

# Molecular-Level Modeling of Phonon Transport: Formulation, Implementation, and Applications

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# Outline

## Session I

1. Introduction (McGaughey)
2. Harmonic lattice dynamics, MD simulation (Ruan)
3. Green Kubo, direct method, spectral methods (Ruan)

## Session II

4. **Anharmonic lattice dynamics, first principles (McGaughey)**
5. Phonon-boundary and phonon-defect scattering (McGaughey)
6. Phonon-electron coupling and non-equilibrium (Ruan)

# Phonon Formula for Thermal Conductivity

Boltzmann transport equation + Fourier law

$$\Rightarrow k_n = \sum_i c_{v,i} v_{g,i,n}^2 \tau_i = \sum_i c_{v,i} v_{g,i,n}^2 \frac{\Lambda_i}{|\mathbf{v}_{g,i}|}$$

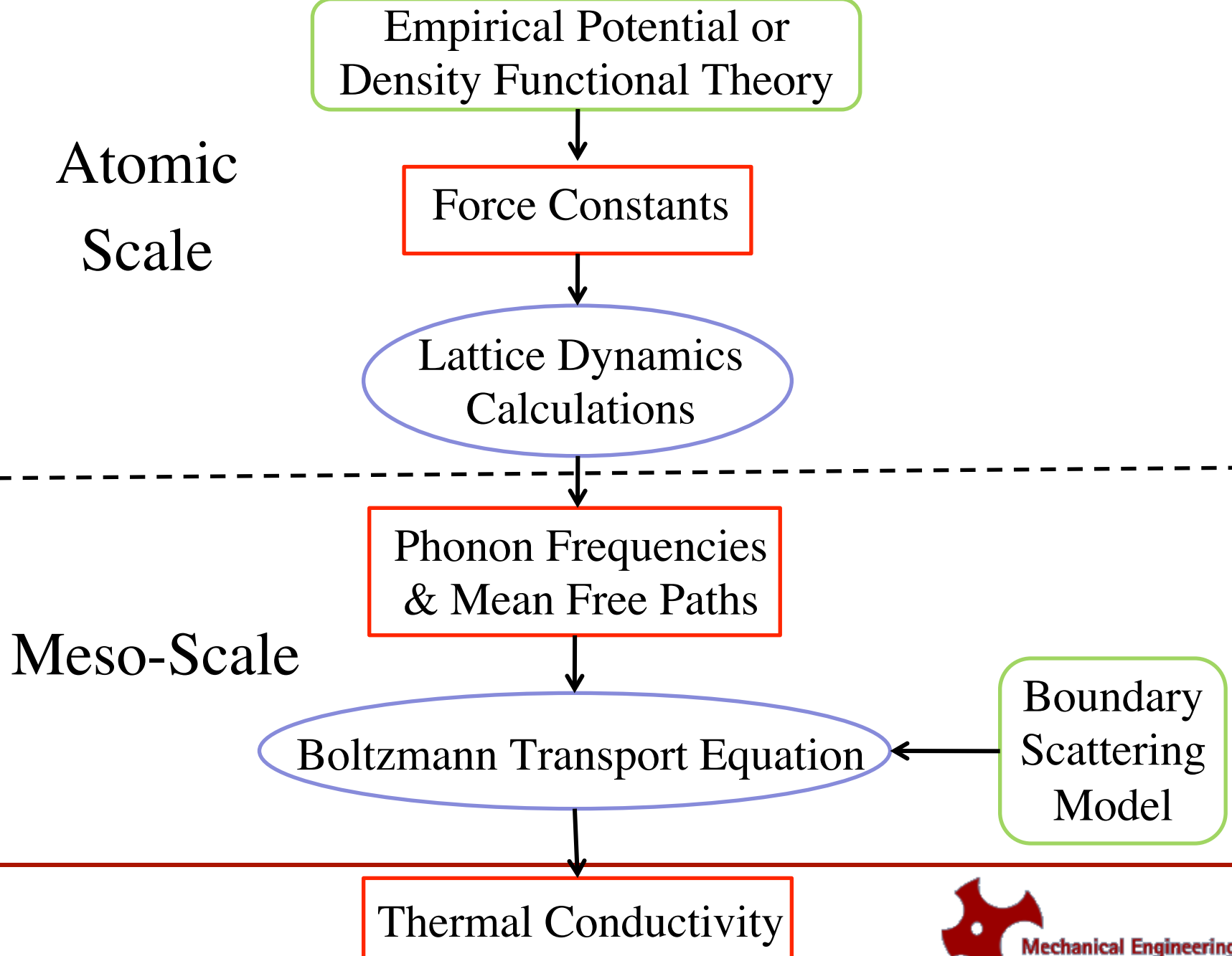
$i$ : indexes over all phonon modes

$c_{v,i}$ : heat capacity

$\mathbf{v}_{g,i}$ : group velocity

$\tau_i$ : lifetime

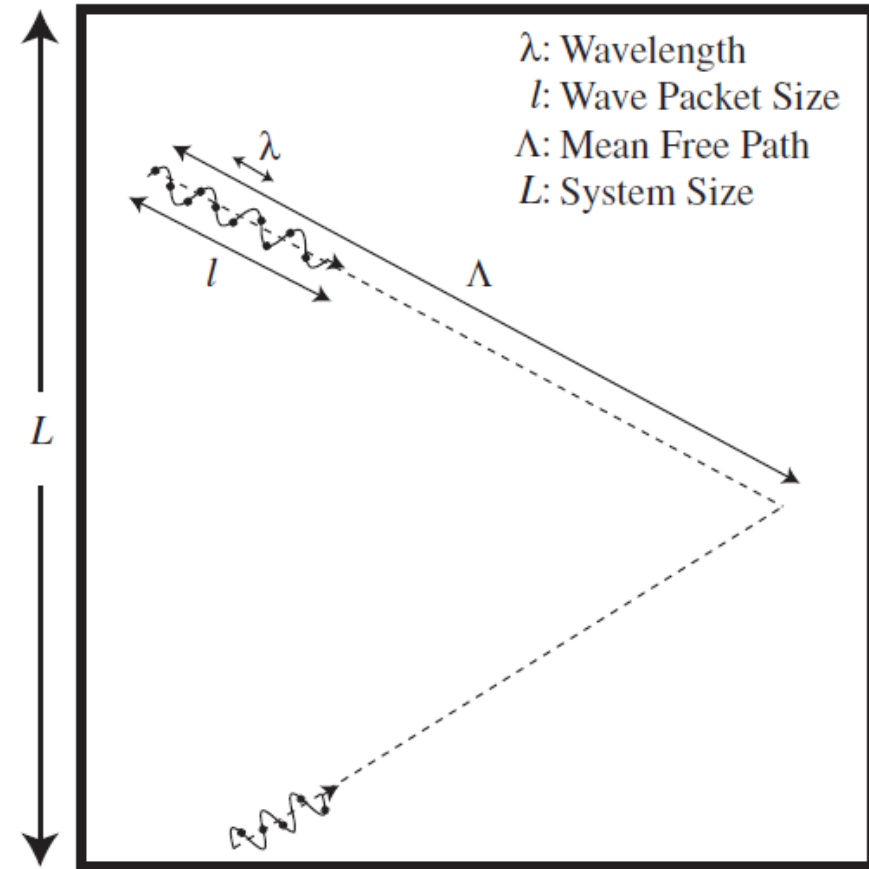
$\Lambda_i$ : mean free path



# Transport and Scattering

Phonons scatter with:

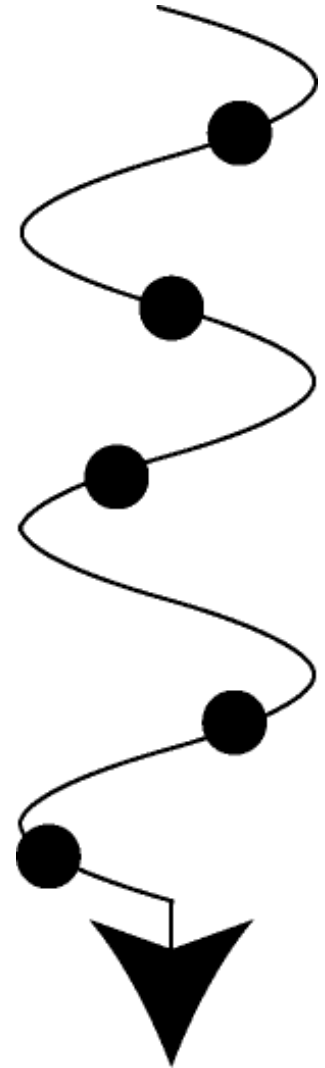
- **other phonons**
- grain boundaries, interfaces, surfaces
- electrons
- defects (isotopes, vacancies, ...)



