

Modeling Transport in Nanostructured Optoelectronics

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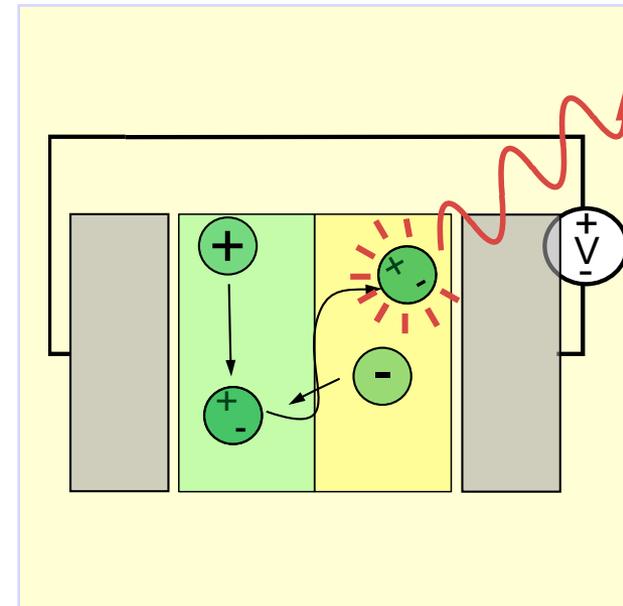
³*Dept. of Materials Science and Engineering*

MOTIVATION:

- Nanostructured optoelectronics: future potential for low-cost, large area, tunable LEDs and solar cells
- Connect macroscopic behavior to microscopic model
- Suggest methods to improve efficiency
- Device physics educational tool

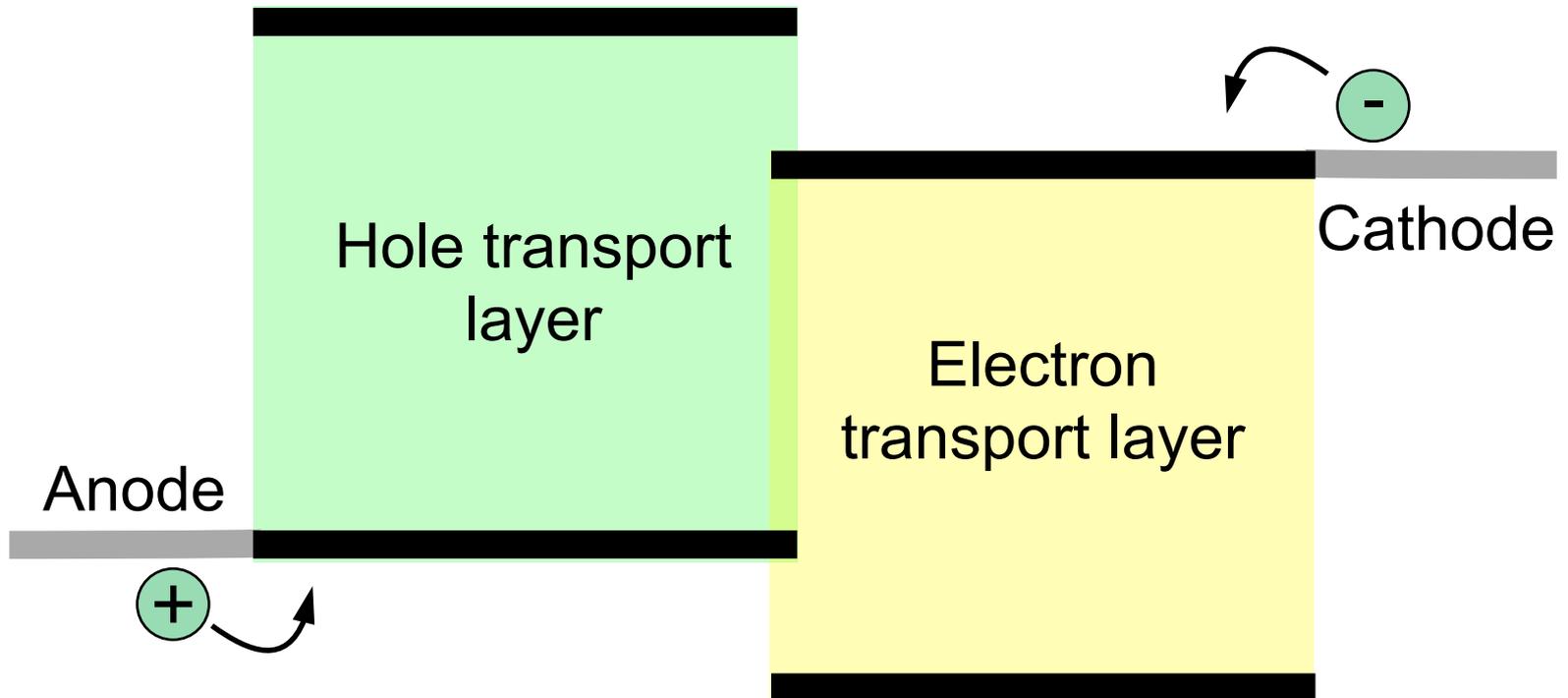
PRESENTATION GOALS:

- Introduce physics in model
- Show reasonable input parameters
- Provide example results



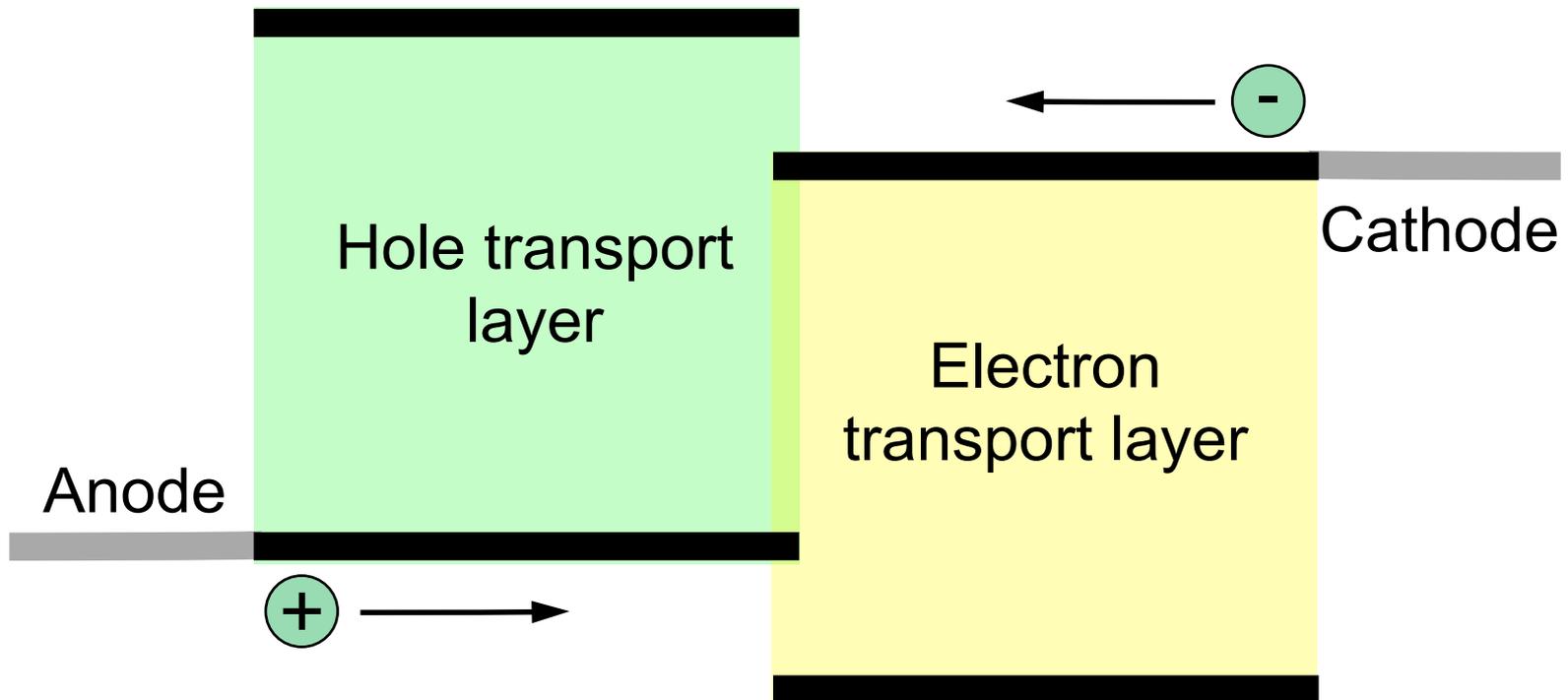
METHOD - MICROSCOPIC PROCESSES:

