ECE-656: Fall 2009

Lecture 34: Ensemble Effects in Non-Local Transport

Professor Mark Lundstrom
Electrical and Computer Engineering
Purdue University, West Lafayette, IN USA

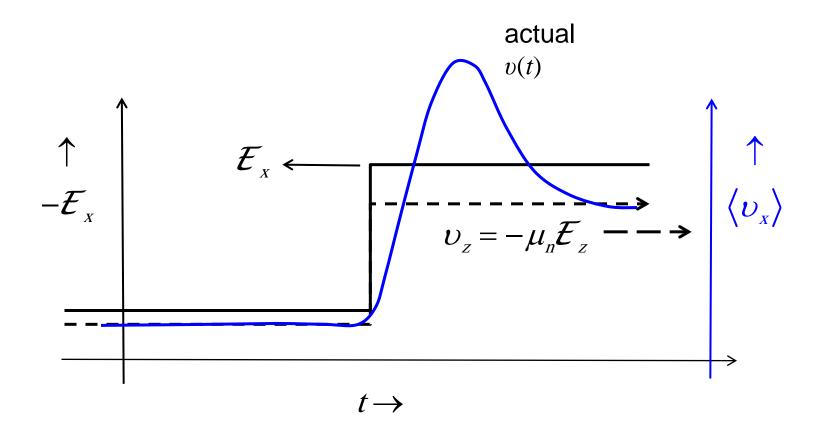




outline

- 1) Review of velocity overshoot
- 2) Steady-state, spatial transients
- 3) Heterojunction launching ramps
- 4) Repeated velocity overshoot
- 5) Questions?

rapidly varying electric fields



Monte Carlo simulation

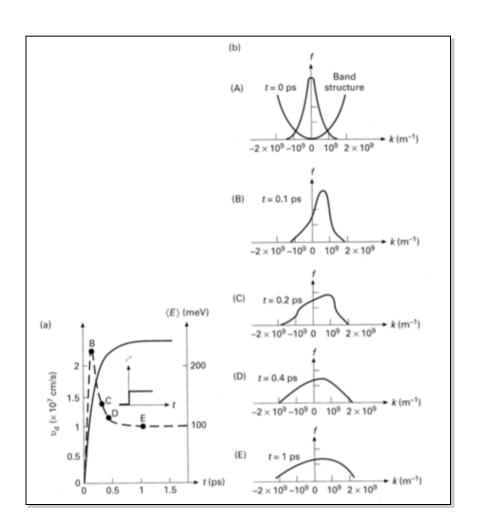
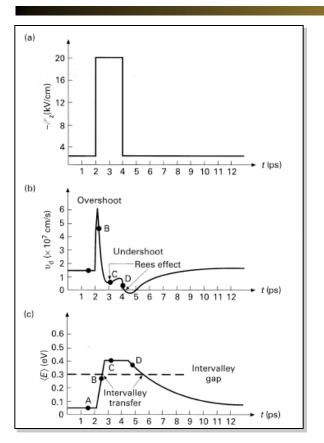


Fig. 8.9 Evolution of the distribution function during a velocity overshoot transient. The average drift velocity and energy are shown in (a), and the evolution of the corresponding distribution function is shown in (b). The results were obtained by Monte Carlo simulation of electron transport in silicon by E. Constant [8.10].

p. 335 of Lundstrom

Rees effect (GaAs)



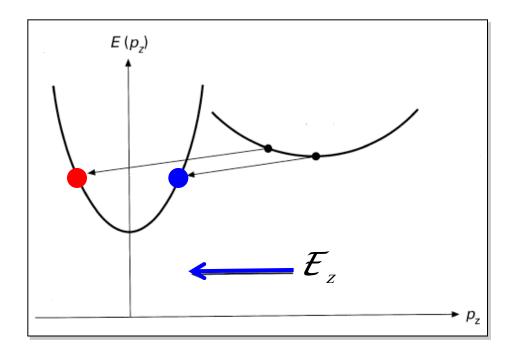


Fig. 8.10 (a) Applied electric field vs. time. (b) Ave. drift velocity vs. time. (c) Ave. electron energy vs. time. (Monte Carlo simulations from E. Constant [8.10].

