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# **Nanometer Scale Patterning and Processing**

Spring 2016

## **Lecture 14**

### **Extreme UV (EUV) Lithography – EUV Source (Hot and Dense Plasma)**

# Extreme UV (EUV) lithography

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1. Overview, why EUV lithography?
2. EUV source (hot and dense plasma).
3. Optics (reflection mirrors).
4. Mask (absorber on mirrors).
5. Resist (sensitivity, LER, out-gassing).
6. Contamination control.

# Source Power vs Throughput

Updated EUV power and wafer throughput:

~~250 W~~  
~~120 W~~

Collectable EUV power  
~~60 W~~

Collectable, in-band, "clean" (no debris, no out-of-band)

EUV Power @ reticle  
3.5 W

Power @ wafer  
140 mW

Illum. time per field  
0.26 s

Illum. time per wafer  
23 s

~~10~~  
~~5~~ mj/cm<sup>2</sup> resist  
300 mm wafers  
89 fields/wafer

120 wafers/hr

Raw wafer throughput  
~~80~~ wafers/hr

Original courtesy of Jos Benschop and Vadim Banine, ASML.

J. Benschop et al., SPIE 3997, 34 (2000),  
V. Banine and R. Moors, SPIE 4343, 203 (2002).

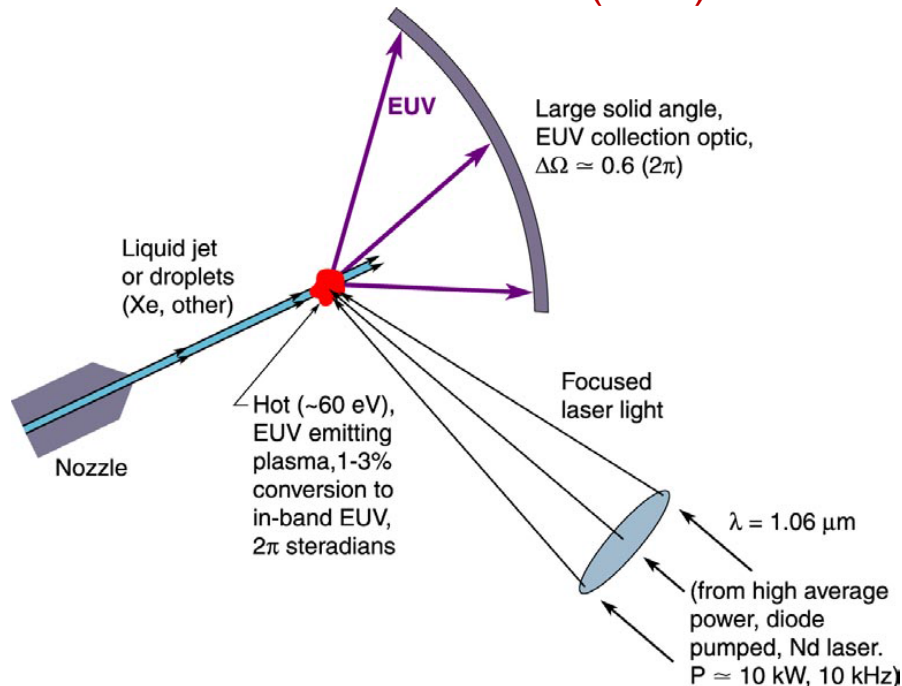
# EUV Radiation Source from Plasma

Must match the wavelengths of 13.5 nm at which Mo/Si multi-layers have high reflectivity

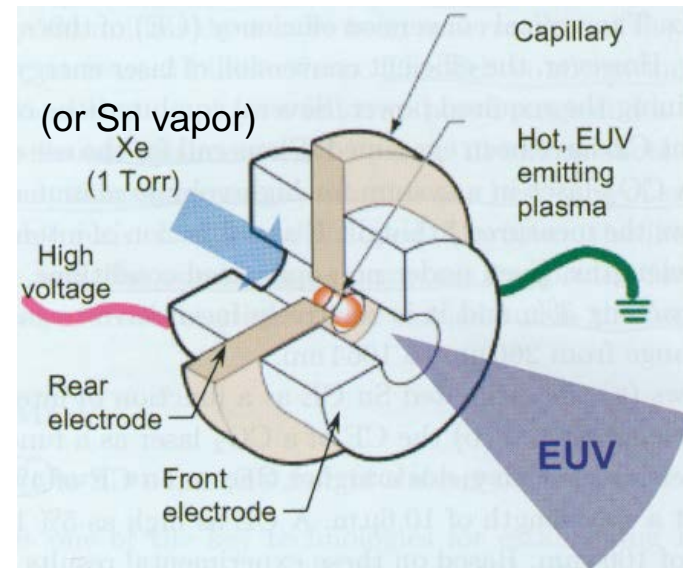
The only viable source for 13.5nm photons is a hot and dense plasma

- temperature of up to 200,000°C, atoms ionized up to +20 state.
- Emit photons by (e - ion) recombination and de-excitation of the ions.
- Plasma must be pulsed: pulse length in pico- to nanosecond range
- produced by powerful pulsed laser or electric arc (discharge) of up to 60,000A peak current.

## Laser Produced Plasma (LPP)



## Discharge Produced Plasma (DPP)



Courtesy of Neil Fornaciari and Glenn Kubiak, Sandia.

# Laser (LPP) and discharge (DPP) produced plasma

**DPP:** change electrical energy directly into EUV light, so high power, high efficiency; but ablation of electrode and more debris.

**LPP:** higher collection efficiency (larger collection angle), high repetition rate (more pulses/sec), more manageable thermal loads and debris, more scalable to HVM (high volume manufacturing).

EUV lithography tools using both DPP and LPP have been built.

## Challenges:

- Radiate from IR to x-ray: need filter.
- Large source size into  $4\pi$  spherical angles: need collector optics.
- Debris and thermal issues: may damage the optics.

## For EUV lithography, ideally:

- Power  $> 110W$
- Maximum in-band emission with narrow bandwidth ( $\leq \pm 2\%$ ).
- Forward directed, no collector

Low ( $\mu W$ ) to mid-power (1W) has application in other fields:

Interference lithography, spectroscopy, microscopy, other metrology, testing EUV resist.

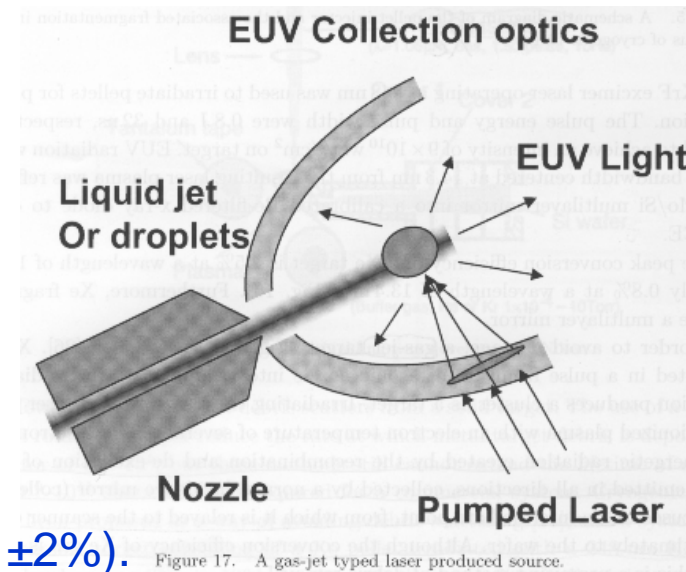


Figure 17. A gas-jet typed laser produced source.