- Relevant processes
 - *Charge injection*: thermionic injection into broadened density of states
 - *Charge transport*: thermally-activated hopping (Miller-Abrahams or Marcus)
 - *Exciton formation*: photon absorption and charge recombination (Langevin)
 - *Exciton transport*: thermally-activated hopping or Förster energy transfer
 - *Exciton decay*: radiative, non-radiative, and dissociation into charge pairs



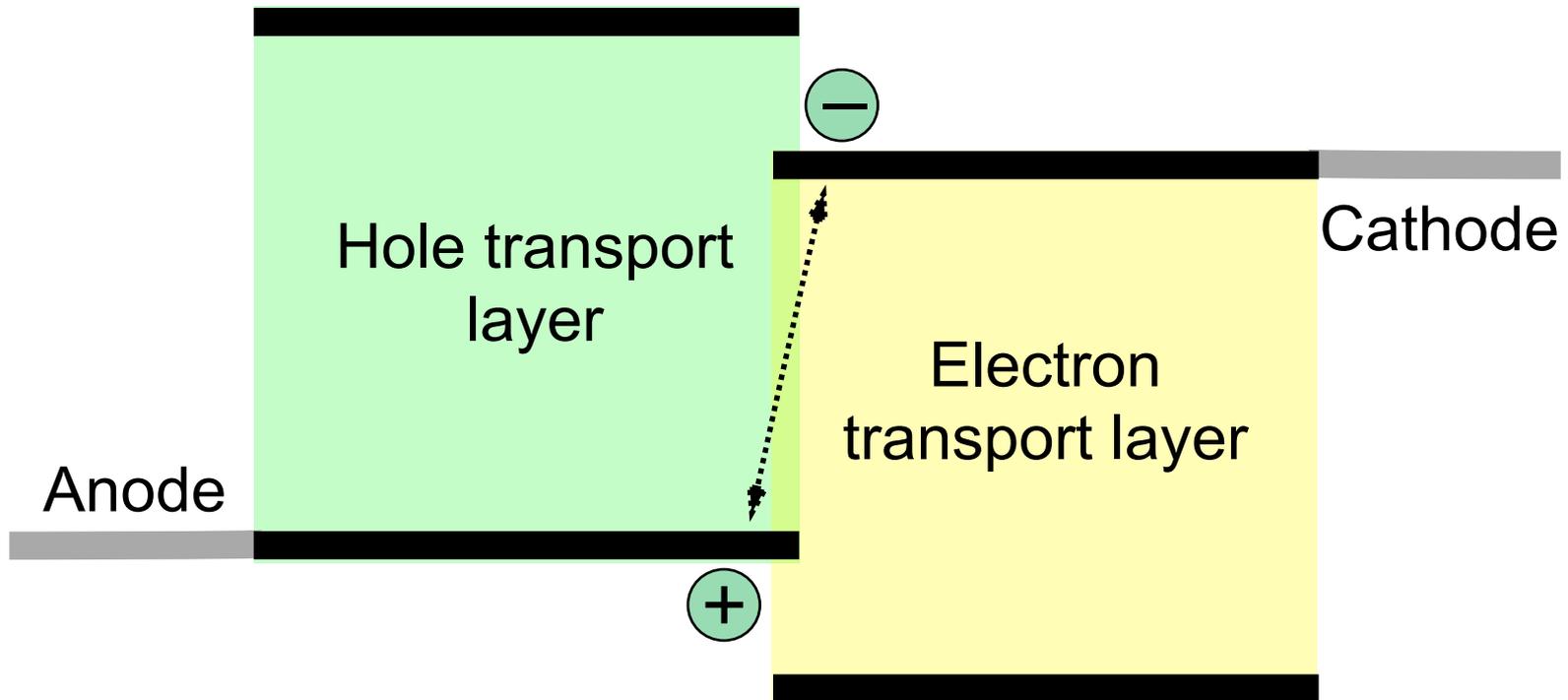
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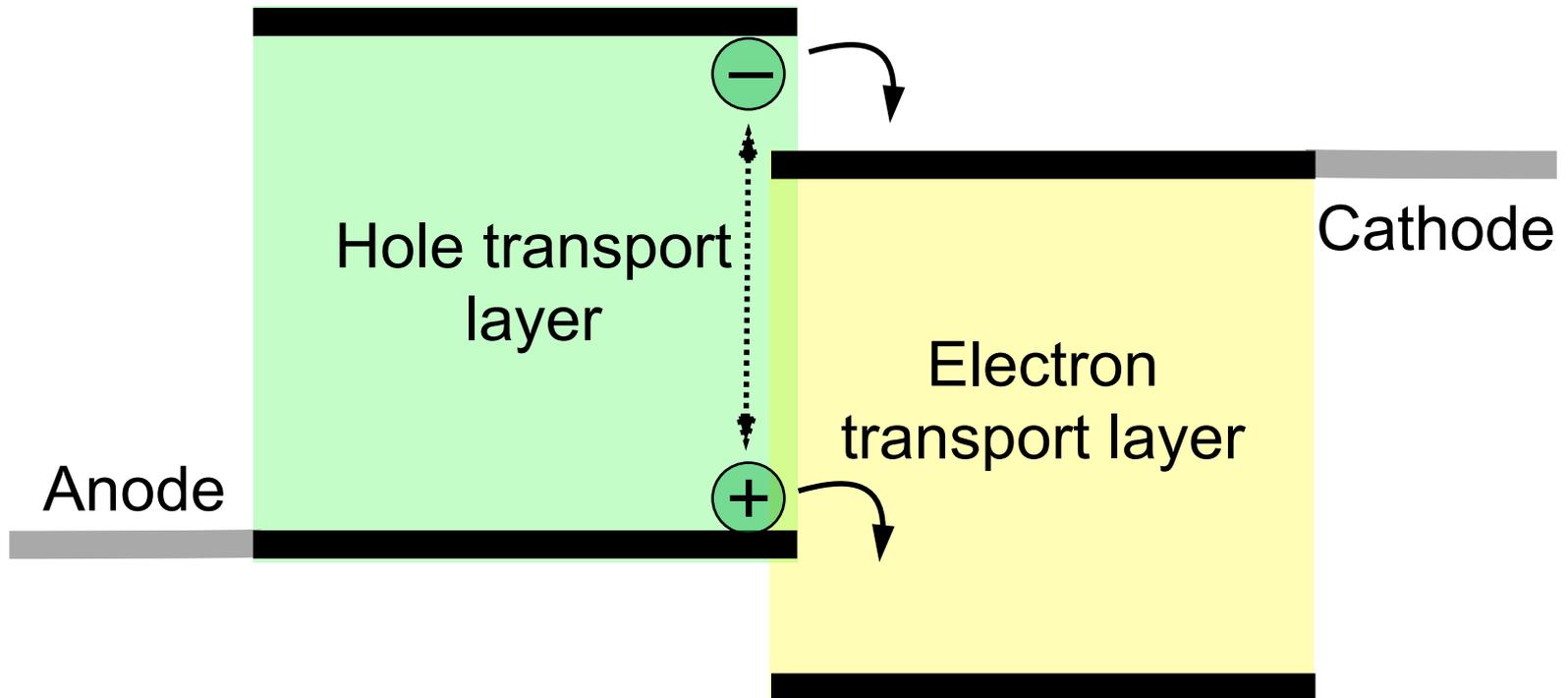
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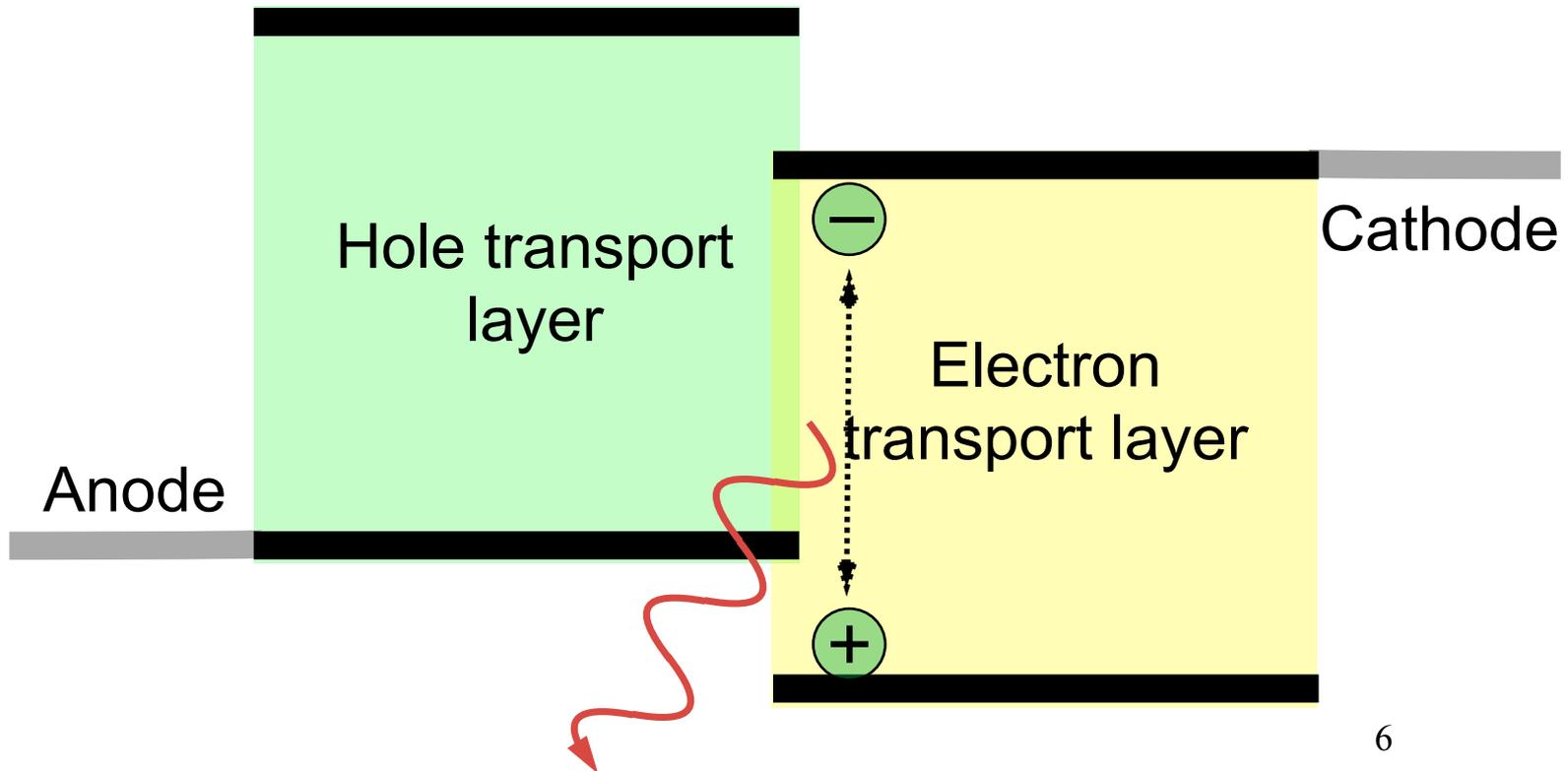
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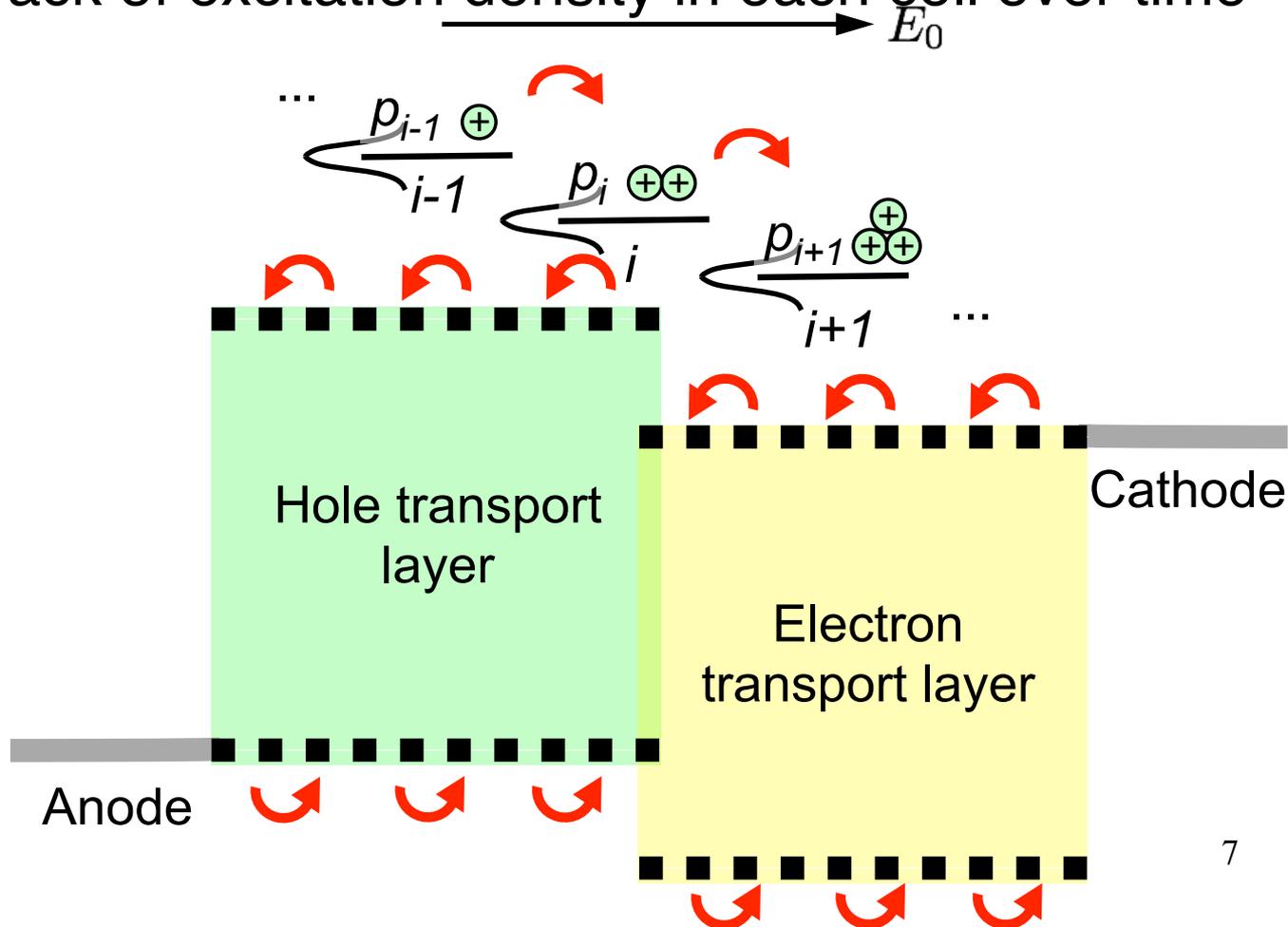
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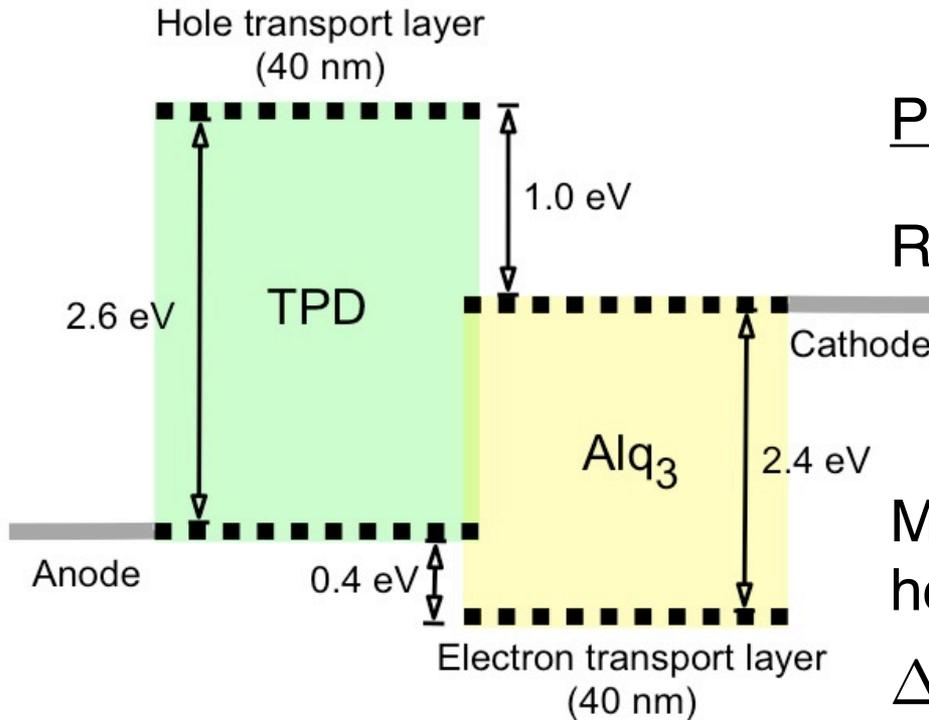
METHOD - NUMERICAL SIMULATION:

