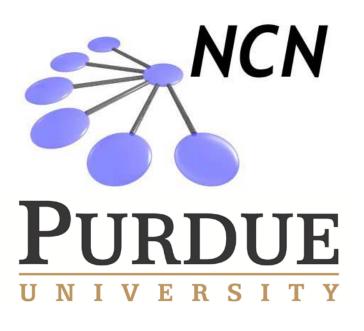


Network for Computational Nanotechnology (NCN)

UC Berkeley, Univ. of Illinois, Norfolk State, Northwestern, Purdue, UTEP

Open 1D Systems: Transmission through & over 1 Barrier



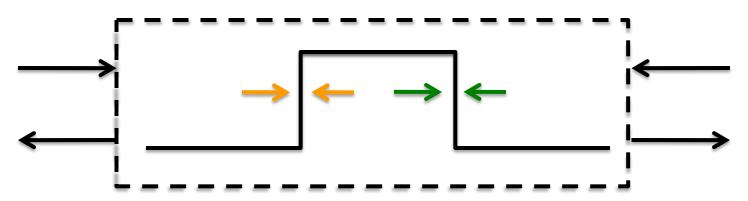
Gerhard Klimeck, Dragica Vasileska, Smartha Agarwal



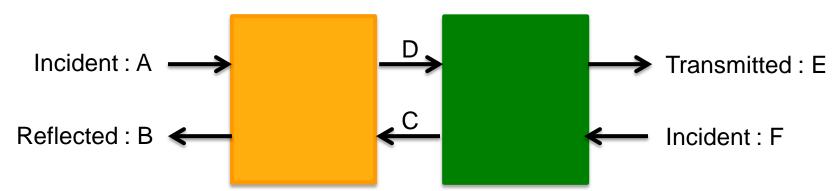


Scattering Matrix approach

Define our system : Single barrier



One matrix each for each interface: 2 S-matrices



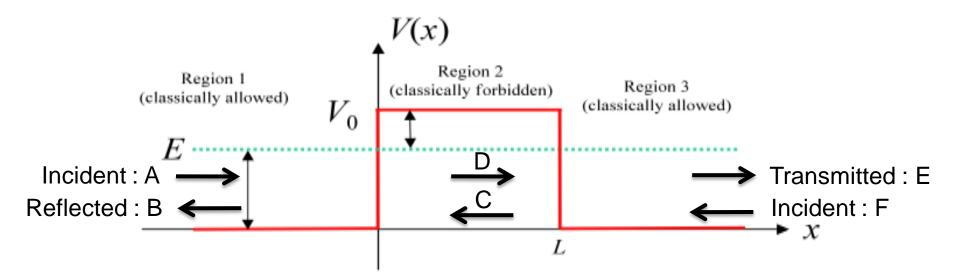
No particles lost! Typically A=1 and F=0.







Tunneling through a single barrier



Wave-function each region,

$$\psi_{1}(x) = Ae^{ikx} + Be^{-ikx}$$

$$\psi_{2}(x) = Ce^{-\gamma_{x}} + De^{\gamma_{x}}$$

$$\psi_{3}(x) = Ee^{ikx} + Fe^{-ikx}$$

$$k = \sqrt{\frac{2mE}{\hbar^2}} \qquad \gamma = \sqrt{\frac{2m(V_0 - E)}{\hbar^2}}$$







Single barrier case

Applying boundary conditions at each interface (x=0 and x=L) gives,

$$\psi_{1}(0) = \psi_{2}(0) \rightarrow A + B = C + D$$

$$\psi_{1}(0) = \psi_{2}(0) \rightarrow ik(A - B) = -\gamma(C - D)$$

$$\psi_{2}(L) = \psi_{3}(L) \rightarrow Ce^{-\gamma L} + De^{\gamma L} = Ee^{ikL} + Fe^{-ikL}$$