# Discharge produced plasma (DPP)

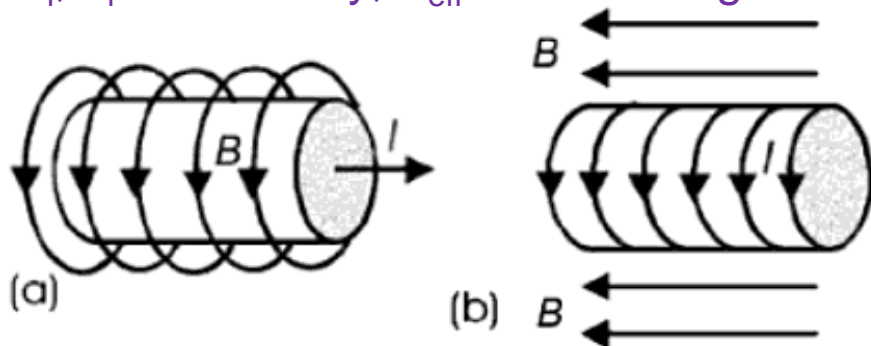
- Hot plasma is created by magnetic compression of low-temperature plasma.
- Plasma is compressed with the magnetic field generated by the current used to heat the plasma, with two common geometries (see figure below).
- Because plasma is compressed by magnetic field  $B$  of current  $I$ , which generates the plasma, plasma is “self-heating”.
- Two forces are present: the magnetic field pressure  $B^2/2\mu_0$  and the plasma pressure.
- When these two forces are equal, the plasma achieves an equilibrium.

**Equilibrium:** 
$$\frac{\mu_0 I^2}{4\pi} = (Z_{eff} + 1)N_i kT_e$$

$I$ : total current ( $B \propto I$ ), supplied by a capacitor in a pulsed mode;

$T_e$ : electron temperature

$N_i = \pi r^2 n_i$ ,  $n_i$ : ion density;  $Z_{eff}$ : mean charge of the ions.



Two plasma compression geometries.

# Laser produced plasma (LPP): overview

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- Focusing a laser beam on a target material.
- Initial **ionization** occurs through photo-ionization, and the electric field of the laser accelerates these electrons (inverse Bremsstrahlung, see next slides).
- Non-elastic collision further ionize the plasma, whereas elastic collision transfers the electron's kinetic energy into ionic kinetic energy.
- As plasma expands, thermal energy is converted into kinetic energy and **charge density decreases**. (Expansion velocity of Sn plasma at  $\sim 30\text{eV}$  is about  $2 \times 10^6\text{cm/s}$ , so at the order of  $100\mu\text{m}$  in front of target for 10 ns pulse.) This decreases both the further absorption of laser energy and energy conversion efficiency (CE).
- Therefore, the laser pulse length should not be very long: **10 ns** is a good time scale.
  
- The corresponding **optimum laser intensity** for maximum CE is  $1 \times 10^{11}\text{ W/cm}^2$  for Nd:YAG laser ( $1.06\ \mu\text{m}$ ), and  $1 \times 10^{10}\text{ W/cm}^2$  for  $\text{CO}_2$  laser ( $10.6\mu\text{m}$ ).
- **Longer wavelength preferred**: Modeling shows that CE depends on laser wavelength -  
- CE ( $10.6\ \mu\text{m}$  laser):CE ( $1.06\ \mu\text{m}$  laser):CE( $0.26\ \mu\text{m}$  laser)=1.9:1.0:0.55.
- This is because: CE depends on a balance between emissivity and opacity; at longer wavelength laser absorption occurs at lower plasma (electron) density that is more transparent for EUV to escape.

# Laser produced plasma: a clean, bright, narrow source

## Plasma temperature:

$T_e \propto Z^{1/5} (\lambda^2 \phi)^{3/5}$ . Longer wavelength is more efficient to heat up plasma, so higher CE.  
 $\lambda$ : laser wavelength;  $\phi$ : laser flux (Watt/cm<sup>2</sup>);  $Z$ : target atomic number.

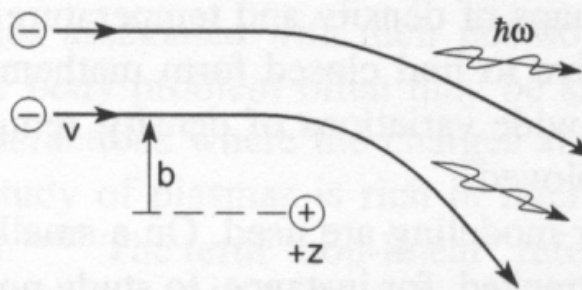
## Energy levels:

Transitions and energy levels are calculated using Hartree-Fock methods. The UTA continuum spectra is due to  $4p^6 4d^n - 4p^5 4d^{n+1} + 4d^{n-1} 4f$  line transitions.

## Atomic processes (routes for radiation):

- Ionization and recombination between successive ion stages.
- UTA line radiation from bound-bound transitions (excitation and de-excitation).
- Continuous Bremsstrahlung occur as well (Bremsstrahlung: electron emits photon when it is accelerated by the positively charged ions).

Same idea as synchrotron radiation where electrons are accelerated radially by magnetic field.