McGaughey and Kaviani, *PRB* **71**, 184305 (2005).

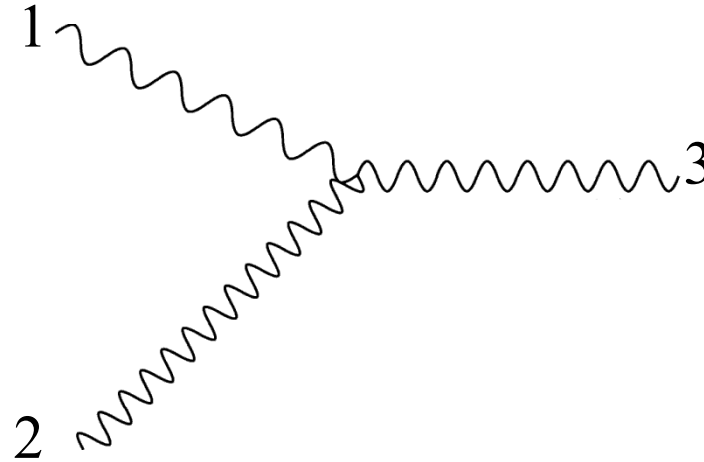
# Outline

1. Force Constants, Phonons, and Thermal Conductivity
2. First Principles Approach
3. Impactful Research



# Phonon-Phonon Scattering

Two phonons can  
combine to form a third  
(and vice-versa)



- Energy conservation:  $\hbar\omega_1 + \hbar\omega_2 = \hbar\omega_3$
- Translational invariance of the lattice:  $\mathbf{K}_1 + \mathbf{K}_2 = \mathbf{K}_3 + \mathbf{G}$   
Reciprocal lattice vector  
↓
- Interactions with  $> 3$  phonons possible,  
rates increase with temperature

# Lattice Dynamics on a Crystal Lattice

- Building from the mass-spring system, include:
  - Motion in three-dimensions
  - More than nearest-neighbor interactions
  - Non-linear interactions
  - Periodic boundary conditions
- Expand the system potential energy as a Taylor series:

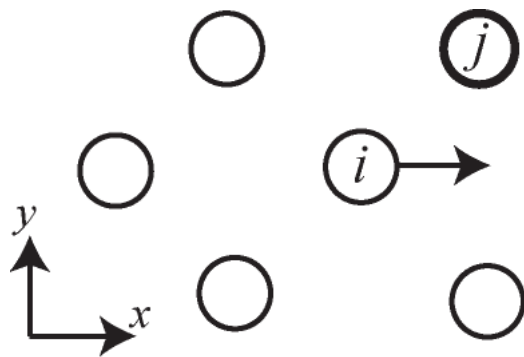
$$E = E_0 + \sum_i \sum_\alpha \frac{\partial E}{\partial u_{i,\alpha}} \Big|_0 u_{i,\alpha} + \frac{1}{2} \sum_{i,j} \sum_{\alpha,\beta} \overbrace{\frac{\partial^2 E}{\partial u_{i,\alpha} \partial u_{j,\beta}} \Big|_0}^{2^{\text{nd}} \text{ Order Force Constant}} u_{i,\alpha} u_{j,\beta} + \frac{1}{6} \sum_{i,j,k} \sum_{\alpha,\beta,\gamma} \underbrace{\frac{\partial^3 E}{\partial u_{i,\alpha} \partial u_{j,\beta} \partial u_{k,\gamma}} \Big|_0}_{3^{\text{rd}} \text{ Order Force Constant}} u_{i,\alpha} u_{j,\beta} u_{k,\gamma} + \dots$$

$E$  = potential energy

$u_{i,\alpha}$  = displacement of atom  $i$  in direction  $\alpha$



# Force Constant Calculation



## Harmonic Force Constant

1. Move  $i$  in  $x$ -direction.
2. Calculate force on  $j$  in  $y$ -direction.

$$\frac{\partial^2 E}{\partial u_{i,x} \partial u_{j,y}} = \frac{\partial}{\partial u_{i,x}} \frac{\partial E}{\partial u_{j,y}} = - \frac{\partial F_{j,y}}{\partial u_{i,x}}$$

- From an empirical interatomic potential
  - Analytically
  - Numerically
- From first-principles calculations
  - Density functional perturbation theory [Baroni et al., *RMP* **73**, 515 (2001)]
  - Numerically

Significant work by David Broido and colleagues.

# Harmonic Force Constants

- Dynamical matrix -> eigenvalue problem
  - Frequencies and mode shapes
  - Group velocities

$$\mathbf{v}_{g,i} = \frac{\partial \omega_i}{\partial \mathbf{\kappa}}$$

$$v_{g,i,n} = \frac{1}{2\omega_i} \left[ \mathbf{e}_i^+ \frac{\partial \mathbf{D}(\mathbf{\kappa})}{\partial \kappa_n} \mathbf{e}_i \right]$$

Wang et al., *Eur. Phys. J. B* **62**, 381 (2008)

# Cubic Force Constants

- Anharmonic + frequencies & mode shapes
  - RTA - > phonon lifetimes

$$\frac{1}{\tau_{q\nu}^{pp}} = \frac{\pi \hbar}{16N} \sum_{q'\nu'} \sum_{q''\nu''} \left| V_{\nu\nu'\nu''}^{qq'q''} \right|^2 \left\{ (n_{q'\nu'}^o + n_{q''\nu''}^o + 1) \delta(\omega_{q\nu} - \omega_{q'\nu'} - \omega_{q''\nu''}) + (n_{q'\nu'}^o - n_{q''\nu''}^o) \right. \\ \left. \times [\delta(\omega_{q\nu} + \omega_{q'\nu'} - \omega_{q''\nu''}) - \delta(\omega_{q\nu} - \omega_{q'\nu'} + \omega_{q''\nu''})] \right\},$$

