

Nanometer Scale Patterning and Processing

Spring 2016

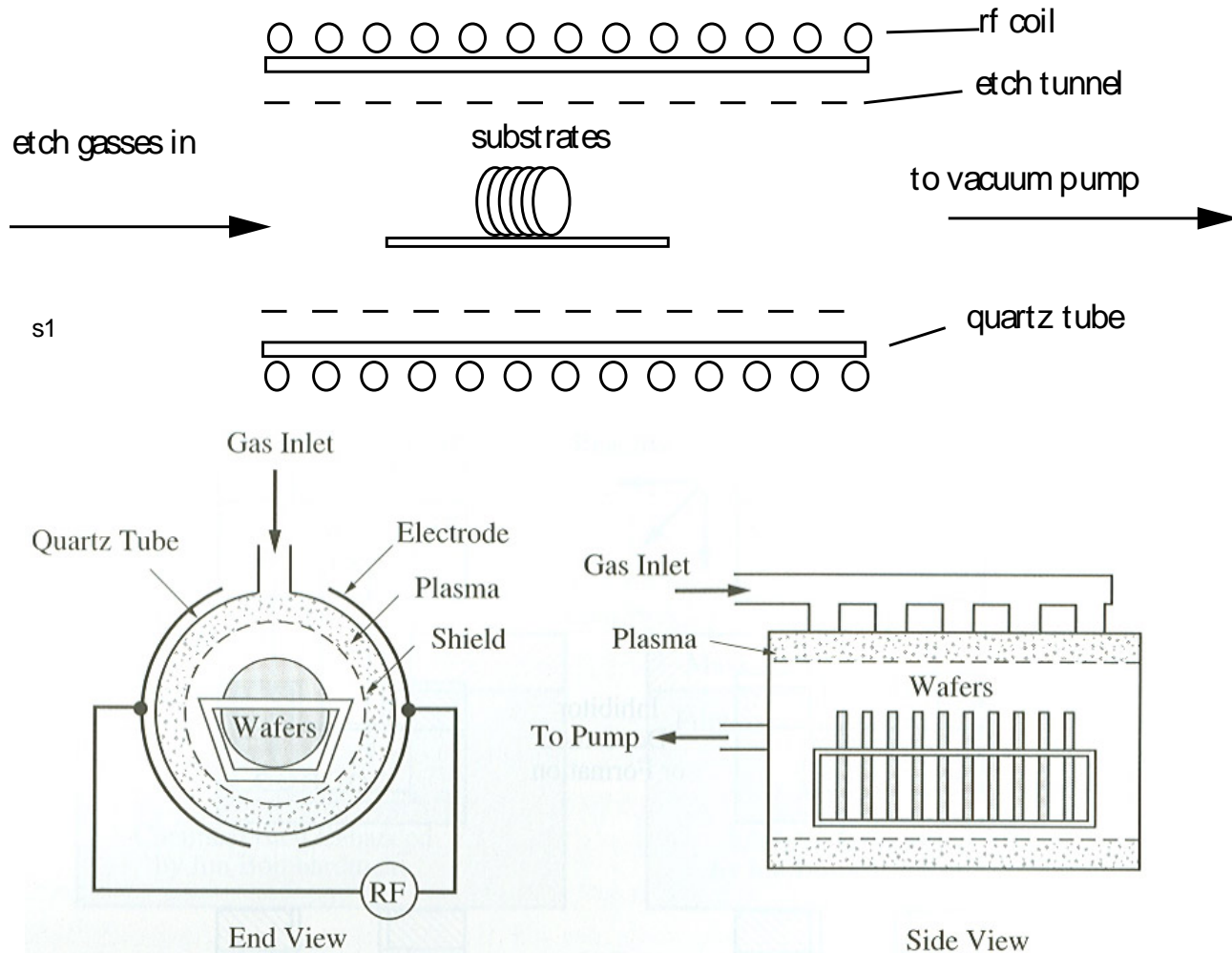
Lecture 45

Dry Etching, Part 1



DRY ETCHING

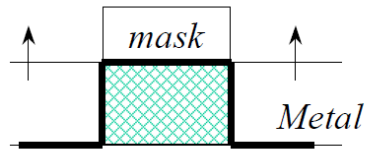
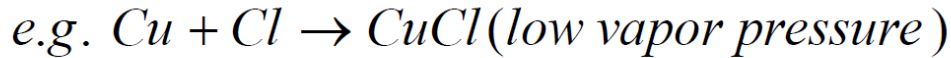
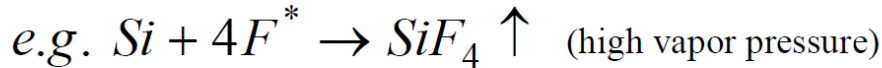
Pure Chemical Etch in Plasmas



