

Modeling of Quantum Cascade Laser Sources with Giant Optical Nonlinearities

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Overview

- Modeling of quantum cascade lasers
- Inclusion of optical cavity field
- THz difference frequency generation in QCL structures
- Mode-locked QCLs and frequency combs
- Conclusion

Quantum Cascade Laser

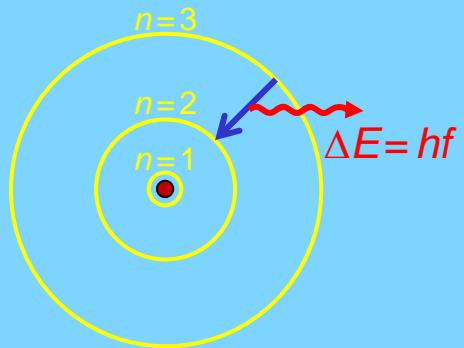
Conventional lasers / light sources

Use optical transitions in atoms, molecules, lattices,...

➤ Usually in infrared, visible or ultraviolet regime

⇒ **Scientifically underdeveloped terahertz gap**

⇒ **No practical compact diode lasers in mid-infrared**

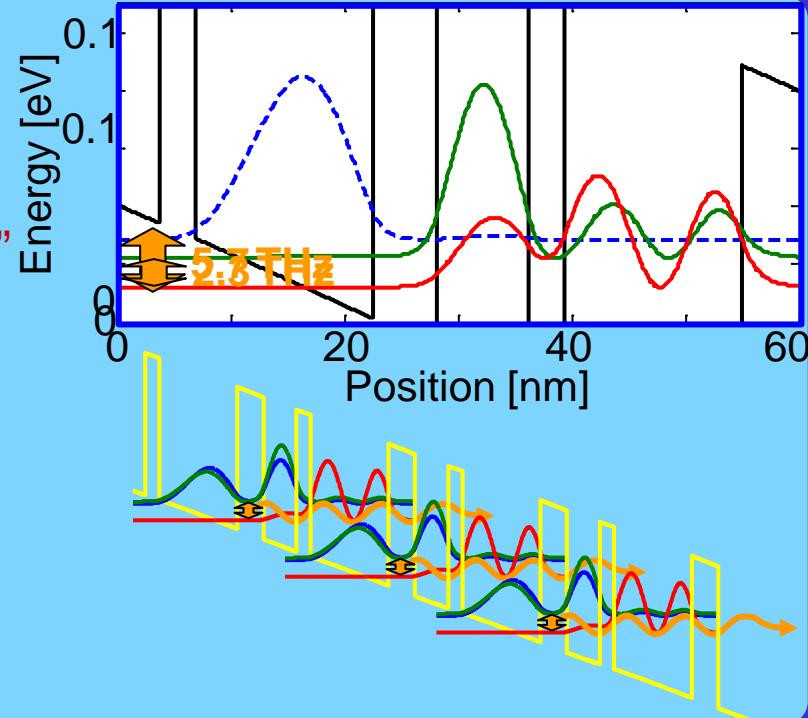


Quantum cascade laser (QCL)

Use nanostructure as “artificial atom”

➤ Wavelength does not depend on material, but can be tailored by “quantum engineering”

⇒ **QCL covers terahertz and mid-infrared**



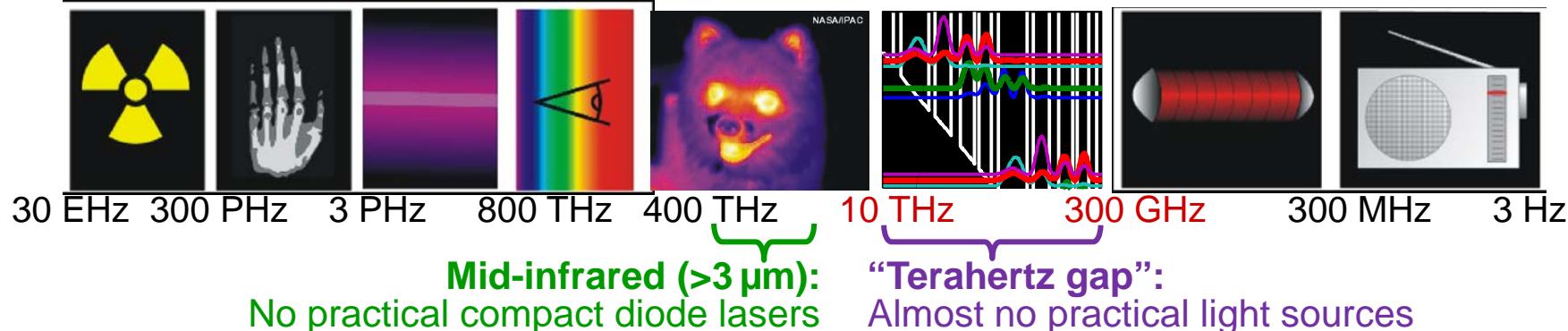
Use many transitions in a series (“cascade”)

➤ A single electron can emit multiple photons

⇒ **Increased optical power and efficiency**

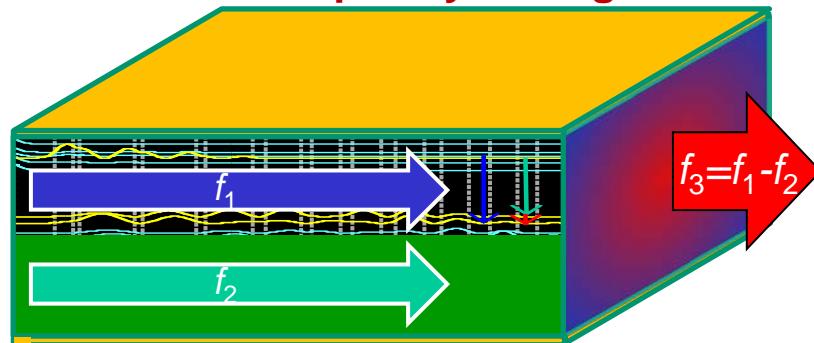
Quantum Engineering of Active Region

- Optical gain characteristics can be custom-tailored
⇒ Mid-infrared and THz ranges become accessible

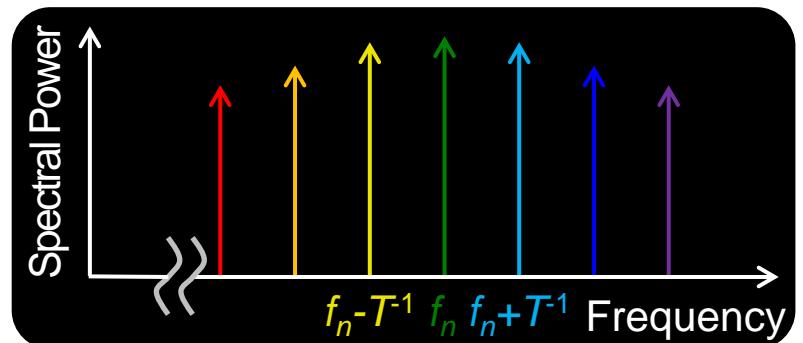


- Artificial giant optical nonlinearities can be integrated

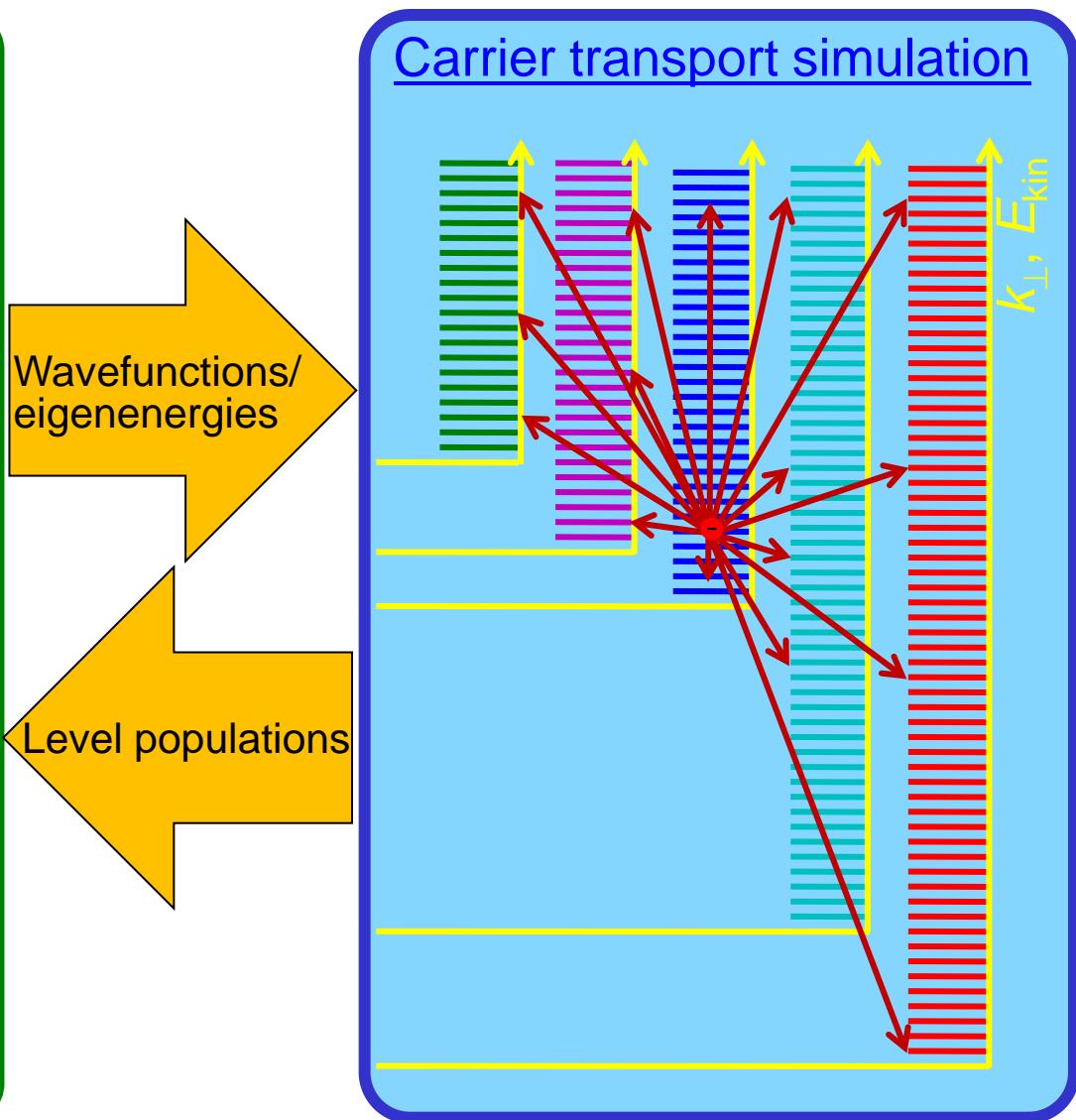
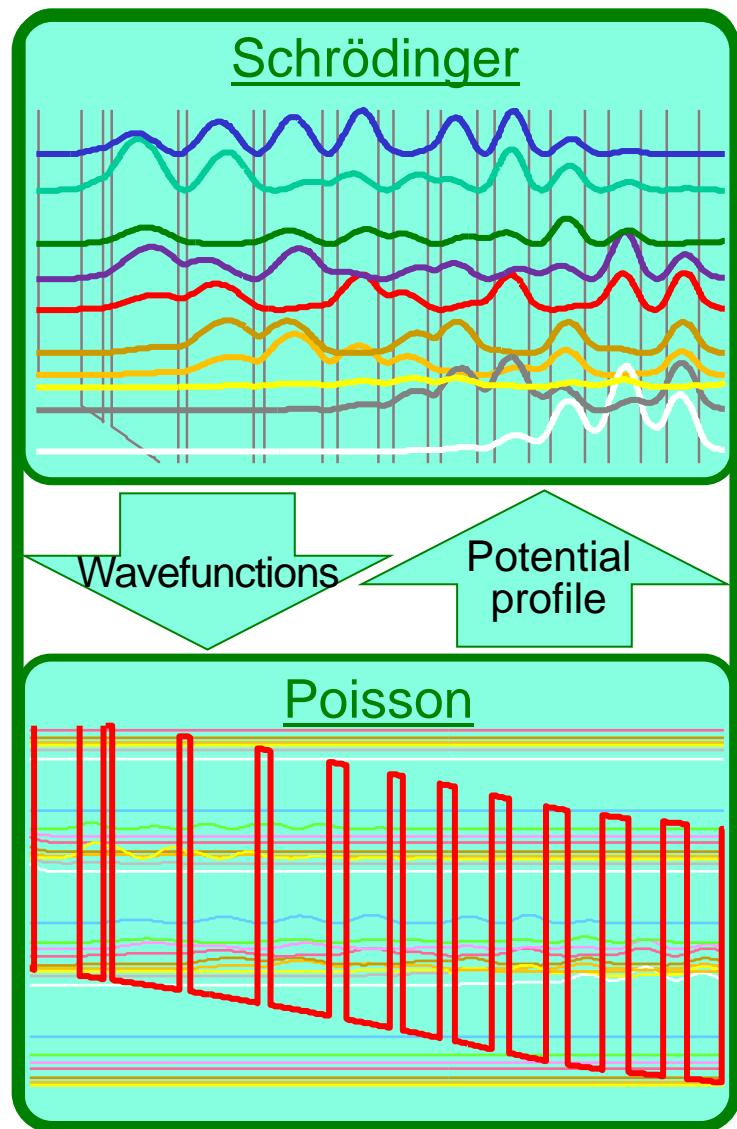
Frequency conversion structures
➤ Based on frequency mixing