velocity overshoot

 $p_{dx} = \langle p_x \rangle$ Let's find an equation for the ave. x-directed momentum.

$$\frac{dp_{dx}}{dt} = -q\mathcal{E}_{x} - \frac{p_{dx}}{\langle \tau_{m} \rangle}$$

 $u = \langle E(p) \rangle$ Let's find an equation for the ave. x-directed momentum.

$$\frac{du}{dt} = -qv_{dx}\mathcal{E}_{x} - \frac{\left(u - u_{0}\right)}{\left\langle\tau_{E}\right\rangle}$$

$$\upsilon_{dx}(t) = -\mu_n \mathcal{E}_x \left(1 - e^{-t/\langle \tau_m \rangle} \right)$$

$$u(t) = u_0 + q \langle \tau_E \rangle \mu_n \mathcal{E}_x^2 \left(1 - e^{-t/\langle \tau_E \rangle} \right)$$

But, we have ignored diffusion (an ensemble effect).

momentum and energy balance

$$\frac{dp_{dx}}{dt} = -q\mathcal{E}_{x} - \frac{p_{dx}}{\langle \tau_{m} \rangle}$$

$$\frac{du}{dt} = -qv_{dx}\mathcal{E}_{x} - \frac{\left(u - u_{0}\right)}{\left\langle\tau_{E}\right\rangle}$$

$$\frac{dP_{x}}{dt} = -\frac{d(2W_{xx})}{dx} - qn\mathcal{E}_{x} - \frac{P_{x}}{\langle \tau_{m} \rangle}$$

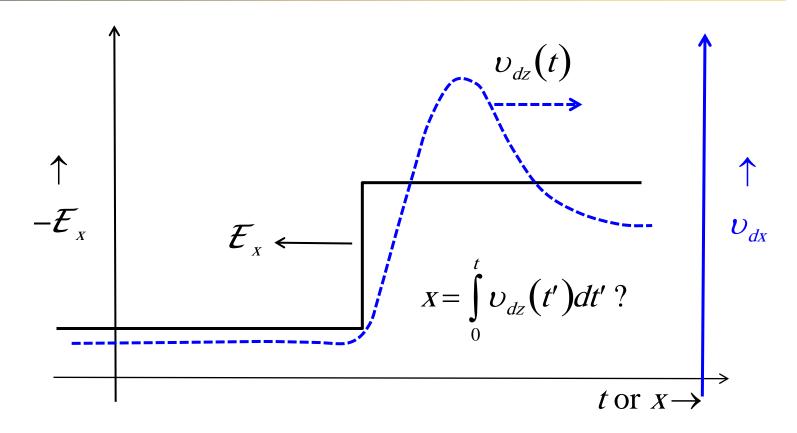
$$P_{X} = n \langle p_{X} \rangle = n p_{dX}$$

$$\frac{dW}{dt} = -\frac{d(F_W)}{dx} + J_X \mathcal{E}_X - \frac{W - W_0}{\langle \tau_E \rangle}$$

$$W = n \langle E(p) \rangle = nu$$

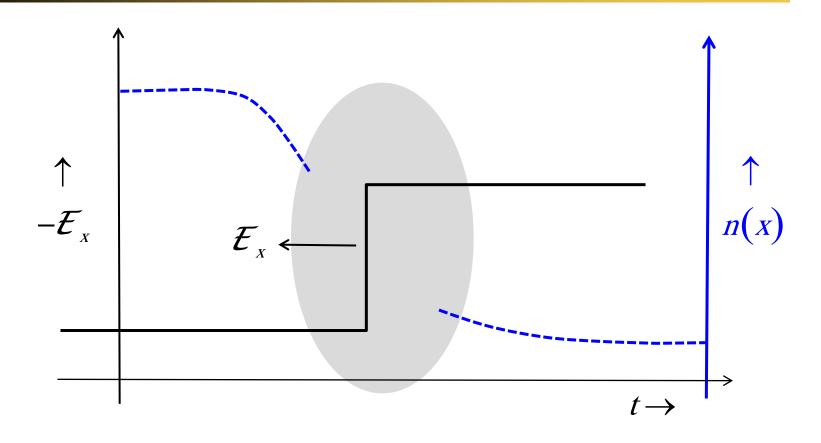
What effect do these spatial gradients have?