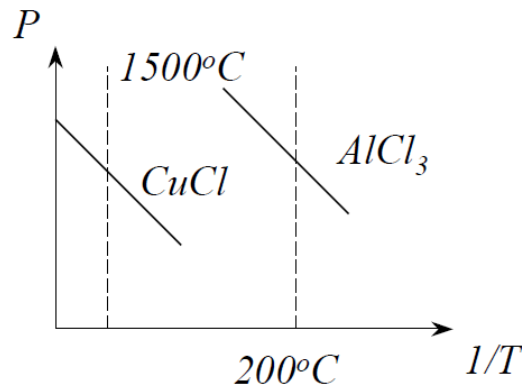
- Isotropic
- Like wet etch, but in gas phase
- Little damage, poor uniformity (e.g. wafer edge to center)
- Suitable for resist stripping
 - $\text{Polymer} + \text{O}_2 \rightarrow \text{CO}_2 + \text{H}_2\text{O}$

Volatility of Etching Product

* Higher vapor pressure \Rightarrow higher volatility



$$P = P_0 e^{-\Delta H_v / kT}$$



Cl_2 as etching gas.

Example

Difficult to RIE Al-Cu alloy with high Cu content

\swarrow 1~2% typical
[Al-Cu alloy]

BOILING POINT (°C) AT 1 atm

Chlorides

$AlCl_3$	177.8
$SiCl_4$	57.6
Cu_2Cl_2	1490
$TiCl_4$	136.4
WCl_6	346.7

Fluorides

AlF_3	1291
SiF_4	-86
Cu_2F_2	1100
TiF_4	284
WF_6	17.5

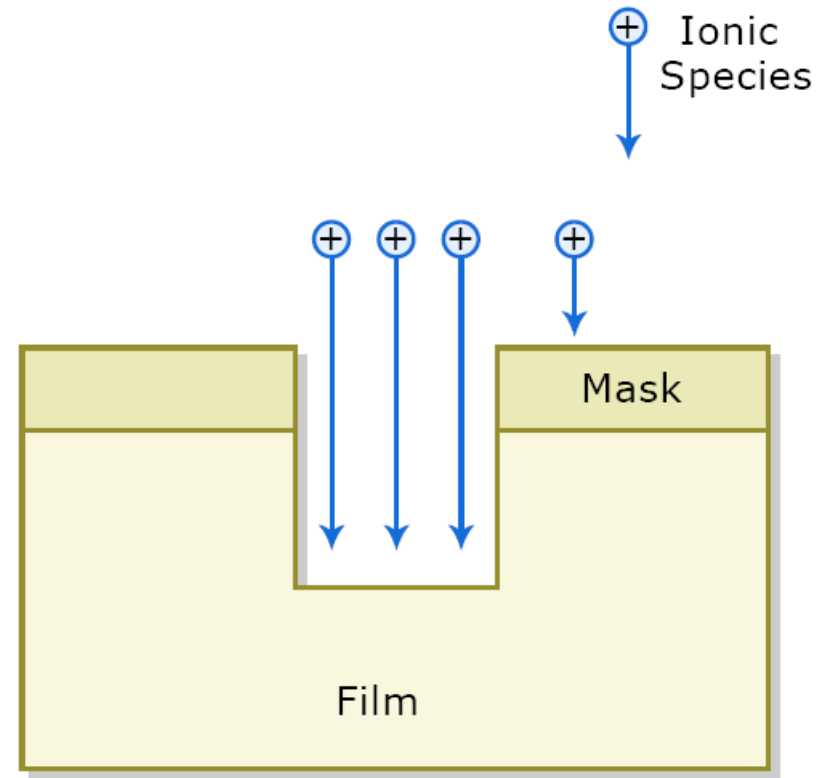
Chemistry	Primary Films Etched
Cl-based (Cl_2 , BCl_3)	Al alloys, Ti, TiN, resist
F-based (SF_6 , CF_4 , CHF_3)	W, TiW, SiO_2 , resist
O-based (O_2 , O_3 , CO_2 , H_2O)	Resist

BOILING POINTS OF TYPICAL ETCH PRODUCTS

ELEMENT	CHLORIDES	BOILING POINT (°C)	FLUORIDES	BOILING POINT (°C)
Al	AlCl ₃	177.8 (subl.)	AlF ₃	1291 (subl.)
Cu	CuCl	1490	CuF	1100 (subl.)
Si	SiCl ₄	57.6	SiF ₄	-86
Ti	TiCl ₃	136.4	TiF ₄	284 (subl.)
W	WCl ₆	347	WF ₆	17.5
	WCl ₅	276	WOF ₄	187.5
	WOCl ₄	227.5		

Physical Etching

- Wafer exposed to inert gas and/or plasma
- Etch is from momentum transfer from accelerated ions
 - Acting anisotropically
 - Good etch definition
 - Low selectivity
 - Rate is mass transport dependent
- Parallel version of Focused Ion beam etching



