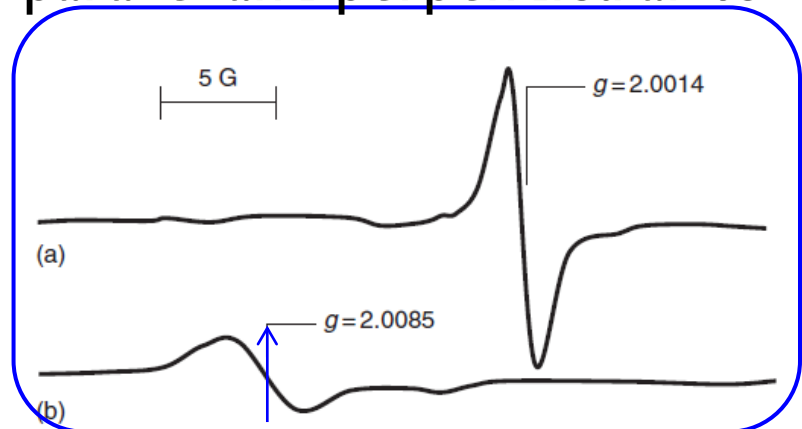
[easyspin.org](http://easyspin.org)



# ESR signature of different defects



Pb parallel and perpendicular to  $\langle 111 \rangle$

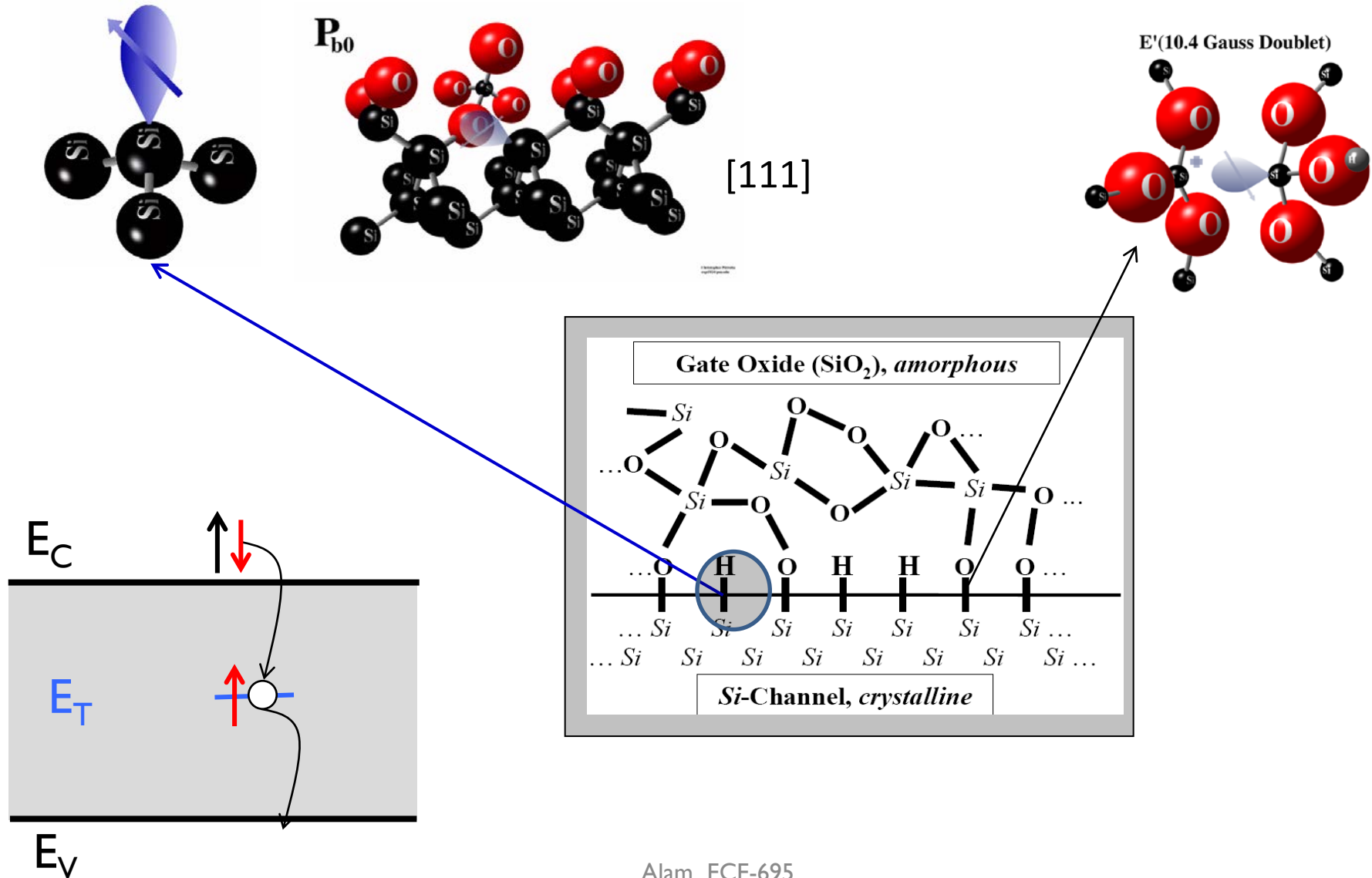


P. M. Lenahan, 2004.

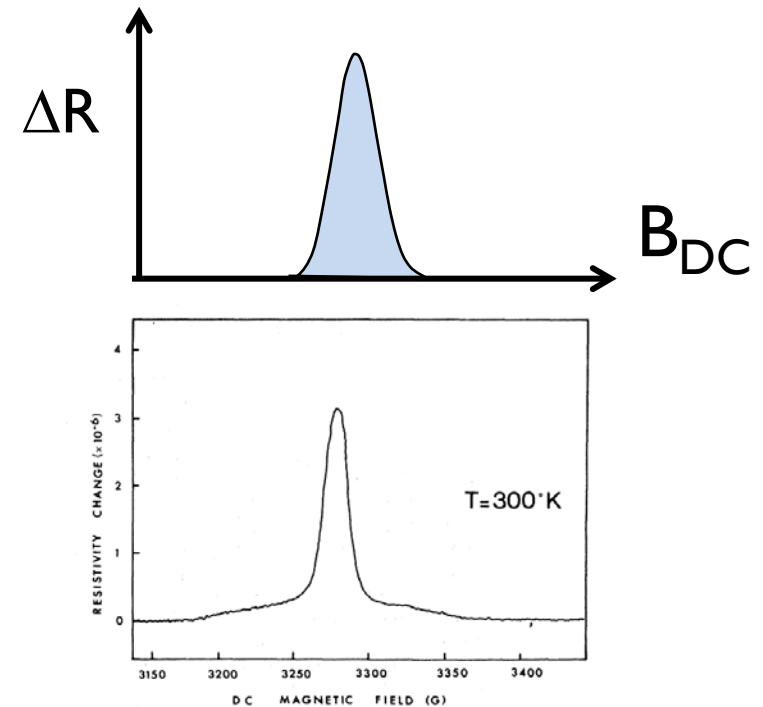
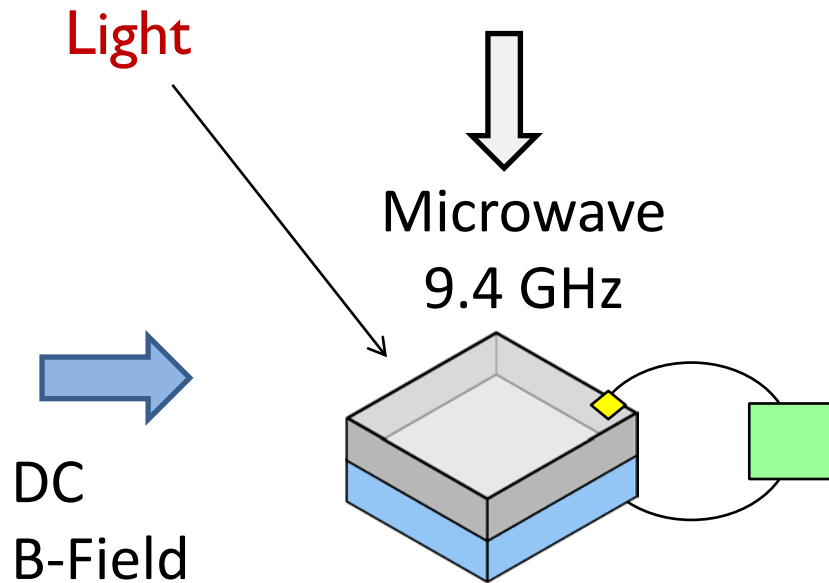
# Outline

1. Importance of measuring interface damage
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# Spin dependent recombination



# Basics of a SDR measurement



$$\frac{d(\Delta n)}{dt} = G - \frac{\Delta n}{\tau(\mathbf{B})} = 0$$

$$\tau(\mathbf{B}) \equiv (\sigma N_{IT} \nu)^{-1}$$

$$\Delta n(B) = G \times \tau(\mathbf{B})$$

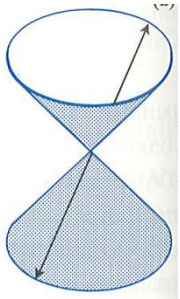
$$\Delta J = J - J_0 = q \Delta n(B) \nu$$