temporal vs. spatial transients



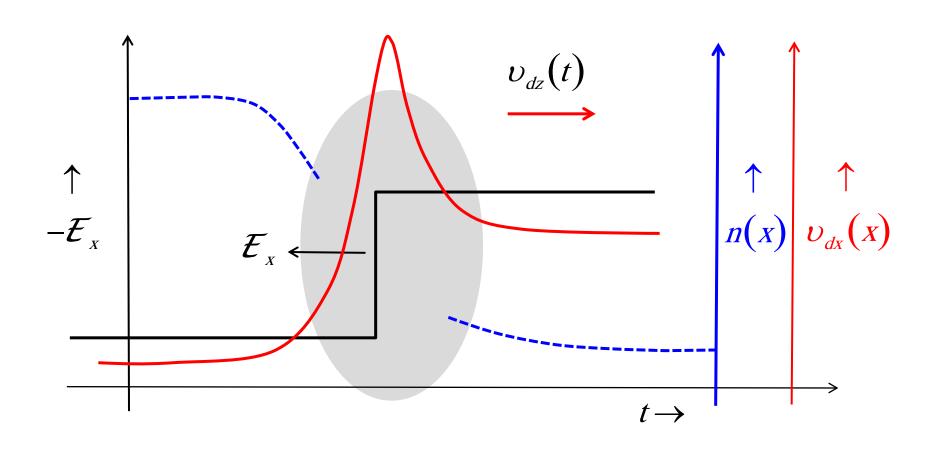
Question: If we change the horizontal axis to distance, what does the steady-state velocity vs. position characteristic look like?

carrier density vs. position



Steady-state current is constant: $J_{nx} = n(x)q\langle v_x(x)\rangle$

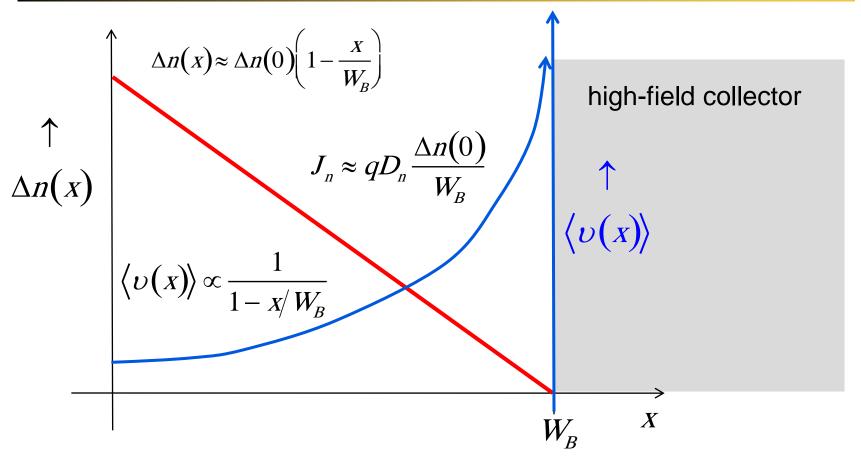
carrier velocity vs. position



outline

- 1) Review of velocity overshoot
- 2) Steady-state, spatial transients
- 3) Heterojunction launching ramps
- 4) Repeated velocity overshoot
- 5) Questions?