- BTE full solution (many flavors) -> scattering rates

# Quantum vs. Classical Statistics

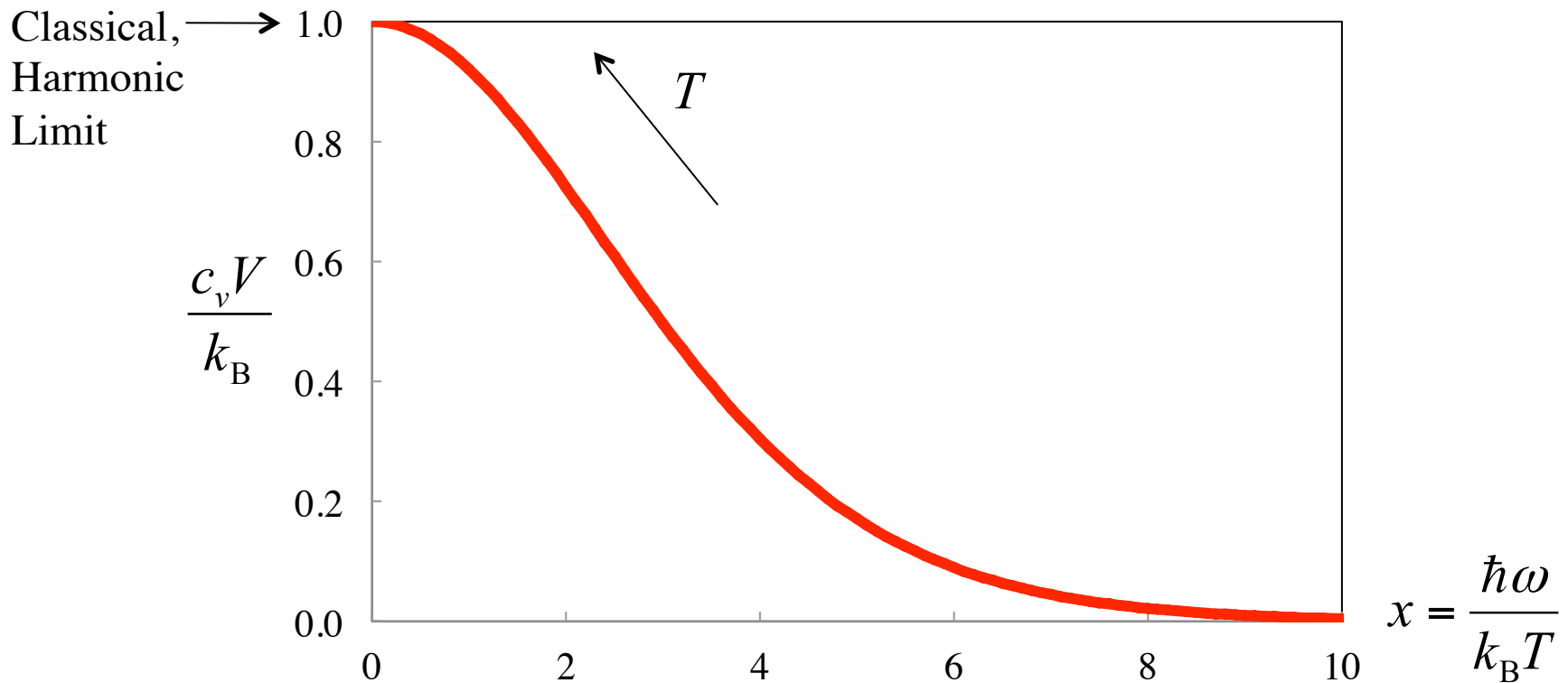
MD simulations are classical

- High temperature limit of Bose-Einstein,  $x = \hbar\omega/k_B T \rightarrow 0$
- Equipartition of energy in a harmonic system

	Quantum	Classical
$f$	$1/(e^x-1)$	$1/x$
$c_v V$	$k_B x^2 e^x / (e^x - 1)^2$	$k_B$

# Heat Capacity

- Phonons are bosons, described by Bose-Einstein statistics
- Energy of phonon mode  $i$  is  $E_i = \hbar\omega_i f_i$
- Heat capacity is  $c_{v,i} = \frac{1}{V} \frac{\partial E_i}{\partial T} = \frac{k_B x_i^2}{V} \frac{\exp(x_i)}{[\exp(x_i) - 1]^2}$



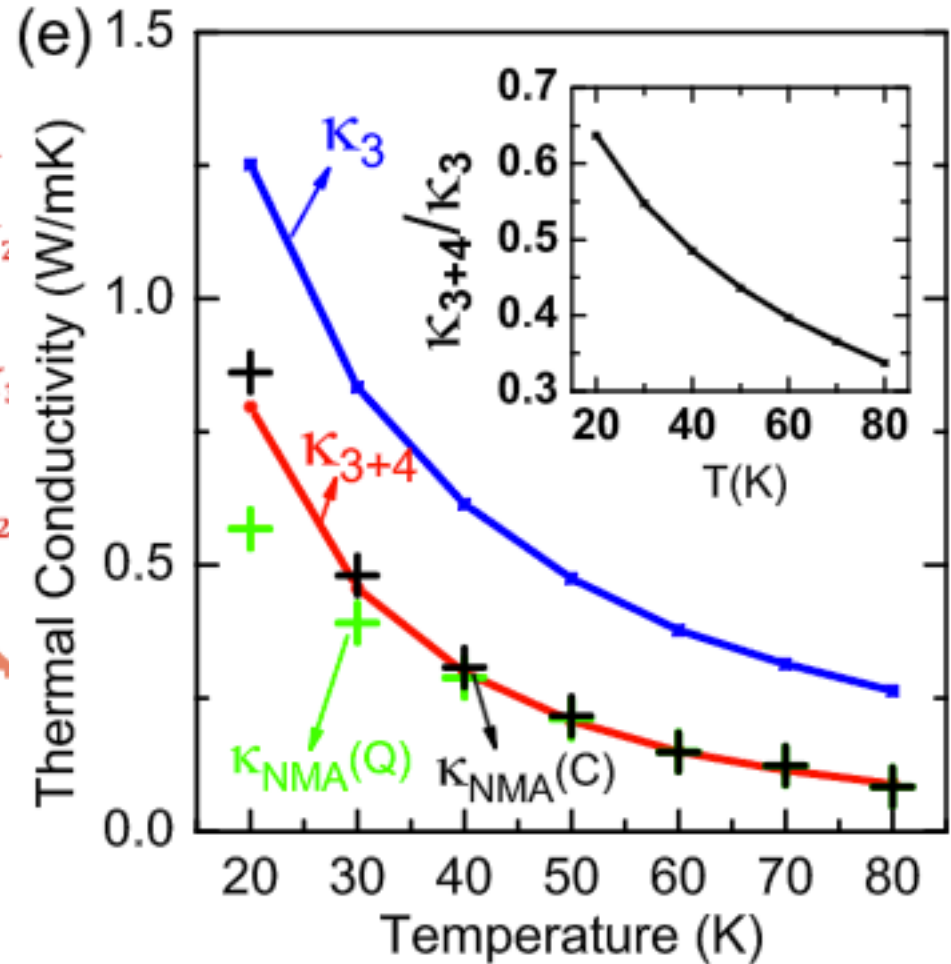
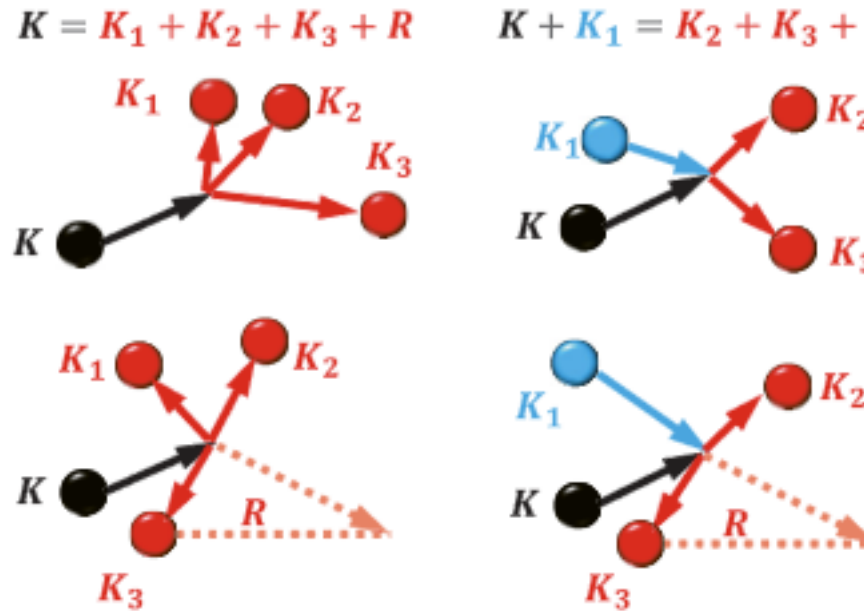
# LJ Crystal Self-Consistent Comparison

$T$ (K)	GK-MD <sup>a</sup>	Direct-MD	BTE-MD	BTE-LD
20	1.2	1.4	1.3	1.4
30	0.72	0.76	0.69	0.89
40	0.47	0.50	0.49	0.63
50	0.32	0.34	0.30	0.49
60	0.26	0.29	0.20	0.38
70	0.20	0.21	0.13	0.31
80	0.16	0.17	0.086	0.26

- Limited to low/medium temperatures
- Cannot include disorder explicitly
- Computational challenges for large unit cells