$$= q G \nu \times \tau(\mathbf{B})$$

# Singlet vs. Triplet States: How to make lifetime B dependent

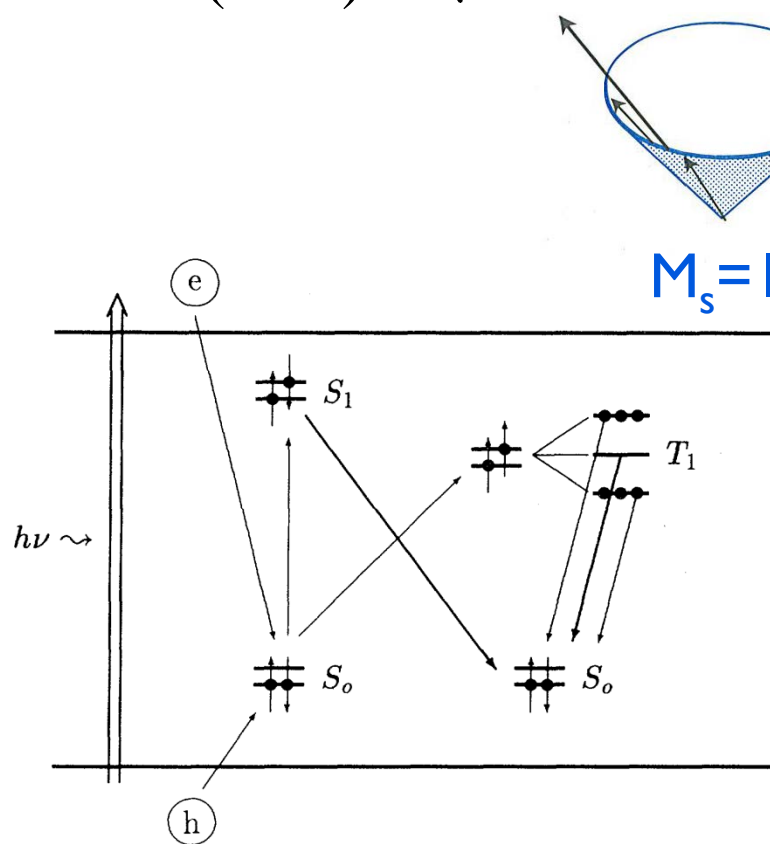
$$\omega_L \propto \frac{\mu_B B_0}{\hbar}$$