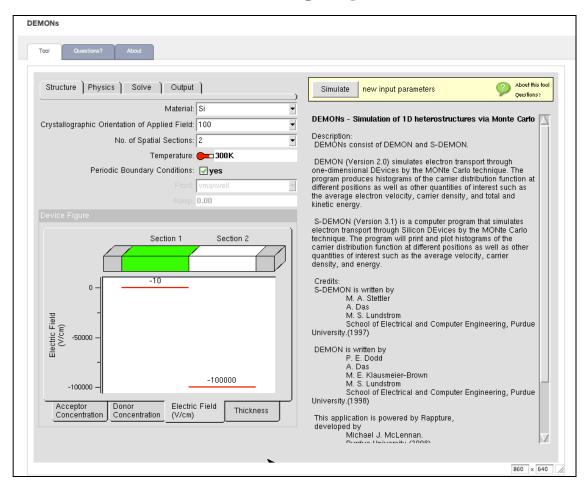
a familiar example



$$J_{nx} = n(x)q\langle \upsilon_{x}(x)\rangle$$

example Monte Carlo simulations

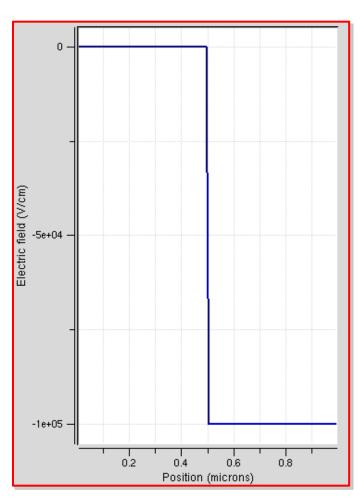
DEMONs



1D, steady-state Monte Carlo simulation for Si and GaAs

Piecewise constant electric field profiles.

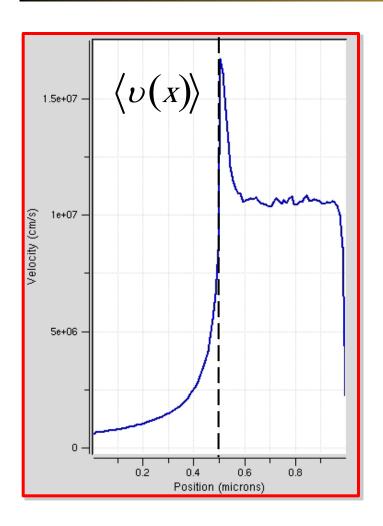
low-high field structure

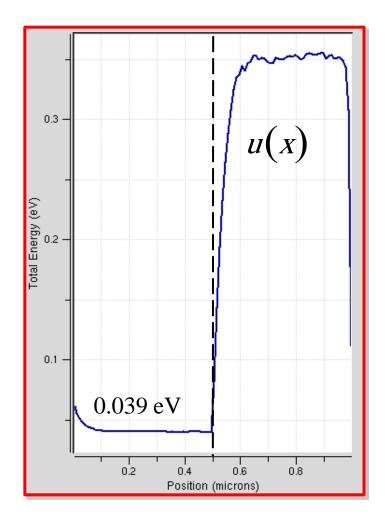


1E16 n(X)Electron Concentration (/cm^3) 1E15 $J_{nx} = n(x)q\langle \upsilon_{x}(x)\rangle$ 1E14 0.2 0.8 0.4 0.6 Position (microns)