Mode-locking & frequency combs
➤ Based on nonlinear coherent interaction



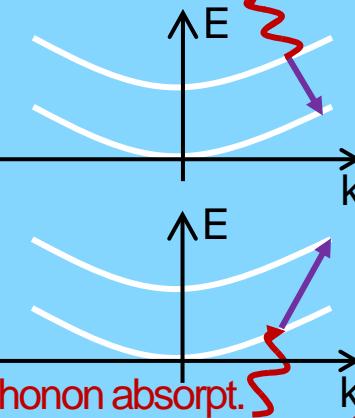
Ensemble Monte Carlo (EMC)



Boltzmann Equation and Scattering

Optical phonons

Phonon emission

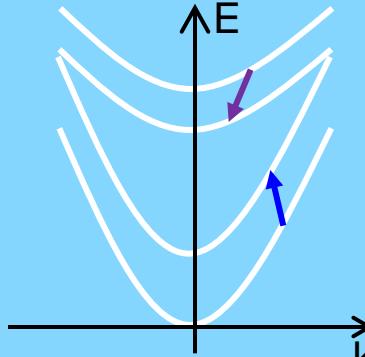


Boltzmann equation for carrier distribution function $f_{n,k}(t)$:

$$\frac{df_{n,k}}{dt} = \sum_s \sum_{m,k'} (W_{mkmk'} f_{m,k'} - W_{nkmk'} f_{n,k})$$

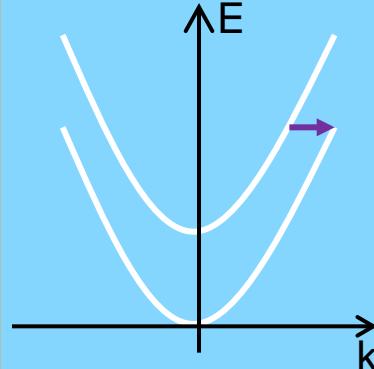
$W_{nkmk'}$

Electron-electron



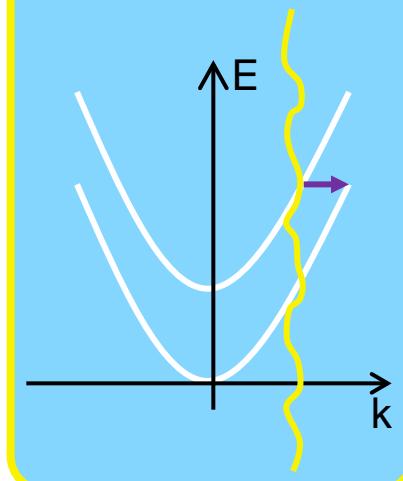
Acoustic phonons

↑E



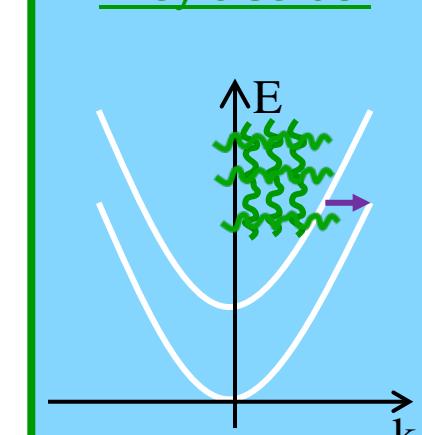
Interface roughness

↑E



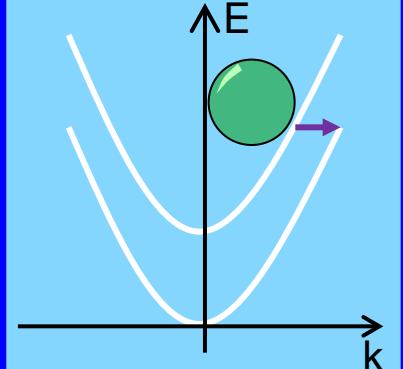
Alloy disorder

↑E



Impurity

↑E

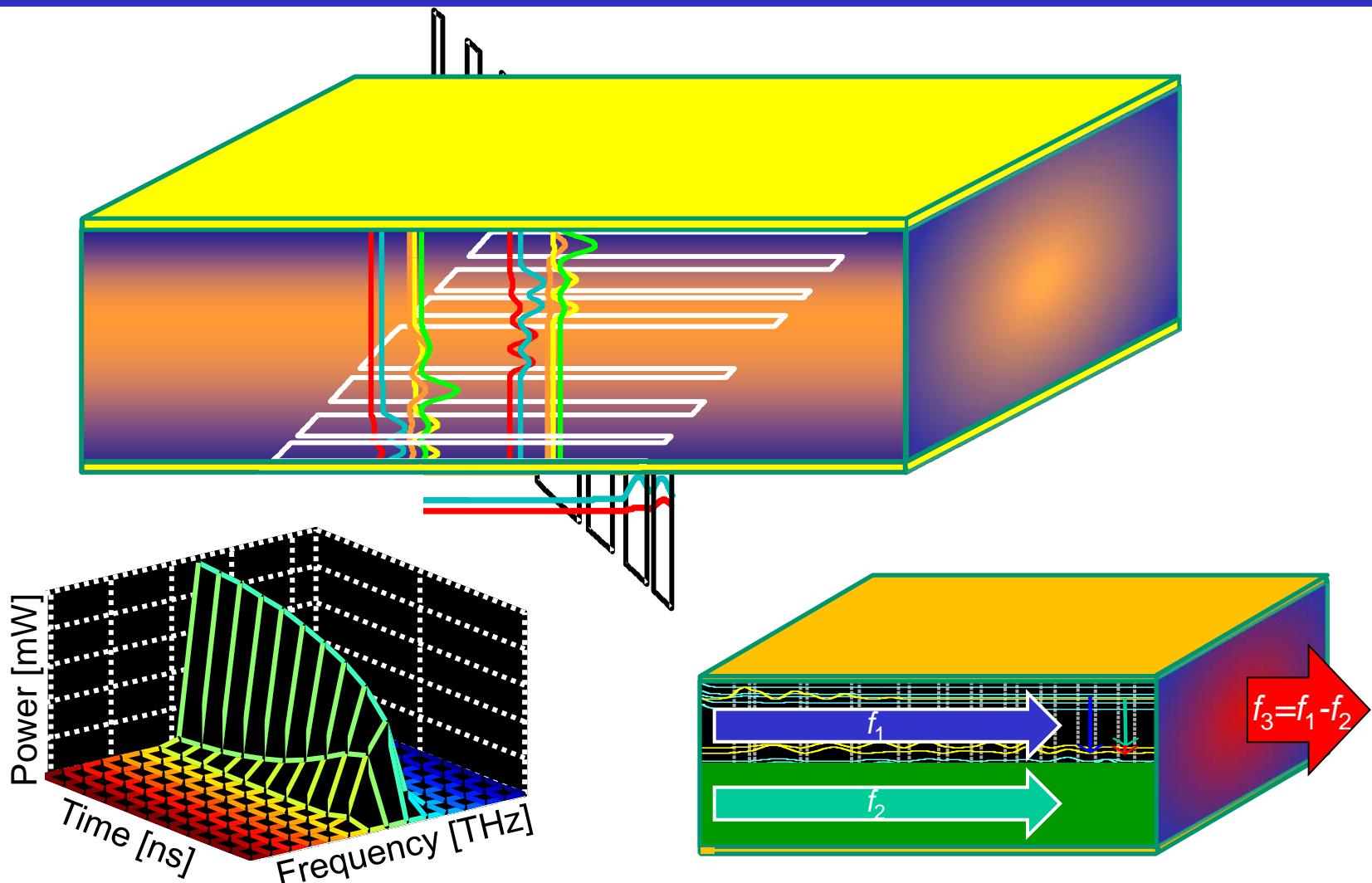


Overview

- Modeling of quantum cascade lasers
- **Inclusion of optical cavity field**
- THz difference frequency generation in QCL structures
- Mode-locked QCLs and frequency combs
- Conclusion



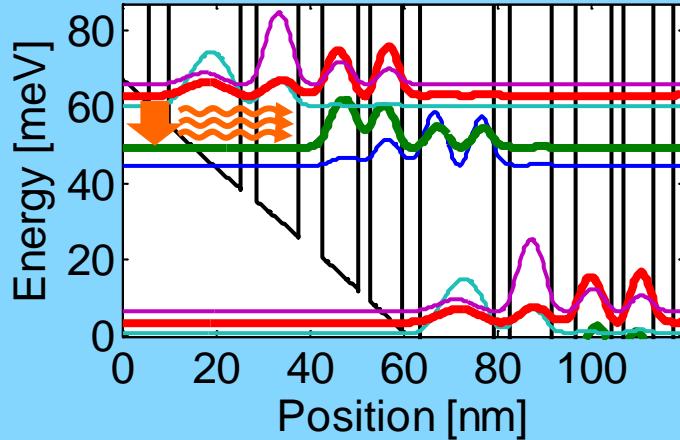
Inclusion of Optical Cavity Field



C. Jirauschek, Appl. Phys. Lett. **96**, 011103 (2010)
A. Mátyás et al., J. Appl. Phys. **110**, 013108 (2011)
C. Jirauschek, Opt. Express **18**, 25922 (2010)