$$s(s+1) = \sqrt{3}/2$$

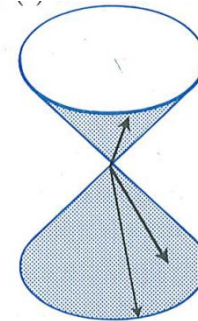


$S=0, M_s=0$

$$s(s+1) = \sqrt{2}$$

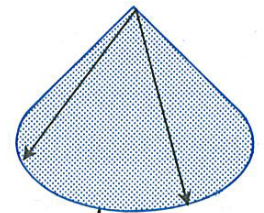


$M_s=1$



$M_s=0$

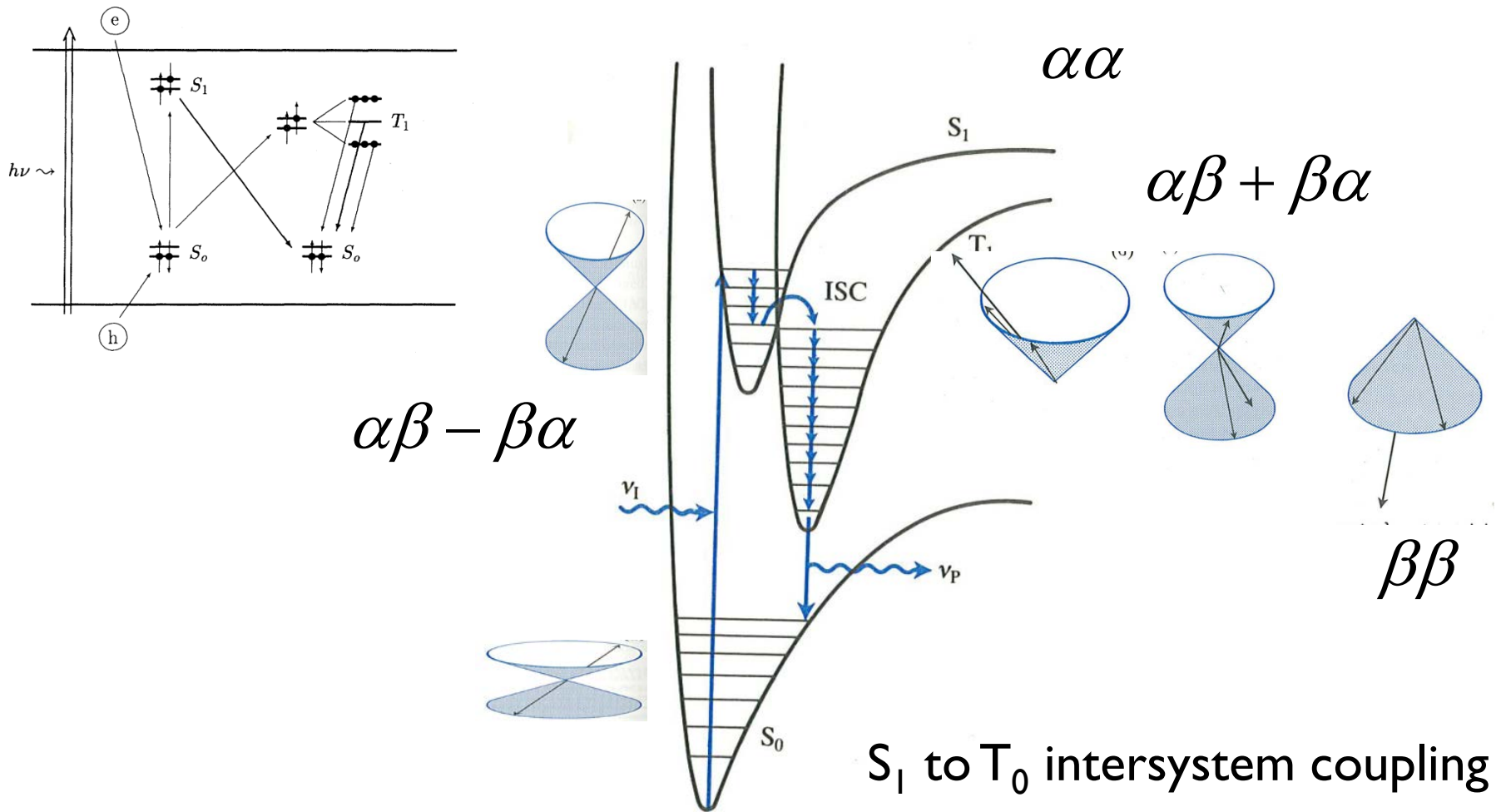
$T_1^{+1}$   
 $T_1^0$   
 $T_1^{-1}$



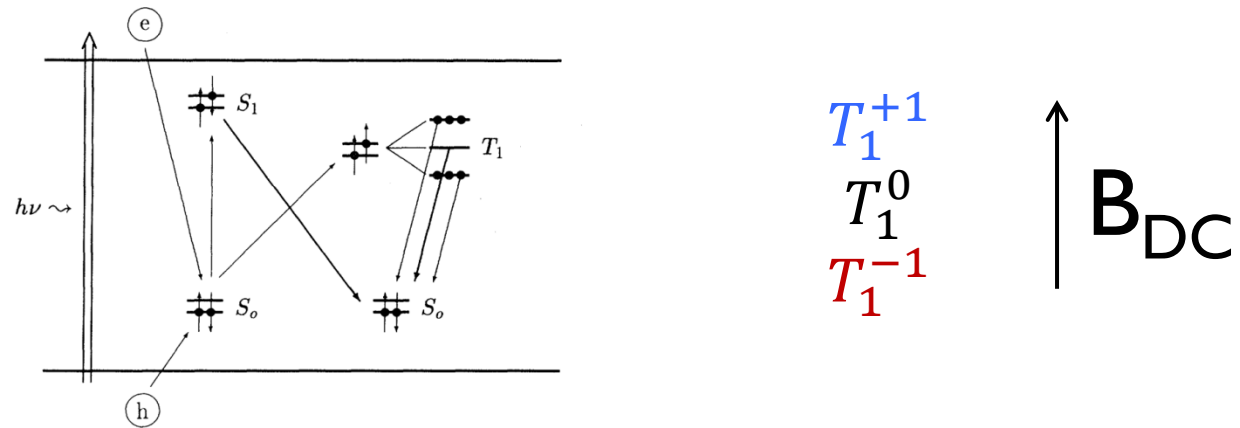
$M_s=-1$

$S=1$

# Another perspective of spin relaxation



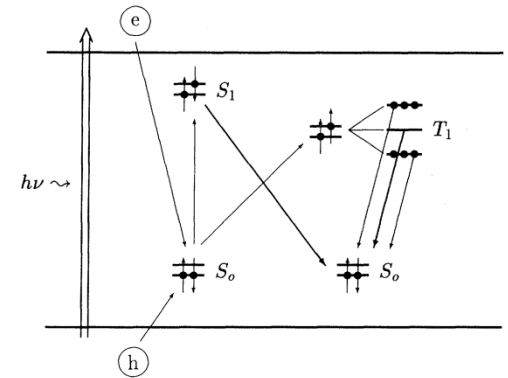
# Rules of spin relaxation



1.  $S_0$  to  $S_1$  transition is allowed both ways.
2.  $S_1$  to  $T_1^0$  transition through intersystem coupling (note  $\Delta S_Z = 0$ )
3. Selection rule allows transition with  $\Delta S_Z = \pm 1$ . Therefore,  $T_1^{+1}$  couples to  $T_1^0$ , and  $T_1^0$  couples to both  $T_1^{+1}$  and  $T_1^{-1}$ , but  $T_1^{+1}$  does not couple to  $T_1^{-1}$ .
3. Coupling between  $T_1^{+1}$  and  $S_0$  is similar to that of  $T_1^{-1}$  to  $S_0$ , but  $T_1^0$  to  $S_0$  is different.
4. The splitting between  $T_1^{+1}$  and  $T_1^0$  is different from that of  $T_1^{-1}$  and  $T_1^0$ . The field couples to only one group. And this restores equilibrium among the states.

# Derivation of the key result

Trap-  
assisted      Direct



$$\frac{dn_e}{dt} = G - \sigma n_e n_{S_0} - R n_e$$

$$\frac{dn^+}{dt} = \frac{\sigma n_e n_{S_0}}{3} - n^+ R - (n^+ - n^0)W - (n^+ - n^0)B$$

$$\frac{dn^0}{dt} = \frac{\sigma n_e n_{S_0}}{3} - n^0 R_0 - (n^0 - n^+)W - (n^0 - n^-)W - (n^0 - n^+)B$$

$$\frac{dn^-}{dt} = \frac{\sigma n_e n_{S_0}}{3} - n^- R - n^- R - (n^- - n^0)W$$