periodic boundary conditions

low-high field structure



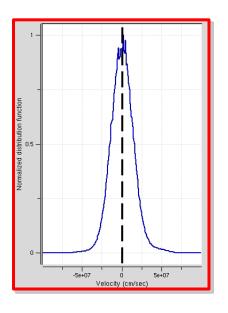


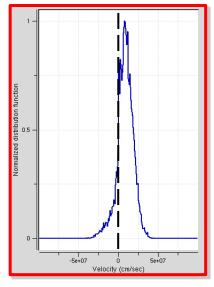
periodic boundary conditions

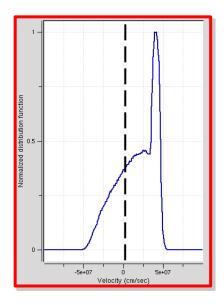
velocity histograms

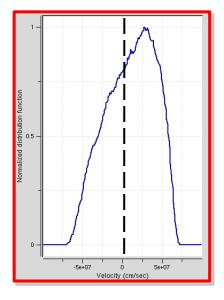
$$\mathcal{E} = -10^5 \text{ V/cm}$$

$$\mathcal{E} = -10 \text{ V/cm}$$





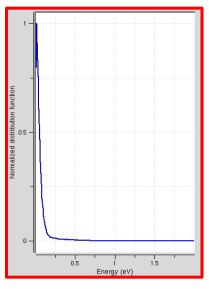


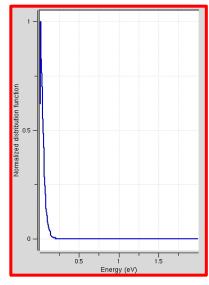


energy histograms

$$n(E) = f(E)D(E)$$

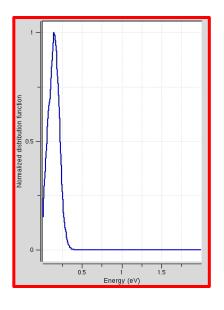
$$\mathcal{E} = -10 \text{ V/cm}$$

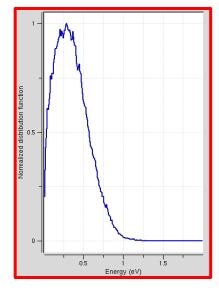




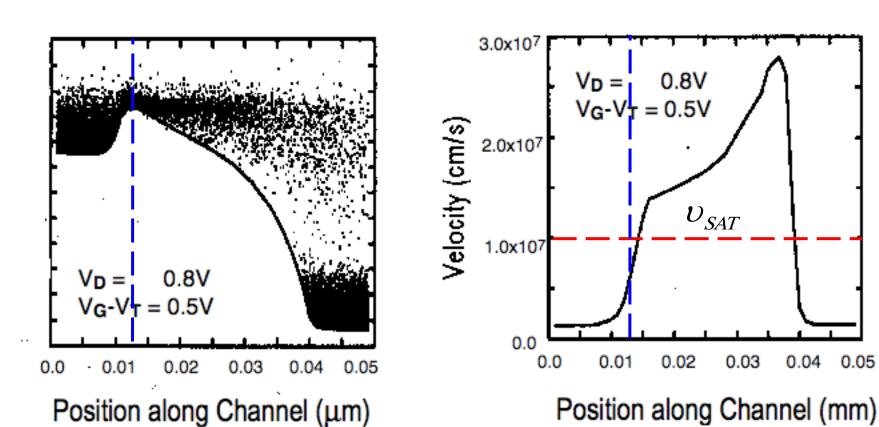
$$n(E) \propto e^{-E/k_BT} \sqrt{E}$$

$$\mathcal{E} = -10^5 \text{ V/cm}$$





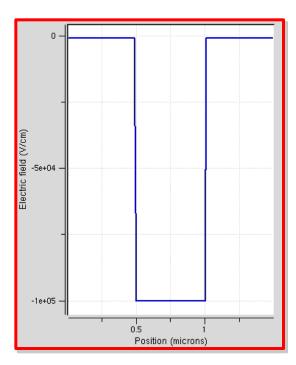
off-equilibrium nanoscale MOSFETs



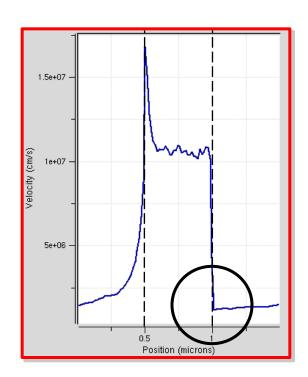
Frank, Laux, and Fischetti, IEDM Tech. Dig., p. 553, 1992

0.05

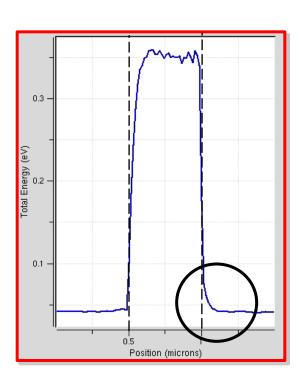
low-high-low field structure



periodic boundary conditions



velocity undershoot



 $u(x) > u_0$

temporal vs. spatial transients

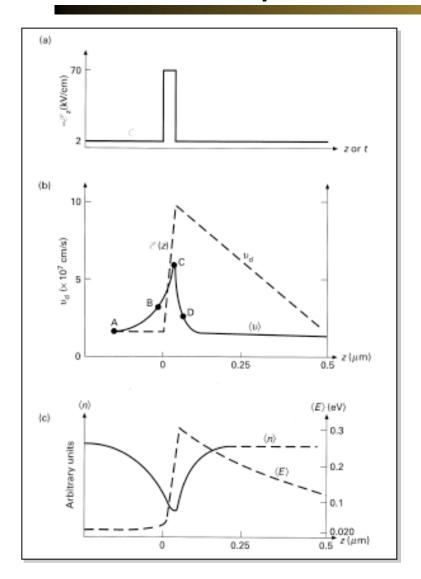


Fig. 8.13 (a) Applied electric field in time and space. (b) Average velocity versus position for a pulse applied in space (solid line) and time (dashed line). (c) Steady-state carrier density (solid line() and energy (dashed line.) The results were obtained by Monte Carlo simulation of electron transport in GaAs by E. Constant [8.10].