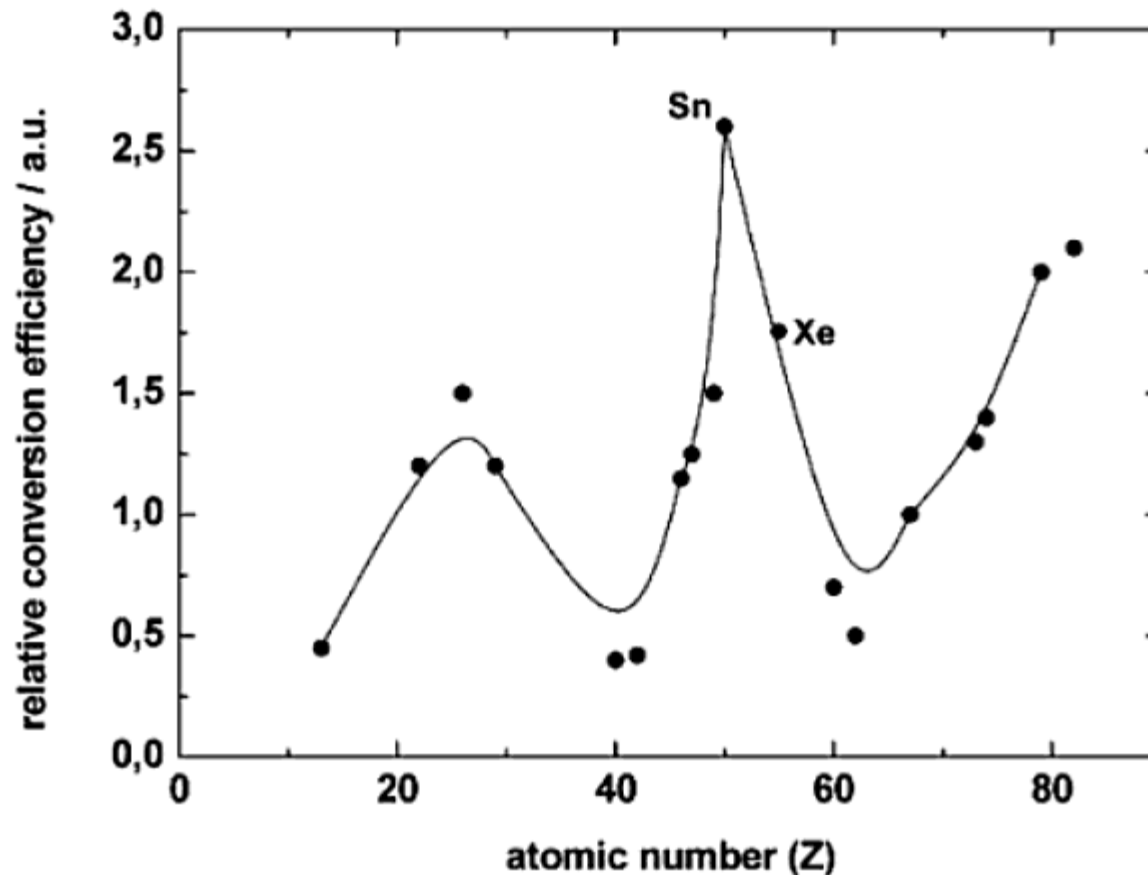


From: Soft x-rays and extreme ultraviolet radiation, written by David Attwood

**FIGURE 6.2.** Bremsstrahlung occurs when a passing electron is accelerated by an ion, causing it to radiate. Because of the wide range of incident electron velocities and the range of distances of closest approach (*impact parameter*  $b$ ) a broad continuum of radiation is generated in a plasma, with a spectrum closely related to the electron velocity distribution, or its characteristic temperature.



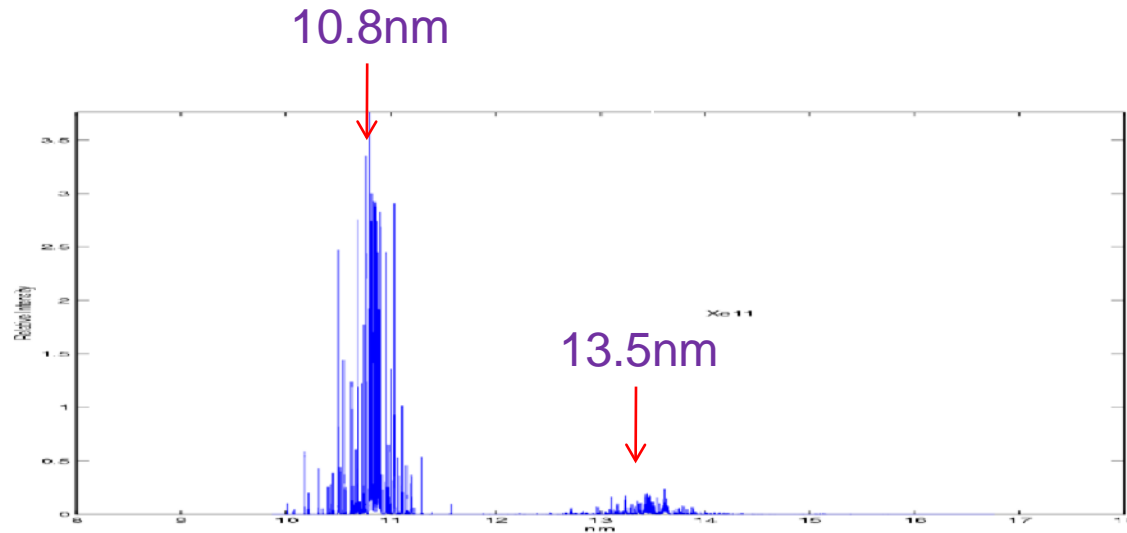
# Energy conversion efficiency (CE) into 13.5nm radiation



**Figure 3.5** Relative CE into 13.5-nm radiation as a function of the atomic number of the emitter. The highly efficient Sn ( $Z = 50$ ) and the frequently used Xe ( $Z = 54$ ) are marked. (Reprinted from Ref. 14.)

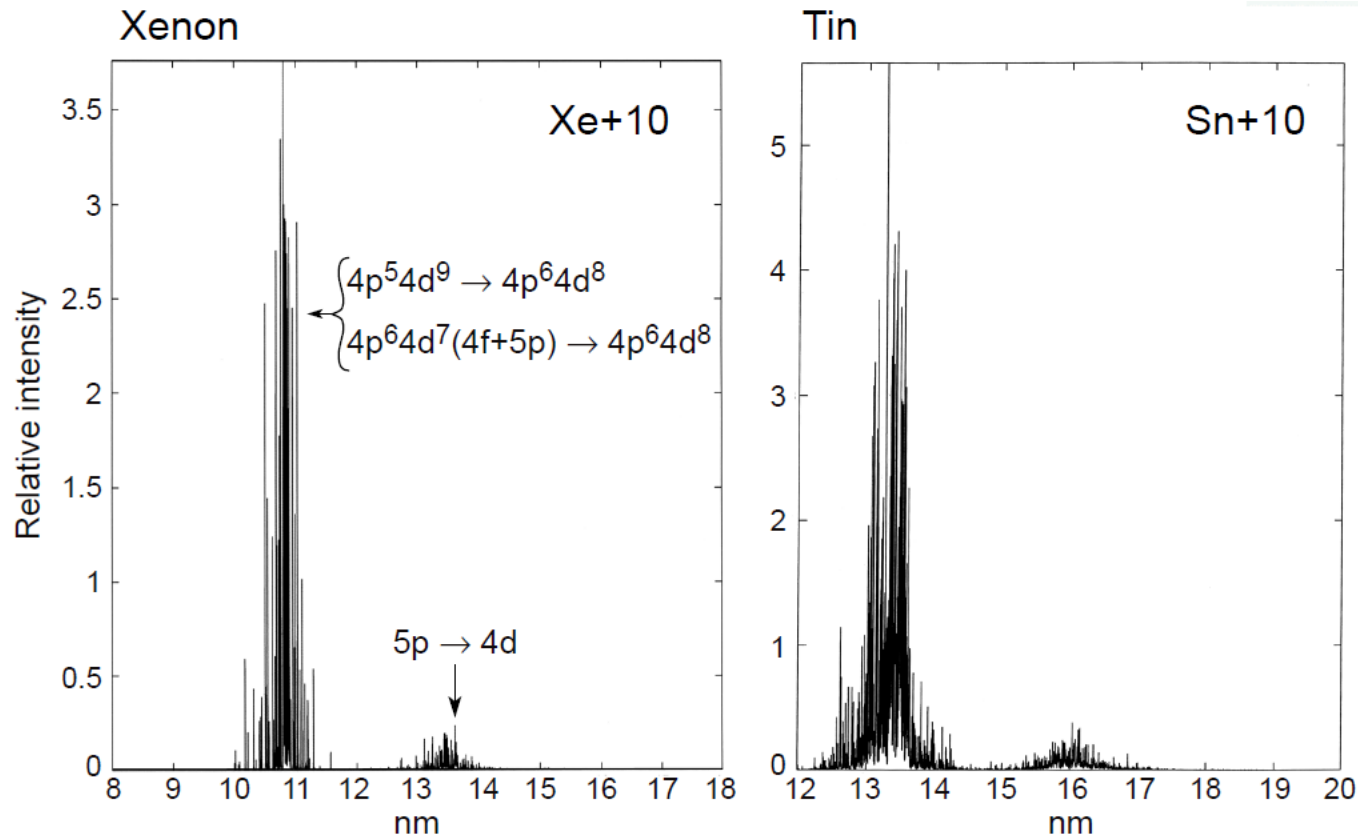
Sn is most efficient, followed by Li (?), then Xe.