Combination of Chemical and Physical etching

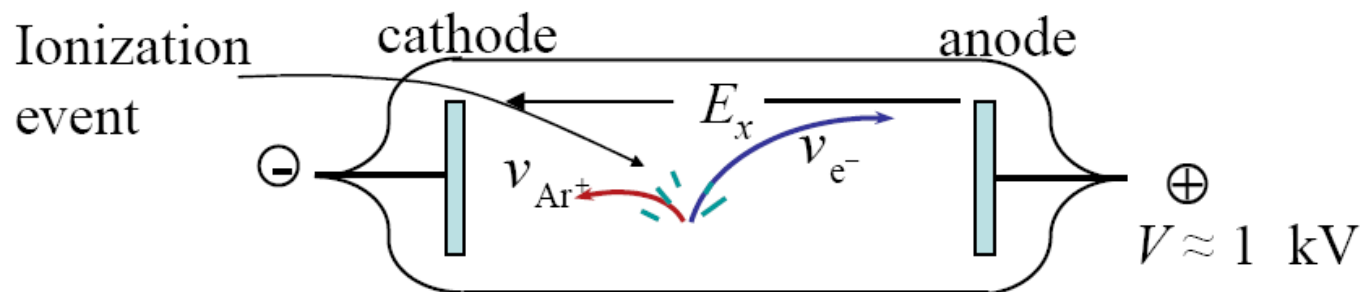
Ion-enhanced or Reactive Ion Etching (RIE)