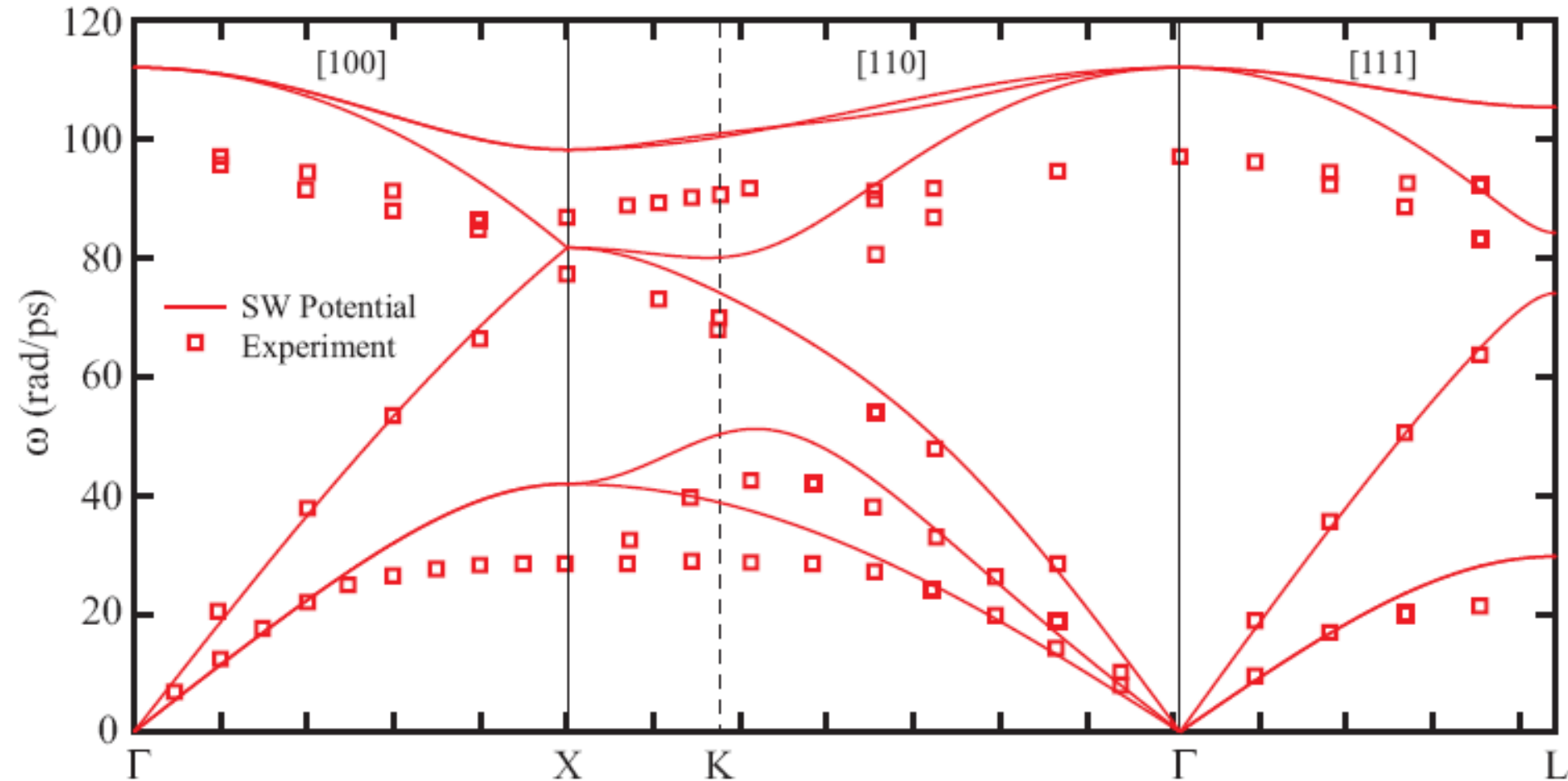
Turney et al., *PRB* **79**, 075316 (2009)

# Inclusion of Four-Phonon Processes



Feng and Ruan, *PRB* **93**, 045202 (2016)

# Failure of Empirical Potentials: Dispersion



Eric Landry PhD thesis (CMU, 2009).



# Failure of Empirical Potentials: Thermal Conductivity

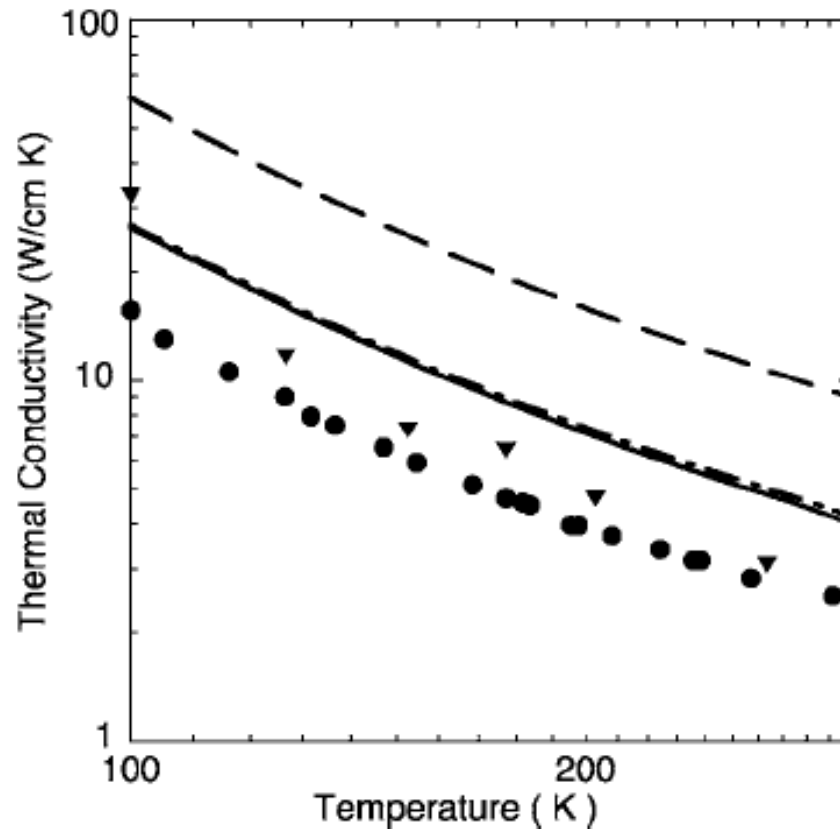
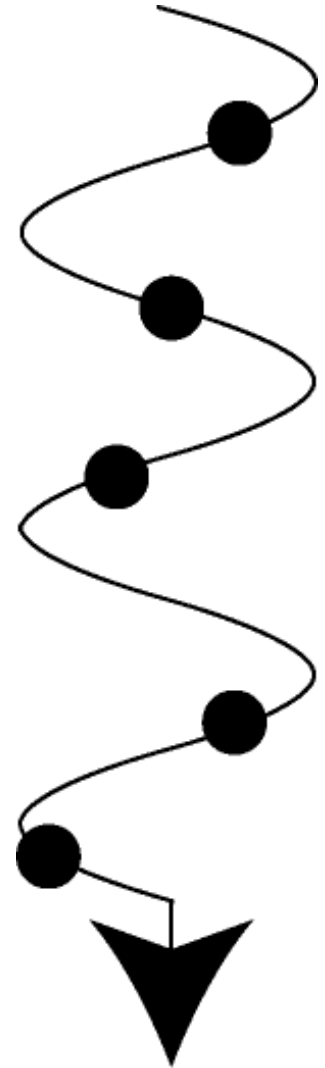


FIG. 5. Calculated lattice thermal conductivity for isotopically enriched Si using SW (dashed line), Tersoff (dashed-dotted line) and ED (solid line) models in the temperature range between 100 K and 300 K compared to measured values.

Broido et al., *PRB* 72 (2005) 014308.