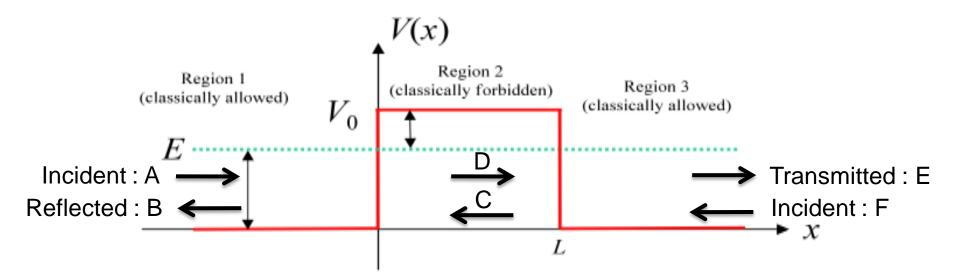
$$\psi_{2}(L) = \psi_{3}(L) \rightarrow -\gamma \left(e^{-\gamma L} - De^{\gamma L} \right) = ik \left(Ee^{ikL} - Fe^{-ikL} \right)$$







Tunneling through a single barrier



Wave-function each region,

$$\psi_{1}(x) = Ae^{ikx} + Be^{-ikx}$$

$$\psi_{2}(x) = Ce^{-\gamma_{x}} + De^{\gamma_{x}}$$

$$\psi_{3}(x) = Ee^{ikx} + Fe^{-ikx}$$

$$k = \sqrt{\frac{2mE}{\hbar^2}} \qquad \gamma = \sqrt{\frac{2m(V_0 - E)}{\hbar^2}}$$







Single barrier case

Applying boundary conditions at each interface (x=0 and x=L) gives,

$$\begin{split} &\psi_1(0) = \psi_2(0) \quad \Rightarrow \quad A + B = C + D \\ &\psi_1(0) = \quad \psi_2(0) \quad \Rightarrow ik(A - B) = -\gamma(C - D) \\ &\psi_2(L) = \psi_3(L) \quad \Rightarrow \quad Ce^{-\gamma L} + De^{\gamma L} = Ee^{ikL} + Fe^{-ikL} \\ &\psi_2(L) = \psi_3(L) \quad \Rightarrow \quad -\gamma \left(Ce^{-\gamma L} - De^{\gamma L} \right) = ik \left(Ee^{ikL} - Fe^{-ikL} \right) \end{split}$$

Which in matrix can be written as,

$$\begin{bmatrix} A \\ B \end{bmatrix} = \begin{bmatrix} \frac{1}{2} \left(1 + i \frac{\gamma}{k} \right) & \frac{1}{2} \left(1 - i \frac{\gamma}{k} \right) \\ \frac{1}{2} \left(1 - i \frac{\gamma}{k} \right) & \frac{1}{2} \left(1 + i \frac{\gamma}{k} \right) \end{bmatrix} \begin{bmatrix} C \\ D \end{bmatrix} = M_1 \begin{bmatrix} C \\ D \end{bmatrix}$$

$$\begin{bmatrix} C \\ D \end{bmatrix} = \begin{bmatrix} \frac{1}{2} \left(1 - i \frac{k}{\gamma} \right) e^{(ik+\gamma)L} & \frac{1}{2} \left(1 + i \frac{k}{\gamma} \right) e^{-(ik-\gamma)L} \\ \frac{1}{2} \left(1 + i \frac{k}{\gamma} \right) e^{(ik-\gamma)L} & \frac{1}{2} \left(1 - i \frac{k}{\gamma} \right) e^{-(ik+\gamma)L} \end{bmatrix} \begin{bmatrix} E \\ F \end{bmatrix} = M_2 \begin{bmatrix} E \\ F \end{bmatrix}$$







Single barrier case

Transmission can be found using the relations between unknown constants,

$$T(E) = \left| \frac{E}{A} \right|^2 = \frac{1}{|m_{11}|^2} \qquad \begin{bmatrix} A \\ B \end{bmatrix} = M_1 \begin{bmatrix} C \\ D \end{bmatrix} = M_1 M_2 \begin{bmatrix} E \\ F \end{bmatrix} = M \begin{bmatrix} E \\ F \end{bmatrix}$$

Case: E<V_o

Case(
$$\gamma L$$
 large): Strong barrier

$$T(E) = \left[1 + \left(\frac{\gamma^2 + k^2}{2k\gamma}\right)^2 sh^2(\gamma L)\right]^{-1} \quad T(E) \approx \left(\frac{4k\gamma}{k^2 + \gamma^2}\right)^2 \exp(-2\gamma L)$$