- Most transport occurs along one-dimension
- Excitations move between discrete unit cells
- Employ rate equations and add applicable rate processes
- Keep track of excitation density in each cell over time



EXAMPLE 1: BILAYER OLED

40 nm TPD/40 nm Alq₃



Processes enabled by default:

Radiative exciton decay:

$$K_x = \frac{1}{\tau_d}$$

Miller-Abrahams thermal hopping* for holes, excitons, and electrons**:

ΔE due to change in HOMO/LUMO/exciton energies, electrostatic, and polarization potentials

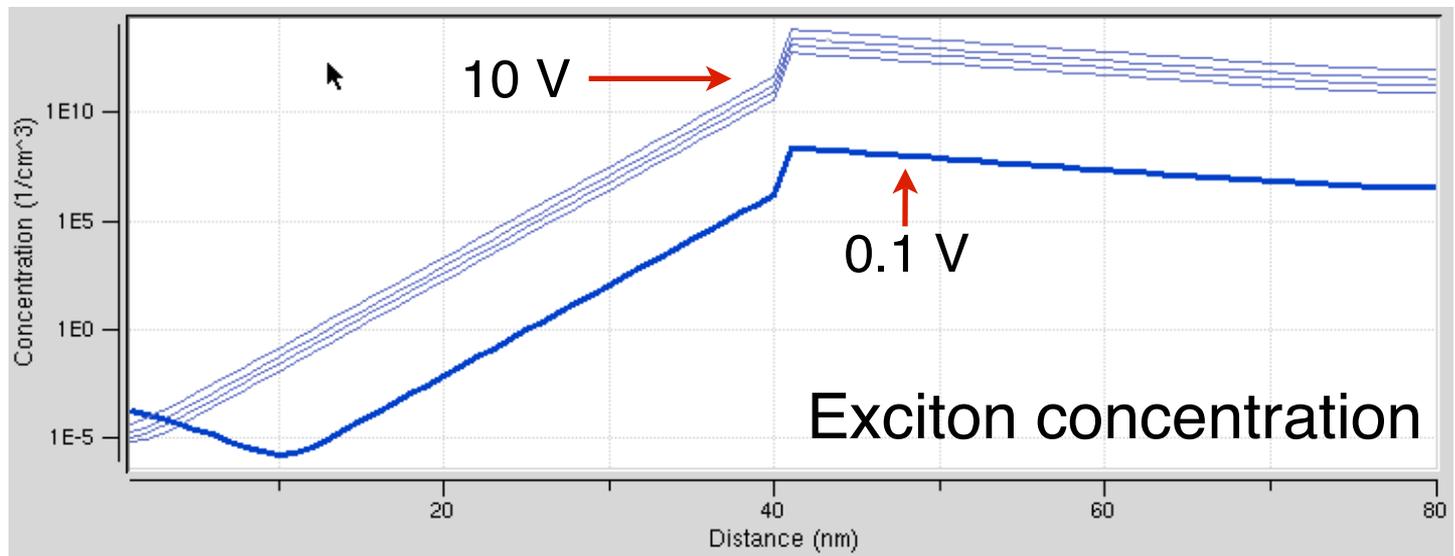
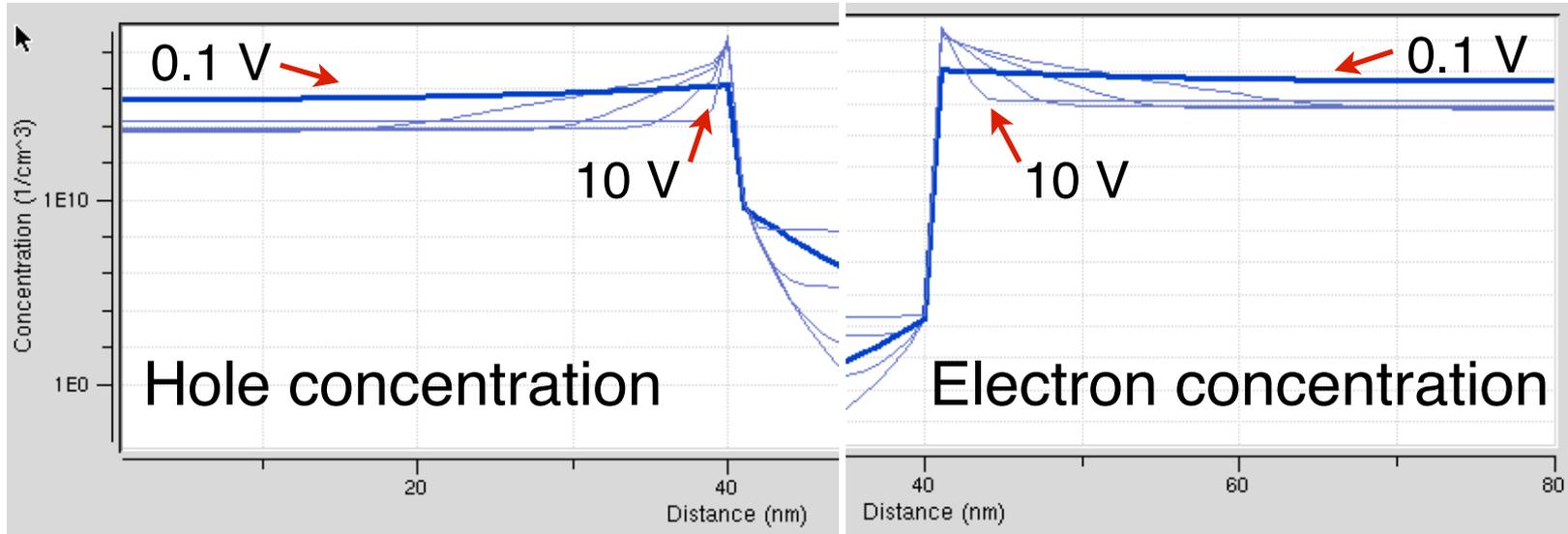
$$K = \mu kT / ed_{\text{mol}}^2 \cdot \begin{cases} 1 & : \Delta E \leq 0 \\ \exp(-\Delta E / kT) & : \Delta E > 0 \end{cases}$$

*reduces to drift-diffusion equations in limit of low field

**between neighboring molecules/nanoparticles in a delta function density of states

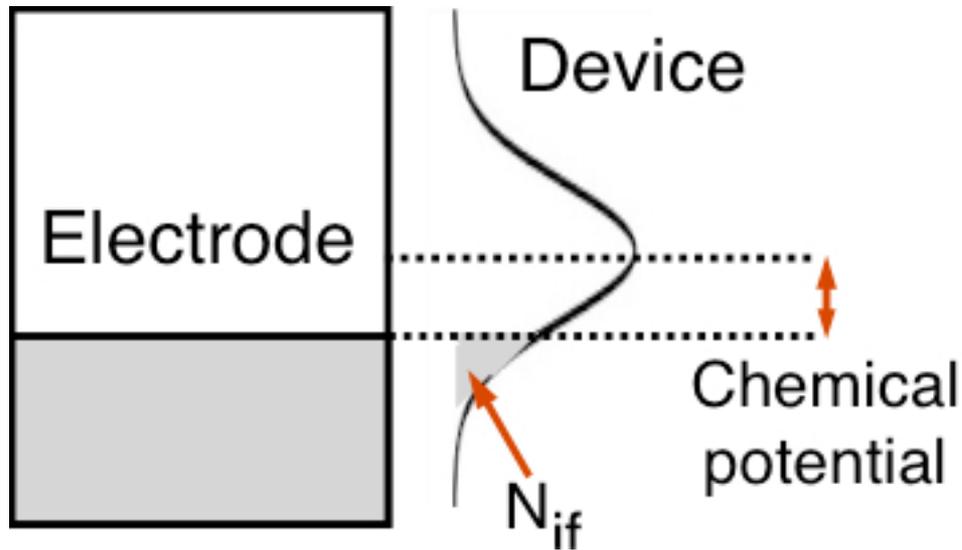
EFFECT OF BIAS

Results: carriers pile up at interface
more excitons created



INJECTION

How charges enter and exit the device



Chemical potential:	-0.38eV
Anode model:	Ohmic
Cathode model:	Marcus transfer
Sigma A:	0.48eV
Sigma B:	0.12eV
Lambda:	0.1eV
Built-in potential:	<input type="checkbox"/> no

Assume Gaussian density of states of width 0.1 eV in first device layer:

$$N_{if} = N_{dos} \operatorname{erf}(\mu/\sigma)$$

Model #1: Ohmic injection

$$J = F \cdot \mu kT / ed_{mol} \cdot \begin{cases} N_{if} & : F > 0 \\ N_0 & : F \leq 0 \end{cases}$$

Where F is the electric field and N_0 is the concentration in the first device layer

INJECTION

How charges enter and exit the device

Model #2: Thermionic injection

Injection over a thermal barrier of height $\Delta E = \frac{4\pi e^2}{\epsilon} N_{if} d_{mol}^2$
with field-dependent recombination current.