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# Density Functional Theory

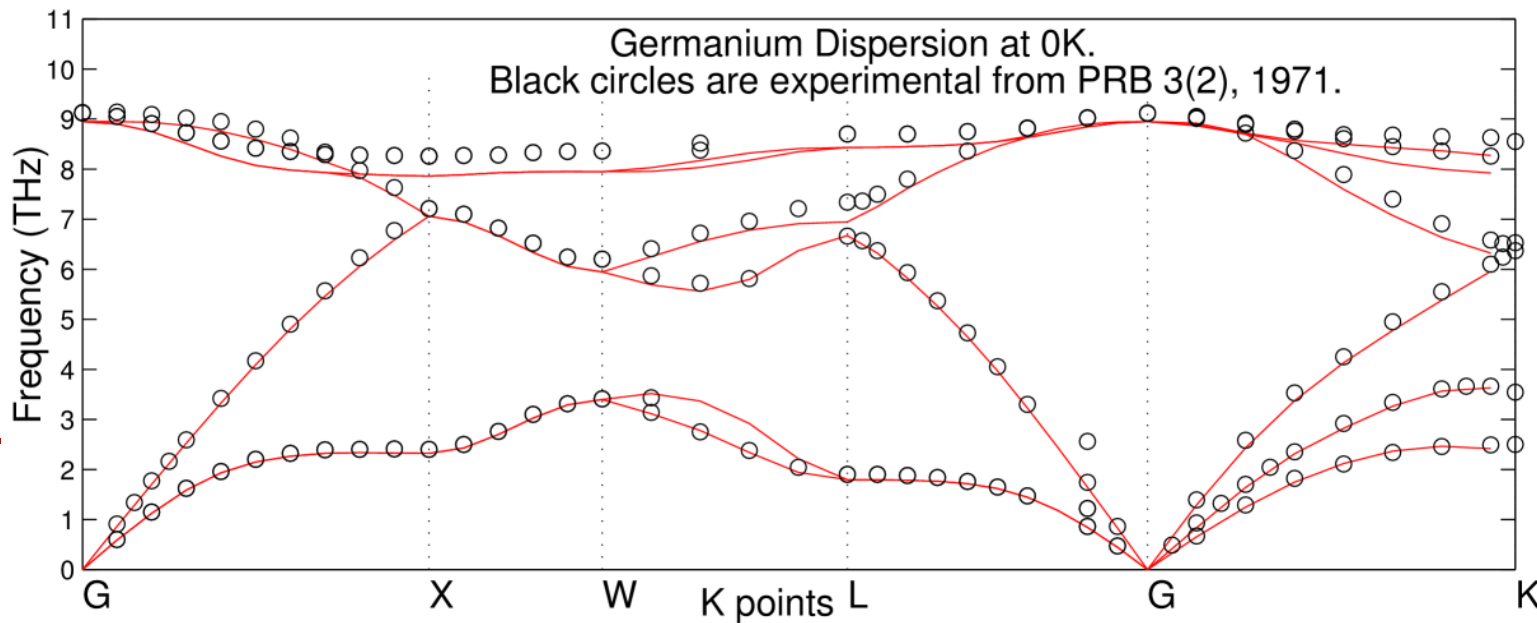
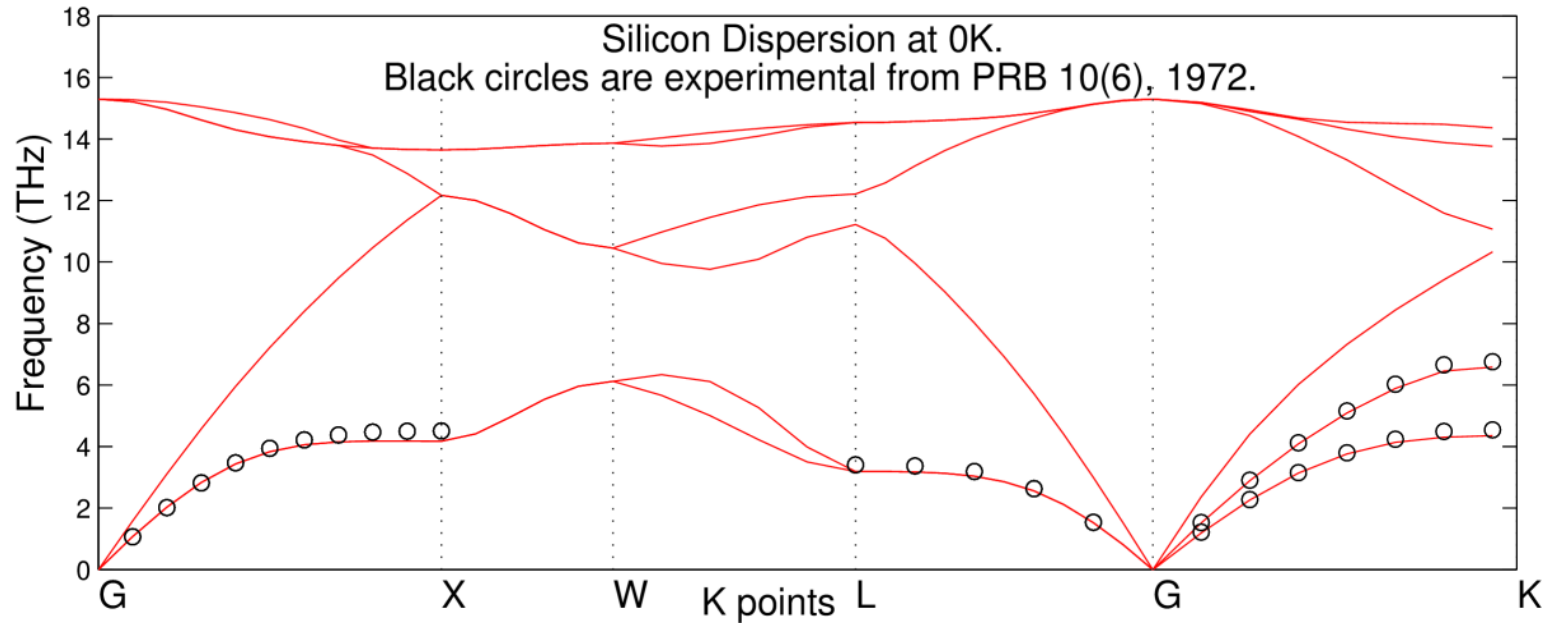
- Approximate numerical solution to the many-body Schrödinger equation  
*(bodies = electrons and ions)*
- Theoretically and computationally complex  
*(large number of degrees of freedom)*
- Expertise is needed to set up and run the calculations.
- Graduate students can write a molecular dynamics or harmonic lattice dynamics code in a one-semester class.  
*(not possible for DFT)*

# DFT -> Force Constants

$$E = E_0 + \sum_i \sum_{\alpha} \left. \frac{\partial E}{\partial u_{i,\alpha}} \right|_0 u_{i,\alpha} + \frac{1}{2} \sum_{i,j} \sum_{\alpha,\beta} \left. \frac{\partial^2 E}{\partial u_{i,\alpha} \partial u_{j,\beta}} \right|_0 u_{i,\alpha} u_{j,\beta} + \frac{1}{6} \sum_{i,j,k} \sum_{\alpha,\beta,\gamma} \left. \frac{\partial^3 E}{\partial u_{i,\alpha} \partial u_{j,\beta} \partial u_{k,\gamma}} \right|_0 u_{i,\alpha} u_{j,\beta} u_{k,\gamma} + \dots$$

- Density functional perturbation theory or finite differences
  - Symmetry reduction, translational invariance
- **Harmonic**: DFPT (standard), FD (convergence issues)
- **Anharmonic**: DFPT (non-standard), **FD (tractable)**