Case($\gamma L <<1$): Weak barrier

$$T(E) \approx \frac{1}{1 + (kL/2)^2}$$

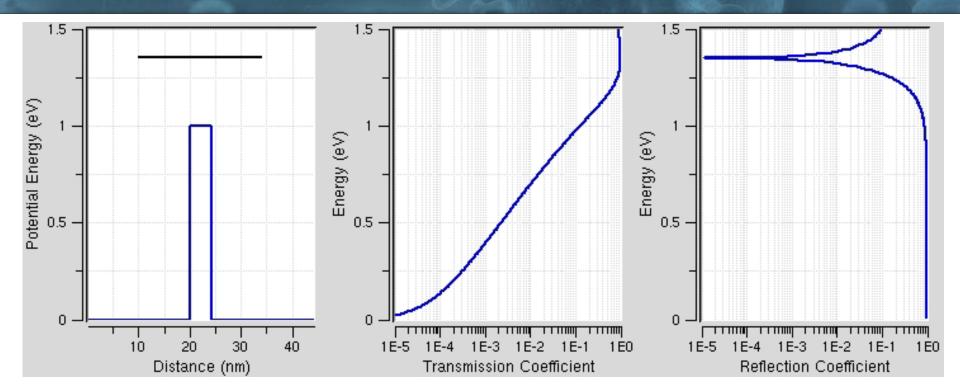
$$T(E) = \left[1 + \left(\frac{k^2 - k_2^2}{2kk_2}\right)^2 \sin^2(k_2 L)\right]^{-1}$$







Single barrier : Concepts



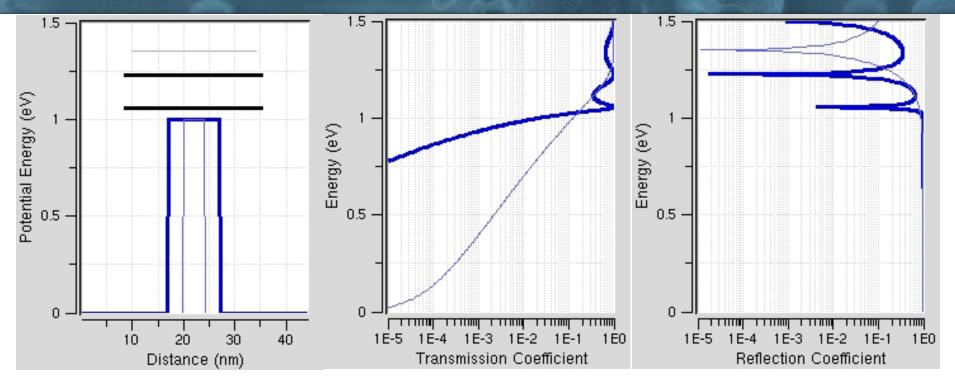
- Transmission is finite under the barrier tunneling!
- Transmission above the barrier is not perfect unity!
- •Quasi-bound state above the barrier. Transmission goes to one.







Effect of barrier thickness below the barrier

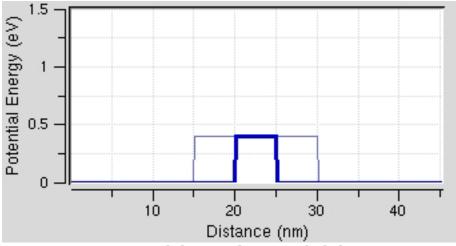


- Increased barrier width reduces tunneling probability
- Thicker barrier increase the reflection probability below the barrier height.
- Quasi-bound states occur for the thicker barrier too.



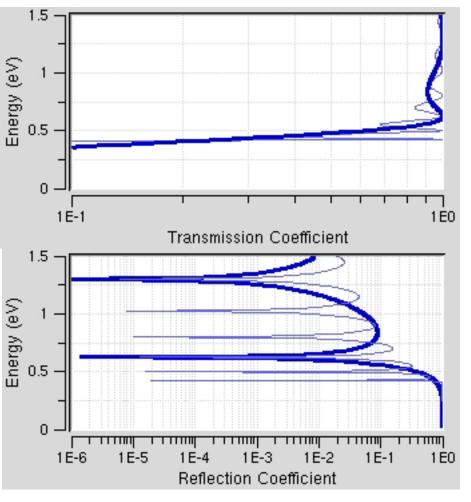


Effect of barrier thickness above the barrier



 Increased barrier width increases oscillation frequency in transmission and reflection.