Blends best of directionality and selectivity

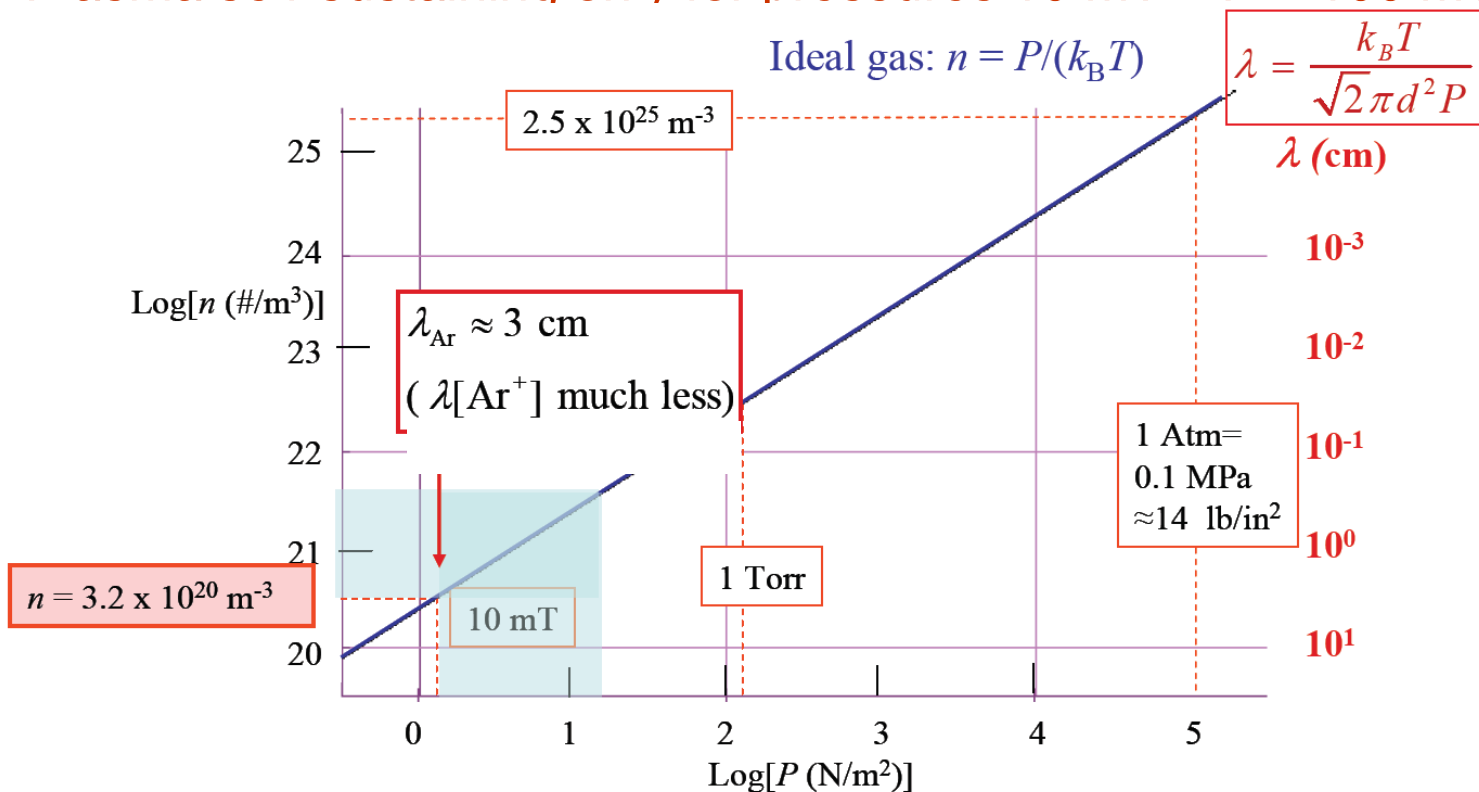
Plasma: A gas of ionized particles

Ar gas discharge

Ionization potential of Ar ≈ 16 eV



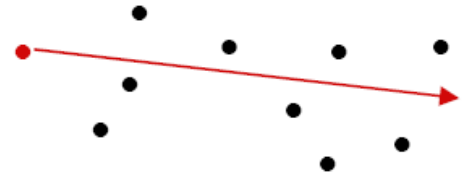
Plasma self-sustaining only for pressures $10 \text{ mT} < P < 100 \text{ mT}$



At 10 mT, molecular spacing $\approx n^{-1/3} = 0.15$ microns

Particle Energy in Plasma

Spacing between molecules $\approx n^{-1/3} = 0.15$ microns.



Thermal velocity of Ar, $v \approx \sqrt{\frac{3k_B T}{m_{\text{Ar}}}} \approx 10^3$ m/s

Velocity of ions at $x = \lambda$:

$$v_f^2 = v_0^2 + 2ax \approx 2 \frac{Eq}{m} \lambda$$

E field accelerates Ar^+ , e^- between collisions.

$$v_{\text{Ar}} \approx 4 \times 10^5 \text{ m/s}$$

$$v_{e^-} \approx 2 \times 10^7 \text{ m/s,}$$

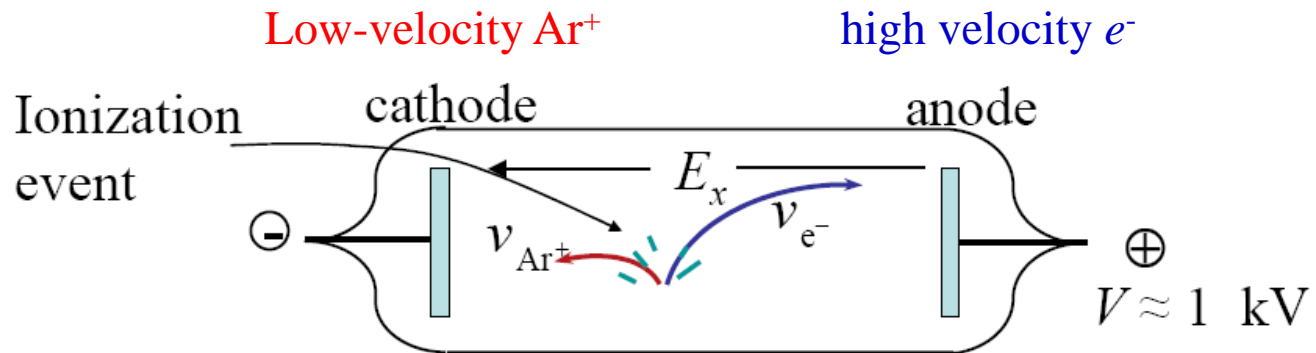
(only 0.1% to 1% of n_{Ar} are ions): **RDUE**
UNIVERSITY

New Picture of Plasma

$$v_{kT} \approx 10^3 \text{ m/s}$$

$$v_{Ar^+} \approx 4 \times 10^5 \text{ m/s}$$

$$v_{e^-} \approx 2 \times 10^7 \text{ m/s}$$



The plasma is highly conducting due to electrons:

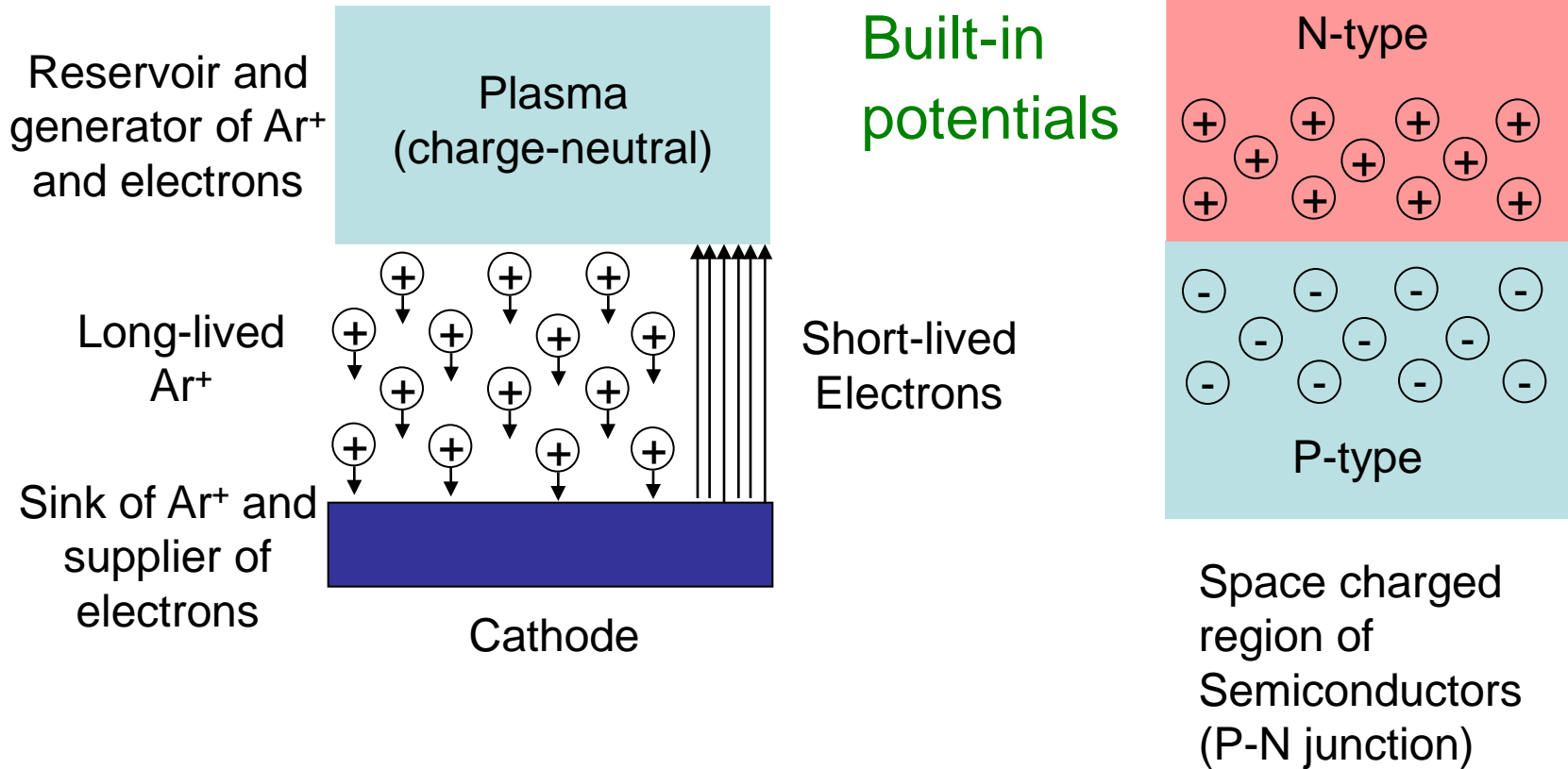
$$J = nq\bar{v}_x$$

Thus $J_{e^-} \gg J_{Ar^+}$

$$J_{e^-} = \sigma E = nq\bar{v}_x \approx nq \frac{at}{2} \approx nq \frac{Eq \lambda}{2m v}$$

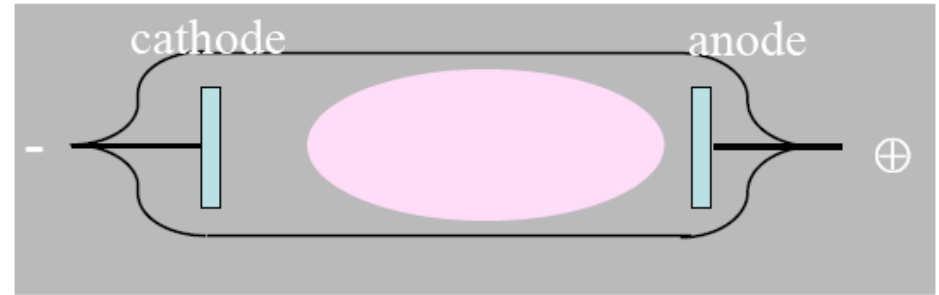
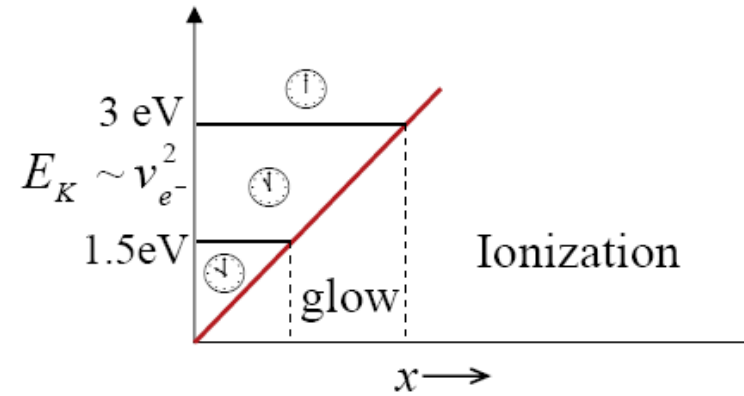
$$\sigma_{e^-} \approx \frac{nq^2 \lambda}{2m v} = \frac{ne^2 \tau}{2m}$$