p. 340 of Lundstrom

$$z = \int_{0}^{t} \upsilon(t') dt'$$

outline

- 1) Review of velocity overshoot
- 2) Steady-state, spatial transients
- 3) Heterojunction launching ramps
- 4) Repeated velocity overshoot?
- 5) Questions?

ballistic launching ramps in HBTs

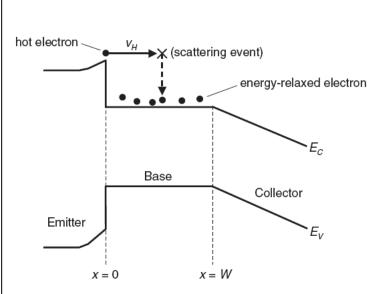


Fig. 1. Modeling of electron transport in the base of InP/InGaAs HBTs. A hot electron having a constant velocity of $v_{\rm H}=1\times10^8$ cm/s is assumed to be thermalized to an energy-relaxed electron after one scattering event. The energy-relaxed electron has a diffusivity $D_{\rm T}$ and a recombination life time $\tau_{\rm T}$.

Hiroki Nakajima, *Jap. J. of Appl. Phys.*, **46**, pp. 485–490, 2007.

Question: How do we experimentally tell whether base transport is ballistic or diffusive?

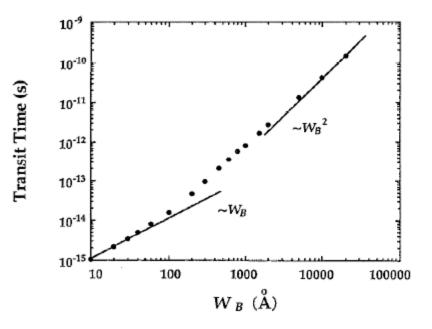
Answer: Look at the base current.

$$I_{B} \propto \frac{t_{t}}{\tau_{n}}$$

ballistic:
$$t_t = W_B/v_{ball}$$

diffusive:
$$t_t = W_B^2 / 2D_n$$

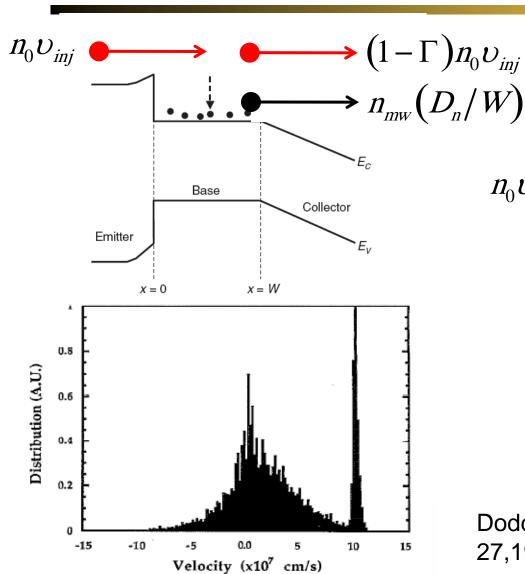
base transit time scaling



~80% traverse a 300A base ballisticaly or quasi-ballisticaly

P.E. Dodd and M.S. Lundstrom, "Minority electron transport in InP/InGaAs heterojunction bipolar transostors," *Appl. Phys. Lett.*, **61**, 27,1992

base transit time scaling



$$n_0 \upsilon_{inj} = (1 - \Gamma) n_0 \upsilon_{inj} + n_{mw} (D_n / W)$$

$$\frac{n_0}{n_{mw}} = \Gamma \frac{v_{inj}}{\left(D_n/W\right)}$$

Dodd and Lundstrom, *APL*, **61**, 27,1992

outline

- 1) Review of velocity overshoot
- 2) Steady-state, spatial transients
- 3) Heterojunction launching ramps
- 4) Repeated velocity overshoot?
- 5) Questions?