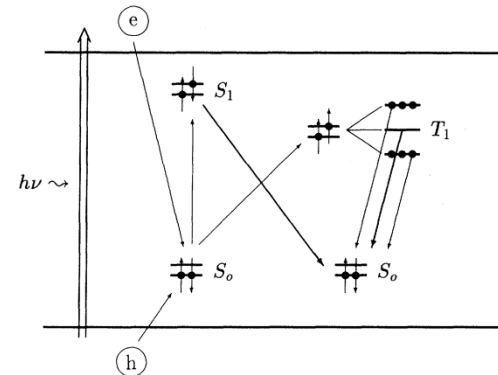
# Derivation ... continued

$$\frac{dn^+}{dt} = \frac{dn^0}{dt} = \frac{dn^-}{dt} = 0$$

$$n_T = n^0 + n^+ + n^- = \frac{\sigma n_e n_{S_0}}{3} F(B)$$

$$f(B) = \frac{(R+W)(2R_0 + R + 9W) + B(R_0 + 5R + 9W)}{(R+W)(RR_0 + 2RW + R_0W) + B(RR_0 + 2RW + R_0W + R^2)}$$

$$G = \frac{\sigma n_e n_D}{1 + \frac{1}{3} \sigma n_e n_D F(B)} + R n_e$$

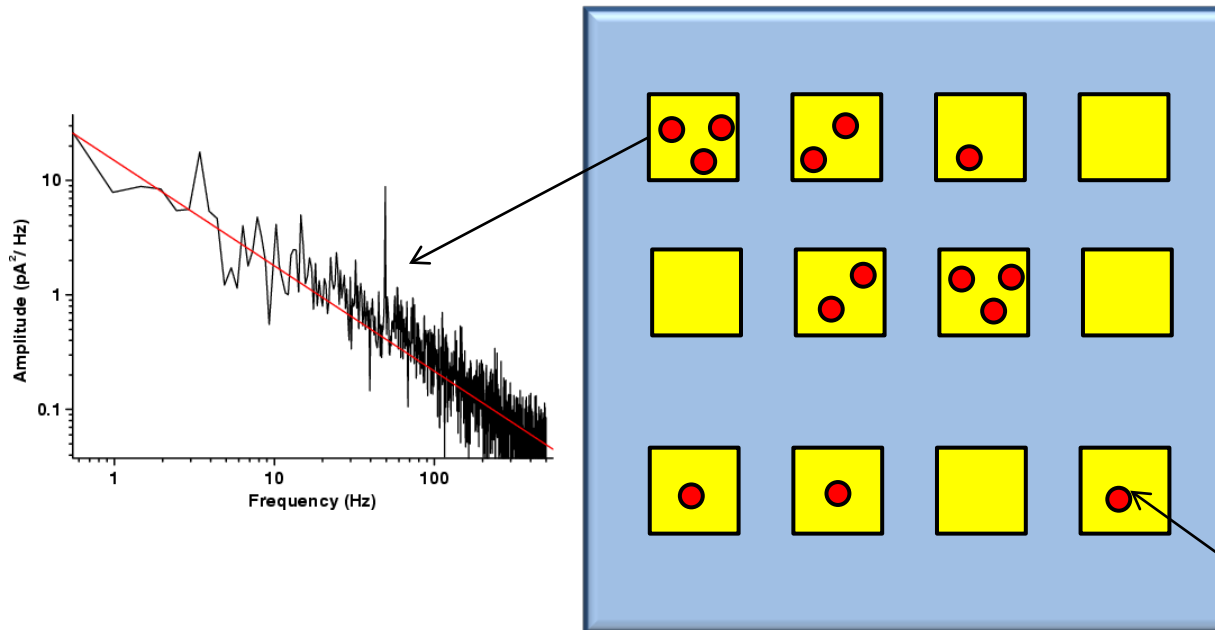


# Outline

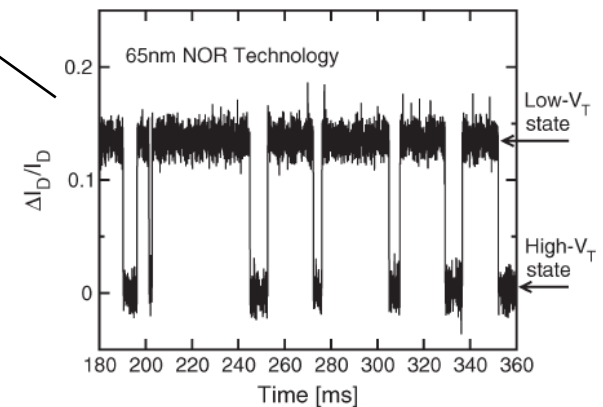
1. Importance of measuring interface damage
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# Recall: Statistics of trapping

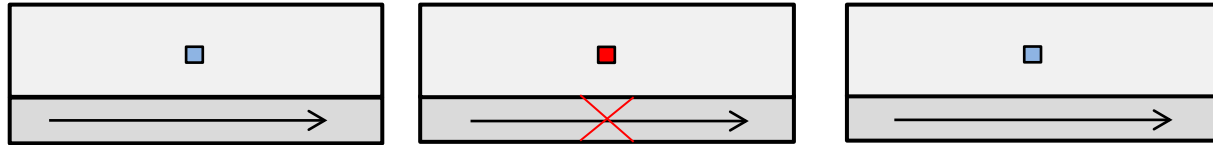
$$10^{17} \text{ cm}^{-3} \times 2\text{nm} \times 100\text{nm} \times 100\text{nm} = \textcolor{red}{2} \text{ traps/device}$$