# Plasma radiation source for 13.5nm: Xe



- 13.5nm photons only generated by one ion state ( $\text{Xe}^{11+}$ ).
- Maximum population of this state is 45%.
- Even this state emits 10 times more at 10.8nm than 13.5nm.
- That is, xenon is inefficient: to produce 100W at 13.5nm, kilowatts of other wavelengths would have to be removed.
- On the plus side, xenon is very clean and easy to work with (no debris).

# Comparative Spectra: Xe and Sn



Courtesy of G. O'Sullivan (Univ. College Dublin)  
R. Faulkner (UCD Ph.D, 1999)  
A. Cummings (Nahond Univ. Ireland)

# Plasma radiation source for 13.5nm: Sn

Peak wavelength emission decreased with increasing atomic number ( $\lambda \propto 1/Z$ ).

## Laser produced Sn plasma

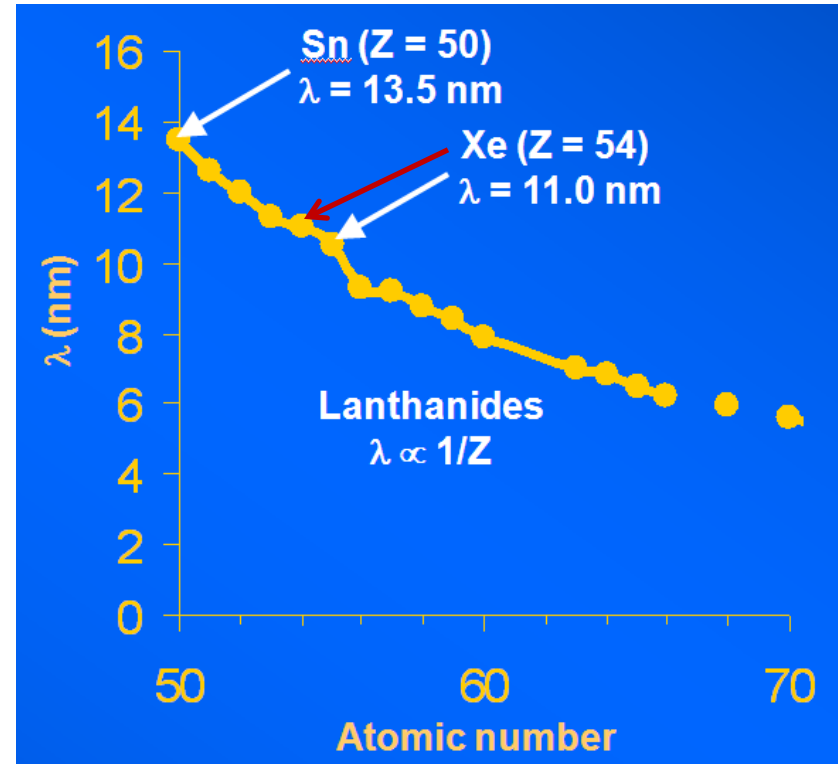
Target, Z=50 [Kr]5s <sup>2</sup> 4d <sup>10</sup> 5p <sup>2</sup>	Sn
Laser wavelength, $\lambda$	1.064 $\mu$ m
Laser flux, $\phi$	$1 \times 10^{11}$ W/cm <sup>2</sup>
Electron temperature, $T_e$	48.8eV
Electron density, $n_e$	$9.88 \times 10^{20}$
cm <sup>-3</sup> (300K is 26meV, 1eV is $1.15 \times 10^4$ K)	

## Ion distribution

Sn X	0.046
Sn XI	0.243
Sn XII	0.306
Sn XIII	0.330
Sn XIV	0.068

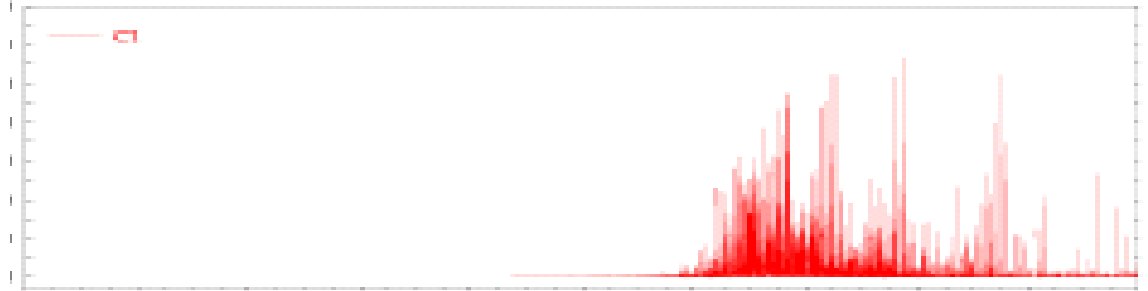
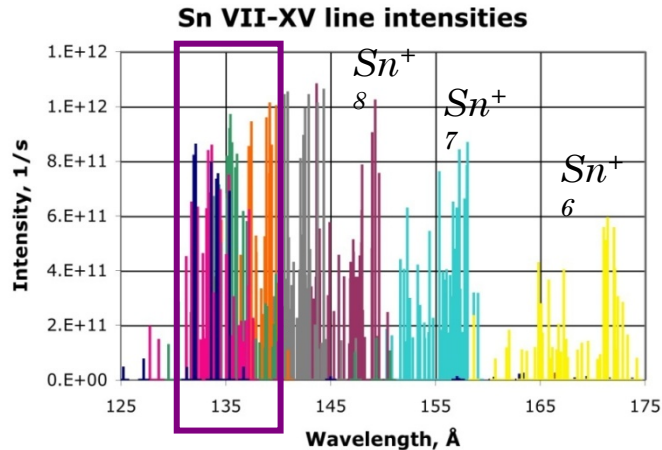
ECE 695 Nanometer Scale Patterning and Processing

## UTA peak wavelength (nm) versus atomic number



UTA: unresolved transition array, consisting of tens of thousands of lines (unresolved, overlapping 4d–4f transitions).