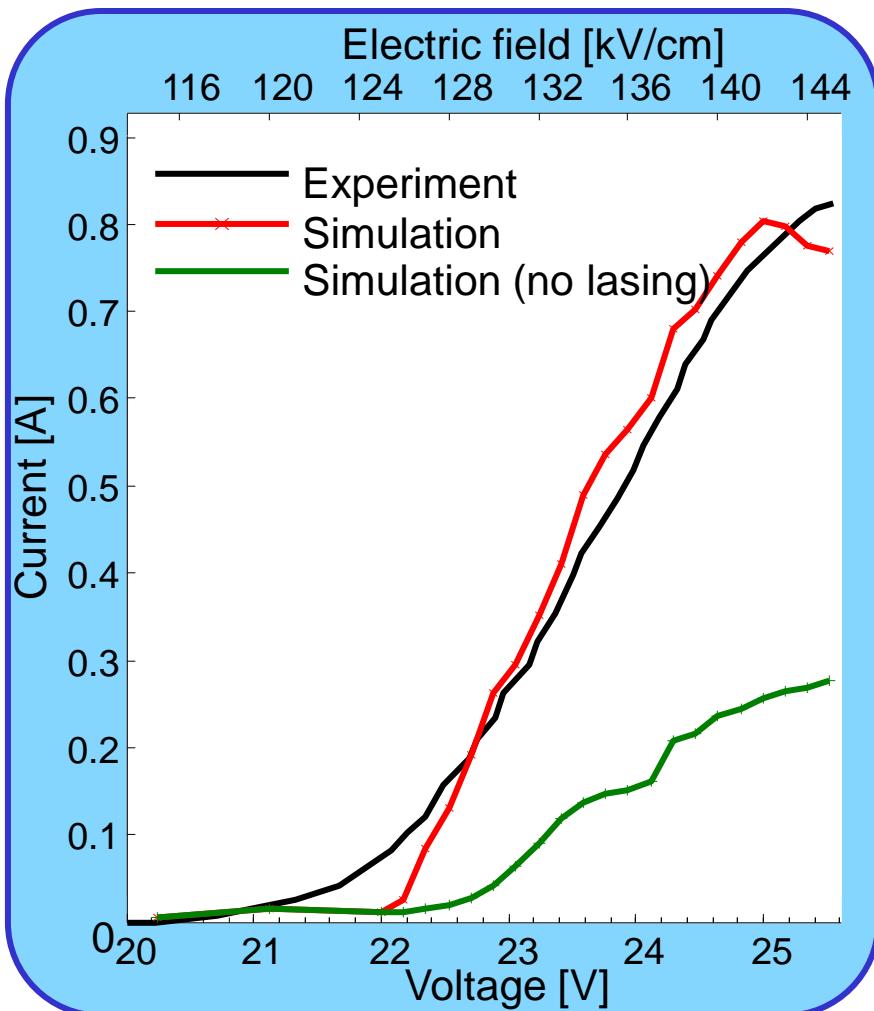
C. Jirauschek et al., Opt. Express **21**, 6180 (2013)
C. Jirauschek et al., Opt. Express **23**, 1670 (2015)

Carrier-Light Coupling in Monte Carlo

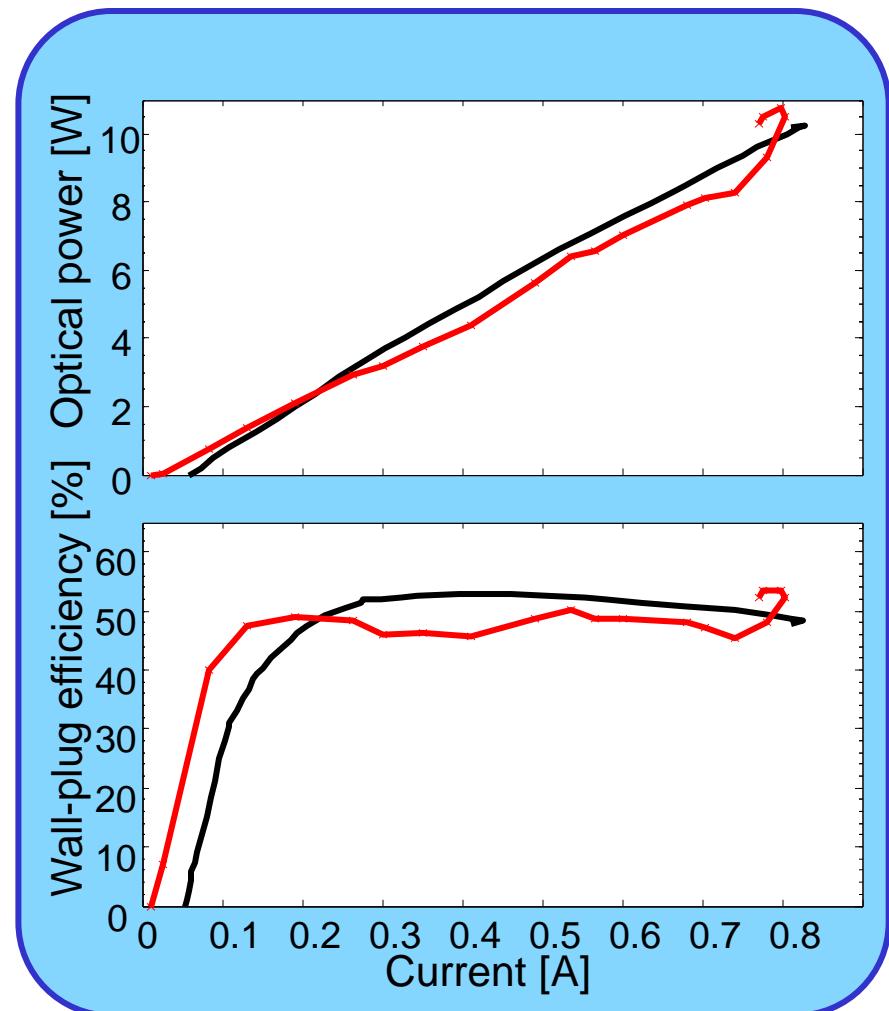


$$g(\omega) = \frac{e^2 \omega}{c \epsilon_0 n_0 \hbar L} n_E^{2D} \sum_{\substack{i,j \\ E_i > E_j}} |z_{ij}|^2$$
$$\times \int_0^\infty d\varepsilon \frac{[f_i(\varepsilon) - f_j(\varepsilon)] \gamma_{ij}(\varepsilon)}{\gamma_{ij}^2(\varepsilon) + [\omega - (E_i - E_j)^2 / \hbar^2]^2}$$

Simulation of High Efficiency Mid-Infrared QCL



Design: Y. Bai et al., Nat. Photonics **4**, 99 (2010)

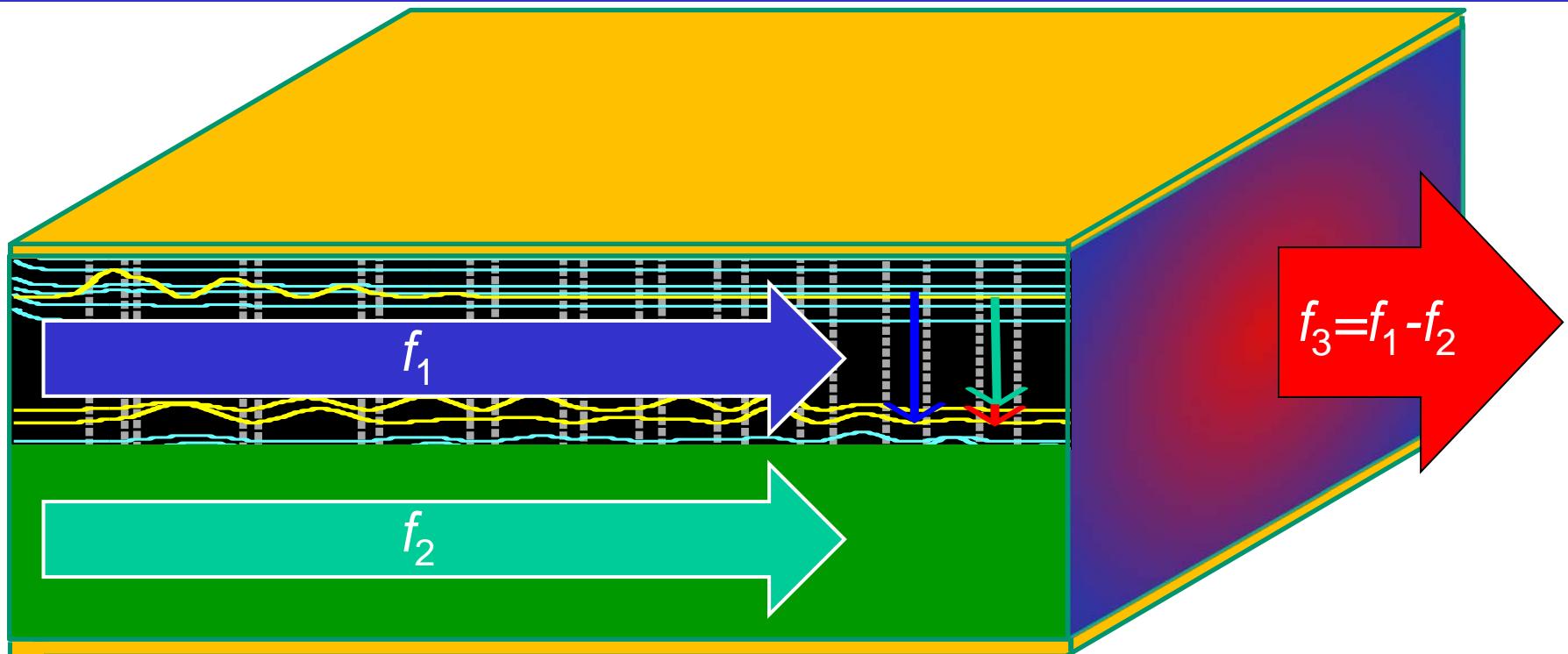


A. Matyas et al., J. Appl. Phys. **110**, 013108 (2011)

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THz Difference Frequency Generation QCL Structure