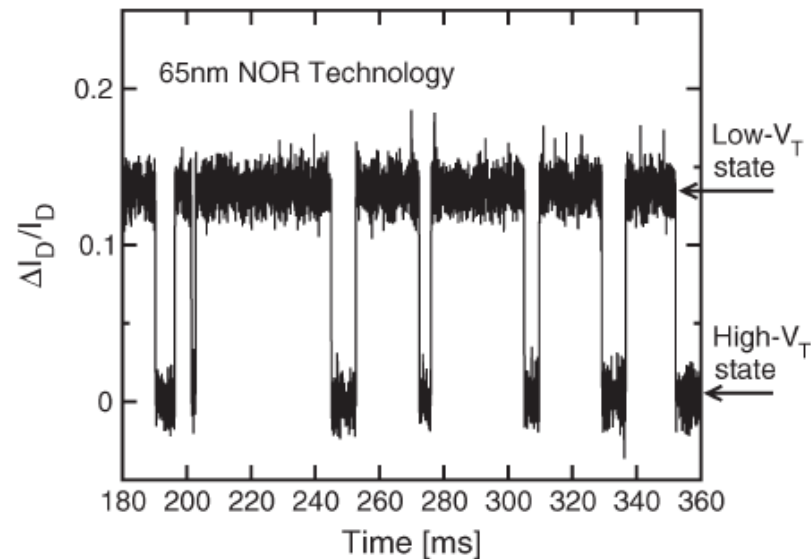
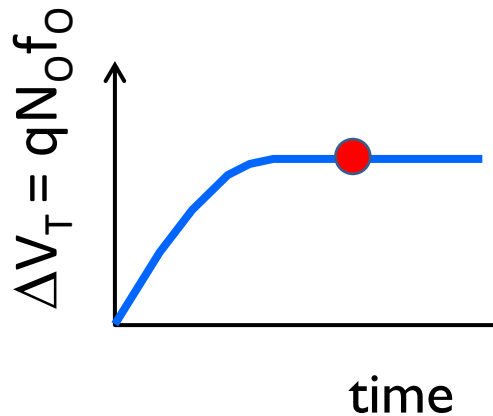
$$P_{\textcolor{red}{n}} = \frac{N^{\textcolor{red}{n}} e^{-N}}{\textcolor{red}{n}!}$$



# Fluctuation in single trap occupation



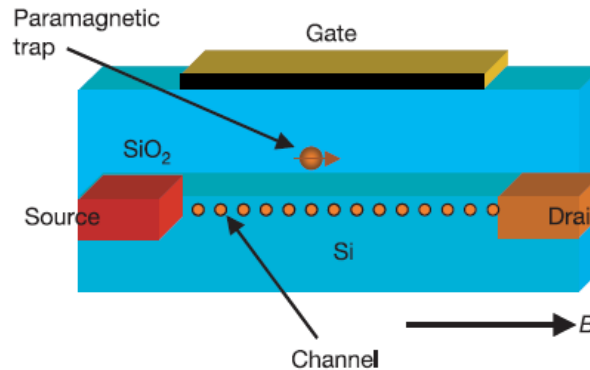
Sodini, IEDM, 2006



$$S(\omega) = \frac{1}{T \rightarrow \infty} \left\langle |F(\omega)|^2 \right\rangle = \left[ \frac{N_0^2 \tau_c}{1 + (\tau_c \omega)^2} \right]$$

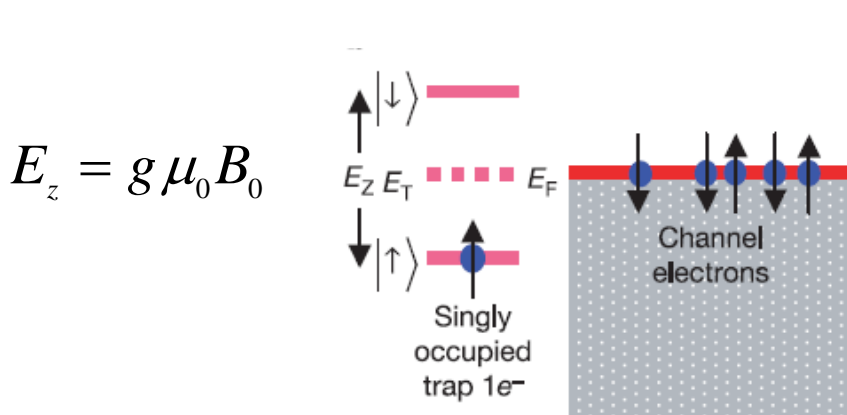
# Electrically detected spin resonance of a single trap in a MOSFET

Ref. Xiao, Nature, 430, 435, 2004.