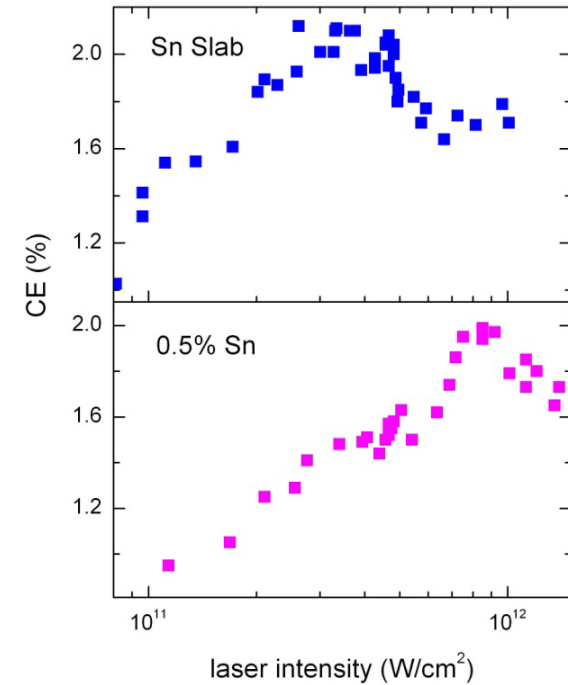
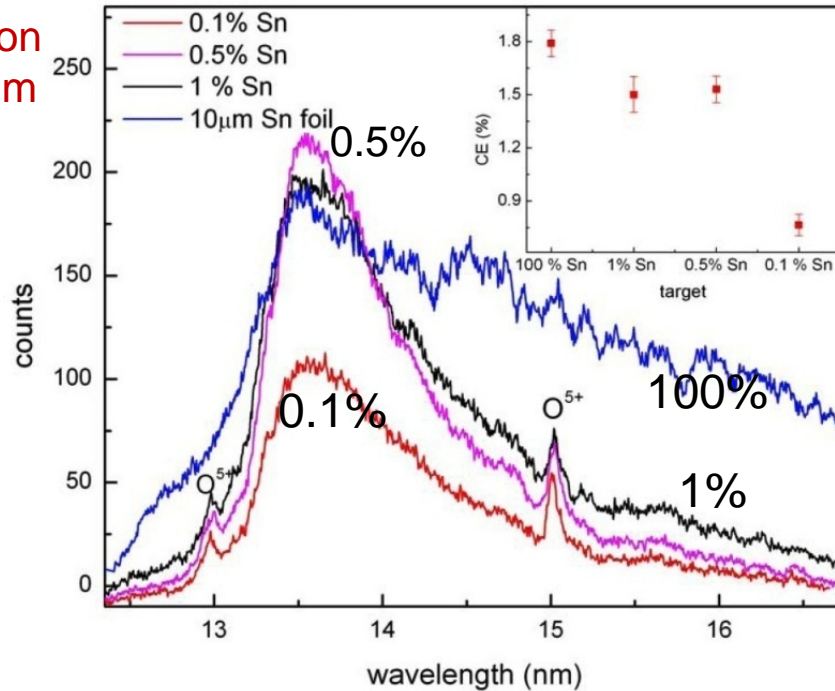
# Plasma compositions for 13.5nm: Tin



- Optimum emission when tin is a low-percentage impurity.
- Light comes from transitions between  $4p^64d^n$  and  $4p^54d^{n+1}$  or  $4d^{n-1}(4f,5p)$ .
- All ion states from  $\text{Sn}^{8+}$  to  $\text{Sn}^{13+}$  can contribute.
- Lighting up these transitions, and *only* these transitions requires exquisite control of laser plasma.
- But tin (debris) tends to condense on optics.
- In summary, tin is great as a 13.5nm source, if one can control the debris (yes).

# Low density (diluted) targets leads to narrow UTA

Radiation spectrum



CE: (energy) conversion efficiency

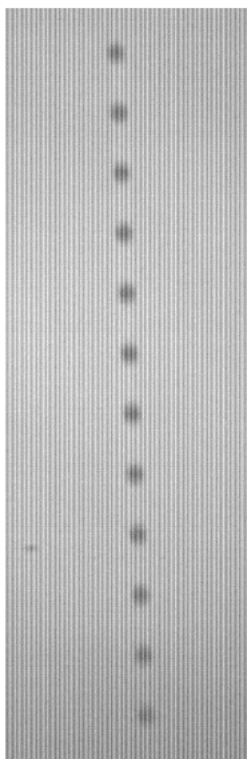
- Low density tin also reduces debris contamination to optics.
- The optical depth at 13.5 nm is only ~7 nm of full density Sn. Beyond that, light is reabsorbed.
- CE (conversion efficiency) slightly reduced with low density Sn.

Tin targets from General Atomics

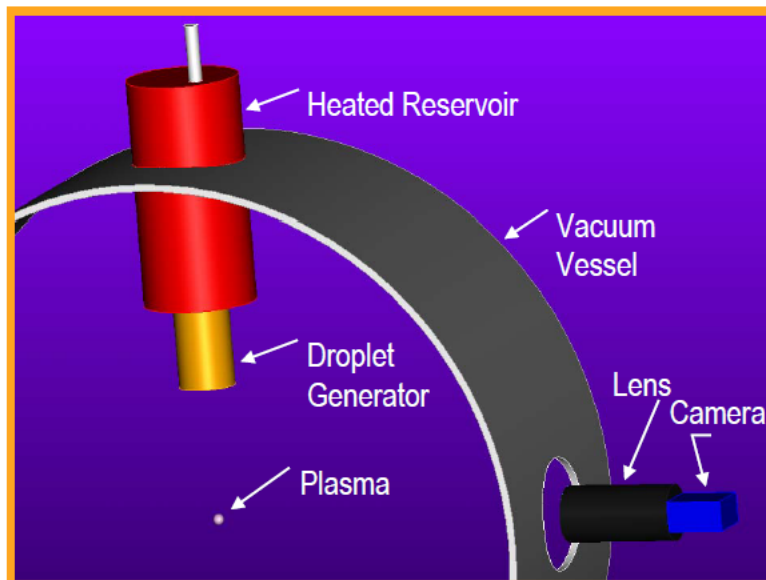
- 100 mg/cc RF foam
- 0.1-1% solid density Sn
- e.g., 0.5% Sn =  $Sn_{1.8}O_{17.2}C_{27}H_{54}$

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# Liquid Metal Droplet Generator Developed



100  $\mu\text{m}$  Sn droplets  
at 36 kHz, captured  
using strobe lighting



- Continuous stimulated droplet generation of liquid metals (Li and Sn) at temperatures up to 250°C
- Droplets diameter  $\leq 100 \mu\text{m}$
- Droplet rates up to 48 kHz
- Working distance of 50mm

# The tin-doped droplet laser plasma EUV source

Laser Plasma Laboratory

College of Optics & Photonics: CREOL & FPCE at UCF

Multi-component 30 -35  $\mu\text{m}$  diameter target  
at 30 kHz -- Location precision 3  $\mu\text{m}$

Modest laser intensities  $I \sim 10^{11} \text{ W/cm}^2$

Mass-limited targets

Target contains only  $10^{13}$  tin atoms

Recently demonstrated 30 kHz  
laser droplet irradiation with  
intelligent feedback beam and  
target control – continuous  
operation for 8 hours

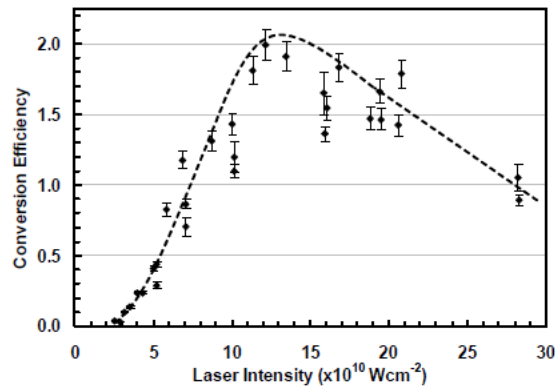


# High CE demonstrated with Droplet Target

Laser Plasma Laboratory

College of Optics & Photonics: CREOL & FPCE at UCF

**CE = 2% at 13.5 nm for tin-doped droplet target source**

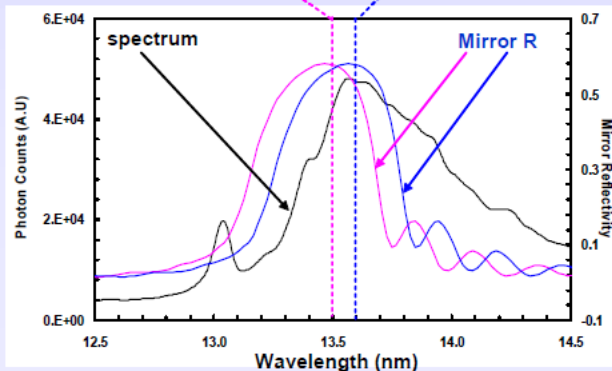
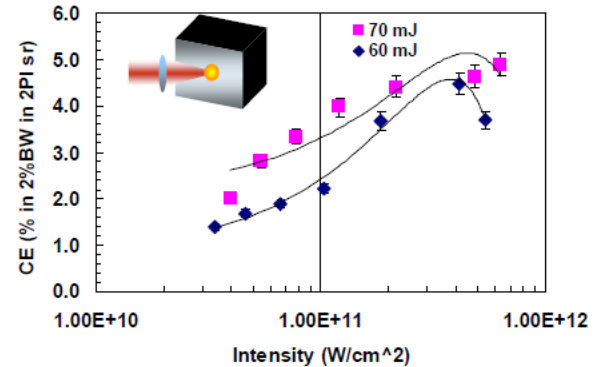