See Scott and Malliaras, CPL 299 (1999) pp 115-119

Model #3: Marcus theory

- Move from first to second device layer is rate-limiting step
- Assume DOS broader at interface (A) than in first layer (B)
- Assumes degenerate carrier distribution and approximates integral for *net* transfer rate*
- Includes recombination current
- Current depends on field and chemical potential
- Broadening introduces a built-in potential

*Limketai and Baldo, PRB 71 (2005) 085207 is the starting point for this model

LANGEVIN RECOMBINATION

hole + electron \Rightarrow exciton

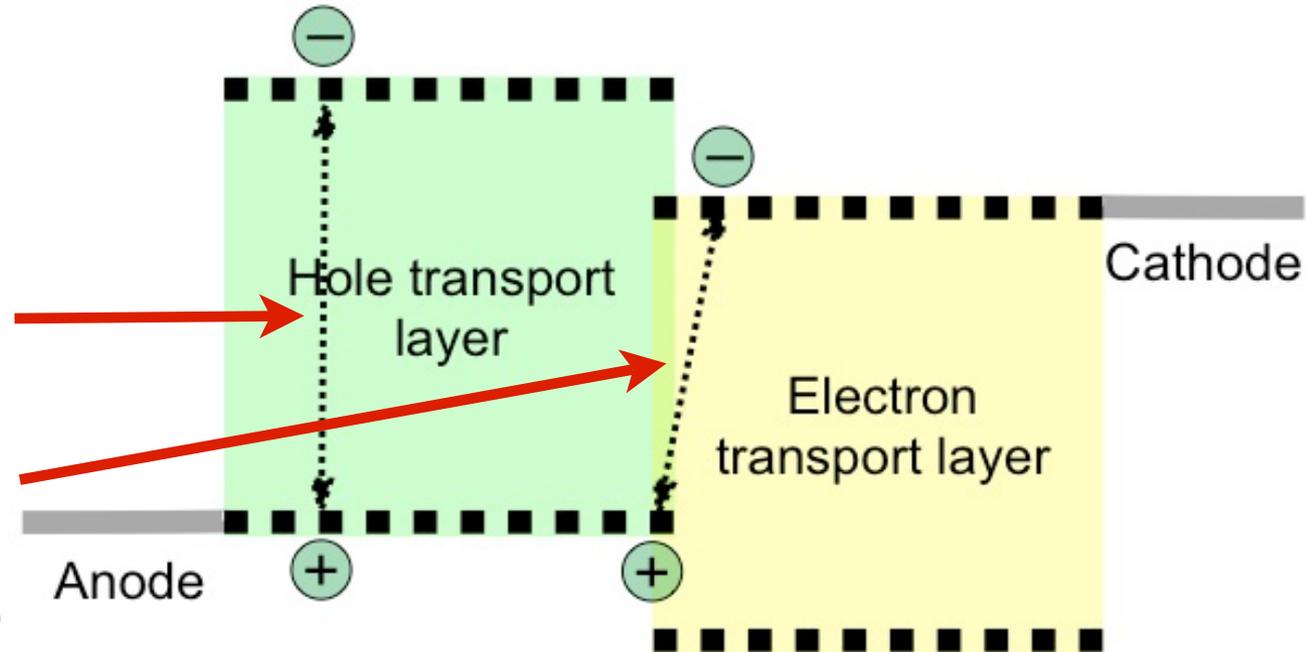
$$R = \gamma np$$

Two types:

Monomolecular
(from mobility)

$$\gamma = \frac{e}{\epsilon} (\mu_p + \mu_n)$$

Bimolecular
(user-specified [cm³/s])



1 Device → 2 Rates → 3 Simulation → 4 Simulate

Chemical potential: **-0.38eV**

Anode model: Ohmic

Cathode model: Ohmic

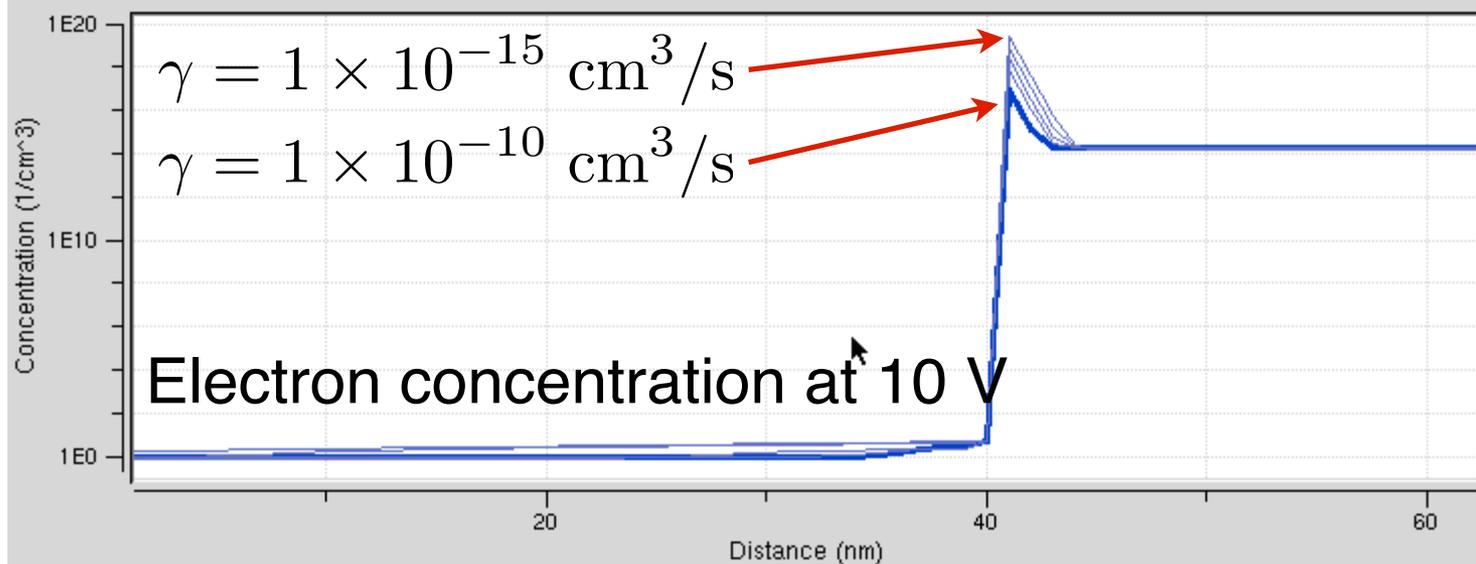
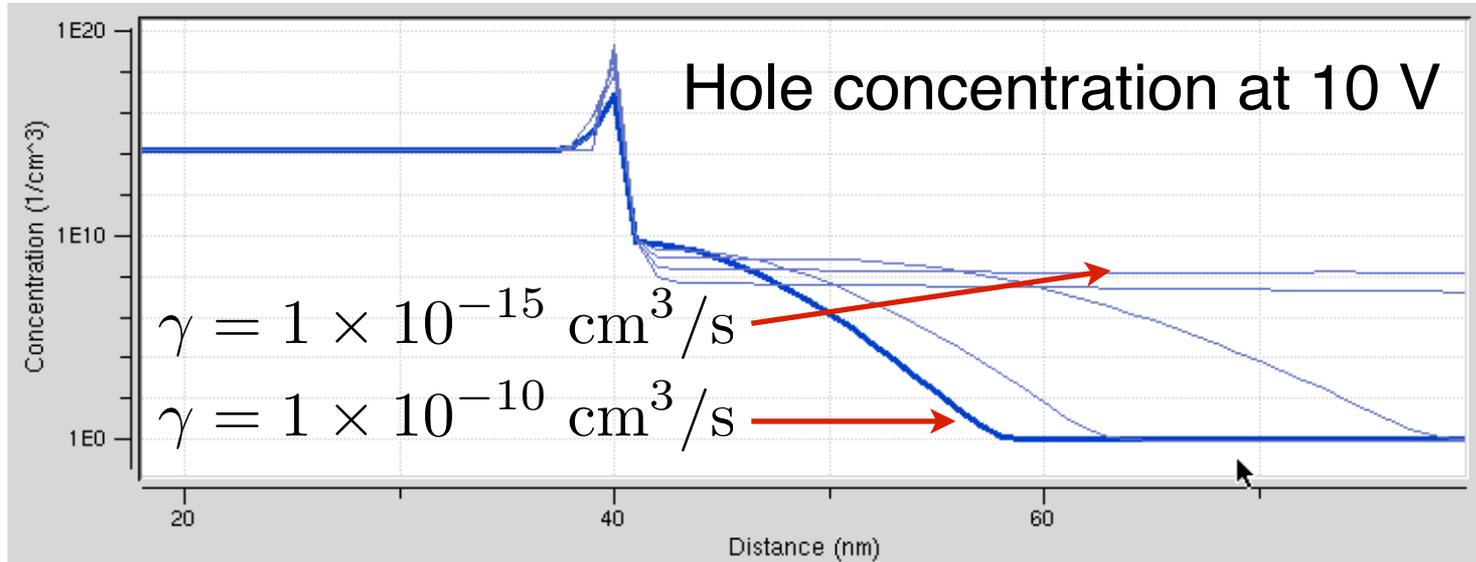
Langevin recombination: **yes**