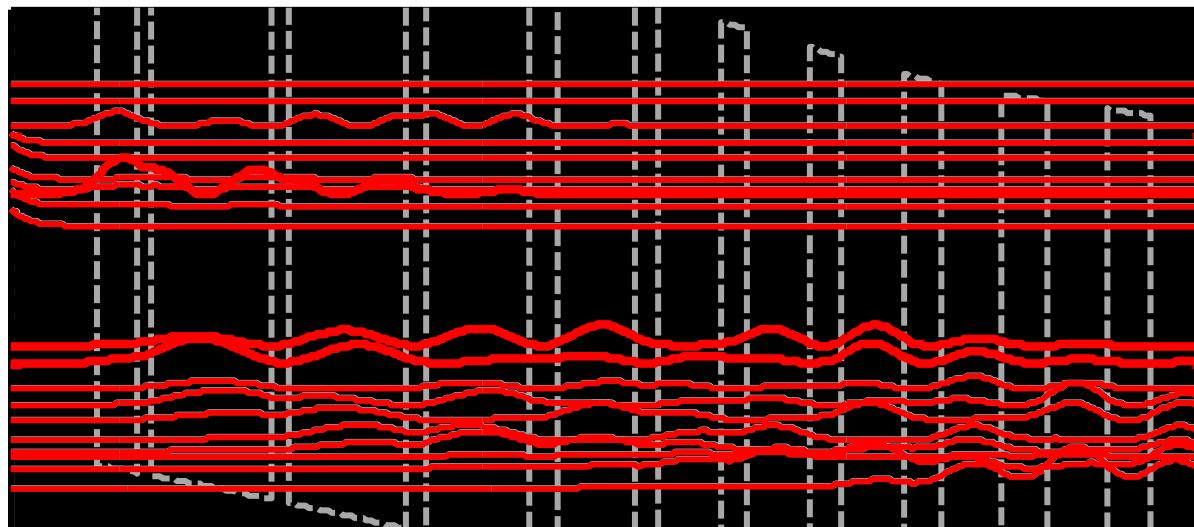
Ideal THz source

- Room temperature operation
- Broadband tunability
- THz output power in mW range

THz DFG QCL source

- M. Belkin et al., Appl. Phys. Lett. **92**, 201101 (2008)
1.0-4.6 THz
Q.Y. Lu et al., Appl. Phys. Lett. **101**, 251121 (2012)
1.7-5.25 THz
K. Vijayraghava et al., Nature Comm. **4**, 2021 (2013)
1.9 mW pulsed, 3 μ W cw (at room temperature)
M. Razeghi et al., Opt. Express **23**, 8462 (2015)

Modeling of Nonlinear Susceptibility



$$\chi^{(2)} = \frac{1}{\hbar^2 \varepsilon_0 L_P} \sum_{\ell,m,n} d_{\ell m} d_{m n} d_{n \ell} n_E^{2D} \int_0^\infty f_\ell(K_{\ell m n} - K_{m \ell n}) d\varepsilon,$$

$$K_{\ell m n} = \left(\frac{1}{\omega_{n \ell} - i\gamma_{n \ell} - \omega} + \frac{1}{\omega_{n m} + i\gamma_{n m} + \omega} \right) \left(\frac{1}{\omega_{m \ell} - i\gamma_{m \ell} + \omega_2} + \frac{1}{\omega_{m \ell} - i\gamma_{m \ell} - \omega_1} \right)$$

L_P :

ε :

$f_m(\varepsilon)$:

$n_E^{2D} = m^*/(\pi\hbar^2)$: 2D density of states

Period length

Kinetic energy

Distribution function

2D density of states

$\gamma_{m n}(\varepsilon)$: Optical linewidth

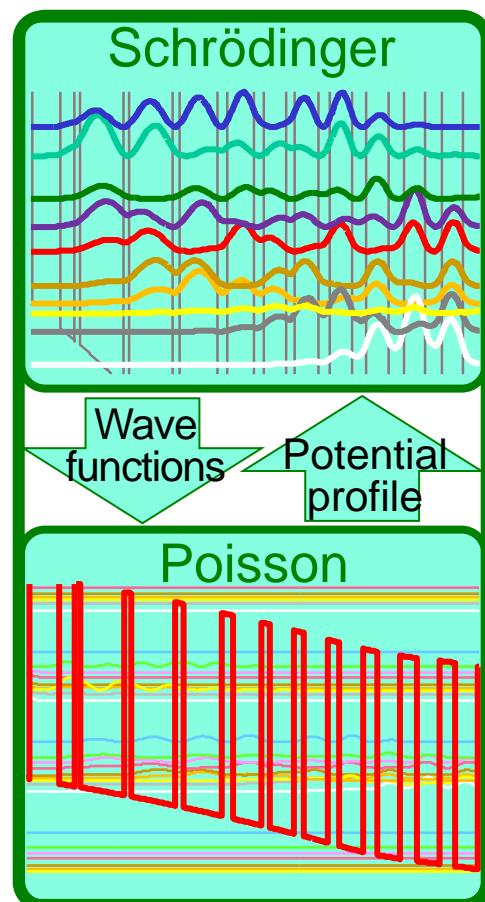
$d_{m n}$: Dipole matrix element

$\omega_{m n}$: Resonance frequency

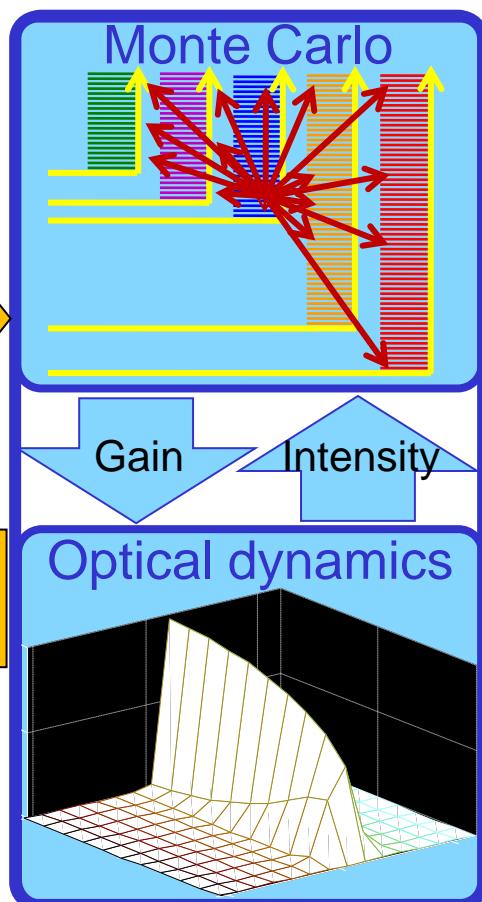
$\omega_{1,2}/\omega$: Mid-IR/THz frequencies