at 13.5nm, CE = 2%

at 13.6nm, CE = 2.25%

**FOM FC2 team**  
*F. Bijkerk S.A. vd Westen C. Bruineman*

**CE = 5.5 % with solid tin!**



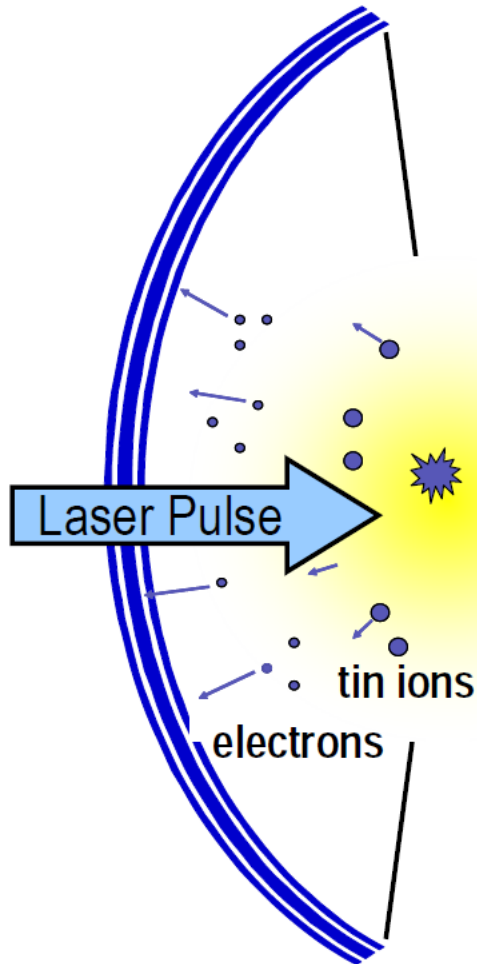
**CE = 3% achievable with droplet source**

--- for 30 kHz, 140 mJ laser

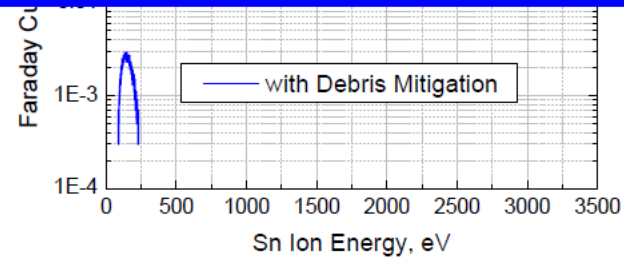
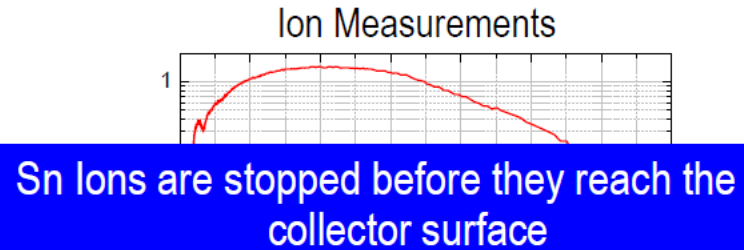


**120 W /  $2\pi$**

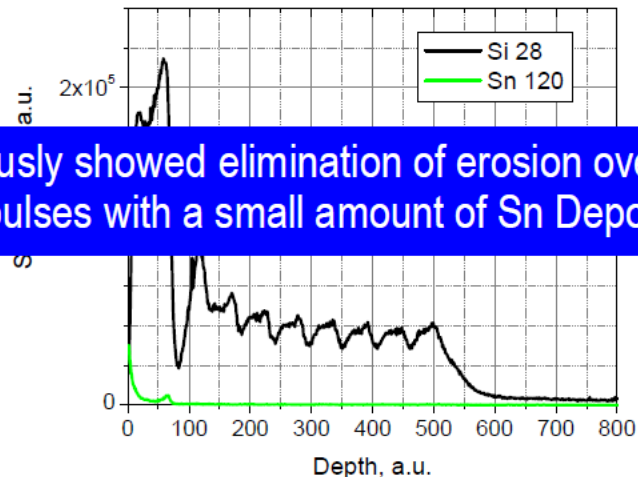
# Debris Mitigation Stops Erosion from Ions



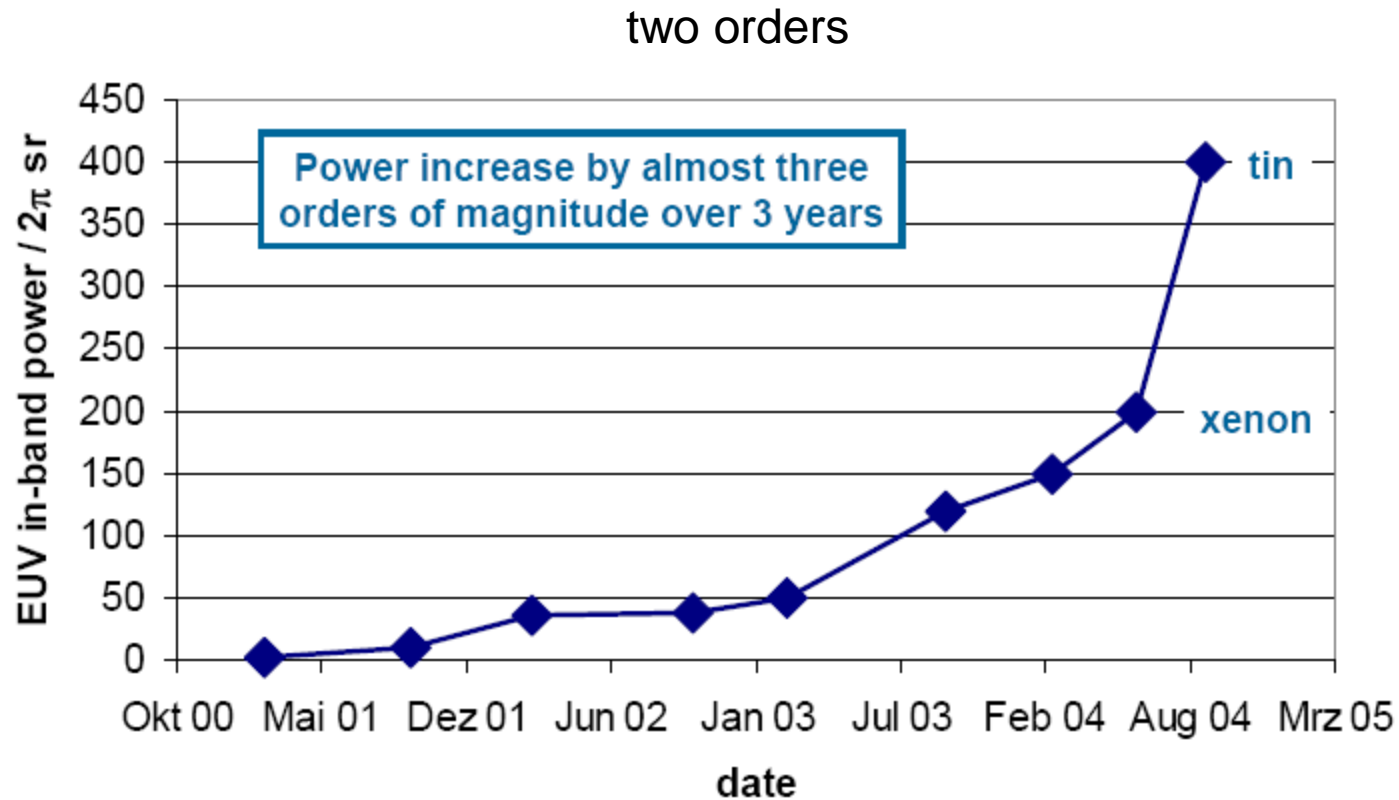
All measurements taken at collector surface