Langevin input: (1,2)->(1e-10,1e-10)
Example: (1,2)->(1e-12,1e-14)

LANGEVIN RECOMBINATION

hole + electron \Rightarrow exciton

Result: Higher recombination \Rightarrow fewer charges at interface

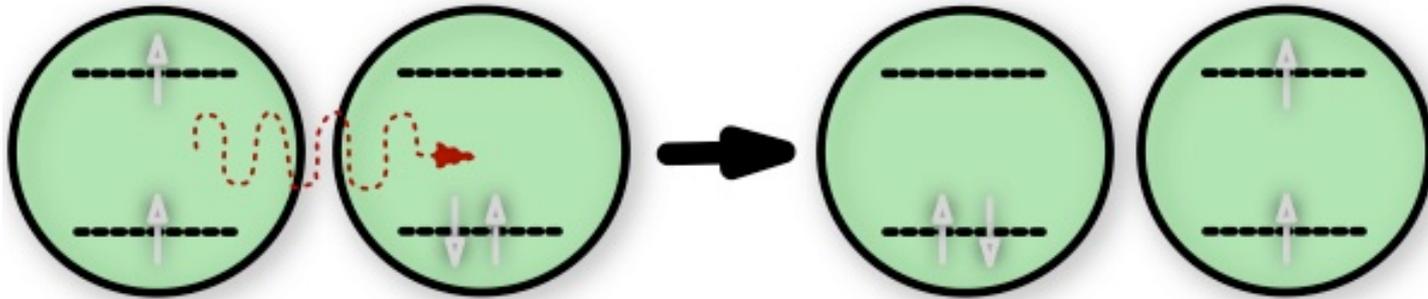


FÖRSTER RESONANT ENERGY TRANSFER

long-range exciton transport

$$K_{D \rightarrow A} = \frac{3c^4}{4\pi\epsilon_D} \frac{1}{\tau_D} \frac{1}{R^6} \int \frac{1}{\omega^4} F_D(\omega) \sigma_A(\omega) d\omega$$

Integral describes overlap between donor emission and acceptor absorbance spectra



Singlet excitons move via non-radiative dipole-dipole coupling

User specifies the effective transfer radius $R_{D \rightarrow A}$ [cm]

τ_D is the exciton natural lifetime

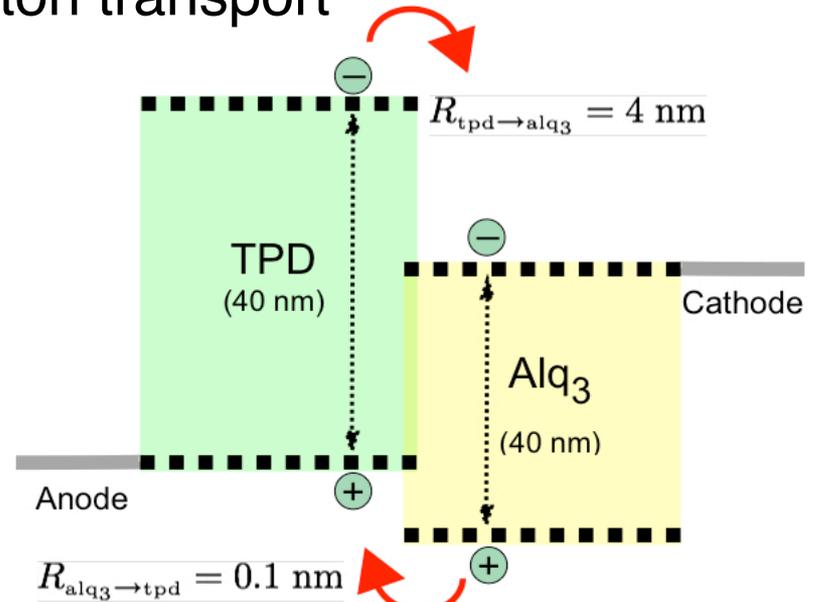
$$K_{D \rightarrow A} = \frac{1}{\tau_D} \left(\frac{R_{D \rightarrow A}}{R} \right)^6$$

FÖRSTER RESONANT ENERGY TRANSFER

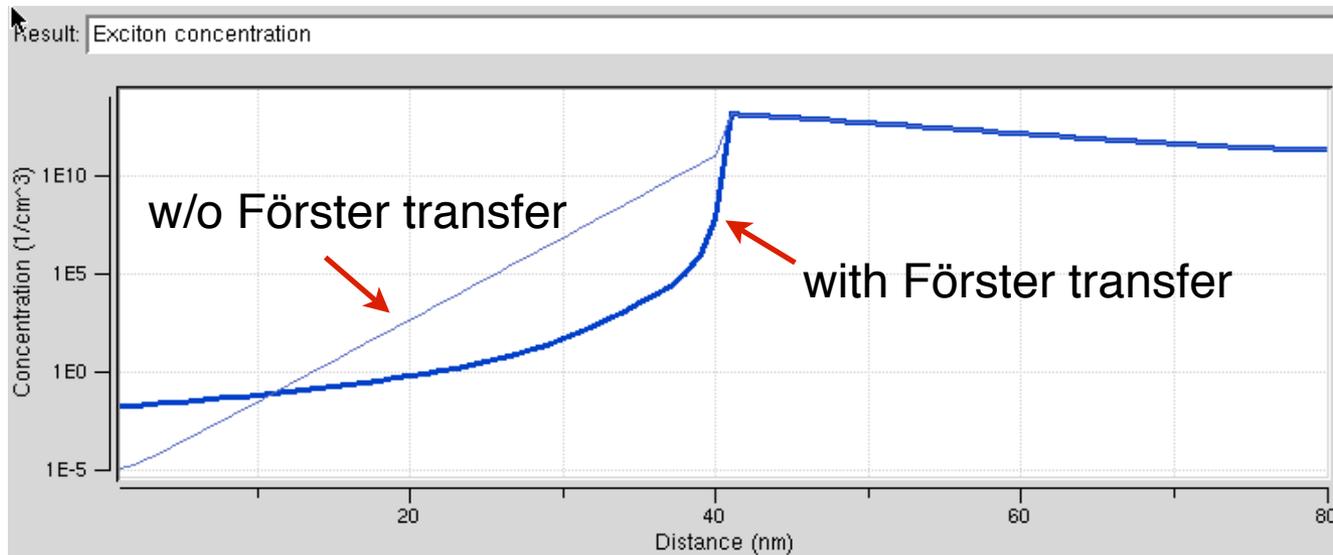
long-range exciton transport

Foerster resonant energy transfer: yes

Foerster transfer input: (1,2)->(4e-7,1e-8)
EXAMPLE: (1,3)->(1e-7,1e-8)