Multi-Domain Simulation Approach

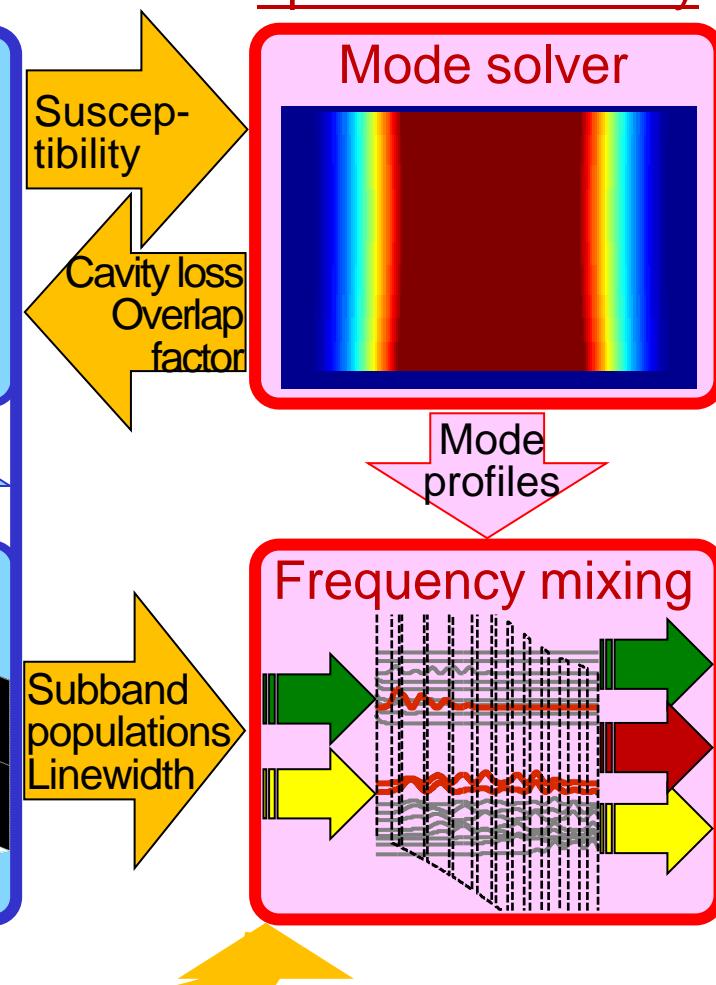
Quantized energy states



Carrier and photon dynamics



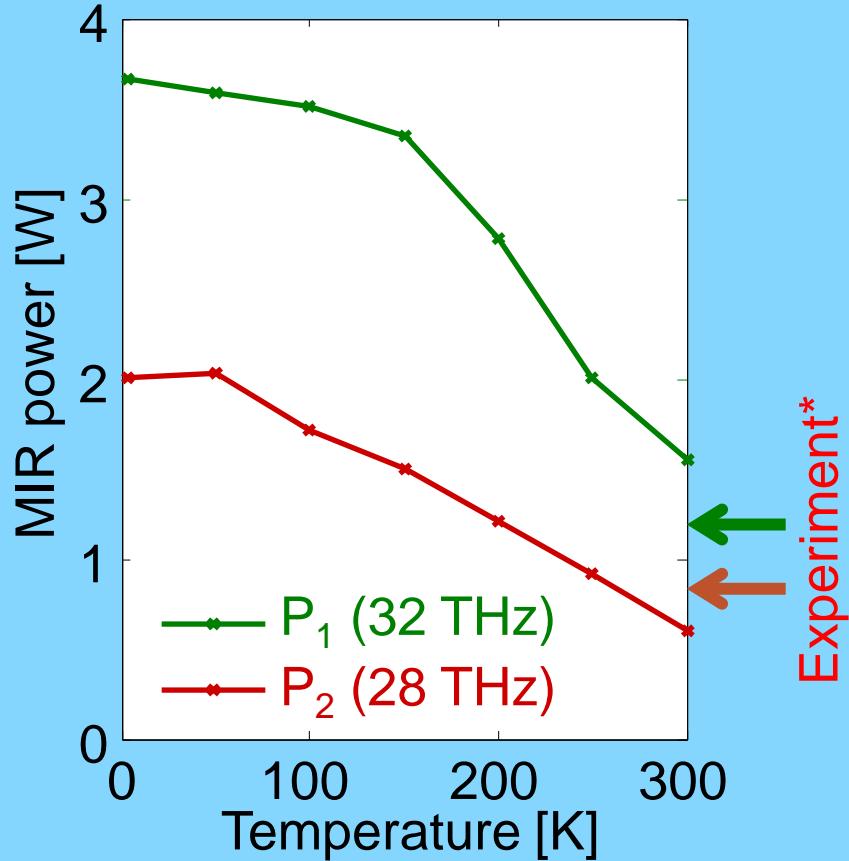
Resonator and optical nonlinearity



Wave functions, energies

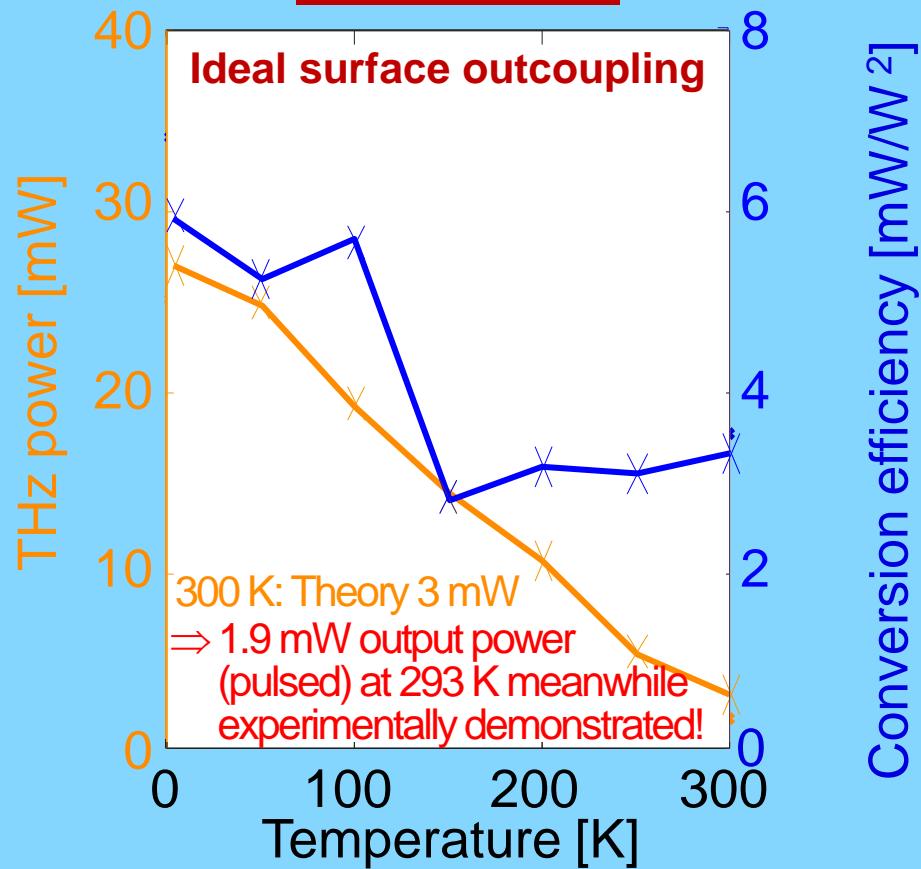
Comparison Simulation - Experiment

Mid-infrared results



Experiment*

THz results



*Q. Y. Lu et al., Appl. Phys. Lett. **99**, 131106 (2011)

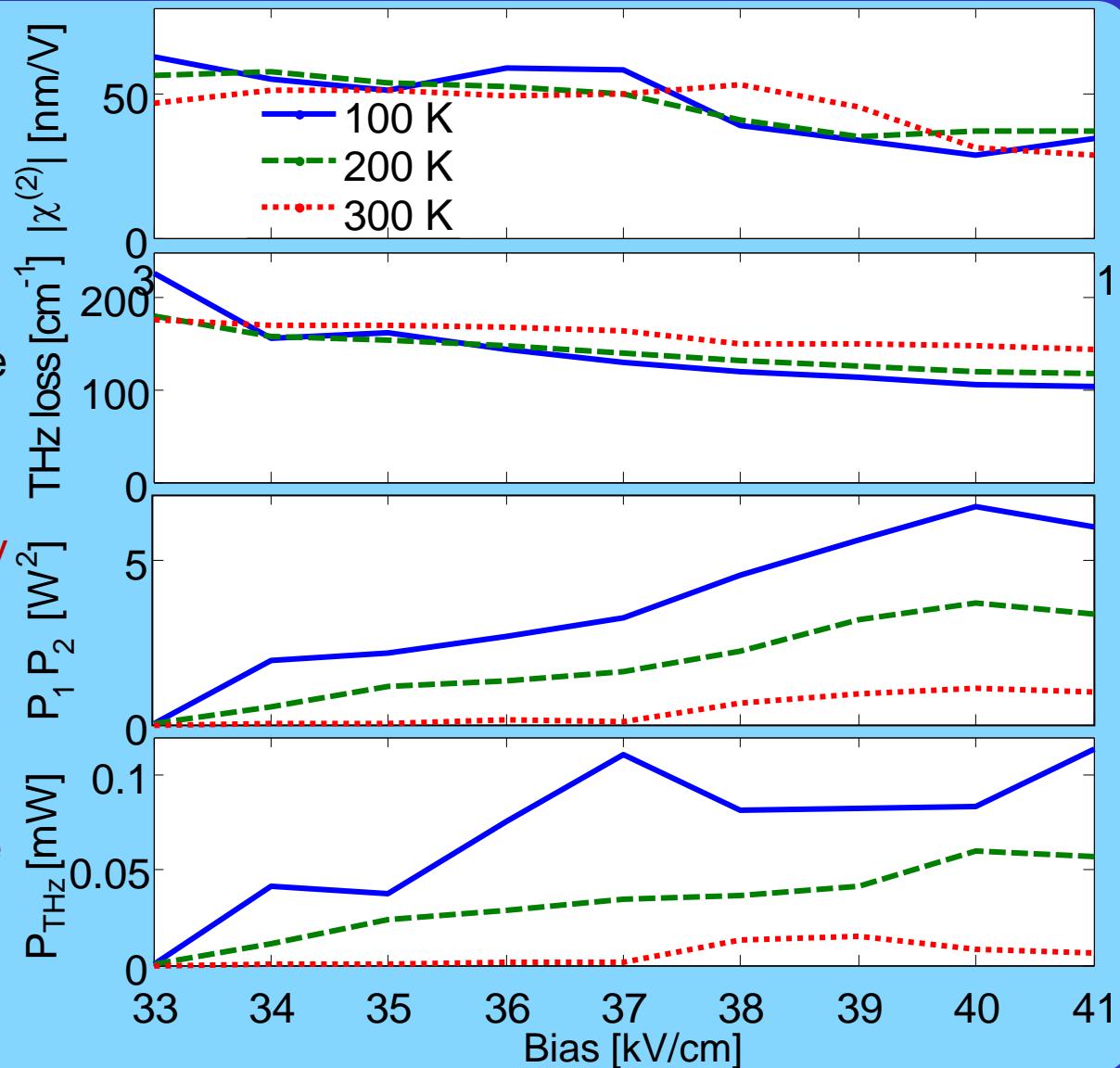
Extracted parameters in good agreement with experimental estimates:

- Nonlinear susceptibility $|\chi^{(2)}| = 44 \text{ nm/V}$
- THz waveguide loss $a = 150 \text{ cm}^{-1}$

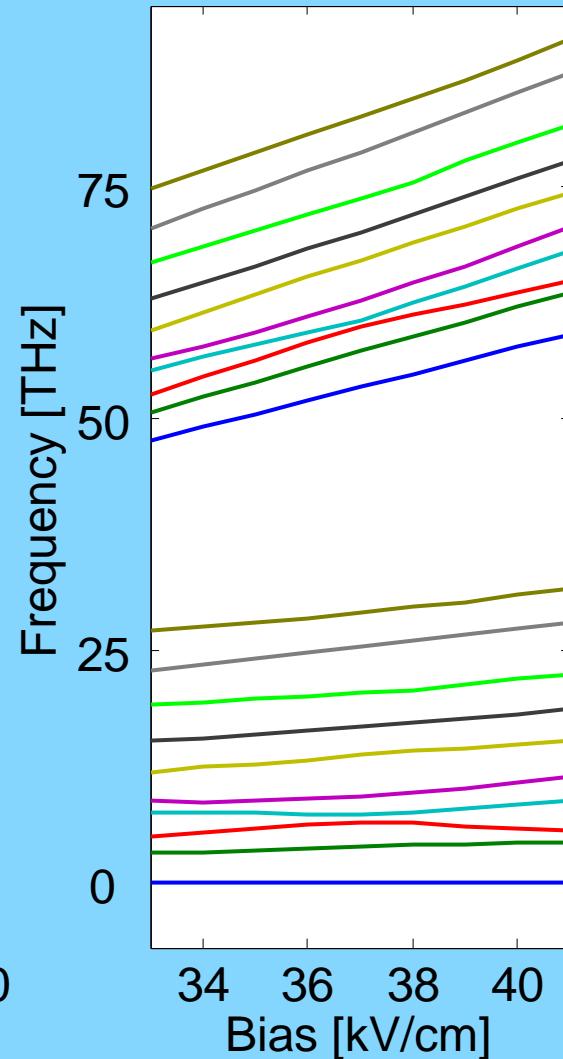
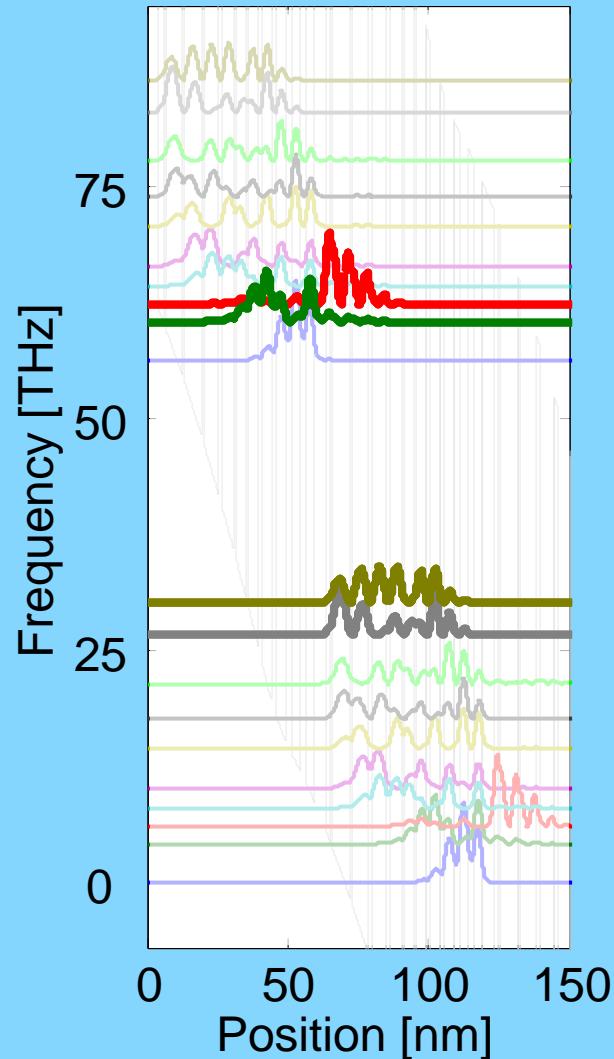
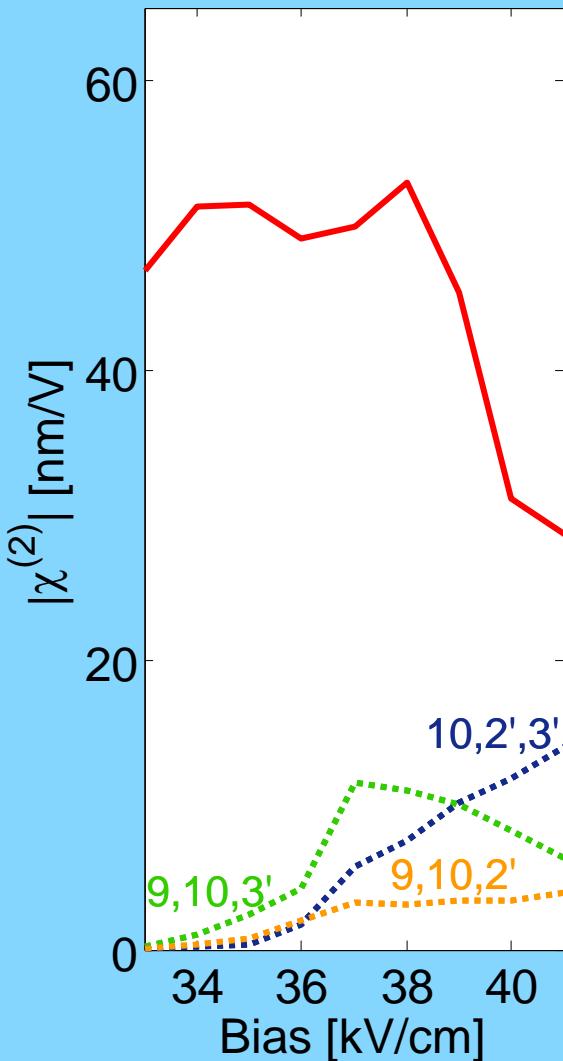
C. Jirauschek et al., Opt. Express **21**, 6180 (2013)