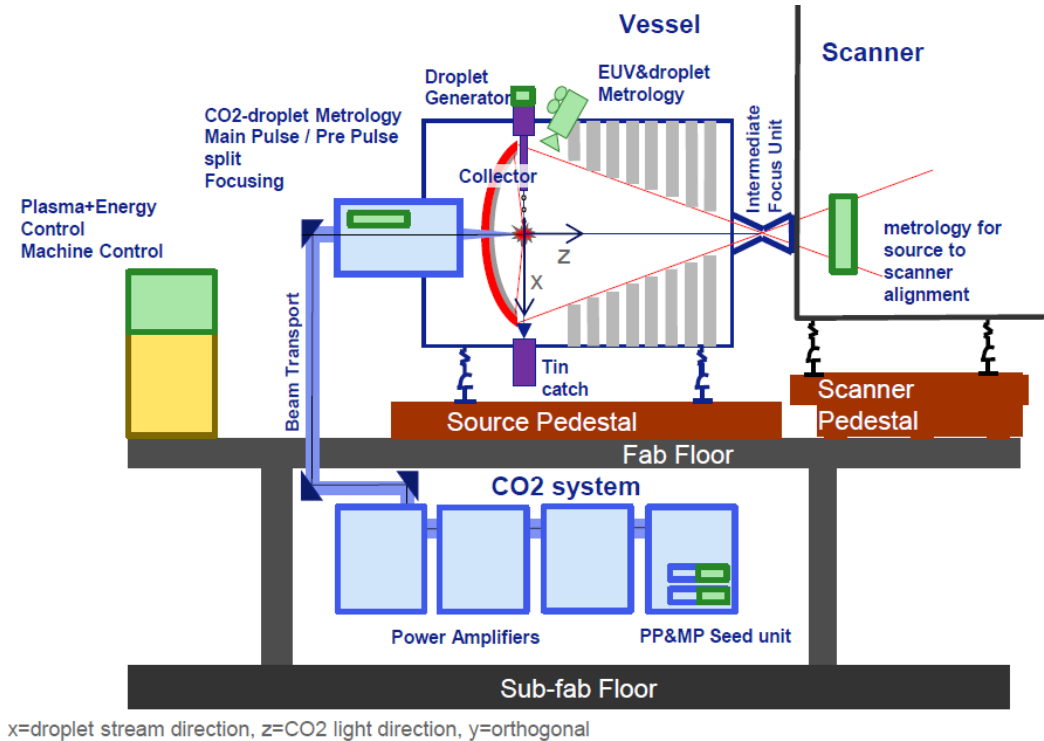
Previously showed elimination of erosion over 3 Mpulses with a small amount of Sn Deposition



# Gas discharge produced plasmas source scaling



# LPP EUV Source System Cross-Section



## Key components:

- Drive Laser
- Collector

Power

- Droplet generator
- Vessel

Availability

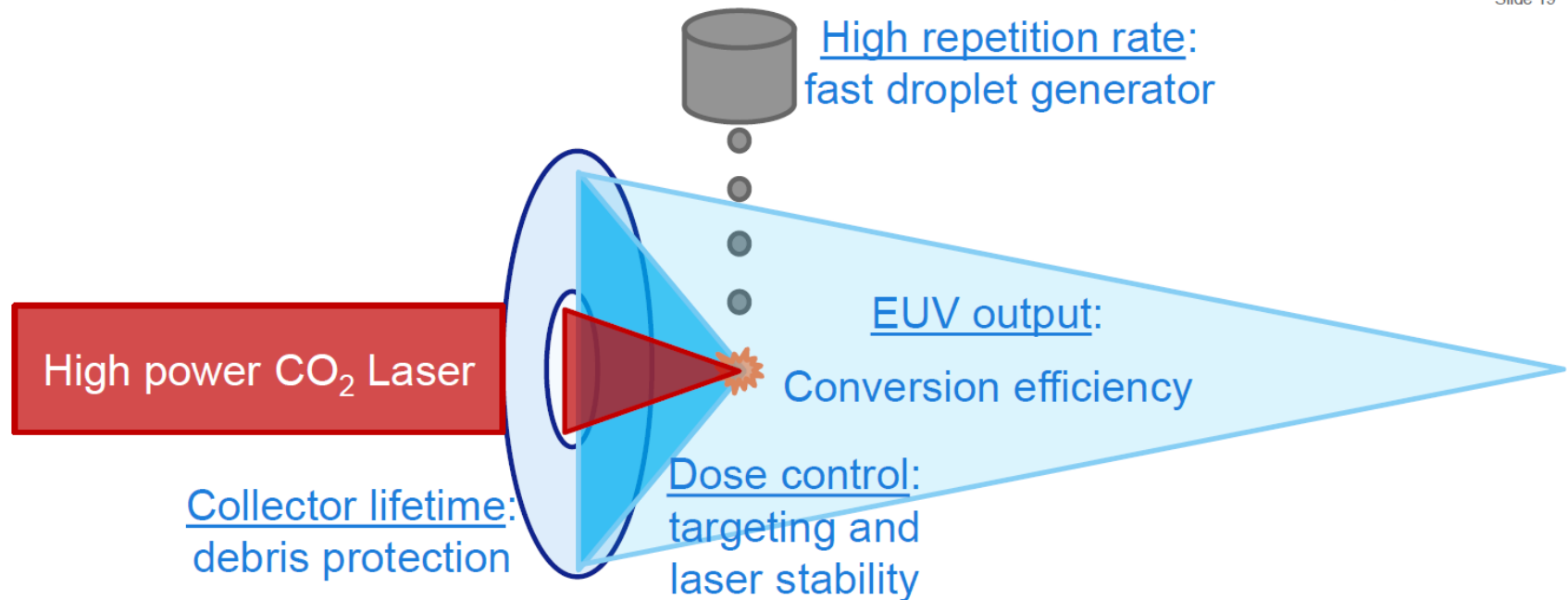
- Controls (E,x,y,z,t)
- Final Focus Assembly