Origin of the Sheath Voltage



Inside a plasma

(D.C. or “cathode sputtering”)

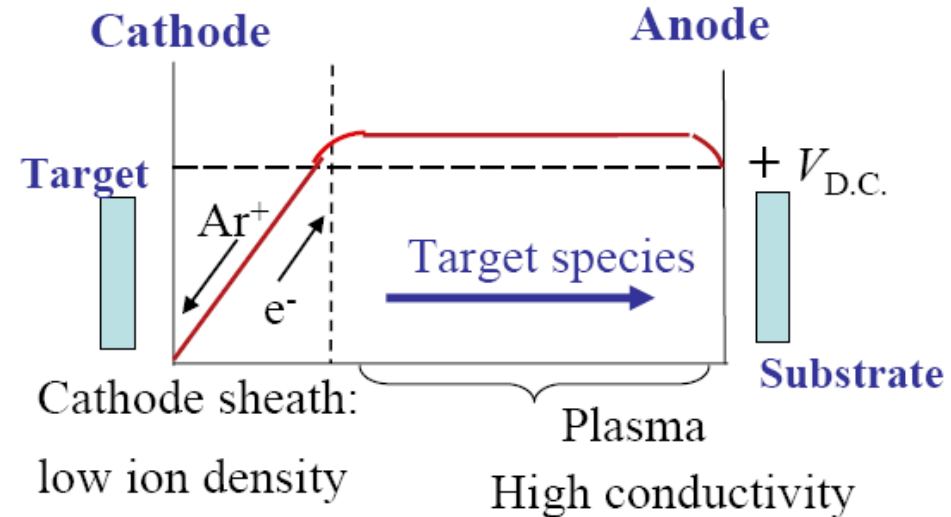


$$J_{e^-}, v_{e^-} \gg J_{Ar^+}, v_{Ar^+} \Rightarrow$$

Surfaces in plasma charge negative, attract Ar^+ repel e^-

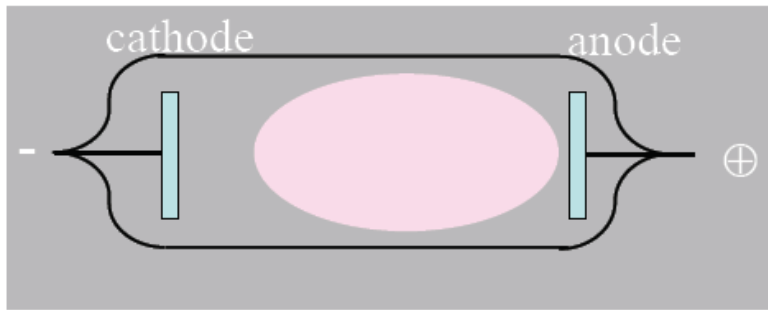
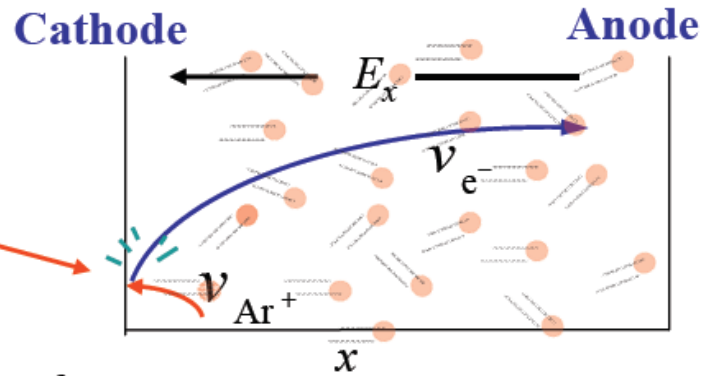
\therefore plasma ≈ 10 V positive

relative to anode

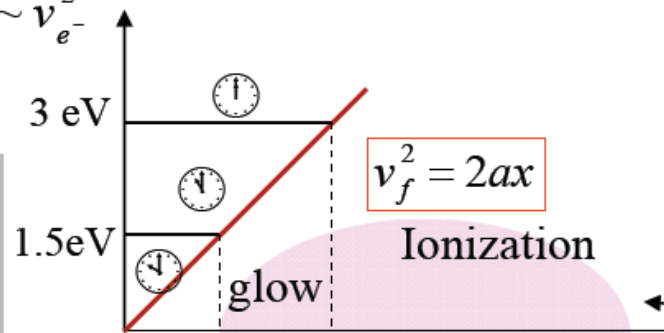


What about Glow?