Temperature Degradation of THz Power

- Susceptibility $\chi^{(2)}$ does not degrade with temperature
- THz loss a_{THz} moderately increases with temperature
- MIR powers P_1, P_2 strongly decrease with temperature
- THz power P_{THz} strongly degrades with temperature

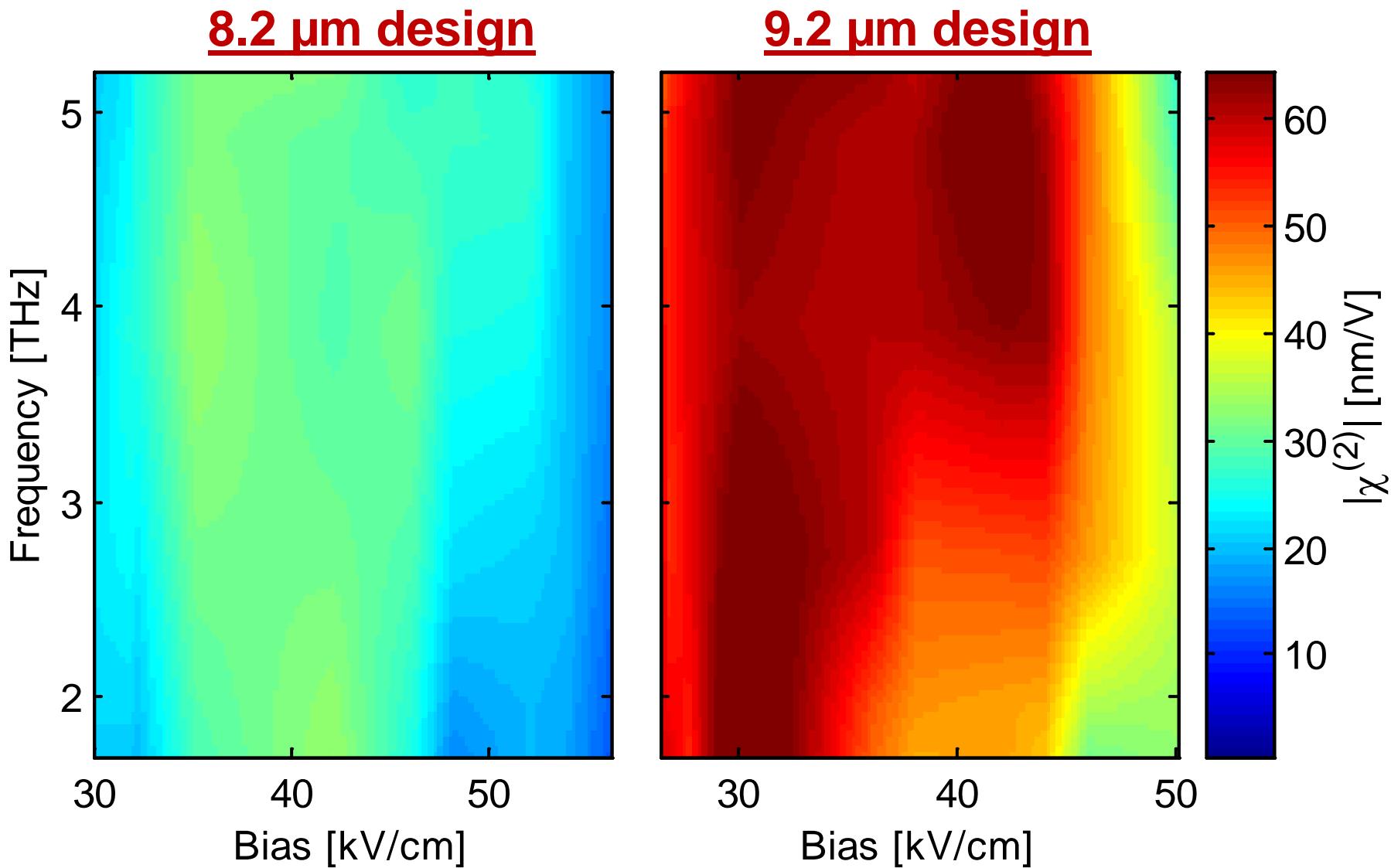


Contributions of Individual Subband Triplets



C. Jirauschek et al., Opt. Express **23**, 1670 (2015)

Susceptibility of Widely Tunable THz DFG Structure



Design: K. Vijayraghava et al., Nature Comm. 4, 2021 (2013)

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Extended Maxwell-Bloch Equations

Field:

$$\frac{n}{c} \partial_t E = -\partial_z E - i \frac{kN\mu\Gamma}{2\epsilon_0 n^2} \eta - \frac{1}{2} \ell(E) E$$

Polarization:

$$\partial_t \eta = \frac{i\mu}{2\hbar} E \Delta - \frac{\eta}{T_2}$$

Inversion:

$$\partial_t \Delta = \frac{\Delta_p - \Delta}{T_1} + \frac{i\mu}{\hbar} (E^* \eta - c.c.)$$

Rate Equations

Field:

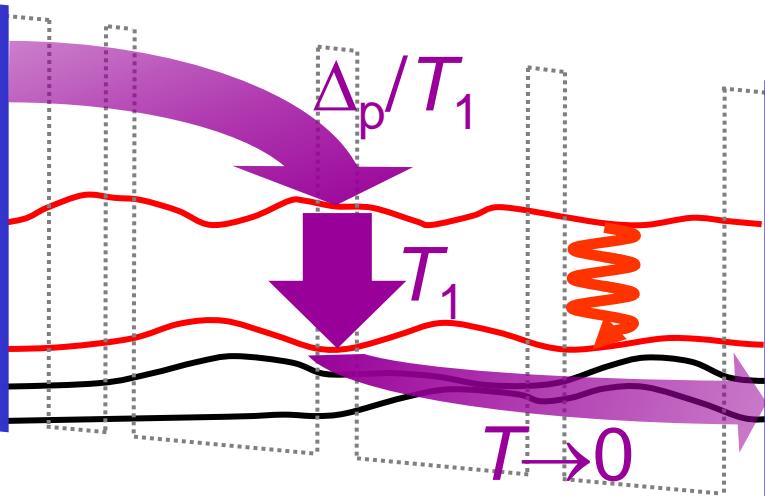
$$\frac{n}{c} \partial_t E = -\partial_z E - i \frac{kN\mu\Gamma}{2\epsilon_0 n^2} \eta - \frac{1}{2} \ell(E) E$$

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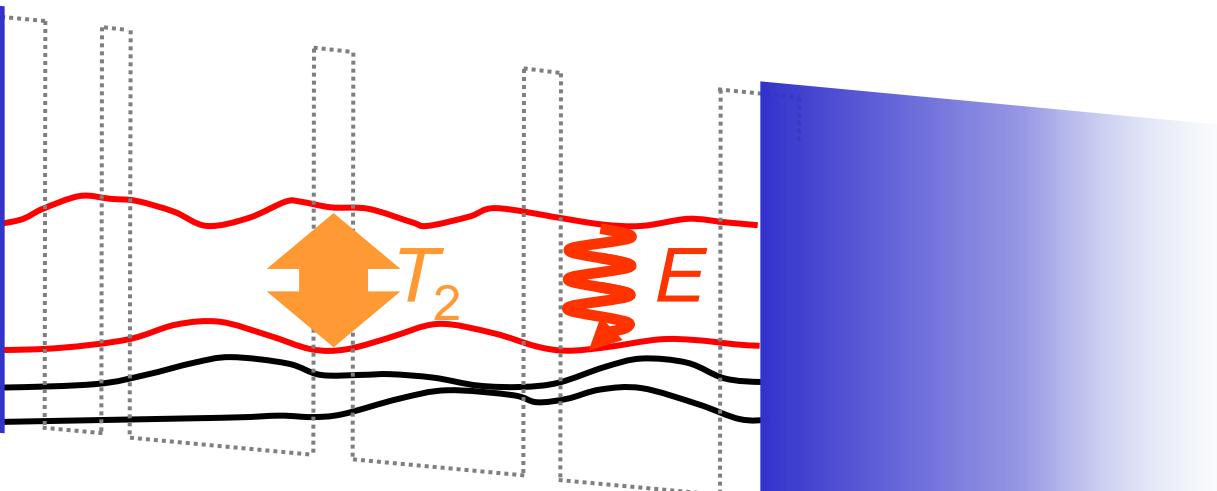
Coherent Effects

Field:

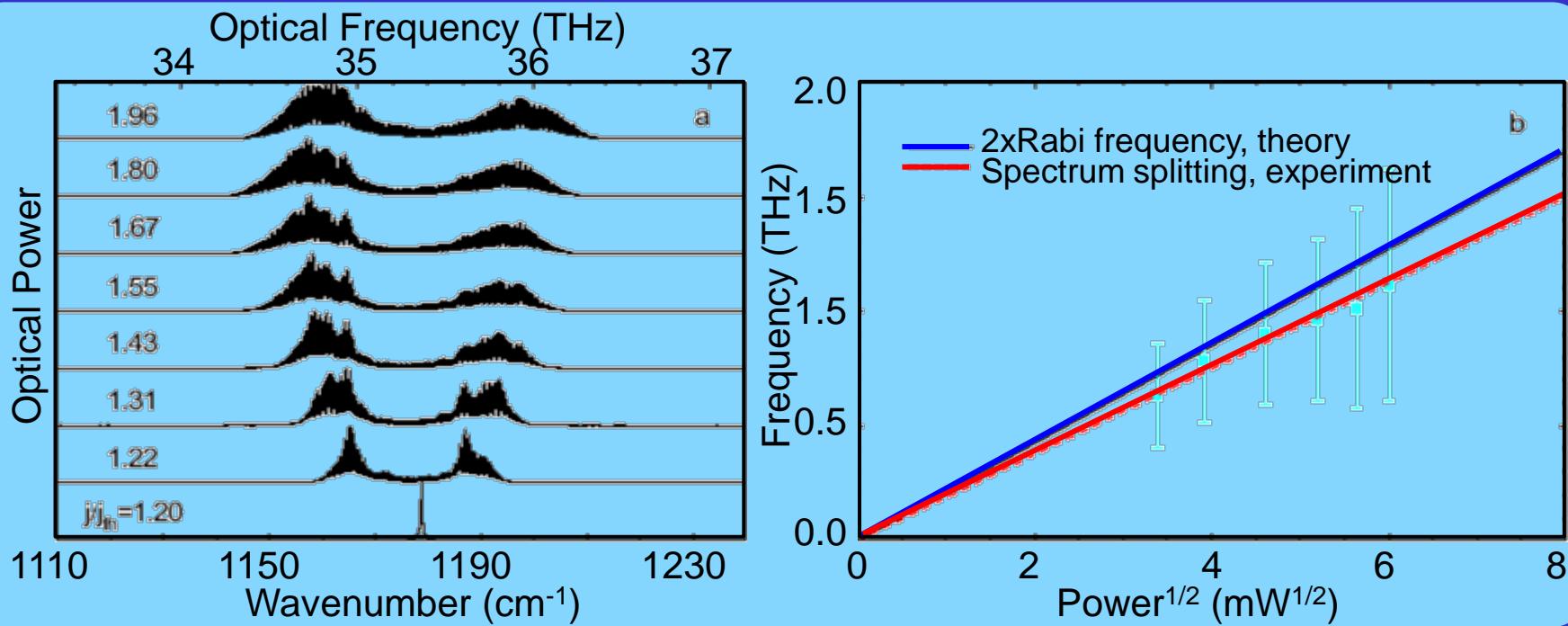
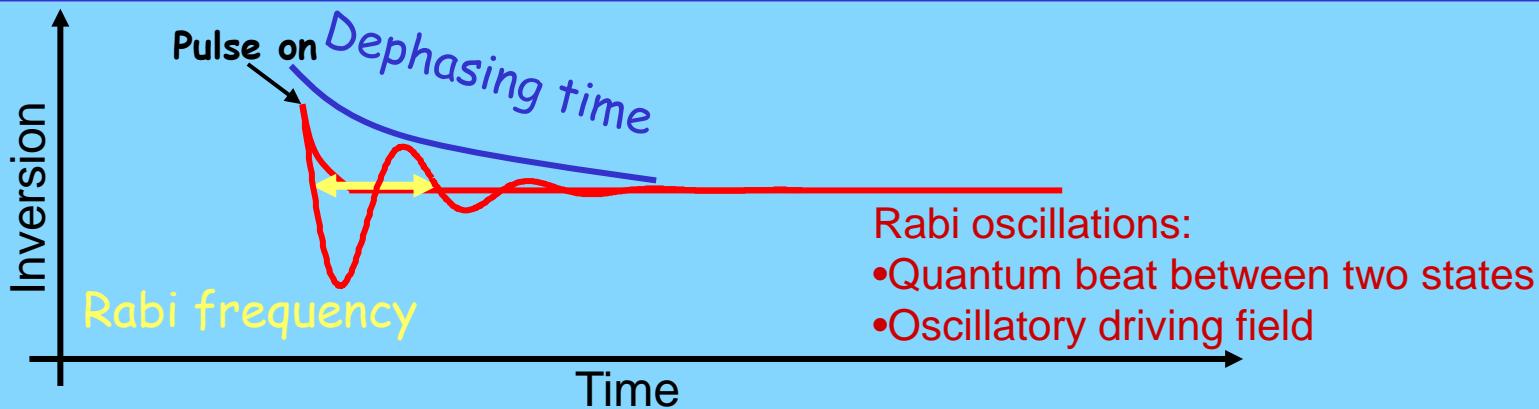
$$\frac{n}{c} \partial_t E = -\partial_z E - i \frac{kN\mu\Gamma}{2\epsilon_0 n^2} \eta - \frac{1}{2} \ell(E) E$$

Polarization: $\partial_t \eta = \frac{i\mu}{2\hbar} E \Delta - \frac{\eta}{T_2}$

Inversion: $\partial_t \Delta = \frac{\Delta_p - \Delta}{T_1} + \frac{i\mu}{\hbar} (E^* \eta - c.c.)$

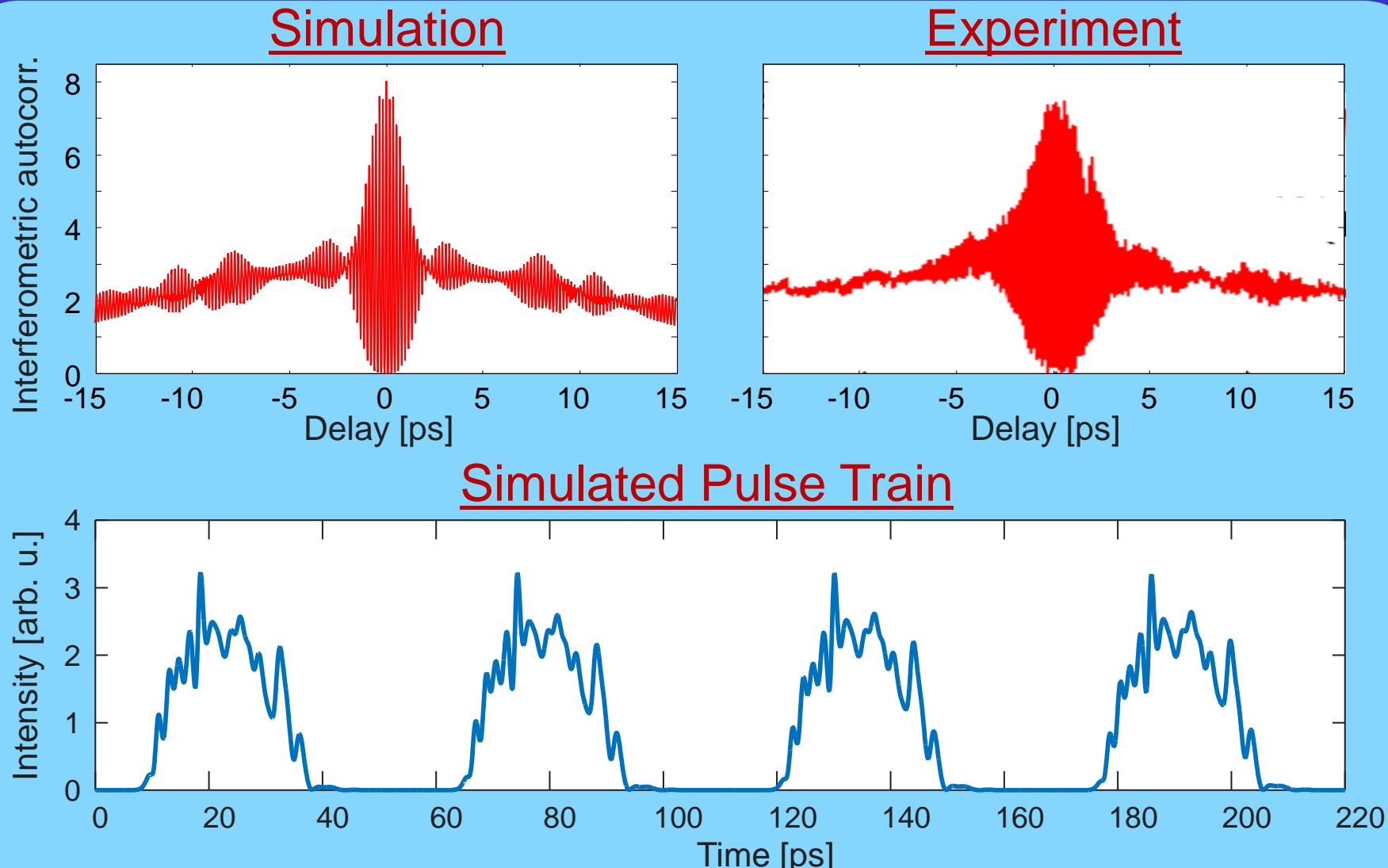


Manifestation of Coherent Effects



C. Y. Wang et al., Phys. Rev. A 75, 031802(R) (2007)

Simulation of Actively Mode-Locked QCLs



V.-M. Gkortsas et al., Opt. Express **18**, 13616 (2010)



Multi-Domain Simulation Approach

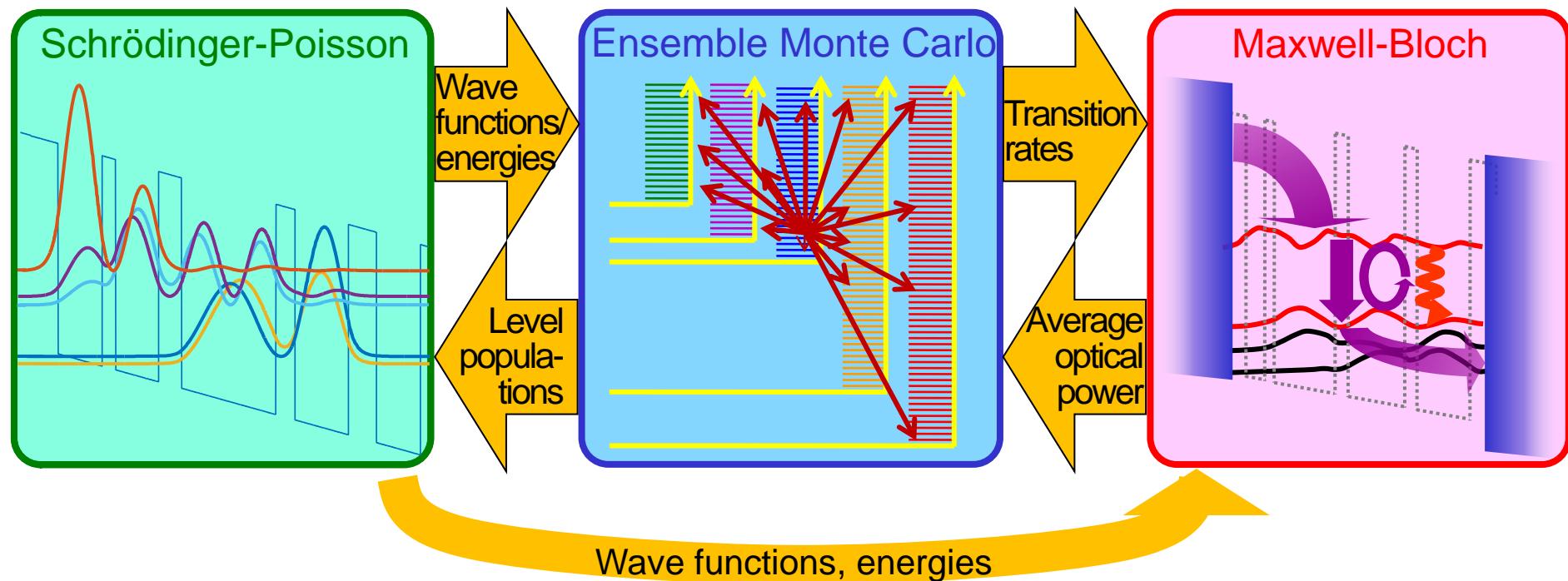
Maxwell-Bloch approach requires lifetimes/transition rates as input parameters

- Empirical or guessed values are used, affecting the quantitative accuracy

EMC self-consistently evaluates scattering based on the corresponding Hamiltonians