 Quasi-bound states above the barrier due to 2 reflections.

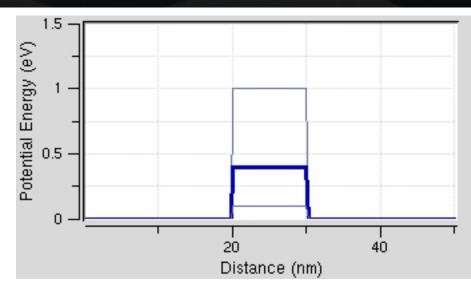




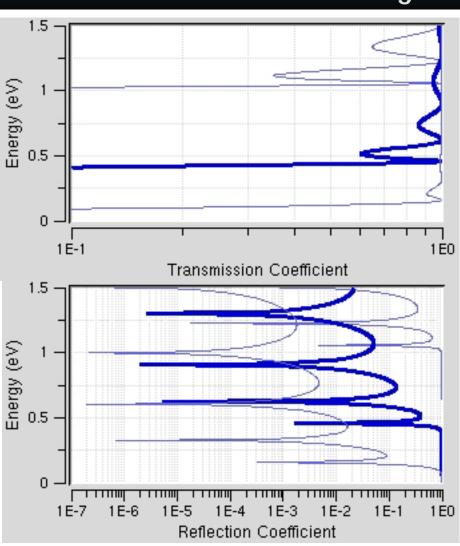




Effect of barrier height



- Increasing the barrier height does not have a significant effect on the modulation frequency above the barrier height.
- Oscillations are strongly related to barrier width but not height!



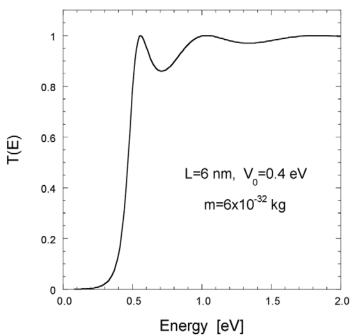




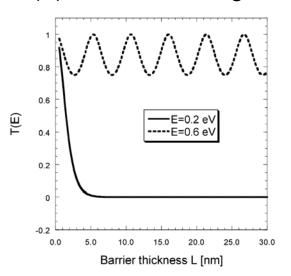


Single barrier case: Discussion

T(E) for a given geometry



T(E) for different lengths



• E<V₀: Classical Physics: T(E)=0,

Quantum Physics: a hyperbolic increase.

•E>V₀: Classical Physics: T(E)=1,

Quantum Physics: total transmission at discrete energies only.

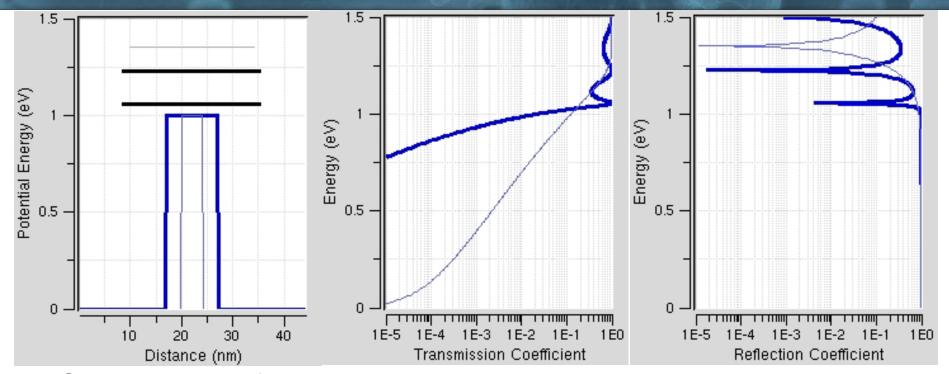
Only barriers of certain width will transmit all particles at a given energies.







Single Barrier - Key Summary



- Quantum wavefunctions can tunnel through barriers
- Tunneling is reduced with increasing barrier height and width
- Transmission above the barrier is not unity
 - » 2 interfaces cause constructive and destructive interference
 - » Quasi bound states are formed that result in unity transmission