Ar⁺ impact on cathode
 => Lots more electrons, plus sputtered atoms



$$E_K \sim v_{e^-}^2$$



⌚ $E_K^{e^-} \lesssim 1.5\text{eV}$

⌚ $1.5 \lesssim E_K^{e^-} \lesssim 3\text{eV}$

⌚ $E_K > 3\text{eV}$

Cathode dark space,
no action

e⁻ - induced
optical excitation
of Ar => visible glow

=> ionization,
High conductivity
plasma

Pressure range for normal glow-discharge plasmas

Necessary conditions:

If pressure is too low, λ is large, too few collisions *in plasma* to sustain energy

$$1) \lambda < L$$

so collisions exchange energy
within plasma

$$\lambda = \frac{k_B T}{\sqrt{2} \pi d^2 P}$$

If pressure is too high, λ is small, very little acceleration between collisions

$$2) E_K > \text{ionization potential of Ar}^+$$

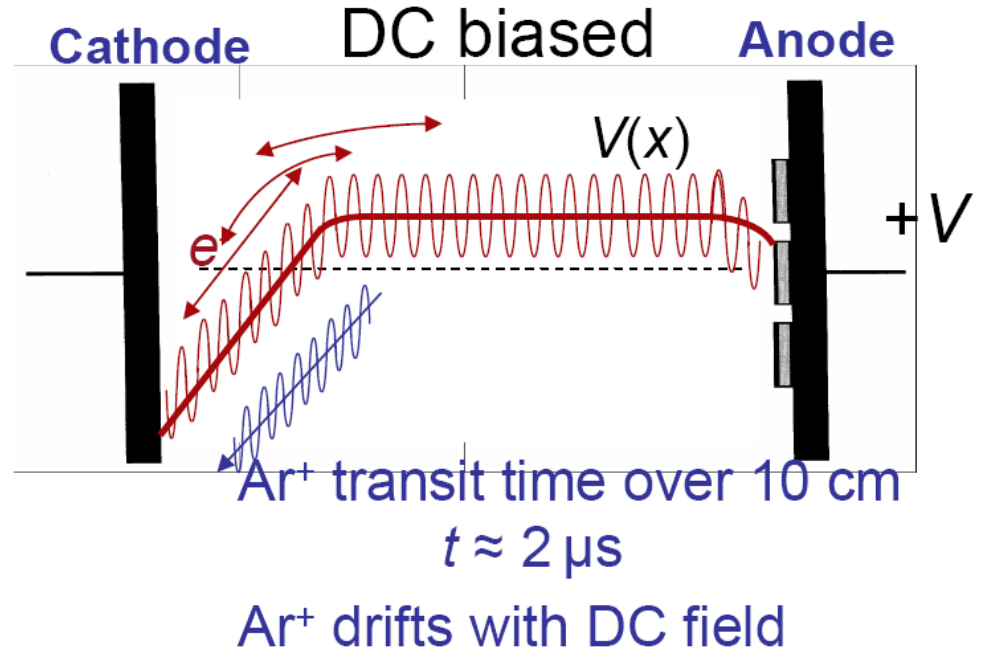
$$\frac{1}{2} m v_f^2 = 2 a x \approx E q \lambda$$

High-density plasmas can sustain at lower pressures

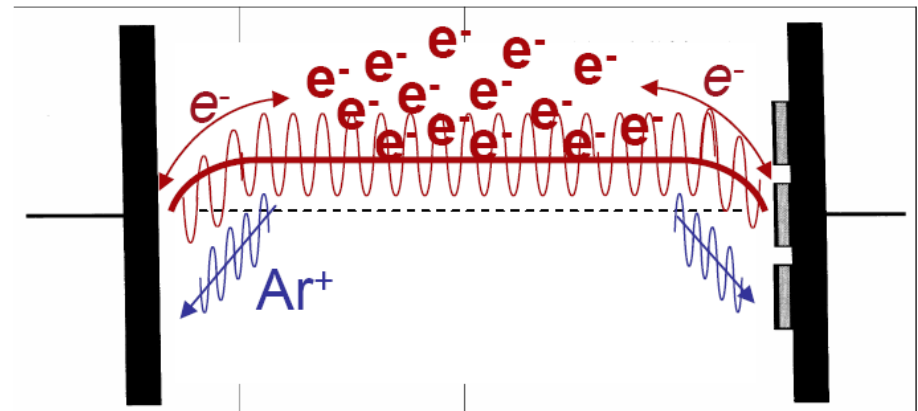
RF Plasma

$f = 13.6 \text{ MHz}$, $\tau \approx 12 \text{ ns}$

e^- transit time over 10 cm:
 $t \approx 10 \text{ ns}$.
 e^- follows RF field



But wait a minute!
If the plasma is
a good conductor,
does the RF field
penetrate it?