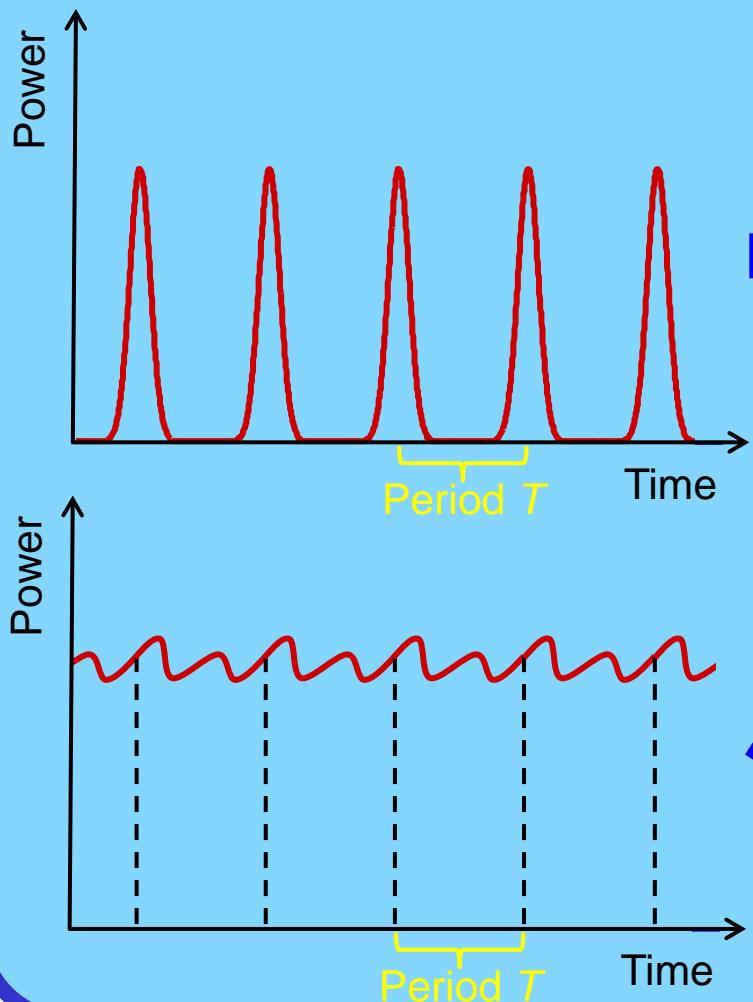
- Computational burden impedes inclusion of the full spatiotemporal dynamics

Coupled Maxwell-Bloch/EMC approach

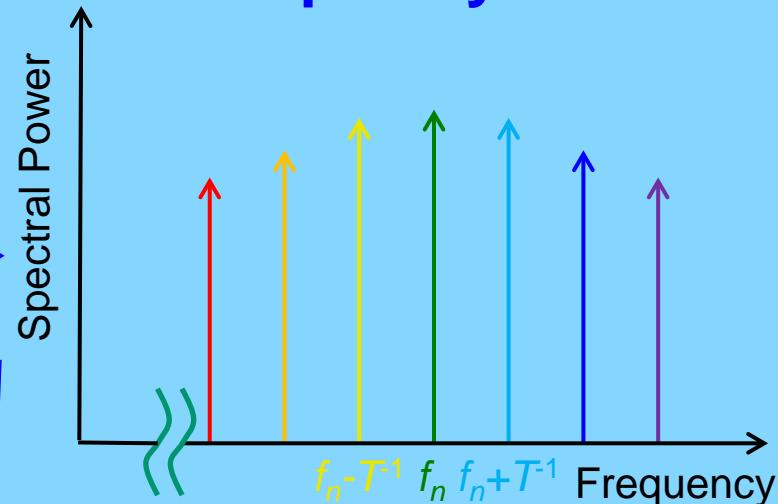


QCL-Based Frequency Combs

Periodic waveform

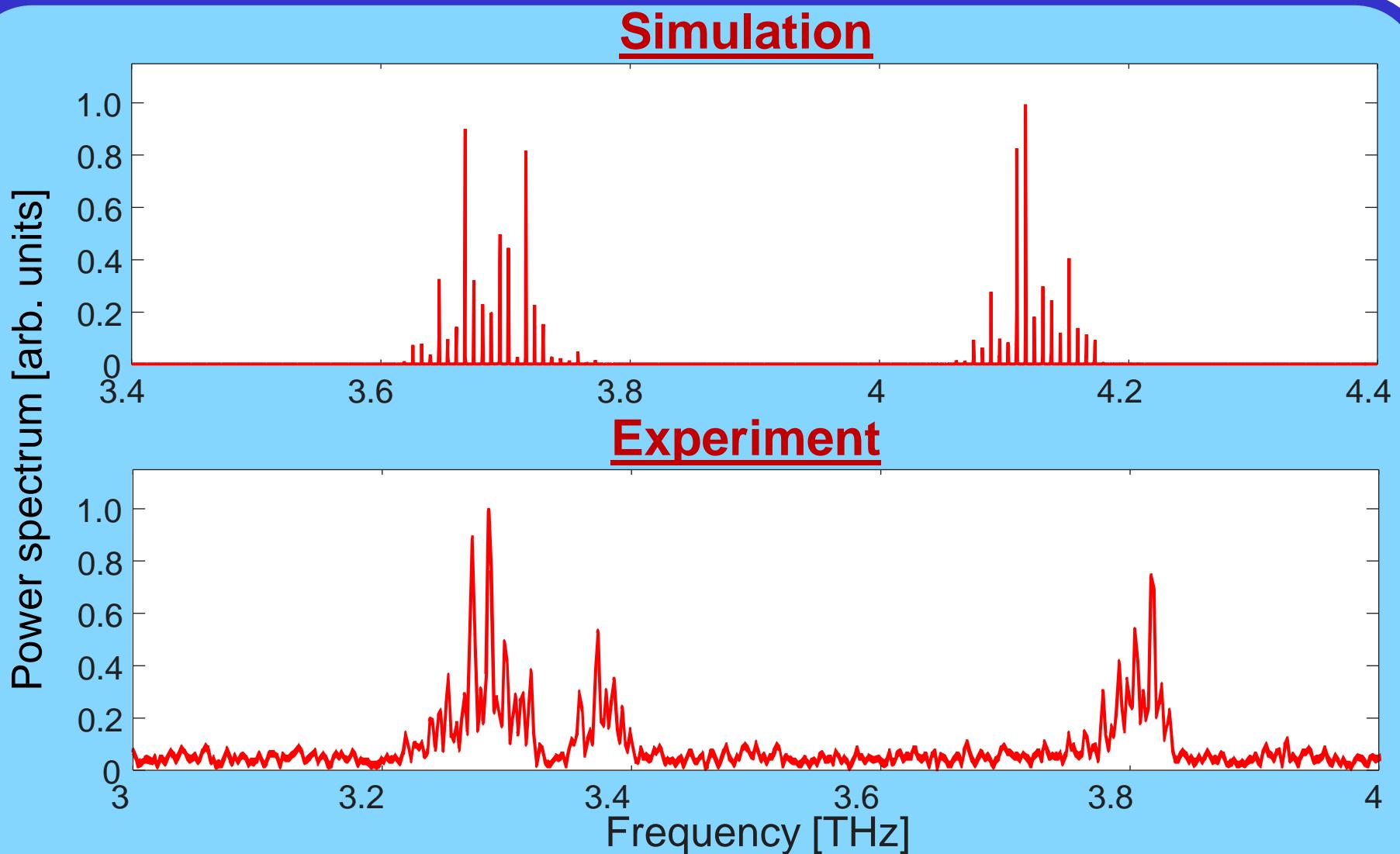


Discrete optical spectrum △ frequency comb



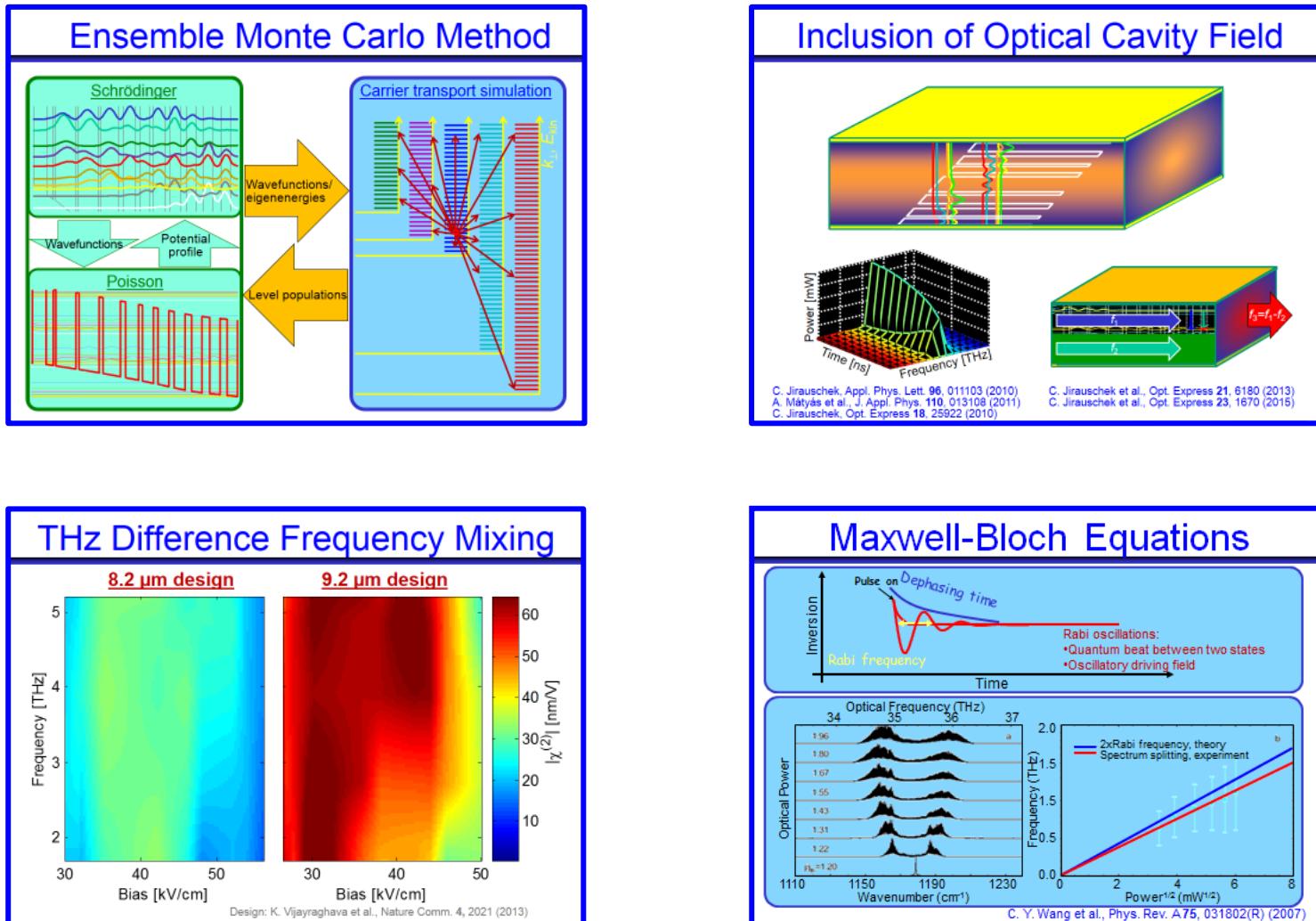
- Compact QCL-based frequency combs in mid-infrared/THz range
- Precision spectroscopy, molecular fingerprint detection (strong rotational-vibrational bands of many molecules at these frequencies)
- Modeling based on Maxwell-Bloch/EMC approach

Terahertz Frequency Comb



Design: D. Burghoff et al., Nature Photon. 8, 462 (2014)

Conclusion



C. Jirauschek and T. Kubis, Appl. Phys. Rev. 1, 011307 (2014)

Acknowledgment

Group members

Alpar Matyas

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Collaborations

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