

Lecture 20: Reciprocal Lattice

Ref. Chapter 5.2



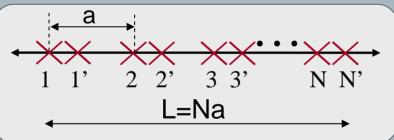
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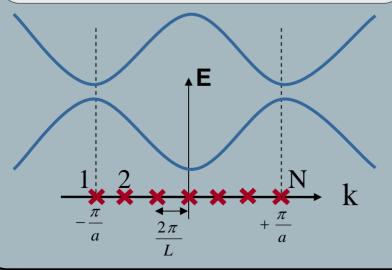
online simulations and more

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Dispersion Relation Quantized k States

 Given any periodic structure we've discussed how to calculate the E-k relationship. For example consider a 1D dimerized solid with two orbitals per unit cell





- The number of allowed values of k can be found as: $\frac{2\pi/a}{2\pi/L} = \frac{L}{a} = N$
- Notice that the spacing between the k states comes from the imposition of periodic boundary conditions. To see this consider:

$$0 \stackrel{1}{\times} 1 \stackrel{2}{\times} 3 \stackrel{\text{n-1 n n+1}}{\times} N$$

$$e^{inka} = e^{i(n+N)ka} \Rightarrow$$

$$e^{iNka} = 1 = e^{i2\pi(\text{integer})} \Rightarrow$$

$$Nka = 2\pi v \quad v : \text{integer} \Rightarrow$$

$$k = \frac{2\pi}{Na} \nu$$

Brillouin Zone

• To see why the k values are bounded between —pi/a and pi/a consider the simple dispersion relation of 1D solid:

$$E = \varepsilon + 2t \cos ka$$
The two k states give us the same wavefunction
$$\frac{1}{-\frac{\pi}{a}} = \frac{2\pi}{L}$$
The two k states give us the same wavefunction

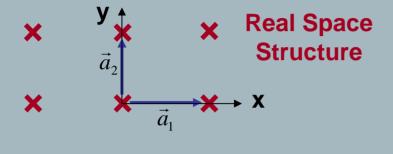
- The point is that if you take any value of k within the range and add 2pi/a to it, you will not get a new independent wavefunction.
- To see this consider the solution...

$$(1)\psi_n = \psi_0 e^{ikna}$$

$$(2) \psi_n = \psi_0 e^{ikna} e^{i2\pi n} = \psi_0 e^{ikna}$$

- You can clearly see that 1 and 2 are the same. This is why we do not need to consider any k values outside of the range −pi/a →2pi/a. The point is that corresponding to any point outside the range there is a point within the range which is an integer multiple of 2pi/a from it. It is the same story for all of them. Add this amount to k and you will get the same answer.
- This symmetric interval around k=0 states that gives us a complete set of k values is called the first Brillouin zone.

Reciprocal Lattice in 2D



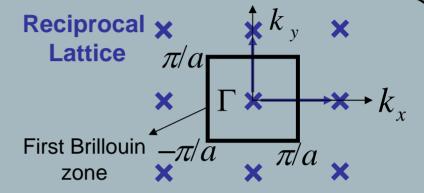
• Any point in the real $\vec{r} = m\hat{x}a + n\hat{y}a$ space can be written as: Where the general solution

is:
$$\psi(\vec{r}) = \psi_0 e^{i\vec{k}\cdot\vec{r}}$$

To construct the reciprocal lattice we need to find a vector K such that:

$$e^{i(\vec{k}+\vec{K})\cdot\vec{r}} = e^{i\vec{k}\cdot\vec{r}} \Longrightarrow$$

the
$$e^{i\vec{K}\cdot\vec{r}}=1\Rightarrow \vec{K}\cdot\vec{r}=(2\pi)\upsilon$$



- Any point in the reciprocal lattice can be written as: $\vec{K} = M\vec{A}_1 + N\vec{A}_2$
- To find K the general procedure is to find the vectors A1 and A2 that satisfy:

$$\vec{a}_1 \cdot \vec{A}_1 = \vec{a}_2 \cdot \vec{A}_2 = 2\pi$$

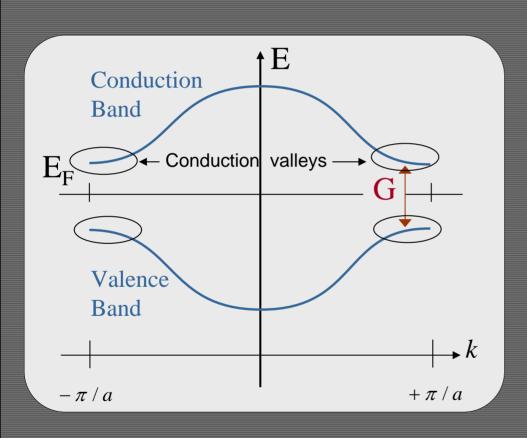
$$\vec{a}_1 \cdot \vec{A}_2 = \vec{a}_2 \cdot \vec{A}_1 = 0$$

$$\Rightarrow \begin{cases} \vec{A}_1 = \hat{x}(2\pi/a) \\ \vec{A}_2 = \hat{y}(2\pi/a) \end{cases}$$

 We can now check to see if we have right answer:

$$\vec{K} \cdot \vec{r} = (m\hat{x}a + n\hat{y}a) \cdot (\hat{x}\frac{2\pi}{a} + \hat{y}\frac{2\pi}{a})$$
$$= 2\pi(Mm + Nn)$$

E-k Diagram and Conduction



• We are studying this course to understand the electrical properties of semiconductors. Looking at this E-k diagram, you can see that the chemical potential lies in the gap between the bands. Ordinary there is no conduction. But if we move the levels relative to Ef via applying a voltage we get conduction.

FCC Lattice to BCC Reciprocal Lattice