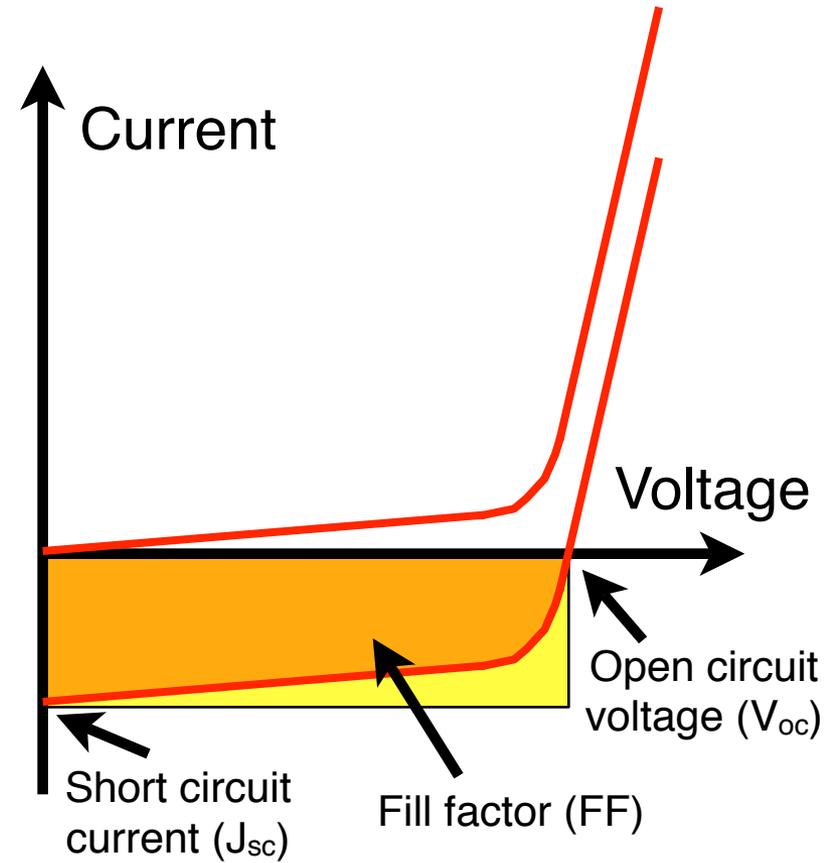
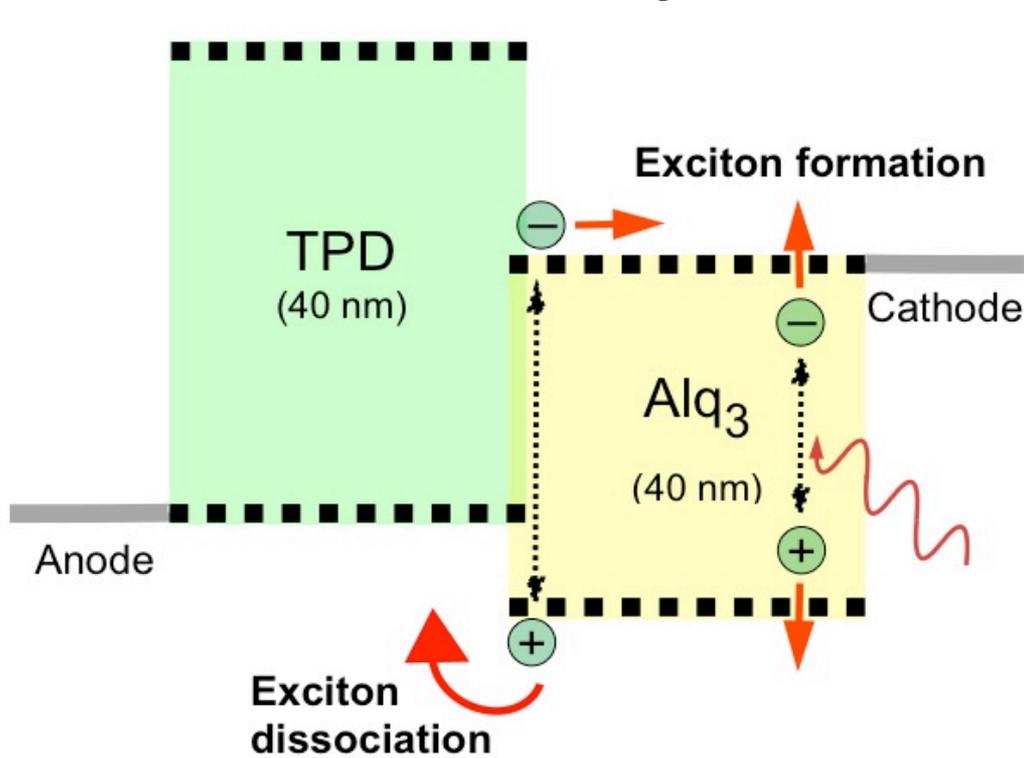


Result: fewer excitons in TPD



EXAMPLE 2: SOLAR CELL

exciton generation and dissociation



$$\text{Efficiency [\%]} = 100 \cdot \frac{J_{sc} \cdot V_{oc} \cdot FF}{I_{sun}}$$

Exciton dissociation: yes

Exciton dissociation input: (1,2)->(1e0,1e10)
 Example: (1,2)->(1e20,1e18); (2,3)->(1e5,0)

Photo generation: yes

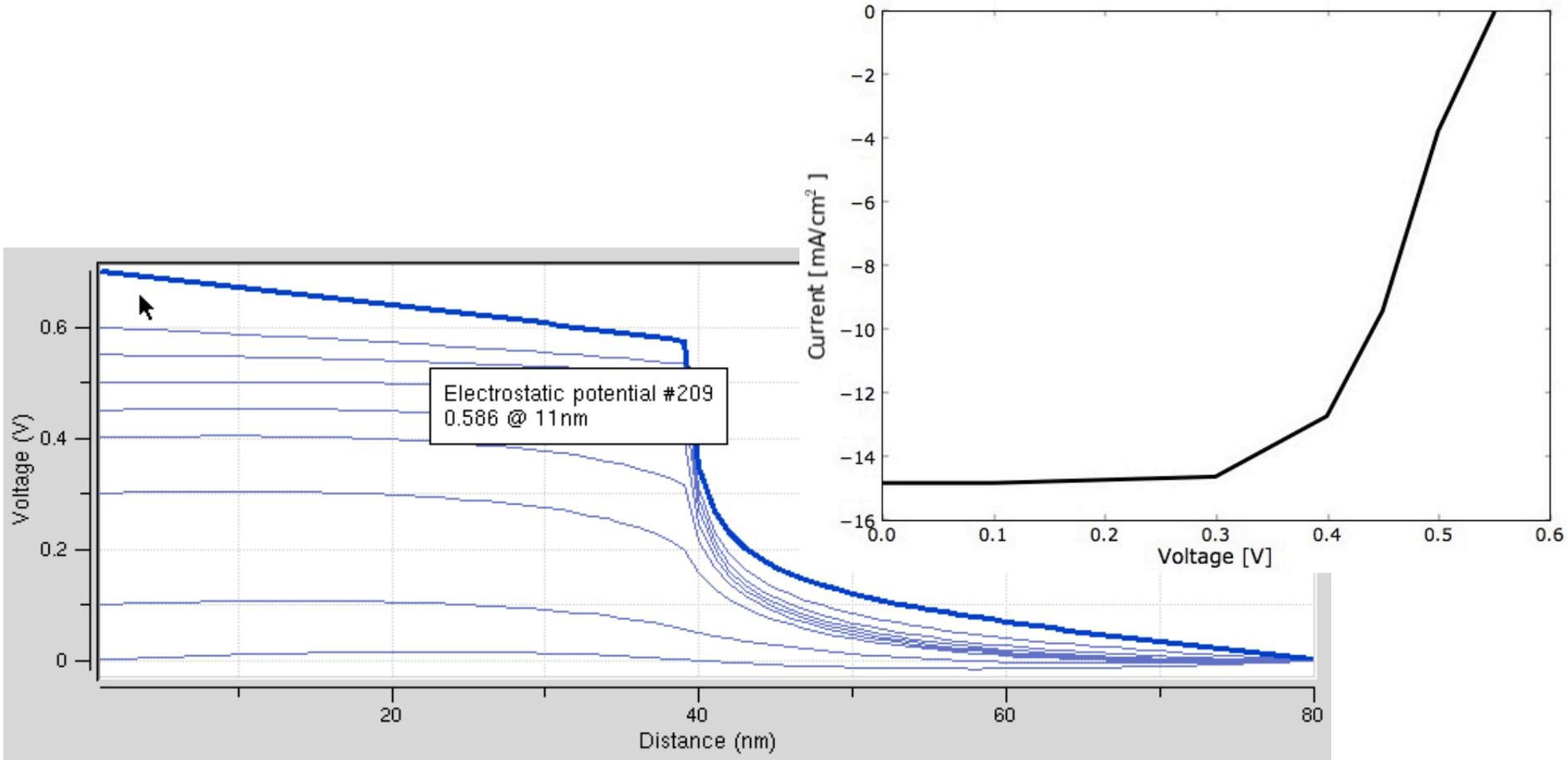
Photo generation input: 2->1e23
 EXAMPLE: 1->1e20; 2->1e24