# Si and Ge Dispersion from First Principles

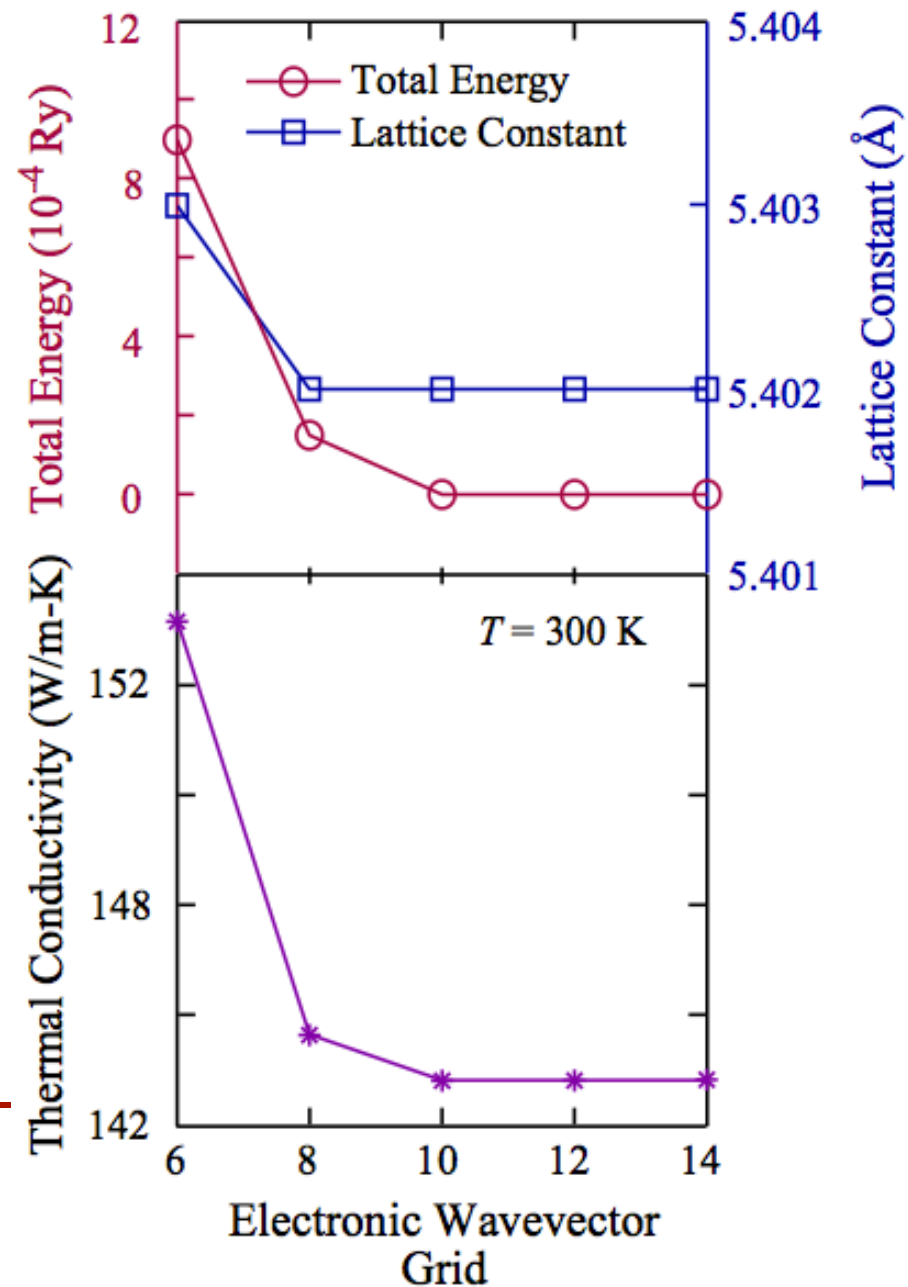


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# Convergence (1)

- DFT and DFPT calculations on isotopically pure silicon
  - Quantum Espresso
- Thermal conductivity < 2%
- Electron wave vector grid
  - Total energy ( < 0.2 mRy),  
lattice constant (< 0.001 Å)
  - Converged at 8x8x8
- Plane wave energy cutoff converged at 60 Ry



Jain and McGaughey,

# Convergence (2)

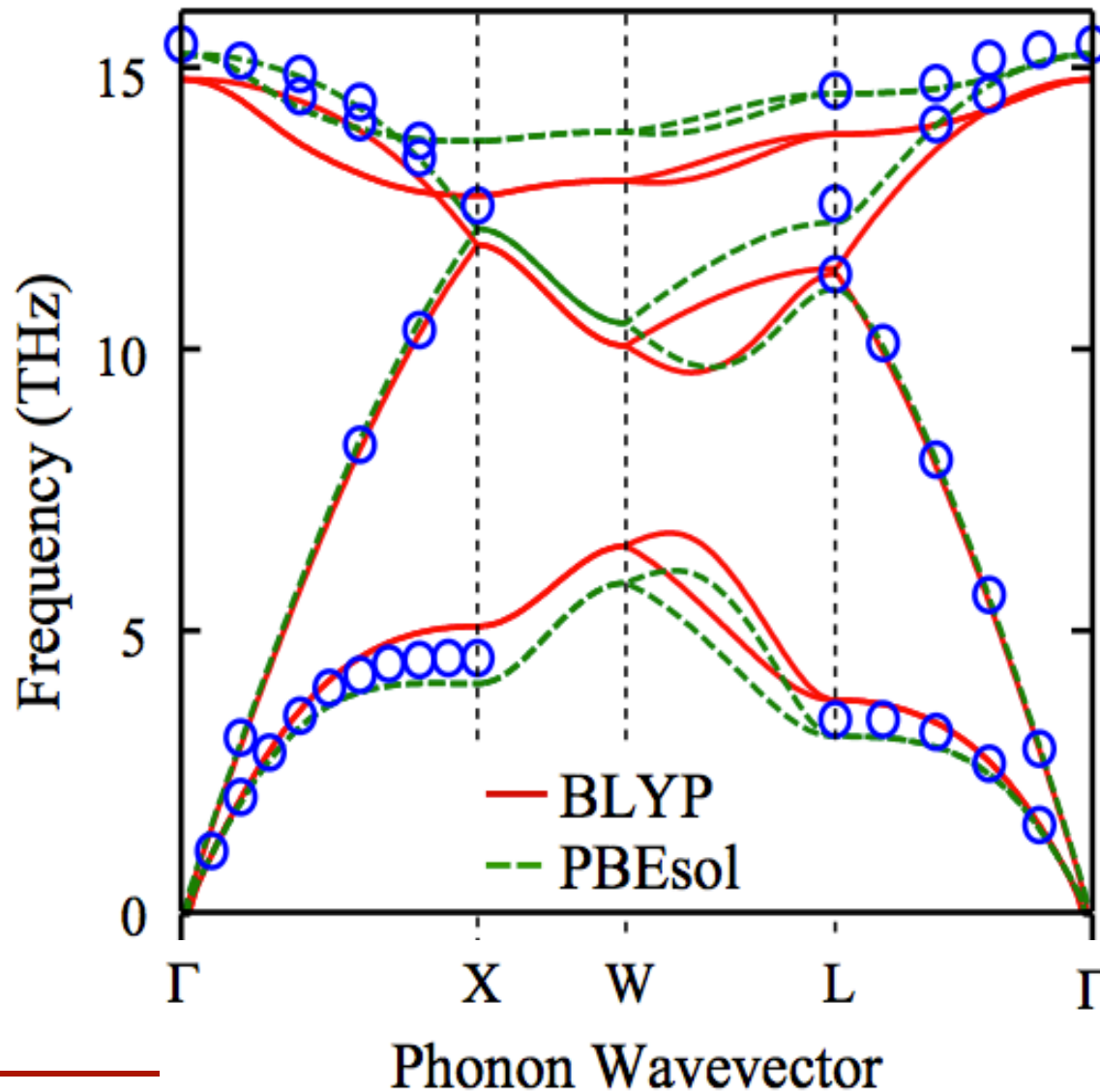
- Phonons
  - 8x8x8 wave vector grid for DFPT (harmonic force constants)
  - 216 atom supercell (cubic force constants)
  - 24x24x24 wave vector grid for thermal conductivity
- Pseudopotential: Core electron model
  - Norm conserving, ultrasoft, PAW
- Exchange correlation: Many-body interaction closure
  - LDA, PBE, PBEsol, PW91, BLYP



# Pseudopotential and Exchange Correlation

Pseudopotential type [24]	Exchange correlation [24]	Lattice constant (Å)	Sound velocity (m/s)
Experiment		5.430 [27]	8430 [30]
Ultrasoft	LDA	5.399	8320
	PBE	5.468	8120
	PBEsol	5.430	8330
	PW91	5.466	5970
Norm-conserving	LDA	5.402	7560
	PBE	5.461	8150
	BLYP	5.505	8510
PAW	LDA	5.400	8340
	PBE	5.466	7830
	PBEsol	5.430	8320