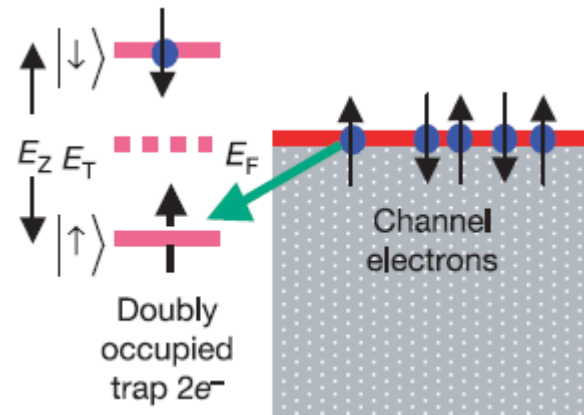


Paramagnetic Trap with single electron

Singly occupied

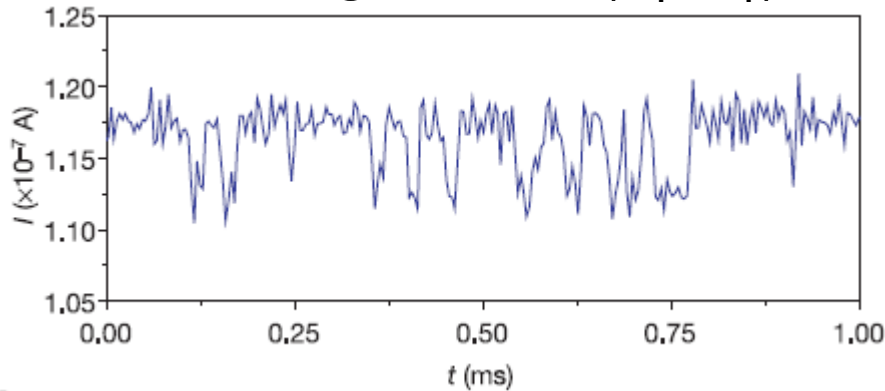


Doubly occupied

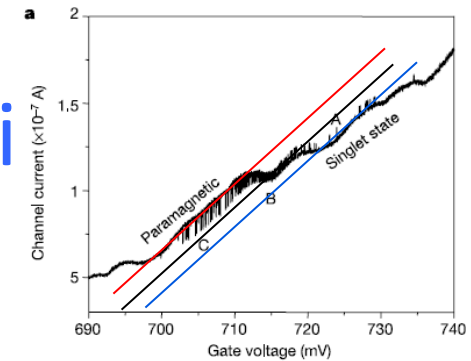
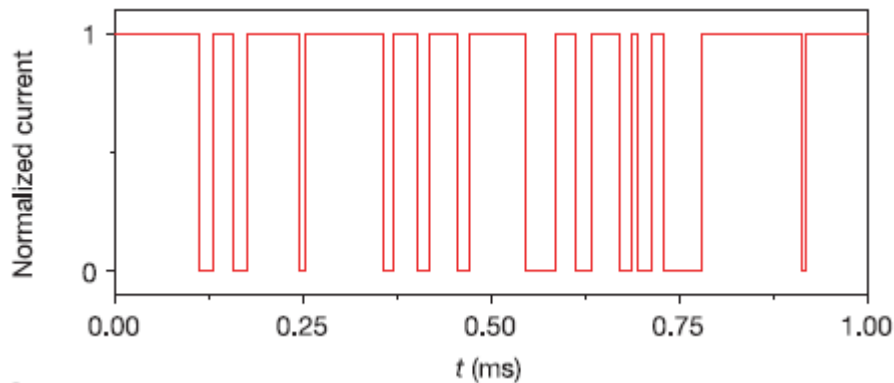


# Energy level splitting of a si

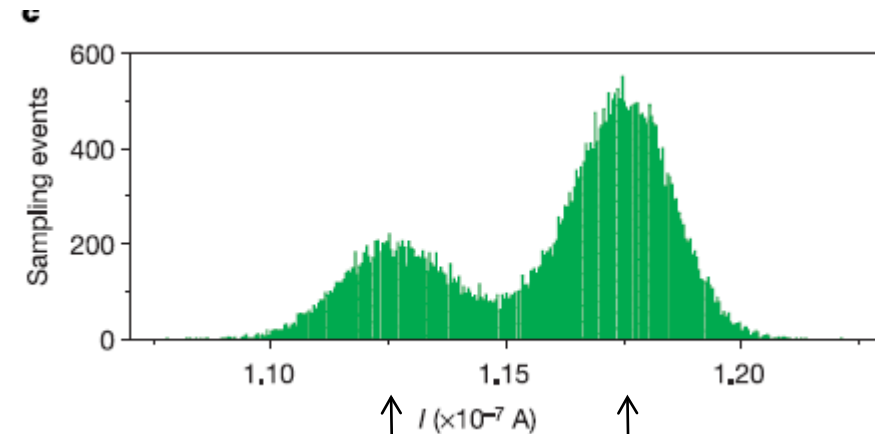
## 1. Original data ( $E_T \sim E_F$ )