Constant rate: exciton \Rightarrow free charge [1/s]

Constant rate: exciton creation [1/cm³-s]

SOLAR CELL

exciton generation and dissociation



Results: Reproduce current-voltage curve

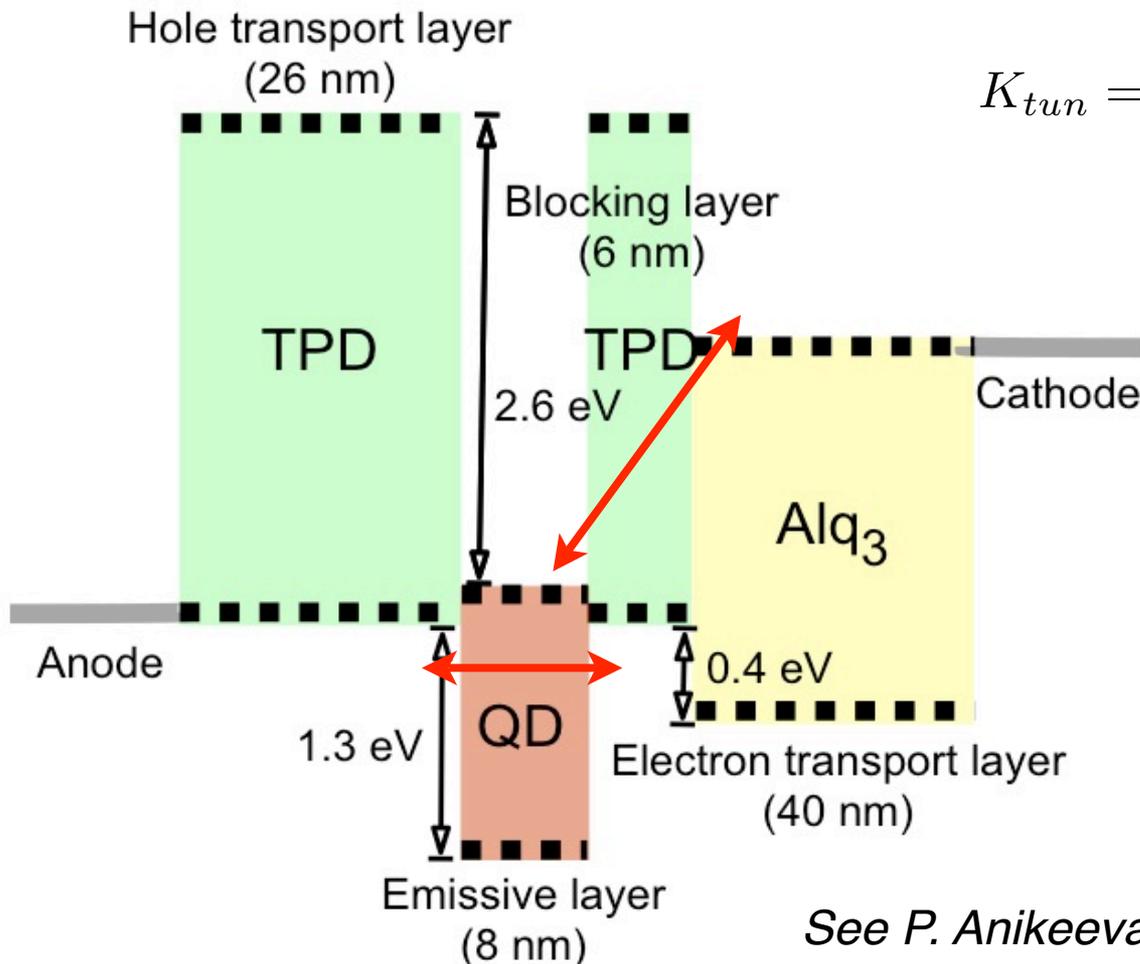
Voltage dropped near interface due to large charge buildup

EXAMPLE 3: “TUNNELING” EFFECTS:

charge movement through imperfections in blocking layers

Idea: thin “blocking” layers in devices may not be complete

Sample device: TPD/quantum dot/TPD/Alq₃ QD-LED



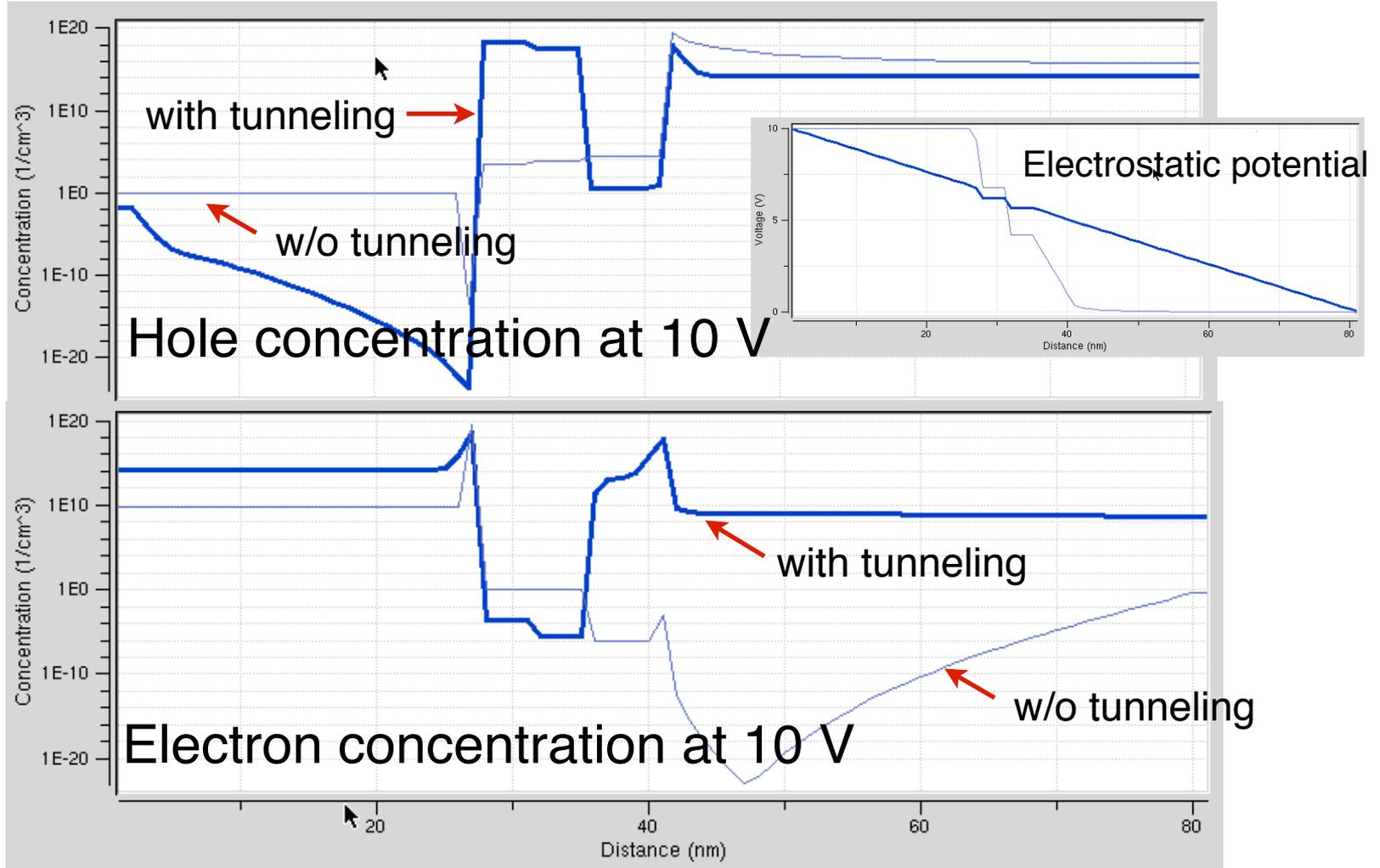
$$K_{tun} = K_{ma} \left(1 - \frac{\pi}{2\sqrt{3}} \right) \exp \left[- \frac{\Delta x}{x_0} \right]$$

- Geometrical factor
- Miller-Abrahams rate
- Exponential decay of wavefunction overlap

See P. Anikeeva's PhD thesis for more information

“TUNNELING” EFFECTS:

charge movement through imperfections in blocking layers



Results: Charges are able to flow \Rightarrow functioning device

Electric potential evenly dropped over device