# Dispersion



# Pseudopotential and Exchange Correlation

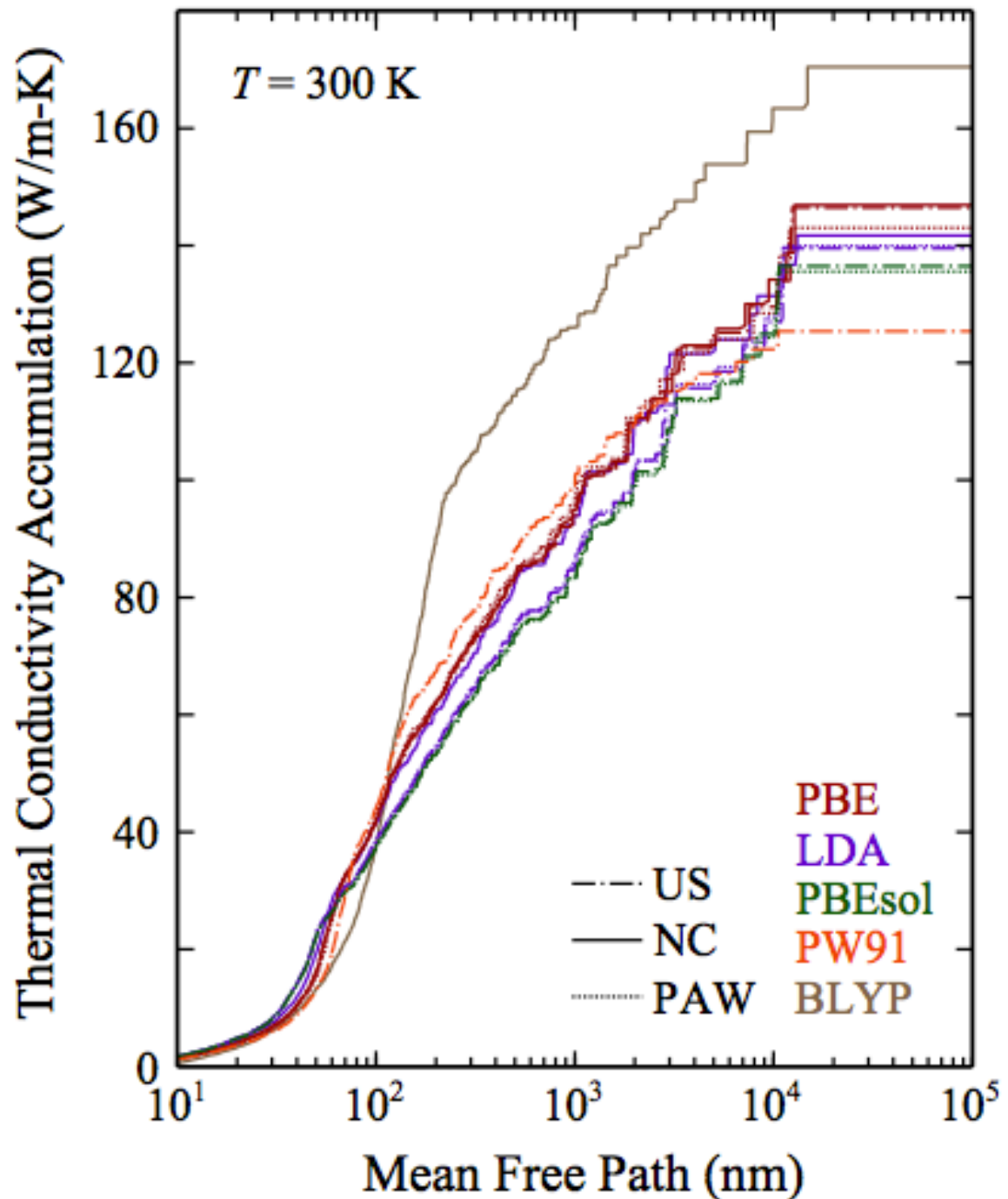
Pseudopotential type [24]	Exchange correlation [24]	Lattice constant (Å)	Sound velocity (m/s)	Thermal conductivity (W/m K) at 300 K
Experiment		5.430 [27]	8430 [30]	153 [19]
Ultrasoft	LDA	5.399	8320	142
	PBE	5.468	8120	148
	PBEsol	5.430	8330	140
	PW91	5.466	5970	127
Norm-conserving	LDA	5.402	7560	144
	PBE	5.461	8150	148
	BLYP	5.505	8510	172
PAW	LDA	5.400	8340	142
	PBE	5.466	7830	145
	PBEsol	5.430	8320	137

Our values: 127-148 W/m-K

Experiment: 153 W/m-K [Inyushkin, *Phys. Stat. Sol. (c)* **1** (2004) 2995]

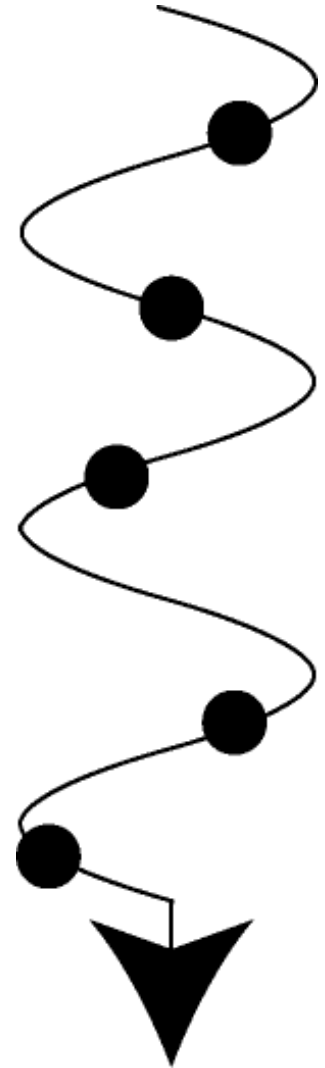
Literature DFT: 132-155, 172 W/m-K

# Thermal Conductivity Accumulation



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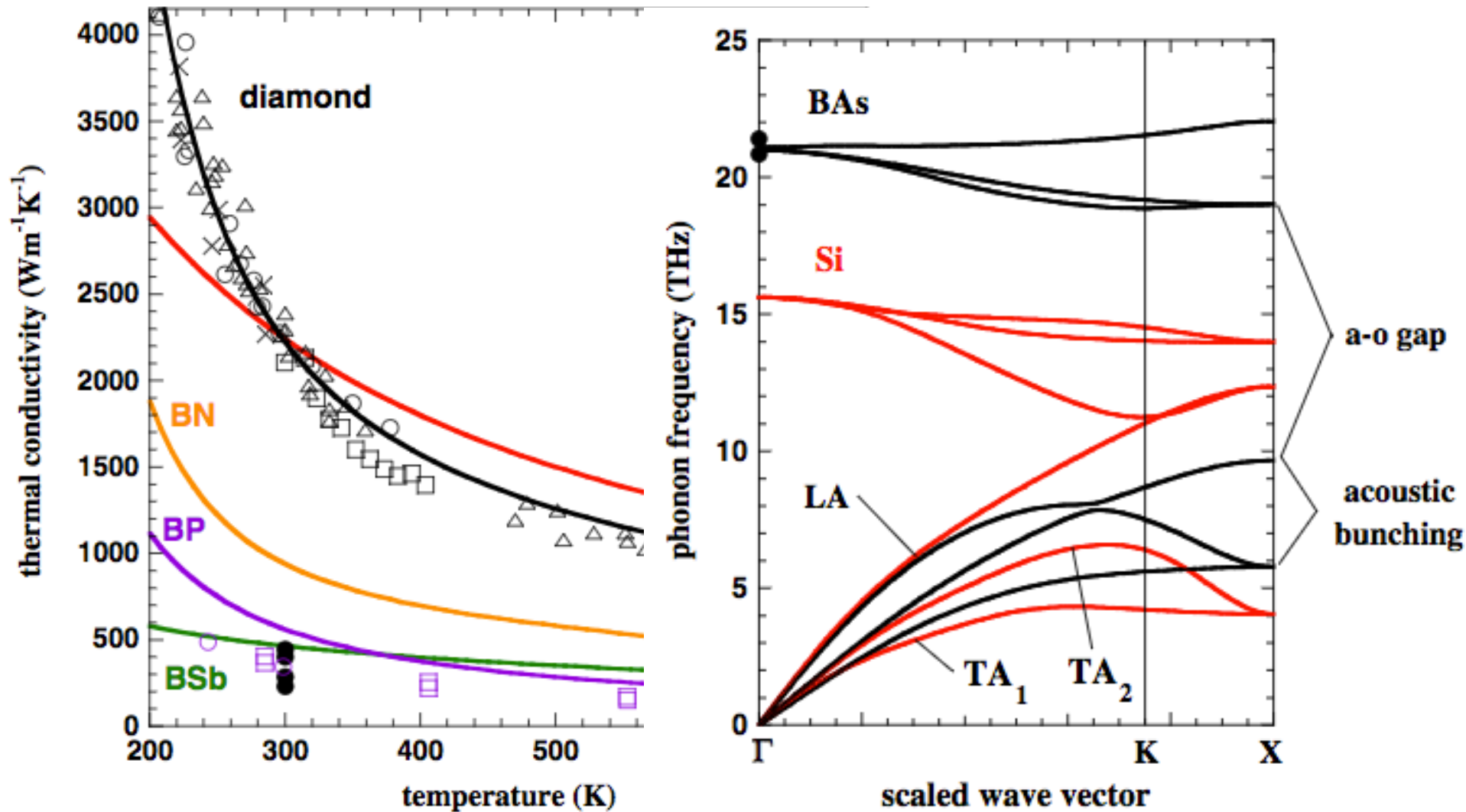