Plasma is conductive but not a good one

τ : average time between collisions

$$\sigma = \frac{ne^2\tau}{m}$$

We estimated $\tau \approx 0.01 \mu\text{s}$, so at 10 mT, $\sigma \approx 300 \text{ s}^{-1}$

Is this a good metal?

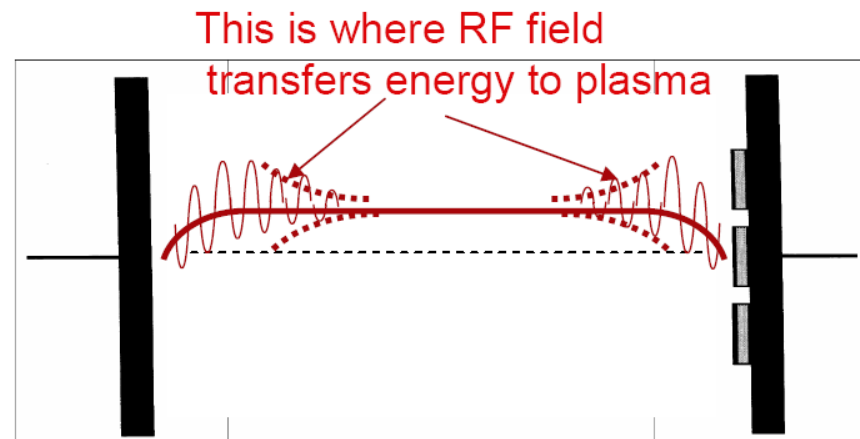
No!

Metals: $\rho_e < 100 \mu\Omega\text{-cm} = 1 \mu\Omega\text{-m}$, $\sigma > 10^6 \text{ s}^{-1}$

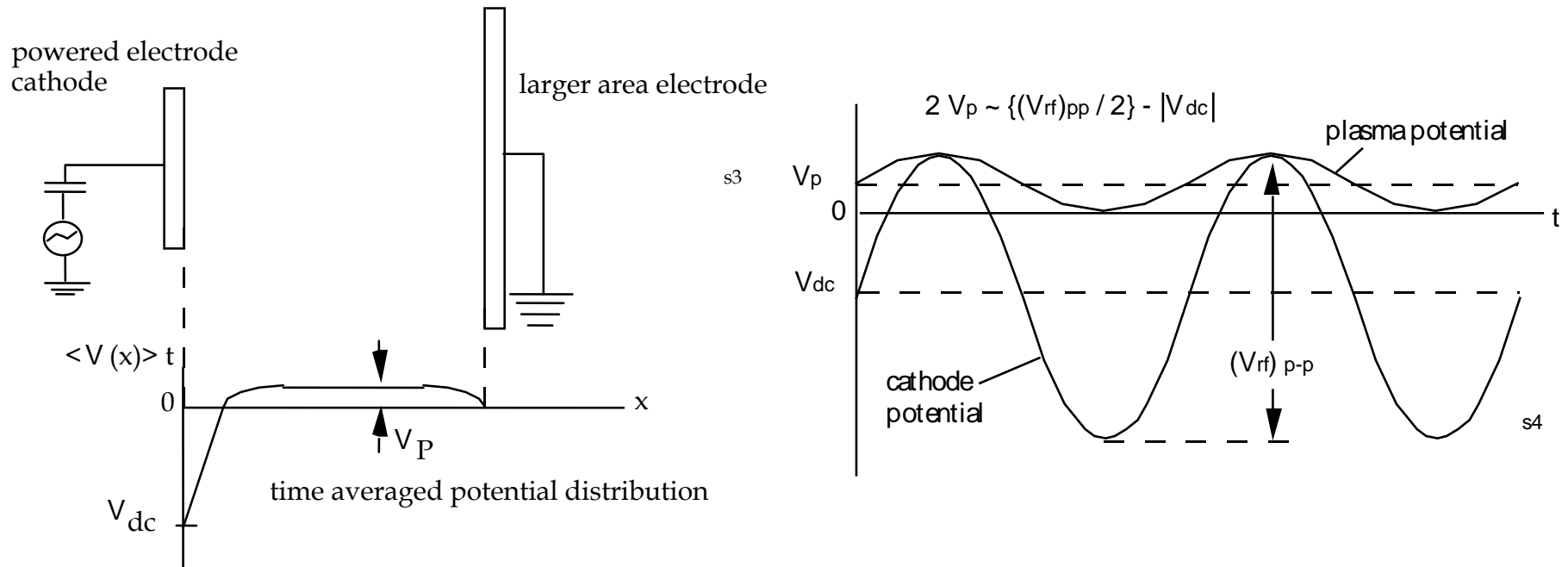
What then is the RF field penetration depth, skin depth?

$$\delta = \frac{1}{\sqrt{\mu\sigma\omega}} \approx 5 \text{ mm}$$

*Energy pumped in
from edges of plasma*



Charge balance in RF plasma



RF plasma allows electrodes to be insulators