Dose control

Source: Cymer ASML company

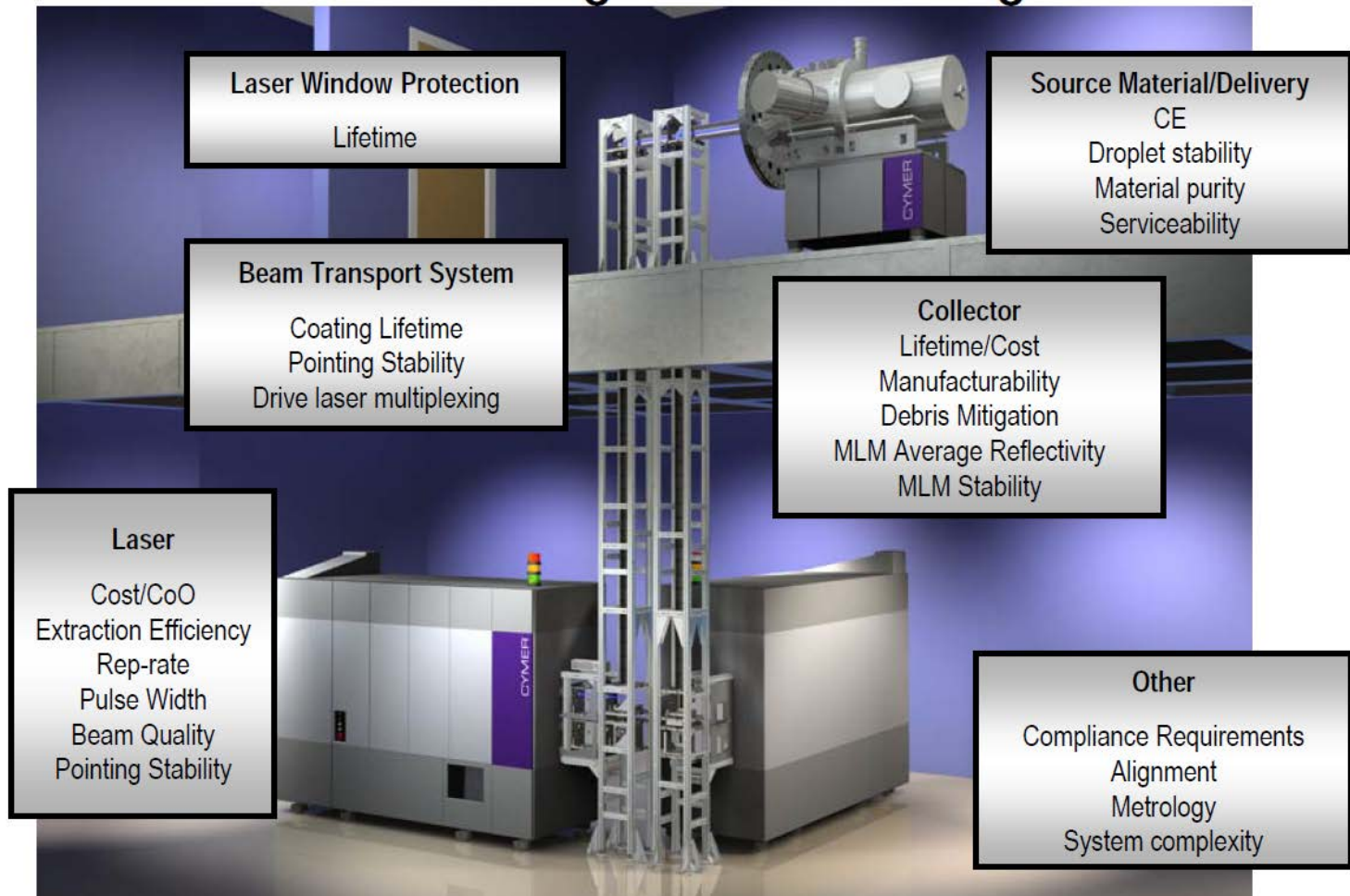
<http://vestige.lmsal.com/TRACE>

# EUV Power Scaling

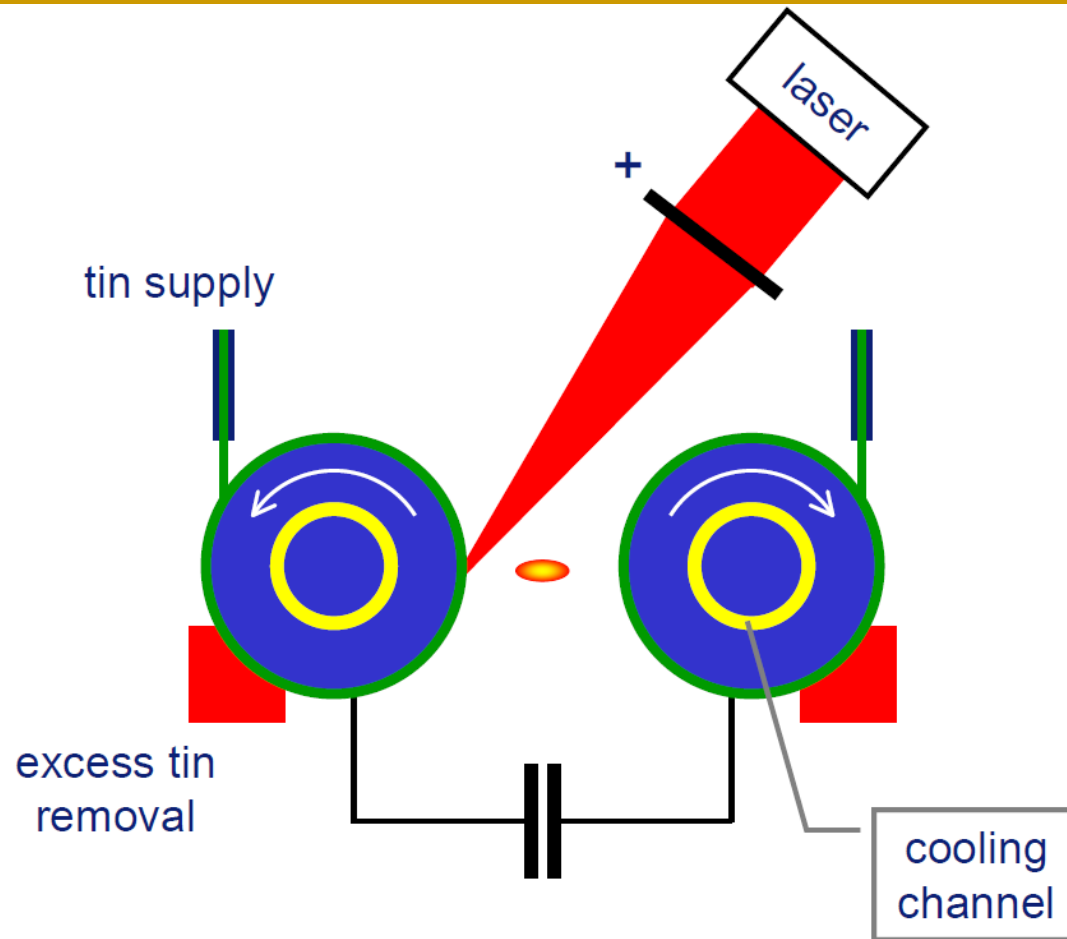


- CO<sub>2</sub> laser power: 47 kW
- Conversion efficiency: 3.3%

# Development of a Laser Produce Plasma EUV Source has Significant Challenges



# Philips EUV source



- CE 2%
- 5 kHz
- 120W continuous operation, 260 W short term

# Basic light generation principle: Sn *D*ischarge *P*roduced *P*lasma using rotating electrodes

- Laser Triggered Vacuum Spark
- Electrical contact through tin
  - Simple power supply to load capacitor bank
- Regenerating liquid tin surface
  - Electrode erosion problem fundamentally solved !
- Liquid metal cooling with tin
  - Very efficient to remove excess heat:  
>>100kW input power