## 2. Cleaned data ( $E_T \sim E_F$ )

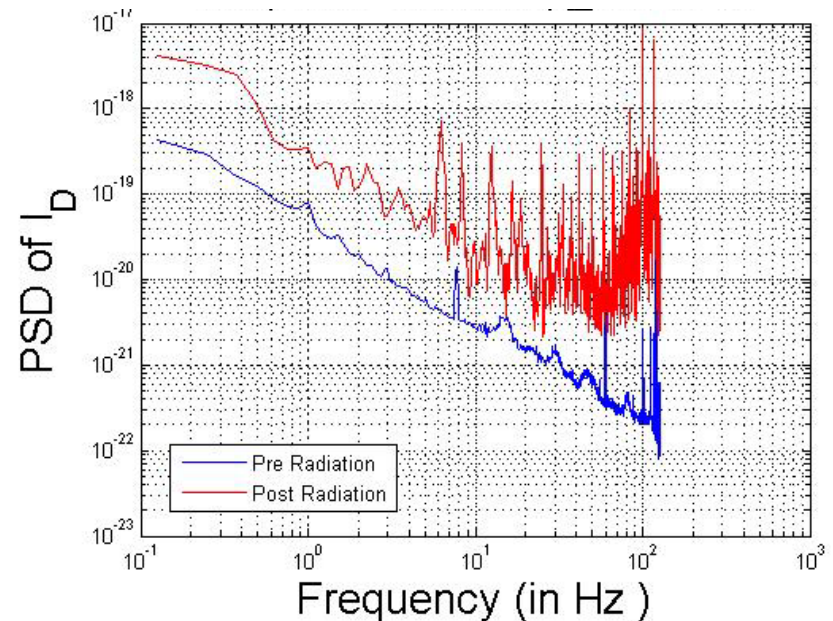
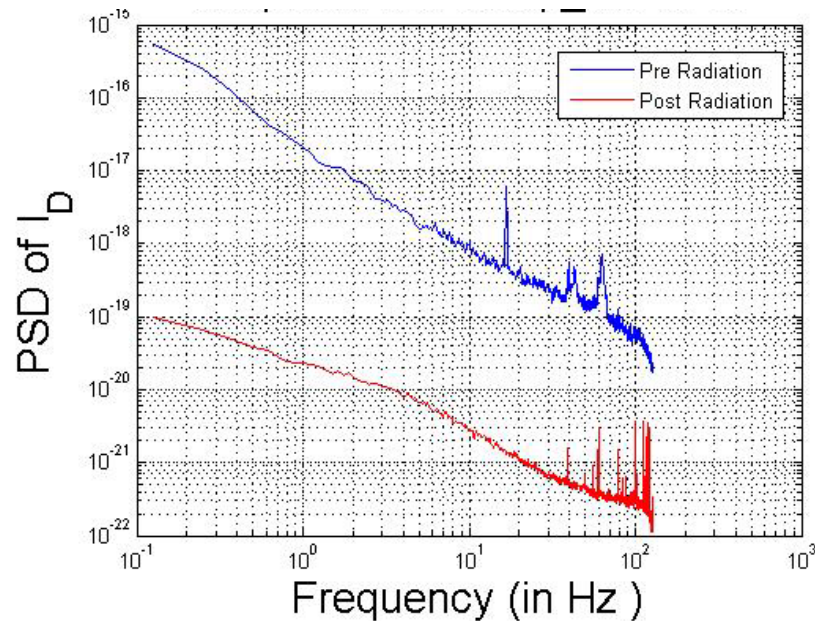


## 2. Statistical Distribution



Doubly occupied      Singly occupied

# Noise as a characterization tool



Overall, noise can be a sensitive monitor of shallow traps

# Conclusions

- ❑ ESR has long been a powerful tool for chemists and physicists, SDR has a long history as well.
- ❑ The modern synthesis of SDR and ESR in the form of EDSR has become a powerful tool.
- ❑ The approach is sensitive, the key challenge is complexity of setup. Cannot be used for routine characterization of wafer.
- ❑ A combination of analytical tools allow fundamental understanding of a given reliability problem. Once this goal is achieved, classical tools are often more convenient.



# References

- *I followed the derivation of L.S. Vlasenko et al. Electron Paramagnetic Resonance vs. spin-dependent recombination: Excited triplet states of structural defects in irradiated silicon, PRB, 52(2), 1995. p. 1144. The original work was done by D. Lepine, Spin-Dependent recombination on silicon surface, PRB, 6(2), 1972. p. 436. The first correct formulation is due to D. Kaplan, I. Solomon, and N. Mott, J. Phys. Lett. 39, 51 ~1978 – although they emphasized direct e-h recombination, not the trap-assisted generalization we discussed.*
- *C. Boehme and K. Lips – Theory of time-domain measurement of spin-dependent recombination with pulse electrically detected magnetic resonance”, PRB, 68, 245105, 2003 provides a more sophisticated quantum mechanical treatment of the problem.*
- *The EDMR (electrically detected magnetic resonance) technique has become sophisticated enough to detect single spin– see, L. Martin, “A scheme for electrical detection of single-electron spin resonance”, PRL, 018301, 2003. T. Wimbauer et al., Defects in planar silicon p-n junctions studied with electrically detected magnetic resonance. APL, 76(16), 2280, 2000. M. Xiao et al., “Electrical Detection of the spin resonance of a single electron in a silicon field-effect transistor, Nature, p. 435, 2004.*
- *K. Hung et al., Random Telegraph Noise in Deep-Submicrometer MOSFETs, 11(2), p. 90, 1990.. B. Gross and Charles G. Sodini. 1/f noise in MOSFETs with ultrathin gate dielectrics. In IEDM Technical Digest, pages 881-884, 1992.*

# Self-Test Questions

- Q1. If a signal disappears from ESR because of negative-U configuration, can it be detected by SDR or EDSR methods?
- Q2. What is the relationship between Gauss and Tesla as units of magnetic field?
- Q3. Was the original SDR method for bulk or interface traps?
- Q4: What is the relationship between RTN noise spectra and EDSR spectra? At what point will they be substantially different?
- Q5. For single spin, we have 2 states. at the ground level. For two spins, we have 4 states (singlet and triplet). How many spin states do you expect for a 3-spin system?
- Q6. What are the advantages and disadvantages of spin-based detection?
- Q7. What is the difference between hyperfine interaction and hyper-polarizability?
- Q8. If ESR experiments are used for NBTI degradation, what type of time exponent would you expect? What about HCI degradation ?