# Extensive Work Across Many Materials Systems

	$a_{\text{calc}} (a_{\text{exp}}) (\text{\AA})$	$g_{\text{cation}} (g_{\text{anion}}) (\times 10^{-4})$	$\kappa_{\text{pure}} (\text{W m}^{-1} \text{K}^{-1})$	$\kappa_{\text{natural}} (\text{W m}^{-1} \text{K}^{-1})$	$P$	$S$
Diamond <sup>a</sup>	3.53 (3.57)	0.75 (0.75)	3450	2290	51	50
Si	5.37 (5.43)	2.01 (2.01)	155	144	8	2.4
Ge	5.61 (5.65)	5.87 (5.87)	74	60	23	6.9
3C-SiC	4.34 (4.36)	2.01 (0.75)	572	479	20	6.7
AlP	5.40 (5.45)	— (—)	90	—	—	3.0
AlAs	5.61 (5.66)	— (—)	105	—	—	0.5
AlSb	6.10 (6.14)	— (0.66)	118	86	36	39
c-GaN <sup>b</sup>	4.42 (4.50)	1.97 (—)	362	215	68	13
GaP <sup>c</sup>	5.34 (5.45)	1.97 (—)	153	131	16	4.8
GaAs <sup>c</sup>	5.55 (5.65)	1.97 (—)	56	54	4	5.6
GaSb <sup>c</sup>	6.00 (6.10)	1.97 (0.66)	48	45	6	3.1
InP <sup>b</sup>	5.79 (5.87)	0.12 (—)	91	89	2	2.4
InAs <sup>b</sup>	5.97 (6.06)	0.12 (—)	36	36	.5	2.8
InSb <sup>b</sup>	6.39 (6.48)	0.12 (0.66)	20	20	2	3.9
w-GaN <sup>b,d</sup>	3.13 (3.19)	1.97 (—)	401 (385) <sup>g</sup>	242 (239) <sup>g</sup>	66	7.2
	5.10 (5.19) <sup>e</sup>					
	0.377 (0.377) <sup>f</sup>					
w-AlN	3.05 (3.11)	— (—)	322 (303) <sup>g</sup>	—	—	14
	4.81 (4.98) <sup>e</sup>					
	0.387 (0.382) <sup>f</sup>					

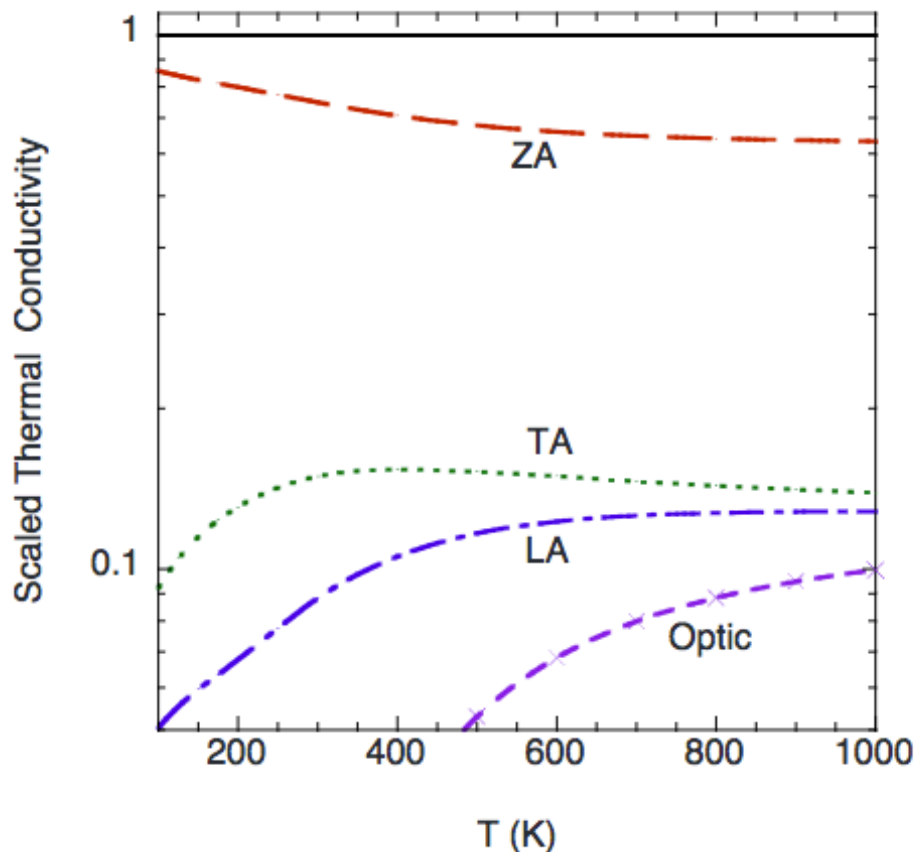
Lindsay et al., *PRB* **87** (2013) 165201.

# Ultra-High Thermal Conductivity of BAs



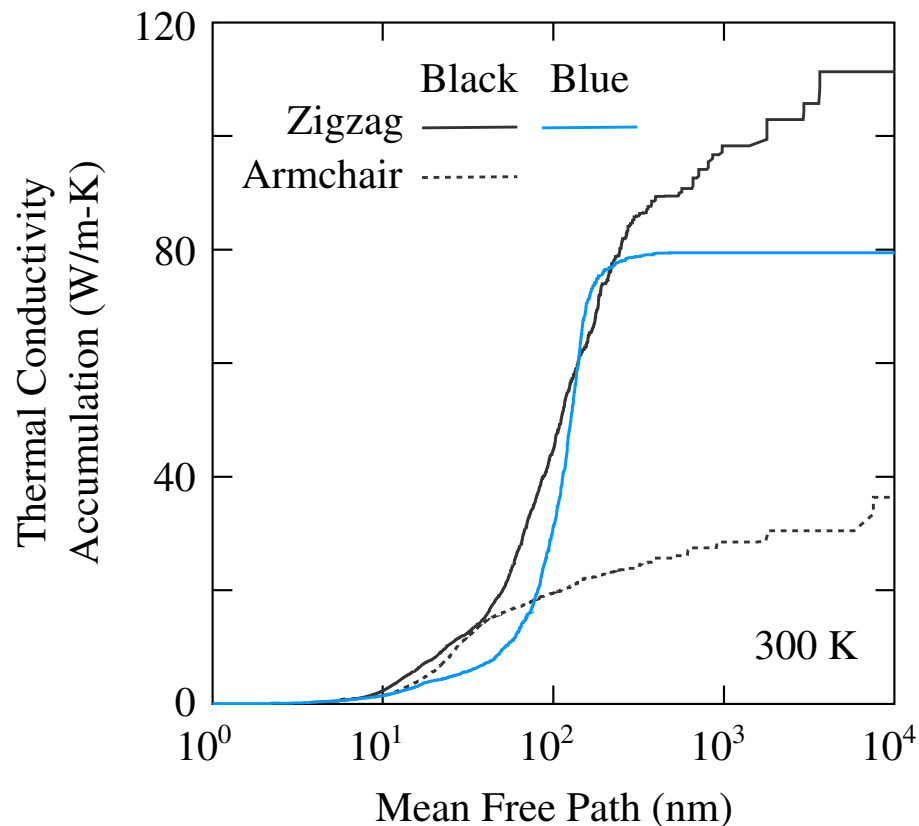
Lindsay et al., *PRL* **111** (2013) 025901.

# Two Dimensional Materials



Origin of graphene's high thermal conductivity.

Lindsay et al. *PRB* **82** (2010) 115427.



Anisotropic thermal conductivity in phosphorene.

Jain and McGaughey, *Sci. Rep.* **5** (2015) 8501.



# Many Others...

- Strain
  - Parrish et al., *PRB* **90** (2014) 235201
  - Mukhopadhyay et al., *PRL* **113** (2014) 025901
- Electrons (next lecture)
  - Liao et al., *PRL* **114** (2015) 115901
  - Jain and McGaughey, *PRB* **93** (2016) 081206(R)
  - Wang et al., *JAP* **119** (2016) 225109.
- BTE Methodology
  - Fugallo et al., *PRB* **88** (2013) 045430
  - Cepellotti and Marzari, *PRX* **109** (2016) 041013

# Advantages and Disadvantages

- Naturally incorporate quantum statistics
  - Integrate with input from first principles calculations
  - Essentially the same framework for all materials
- 
- Thermal conductivity prediction limited by unit cell size,  $N$ 
    - Scales as  $N^4$ ,  $N < 20$
  - Anharmonicity typically to 3<sup>rd</sup> order

# Recommendations

- Some codes freely available, but be careful!
  - GULP, Phonopy, ShengBTE, ...
  - We write our own codes
- Many subtle decisions to make
  - How to enforce energy conservation
  - Convergence/size effects

# Summary of Lecture 4

- Phonon-phonon mean lifetimes from anharmonic lattice dynamics
  - Theoretical and computational challenges
- Force constants from first principles
  - Better agreement with experiments compared to empirical potentials
  - But, there is still uncertainty!